## Generic Potential

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Let us work on the case for a generic potential  $V(\phi)$  for scalars  $\phi$ . For the simplicity, start with the single scalar action:

$$S = \int d^d x \{ \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + V(\phi) \} \tag{1}$$

Then, the quadratic term is:

$$S^{(2)} = \frac{1}{2} \int d^d x \eta(-\Box + V''(\phi))\eta \tag{2}$$

So one may obtain:

$$\frac{\delta^2 S}{\delta \phi \delta \phi} = -\Box + V''(\phi) := D \tag{3}$$

Therefore:

$$W = \frac{1}{2} \operatorname{Tr} \ln D$$

$$= -\frac{1}{2} \int_0^\infty \frac{dt}{t} \operatorname{Tr} e^{-tD}$$
(4)

But, now one can use the expansion for the trace of the heat kernel:

$$Tre^{-tD} = \sum_{k>0} t^{(k-d)/2} a_k \tag{5}$$

where  $a_k$  are heat kernel coefficients which are, for instance, given in Vassilevich. Thus, introducing a cut-off at  $t = \Lambda^{-2}$ :

$$\mathcal{W} = -\frac{1}{2} \int_{\Lambda^{-2}}^{\infty} dt \sum_{k=0}^{\infty} t^{\frac{k-d}{2}-1} a_k 
= -\frac{1}{2} \int_{\Lambda^{-2}}^{\infty} dt \{ t^{-d/2-1} a_0 + t^{-d/2} a_2 + t^{-d/2+1} a_4 \cdots \}$$
(6)

Notice that only finitely many terms involve the divergent one would need to care, i.e. only the coefficients with  $k \le d$  are necessary. Therefore, defining  $\tau \equiv \ln \frac{\Lambda}{\Lambda_0}$ :

$$\partial_{\tau} V_{\Lambda} = -\frac{1}{2} \partial_{\tau} \left\{ \int_{\Lambda^{-2}}^{\infty} dt \sum_{k < d} t^{\frac{k-d}{2} - 1} a_k \right\} \tag{7}$$

From this "beta functional", each beta-function corresponding to each coupling can be extracted. For a theory with multiple scalars  $\phi^i$ , this could be generalized. In this case, the second functional derivative is in the form:

$$\frac{\delta S}{\delta \phi^i \delta \phi^j} = -\Box \delta^{ij} + U(\phi) \tag{8}$$

where  $U(\phi)$  includes the additional terms in its potential which involve  $i \neq j$ .

Thus, specifying terms involving the divergence with heat kernel coefficients with same procedure, eventually one may obtain a "beta-functional" depending on scalars from which all beta-functions of all couplings are extracted. Ultimately, once the form of the potential of the theory is determined, plugging them into the heat kernel coefficients, all beta-functions can be specified.