

# Assignment3

Shufan Xia  
(Dated: Feb 2020)

## 1. 5.3 COMPUTATION OF ERROR FUNCTION

I used Simpson's method to evaluate the integration. To find out what the suitable number of slices  $N$  is, I evaluated the estimated error using Eq 5.24:

$$\epsilon = \frac{1}{90}h^4[f'''(a) - f'''(b)] \quad (1)$$

$$\text{where } a = 0, b = 3 \text{ and } h = \frac{b-a}{N} \quad (2)$$

See 1 for the result of fractional error as a function of number of bins  $N$  plotted in logarithmic scale. The precision limit of my computer is at  $10^{-15}$ . From fig. 1, after about  $N = 5000$ , the fractional error falls down below the  $10^{-14}$  and the trend becomes flatter as  $N$  increases. Therefore I chose 5000 bins for computing the integrals.

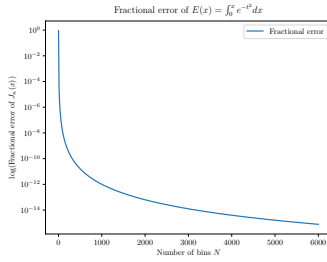


FIG. 1: fractional error as a function of number of bins  $N$

Given the definition of **Erf** function,  $E(x) = \int_0^x e^{-t^2} dx$  can be rewritten as  $\frac{\sqrt{\pi}}{2} \text{erf}(x)$ . Fig 2 shows the two methods using Simpson's integration method with 1000 bins and using **math.Erf(x)** from python library give very close results.

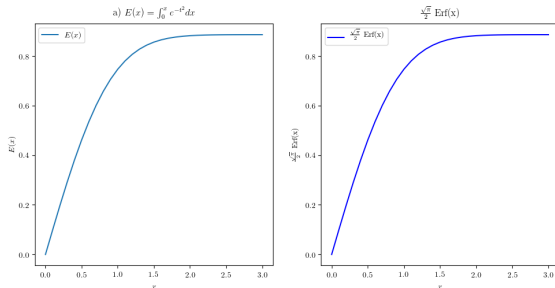


FIG. 2: Comparison of the result for  $E(x) = \int_0^x e^{-t^2} dx$  using Simpson's integration method and **math.Erf(x)**

## 2. 5.4 DIFFRACTION LIMIT OF A TELESCOPE

a) I first defined the integrand as a function of  $m$  and  $x$ ,  $f(m, x)$ . The output of the function is a function  $f(\theta) = \cos(m\theta - x \sin \theta)$  for a specific pair of  $m$  and  $x$  input values. Then I used Simpson's integration method with 1000 slices to evaluate the integration form of Bessel function:

$$J_m(m, x) = \frac{1}{\pi} \int_0^\pi \cos(m\theta - x \sin \theta) d\theta. \quad (3)$$

To plot  $J_0(x)$ ,  $J_1(x)$  and  $J_2(x)$  as a function of  $X$  from 0 to 20, I created a list to store the values of  $J_m(m, x)$  for each  $m$  value at specific  $x$  values. I used 201 equally spaced  $x$  value from 0 to 20. In another word, I have three lists of 201  $J_m$  values.

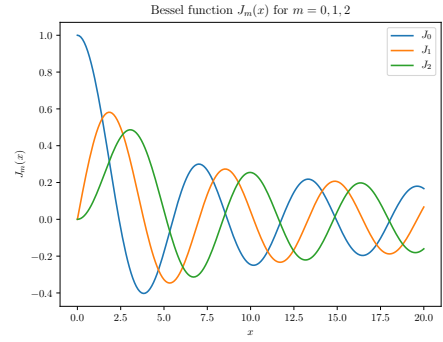


FIG. 3: Bessel function  $J_0, J_1, J_2$  as a function of  $x$  from  $x = 0$  to  $x = 20$

b) To make a 2D density plot, I generated a 2D array of  $I(r)$  at each  $(x, y)$  value. I created a 101 by 101 matrix to hold all coordinates from  $-1\mu\text{m}$  to  $1\mu\text{m}$  along both  $x$  and  $y$  axis.

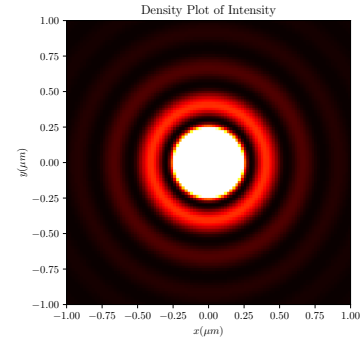


FIG. 4: Caption

c) We can rewrite the recursion definition as:

$$J_n(x) = \frac{2n}{x} J_{n-1}(x) - J_{n-2}(x). \quad (4)$$

I defined a recursion function  $J(n)$ . The base cases for this recursion are  $J_0$  and  $J_1$ . It means if the input of the function  $n$  is 0 or 1, the function returns the results evaluated by integration. For other input value of  $n$ , the functions calls  $J(n-1)$  and  $J(n-2)$ , and evaluated  $J(n)$  in the recursion definition.

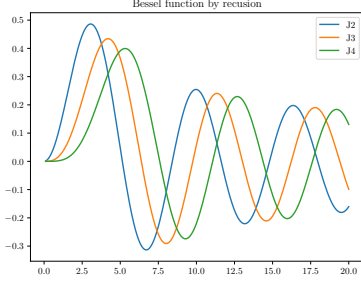


FIG. 5: Bessel function  $J_n$  evaluated using recursion definition for  $n = 2, 3, 4$ .

The error between using integration and recursion methods to find  $J_m(x)$  from  $x=0$  to  $x=20$  for  $J_2, J_3$  and  $J_4$  is shown in Fig 6. It is very noisy. Although the error is relatively large for small  $x$  value, the error drops down below the precision limit of my computer very soon.

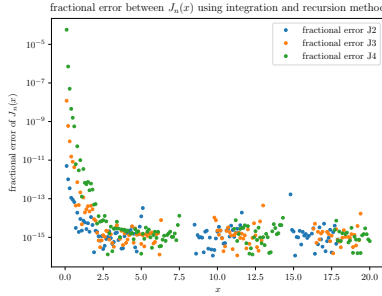


FIG. 6: Fractional error computed from  $\frac{J_{integration} - J_{recursion}}{J_{integration}}$  plotted in logarithmic scale.

### 3. 5.10 ANHARMONIC OSCILLATOR

a) Find the period  $T$ : At  $t = 0, x = a$  and  $v = \frac{dx}{dt} = 0$ . And since energy is conserved over time:

$$E = V(a) = \frac{1}{2}m\left(\frac{dx}{dt}\right)^2 + V(x) \quad (5)$$

We can solve this differential equation. Rearrange the equation first :

$$dt = \sqrt{\frac{m}{2}} \frac{dx}{\sqrt{V(a) - V(X)}} \quad (6)$$

When the particle reaches  $x = 0$ , it has gone through  $\frac{1}{4}T$ . Integrate both sides of the function from  $x = 0$  to  $a$  and from  $t = 0$  to  $\frac{1}{4}T$ :

$$\int_0^{\frac{1}{4}T} dt = \sqrt{\frac{m}{2}} \int_0^a \frac{dx}{\sqrt{V(a) - V(X)}} \quad (7)$$

$$T = \sqrt{8m} \int_0^a \frac{dx}{\sqrt{V(a) - V(X)}} \quad (8)$$

b) I used Simpson's integration with 1000 bins. See 4.

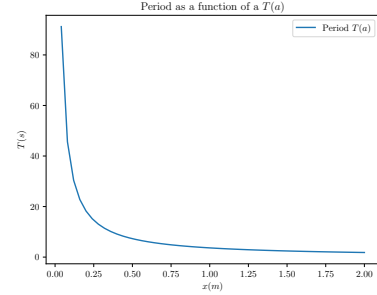


FIG. 7: The period for amplitudes ranging from  $a = 0$  to  $a = 2$  for a particle with mass = 1, in a  $V(x) = x^4$  potential well

c) The restoring force in this anharmonic oscillator is nonlinear:  $F_{restore} = -\frac{V(x)}{dx} = -4x^3$ . As a result, increasing the amplitude increases the acceleration cubically. Plugging  $V(x) = x^4$ , and evaluate Eq.8 in **Mathematica**, I found:

$$T = \frac{Constant}{a} \quad (9)$$

The period is inverse proportional to the amplitude. Therefore, as the amplitude increases, the oscillator gets faster. As  $a$  approaches 0, the potential well is approximately flat, and the oscillator hardly oscillates for small amplitude. Thus, the period approaches inf as the amplitude goes to zero.

## 4. SURVEY QUESTION

This homework took 12 hours. I learned to use my own python library, plotting 2D density plot from a 2D array and recursion. I also learned nested *for* loops can take a long time to run, but I was not able to improve the run time this time. I will try to keep this problem in consideration in the future. The most interested problem on is one on diffraction limit of a telescope because it is I found it is important to know the general "flow chart" of the program before coding, and I did some recursion as well. Moreover, it was a relatively long program, I found some tips on the process of debugging: checking variables in certain points in the program helps me narrow down the issues.