## Notes on Feller Condition

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## 1 Fokker-Planck Equation

Consider the following parabolic partial differential equation (PDE) [1]

$$\frac{\partial u(t,x)}{\partial t} = -\frac{\partial}{\partial x} \left( (bx+c)u(t,x) \right) + \frac{\partial^2}{\partial x^2} \left( axu(t,x) \right), \tag{1}$$

with x > 0. This can be viewed as the corresponding Fokker-Planck equation for the Cox-Ingersoll-Ross (CIR) process

$$dx_t = \kappa(\theta - x_t)dt + \sigma\sqrt{x_t}dW_t, \tag{2}$$

with  $a = \sigma^2/2$ ,  $b = -\kappa$ , and  $c = \kappa \theta$ .

We want to find the solution of the PDE (1), with certain initial condition. In particular, we are interested in the fundamental solution of the PDE (1), *i.e.*, the initial condition is given by

$$u(0,x) = \delta(x-y),\tag{3}$$

where  $\delta(x)$  is the Dirac delta function. To this end, introduce the Laplace transform of u(t,x) as

$$v(t,\lambda) = \int_0^{+\infty} e^{-\lambda x} u(t,x) dx,$$
 (4)

for  $\lambda > 0$ . We denote the Laplace transform of the initial condition as

$$\pi(\lambda) = \int_0^{+\infty} e^{-\lambda x} u(0, x) dx,\tag{5}$$

and for the fundamental solution,  $\pi(\lambda) = e^{-\lambda y}$ . Consider the Laplace transform of the right hand side of Eq. (1),

$$\int_{0}^{+\infty} e^{-\lambda x} \left[ (axu(t,x))_{xx} - ((bx+c)u(t,x))_{x} \right] dx$$

$$= e^{-\lambda x} \left[ (axu(t,x))_{x} - ((bx+c)u(t,x)) \right] \Big|_{0}^{+\infty}$$

$$+\lambda \int_{0}^{+\infty} e^{-\lambda x} \left[ (axu(t,x))_{x} - ((bx+c)u(t,x)) \right] dx$$

$$= f(t) + \lambda (b - \lambda a) v_{\lambda} - c\lambda v, \tag{6}$$

where

$$f(t) = \lim_{x \to 0} \left[ (bx + c)u - (axu)_x \right],\tag{7}$$

is the flux at x = 0 and cannot be arbitrarily specified.

After the Laplace transform, Eq. (1) becomes

$$v_t + \lambda(\lambda a - b)v_\lambda = f(t) - c\lambda v. \tag{8}$$

This first order PDE can be solved by the method of characteristics. The characteristics can be determined by integrating

$$dt = \frac{d\lambda}{\lambda(a\lambda - b)},\tag{9}$$

which leads to

$$e^{-bt}\frac{a\lambda - b}{\lambda} = C_1,\tag{10}$$

with  $C_1 = a - b/\lambda$ . Then, Eq. (8) can be integrated.

$$v(t,\lambda) = \left| C_1 - ae^{-bt} \right|^{c/a} \left\{ C_2 + \int_0^t \frac{f(\tau)}{|C_1 - ae^{-bt}|^{c/a}} \right\}.$$
 (11)

Together with the initial condition, the general solution becomes

$$v(t,\lambda) = \left(\frac{b}{\lambda a(e^{bt}-1)+b}\right)^{c/a} \pi \left(\frac{\lambda b e^{bt}}{\lambda a(e^{bt}-1)+b}\right) + \int_0^t f(\tau) \left(\frac{b}{\lambda a(e^{b(t-\tau)}-1)+b}\right)^{c/a} d\tau. \tag{12}$$

## References

[1] W. Feller, Two Singular Diffusion Problems, Annals of Mathematics 54, 173 (1951).