

Notes on CEV Model

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1 Fokker-Planck Equation

Assume the forward process follows the CEV process,

$$dS_t = \sigma S_t^\beta dW_t, \quad (1)$$

with $\beta < 1$, the corresponding Fokker-Planck equation is given by

$$-\frac{\partial p(0, s_0; t, s)}{\partial t} + \frac{1}{2}\sigma^2 \frac{\partial^2}{\partial s^2} (s^{2\beta} p(0, s_0; t, s)) = 0, \quad (2)$$

with the initial condition

$$p(0, s_0; 0, s) = \delta(s - s_0), \quad (3)$$

where $\delta(x)$ is the Dirac delta function. Since the CEV process is time homogeneous, we do not need to make the initial time explicit, and we will write $p(t, s) \equiv p(0, s_0; t, s)$.

Define the Laplace transform of the forward transition density $p(t, s)$ as [1]

$$p(t, s) = \int_0^{+\infty} e^{-\lambda t} g(\lambda, s) d\lambda, \quad (4)$$

the Fokker-Planck equation (2) becomes

$$\frac{\partial^2 g(\lambda, s)}{\partial s^2} + \frac{4\beta}{s} \frac{\partial g(\lambda, s)}{\partial s} + \left(\frac{2\lambda}{\sigma^2 s^{2\beta}} + \frac{2\beta(2\beta-1)}{s^2} \right) g(\lambda, s) = 0. \quad (5)$$

Following [2], if $u(x)$ satisfies the Bessel equation of order ν ,

$$u'' + \frac{1}{x}u' + \left(1 - \frac{\nu^2}{x^2}\right)u = 0, \quad (6)$$

and let $y(x) = x^a u(bx^c)$, then $y(x)$ satisfies

$$y'' + \frac{1-2a}{x}y' + \left(b^2c^2x^{2c-2} + \frac{a^2 - \nu^2c^2}{x^2}\right)y = 0. \quad (7)$$

Comparing the coefficients with Eq. (5), we have

$$a = \frac{1}{2}(1 - 4\beta), \quad b = \frac{2\nu\sqrt{2\lambda}}{\sigma} = \frac{\sqrt{2\lambda}}{\sigma(1 - \beta)}, \quad c = \frac{1}{2\nu} = 1 - \beta, \quad \nu = \frac{1}{2(1 - \beta)}. \quad (8)$$

Then, the general solution to Eq. (5) is given by

$$g(\lambda, s) = s^{(1-4\beta)/2} \left[A_\nu(\lambda) J_\nu \left(\frac{2\nu\sqrt{2\lambda}}{\sigma} s^{\frac{1}{2\nu}} \right) + B_\nu(\lambda) Y_\nu \left(\frac{2\nu\sqrt{2\lambda}}{\sigma} s^{\frac{1}{2\nu}} \right) \right], \quad (9)$$

where $J_\nu(x)$ and $Y_\nu(x)$ are the Bessel functions of the first and second kind, respectively. To determine the coefficients $A_\nu(\lambda)$ and $B_\nu(\lambda)$, and subsequently the forward transition density $p(t, s)$, boundary condition at $s = 0$ is needed. To this end, we can write down the asymptotic expansion of Eq. (9),

$$\begin{aligned} s^{2\beta} g(\lambda, s) \sim & \left((A_\nu(\lambda) + B_\nu(\lambda) \cot(\nu\pi)) \left(\frac{\nu\sqrt{2\lambda}}{\sigma} \right)^\nu s \right) \left[\frac{1}{\Gamma(1+\nu)} - \frac{1}{\Gamma(2+\nu)} \left(\frac{\nu\sqrt{2\lambda}}{\sigma} \right)^2 s^{\frac{1}{\nu}} \right] \\ & - \frac{B_\nu(\lambda)}{\sin(\nu\pi)} \left(\frac{\nu\sqrt{2\lambda}}{\sigma} \right)^{-\nu} \left[\frac{1}{\Gamma(1-\nu)} - \frac{1}{\Gamma(2-\nu)} \left(\frac{\nu\sqrt{2\lambda}}{\sigma} \right)^2 s^{\frac{1}{\nu}} \right]. \end{aligned} \quad (10)$$

Here, we have used the series representation of the Bessel function of the first and second kind

$$J_\nu(x) = \sum_{m=0}^{+\infty} \frac{(-1)^m}{m! \Gamma(m + \nu + 1)} \left(\frac{x}{2} \right)^{2m+\nu}, \quad (11)$$

and

$$Y_\nu(x) = \frac{J_\nu(x) \cos(\nu\pi) - J_{-\nu}(x)}{\sin(\nu\pi)}. \quad (12)$$

The asymptotics of $p(t, s)$ for small s have the form of $p_A(t, s) \sim s^{1-2\beta}$ and $p_R(t, s) \sim s^{-2\beta}$. The latter exists only $\beta < 1/2$. Otherwise, the norm will diverge around 0. Consider the Fokker-Planck equation (2), it can be expressed in the following current conservation form,

$$\frac{\partial p(t, s)}{\partial t} + \frac{\partial j(t, s)}{\partial s} = 0, \quad (13)$$

with the current given by

$$j(t, s) = -\frac{1}{2}\sigma^2 \frac{\partial}{\partial s} (s^{2\beta} p(t, s)). \quad (14)$$

With the asymptotics of $p(t, s)$, we can see that the current is a constant for $p_A(t, s)$, and vanishes for $p_R(t, s)$. These correspond to an absorbing boundary condition for $p_A(t, s)$, and a reflecting boundary condition for $p_R(t, s)$.

Now, consider an arbitrary function $f(s)$. From the Fokker-Planck equation (2), we have

$$\frac{\partial}{\partial t} \left(\int f(s) p(t, s) ds \right) = \frac{1}{2}\sigma^2 \int f(s) \frac{\partial^2}{\partial s^2} (s^{2\beta} p(t, s)) ds. \quad (15)$$

Integrating the right hand side by parts twice, we have

$$\begin{aligned} \frac{\partial}{\partial t} \left(\int f(s) p(t, s) ds \right) &= \frac{1}{2} \sigma^2 \int f''(s) s^{2\beta} p(t, s) ds \\ &+ \frac{1}{2} \sigma^2 \left[f'(s) s^{2\beta} p(t, s) - f(s) \frac{\partial}{\partial s} (s^{2\beta} p(t, s)) \right] \Big|_{s=0}. \end{aligned} \quad (16)$$

Let $f(s) \equiv 1$, it can be seen that the right hand side of Eq. (16) does not vanish for absorbing boundary condition, and is zero for reflecting boundary condition. Therefore, the total probability is not conserved for the absorbing boundary condition, and is conserved for the reflecting boundary condition.

Now consider $f(s) \equiv s$. The right hand side of Eq. (16) will vanish for absorbing boundary condition, implying that S_t is a martingale under the absorbing boundary condition, but not for reflecting boundary condition.

Finally, let us consider the call payoff, $f(s) = (s - K)^+$. Then, Eq. (16) becomes

$$\frac{\partial}{\partial t} \left(\int f(s) p(t, s) ds \right) = \frac{1}{2} \sigma^2 K^{2\beta} p(t, K), \quad (17)$$

irrespective of the boundary condition, which gives the call option price as

$$C(t, S_t, T, K) = \int (s - K)^+ p(T, s) ds = (S_t - K)^+ + \frac{1}{2} \sigma^2 \int_t^T K^{2\beta} p(u, K) du. \quad (18)$$

However, for the put payoff, there will be an additional boundary term for the reflecting boundary condition.

2 Absorbing boundary condition

The absorbing boundary condition implies that $B_\nu(\lambda) = 0$, which leads to

$$g(\lambda, s) = s^{(1-4\beta)/2} A_\nu(\lambda) J_\nu \left(\frac{2\nu\sqrt{2\lambda}}{\sigma} s^{\frac{1}{2\nu}} \right). \quad (19)$$

To determine the coefficient $A_\nu(\lambda)$, initial condition will be needed. From Eq. (4), the transition density is given by

$$p_A(t, s) = \int_0^{+\infty} e^{-\lambda t} g(\lambda, s) d\lambda = s^{(1-4\beta)/2} \int_0^{+\infty} e^{-\lambda t} A_\nu(\lambda) J_\nu \left(\frac{2\nu\sqrt{2\lambda}}{\sigma} s^{\frac{1}{2\nu}} \right) d\lambda. \quad (20)$$

Take into account of the initial condition (3), we have

$$\delta(s - s_0) = s^{(1-4\beta)/2} \int_0^{+\infty} A_\nu(\lambda) J_\nu \left(\frac{2\nu\sqrt{2\lambda}}{\sigma} s^{\frac{1}{2\nu}} \right) d\lambda. \quad (21)$$

Now, multiply both sides by

$$\sqrt{s}J_\nu\left(\frac{2\nu\sqrt{2\lambda}}{\sigma}s^{\frac{1}{2\nu}}\right),$$

and integrate over s , we will have

$$\begin{aligned}\sqrt{s_0}J_\nu\left(\frac{2\nu\sqrt{2\lambda}}{\sigma}s_0^{\frac{1}{2\nu}}\right) &= \int_0^{+\infty} d\lambda' A_\nu(\lambda') \int_0^{+\infty} ds s^{1-2\beta} J_\nu\left(\frac{2\nu\sqrt{2\lambda'}}{\sigma}s^{\frac{1}{2\nu}}\right) J_\nu\left(\frac{2\nu\sqrt{2\lambda}}{\sigma}s^{\frac{1}{2\nu}}\right) \\ &= 2\nu \int_0^{+\infty} d\lambda' A_\nu(\lambda') \int_0^{+\infty} du u J_\nu\left(\frac{2\nu\sqrt{2\lambda'}}{\sigma}u\right) J_\nu\left(\frac{2\nu\sqrt{2\lambda}}{\sigma}u\right) \\ &= 2\nu \int_0^{+\infty} d\lambda' A_\nu(\lambda') \frac{\delta\left(\frac{2\nu\sqrt{2\lambda}}{\sigma} - \frac{2\nu\sqrt{2\lambda'}}{\sigma}\right)}{\frac{2\nu\sqrt{2\lambda}}{\sigma}} \\ &= \frac{\sigma^2}{2\nu} \int_0^{+\infty} d\lambda' A_\nu(\lambda') \delta(\lambda - \lambda') \\ &= \frac{\sigma^2}{2\nu} A_\nu(\lambda).\end{aligned}\tag{22}$$

Here, we have used the Hankel transform

$$\int_0^{+\infty} r J_\nu(kr) J_\nu(k'r) = \frac{\delta(k - k')}{k},\tag{23}$$

and the property of delta function,

$$\delta(f(x)) = \frac{\delta(x)}{|f'(0)|}.\tag{24}$$

Now, we have

$$A_\nu(\lambda) = \frac{2\nu}{\sigma^2} \sqrt{s_0} J_\nu\left(\frac{2\nu\sqrt{2\lambda}}{\sigma}s_0^{\frac{1}{2\nu}}\right),$$

the transition density can be calculated,

$$\begin{aligned}p_A(t, s) &= \frac{2\nu}{\sigma^2} (s_0 s^{1-4\beta})^{1/2} \int_0^{+\infty} e^{-\lambda t} J_\nu\left(\frac{2\nu\sqrt{2\lambda}}{\sigma}s^{\frac{1}{2\nu}}\right) J_\nu\left(\frac{2\nu\sqrt{2\lambda}}{\sigma}s_0^{\frac{1}{2\nu}}\right) d\lambda \\ &= \frac{4\nu}{\sigma^2} (s_0 s^{1-4\beta})^{1/2} \int_0^{+\infty} e^{-tu^2} u J_\nu\left(\frac{\sqrt{8\nu}}{\sigma}s^{\frac{1}{2\nu}}u\right) J_\nu\left(\frac{\sqrt{8\nu}}{\sigma}s_0^{\frac{1}{2\nu}}u\right) du \\ &= \frac{2\nu}{\sigma^2 t} (s_0 s^{1-4\beta})^{1/2} \exp\left(-\frac{2\nu^2}{\sigma^2 t} \left(s_0^{\frac{1}{2\nu}} + s^{\frac{1}{2\nu}}\right)\right) I_\nu\left(\frac{4\nu^2}{\sigma^2 t} (s_0 s)^{\frac{1}{2\nu}}\right),\end{aligned}\tag{25}$$

where we have used the result [3]

$$\int_0^{+\infty} e^{-tu^2} u J_\nu(au) J_\nu(bu) du = \frac{1}{2t} \exp\left(-\frac{a^2 + b^2}{4t}\right) I_\nu\left(\frac{ab}{2t}\right), \quad (26)$$

and $I_\nu(x)$ is the modified Bessel function of the first kind.

Finally, the transition density is given by

$$p_A(t, s) = \frac{(s_0 s^{1-4\beta})^{1/2}}{(1-\beta)\sigma^2 t} \exp\left(-\frac{s_0^{2(1-\beta)} + s^{2(1-\beta)}}{2(1-\beta)^2 \sigma^2 t}\right) I_\nu\left(\frac{(s_0 s)^{1-\beta}}{(1-\beta)^2 \sigma^2 t}\right). \quad (27)$$

3 Reflecting boundary condition

Consider the reflecting boundary condition. It can be shown, to impose this boundary condition, we must have

$$A_\nu(\lambda) + B_\nu(\lambda) \cot(\nu\pi) = 0, \quad (28)$$

and

$$\frac{1}{\nu} - 1 \geq 0, \quad (29)$$

which is equivalent to

$$\beta \leq \frac{1}{2}. \quad (30)$$

Now, we have

$$g(\lambda, x) = -s^{(1-4\beta)/2} \frac{B_\nu(\lambda)}{\sin(\nu\pi)} J_{-\nu}\left(\frac{2\nu\sqrt{2\lambda}}{\sigma} s^{\frac{1}{2\nu}}\right). \quad (31)$$

Proceed in the same way as the absorbing boundary condition, the transition density is then given by

$$p_R(t, s) = \frac{(s_0 s^{1-4\beta})^{1/2}}{(1-\beta)\sigma^2 t} \exp\left(-\frac{s_0^{2(1-\beta)} + s^{2(1-\beta)}}{2(1-\beta)^2 \sigma^2 t}\right) I_{-\nu}\left(\frac{(s_0 s)^{1-\beta}}{(1-\beta)^2 \sigma^2 t}\right). \quad (32)$$

4 Integral representation of European call option price

To preclude arbitrage opportunity, we are going to consider the European call option price under the absorbing boundary condition. The price for a European call option with expiry T , forward S , and strike K is given by

$$C_A(T, S, K) = (S - K)^+ + \frac{\sigma^2}{2} K^{2\beta} \int_0^T p_A(0, S; \tau, K) d\tau, \quad (33)$$

with the density $p(0, S; t, K)$ given by Eq. (25). To this end, we can use the integral representation of the modified Bessel function [4],

$$I_\nu(z) = \frac{1}{\pi} \int_0^\pi e^{z \cos \phi} \cos(\nu \phi) d\phi - \frac{\sin(\nu\pi)}{\pi} \int_0^{+\infty} e^{-z \cosh \psi - \nu \psi} d\psi \quad (34)$$

Define

$$q_S = \frac{S^{1-\beta}}{1-\beta}, \quad q_K = \frac{K^{1-\beta}}{1-\beta}, \quad (35)$$

the integration term of Eq. (33) becomes

$$\begin{aligned} & \nu\sqrt{SK} \int_0^T \frac{d\tau}{\tau} \exp\left(-\frac{q_K^2 + q_S^2}{2\sigma^2\tau}\right) \left(\frac{1}{\pi} \int_0^\pi \exp\left(\frac{q_K q_S \cos \phi}{\sigma^2\tau}\right) \cos(\nu\phi) d\phi \right. \\ & \quad \left. - \frac{\sin(\nu\pi)}{\pi} \int_0^{+\infty} \exp\left(-\frac{q_K q_S \cosh \psi}{\sigma^2\tau}\right) e^{-\nu\psi} d\psi \right) \\ = & \nu\sqrt{SK} \int_0^{\sigma^2 T} \frac{d\tau}{\tau} \exp\left(-\frac{q_K^2 + q_S^2}{2\tau}\right) \left(\frac{1}{\pi} \int_0^\pi \exp\left(\frac{q_K q_S \cos \phi}{\tau}\right) \cos(\nu\phi) d\phi \right. \\ & \quad \left. - \frac{\sin(\nu\pi)}{\pi} \int_0^{+\infty} \exp\left(-\frac{q_K q_S \cosh \psi}{\tau}\right) e^{-\nu\psi} d\psi \right) \end{aligned} \quad (36)$$

Now, consider the first integration of the above expression. Change the variable $y = q_K q_S / \tau$, define $b = (q_K^2 + q_S^2) / 2q_K q_S$, and integrate by parts, we have

$$\begin{aligned} & \frac{\nu\sqrt{SK}}{\pi} \int_0^{\sigma^2 T} \frac{d\tau}{\tau} \exp\left(-\frac{q_K^2 + q_S^2}{2\tau}\right) \int_0^\pi \exp\left(\frac{q_K q_S \cos \phi}{\tau}\right) \cos(\nu\phi) d\phi \\ = & \frac{\sqrt{SK}}{\pi} \int_{\frac{q_K q_S}{\sigma^2 T}}^{+\infty} \frac{dy}{y} \int_0^\pi e^{-(b - \cos \phi)y} d \sin(\nu\phi) \\ = & \frac{\sqrt{SK}}{\pi} \int_{\frac{q_K q_S}{\sigma^2 T}}^{+\infty} dy \left[\frac{\sin(\nu\phi)}{y} e^{-(b - \cos \phi)y} \Big|_0^\pi + \int_0^\pi \sin \phi \sin(\nu\phi) e^{-(b - \cos \phi)y} d\phi \right] \\ = & \frac{\sqrt{SK}}{\pi} \int_{\frac{q_K q_S}{\sigma^2 T}}^{+\infty} \frac{\sin(\nu\pi)}{y} e^{-(b+1)y} dy \\ & + \frac{\sqrt{SK}}{\pi} \int_0^\pi \frac{\sin \phi \sin(\nu\phi)}{b - \cos \phi} \exp\left(-\frac{q_K q_S}{\sigma^2 T} (b - \cos \phi)\right) d\phi. \end{aligned} \quad (37)$$

Similarly, the second integration can be manipulated to give

$$\begin{aligned} & -\frac{\sqrt{SK}}{\pi} \int_{\frac{q_K q_S}{\sigma^2 T}}^{+\infty} \frac{\sin(\nu\pi)}{y} e^{-(b+1)y} dy \\ & + \frac{\sqrt{SK} \sin(\nu\pi)}{\pi} \int_0^{+\infty} \frac{e^{-\nu\psi} \sinh \psi}{b + \cosh \psi} \exp\left(-\frac{q_K q_S}{\sigma^2 T} (b + \cosh \psi)\right) d\psi. \end{aligned} \quad (38)$$

Finally, the call option price is

$$\begin{aligned} C_A(T, S, K) = & (S - K)^+ \\ & + \frac{\sqrt{SK}}{\pi} \left(\int_0^\pi \frac{\sin \phi \sin(\nu\phi)}{b - \cos \phi} \exp\left(-\frac{q_K q_S}{\sigma^2 T} (b - \cos \phi)\right) d\phi \right. \\ & \quad \left. + \sin(\nu\pi) \int_0^{+\infty} \frac{e^{-\nu\psi} \sinh \psi}{b + \cosh \psi} \exp\left(-\frac{q_K q_S}{\sigma^2 T} (b + \cosh \psi)\right) d\psi \right). \end{aligned} \quad (39)$$

For completeness, we also record the result for reflecting boundary condition as

$$\begin{aligned}
C_R(T, S, K) &= (S - K)^+ \\
&+ \frac{\sqrt{SK}}{\pi} \left(\int_0^\pi \frac{\sin \phi \sin(\nu \phi)}{b - \cos \phi} \exp \left(-\frac{q_K q_S}{\sigma^2 T} (b - \cos \phi) \right) d\phi \right. \\
&\quad \left. + \sin(\nu \pi) \int_0^{+\infty} \frac{e^{\nu \psi} \sinh \psi}{b + \cosh \psi} \exp \left(-\frac{q_K q_S}{\sigma^2 T} (b + \cosh \psi) \right) d\psi \right). \quad (40)
\end{aligned}$$

5 SABR

The SABR model is characterized by the following stochastic process,

$$\begin{aligned}
dS_t &= \sigma_t S_t^\beta dW_t, \\
d\sigma_t &= \gamma \sigma_t dZ_t, \\
E[dW_t dZ_t] &= \rho dt,
\end{aligned} \quad (41)$$

where the volatility of the forward process is now stochastic, γ is the volatility of volatility (volvol), and ρ is the correlation. The initial volatility of σ_t is assumed to be σ_0 . We also assume that the absorbing boundary condition is imposed for the forward process to preserve the martingale property. In the following, we will first consider the zero correlation model, then map the non-zero correlation model to the zero correlation one.¹

5.1 Zero correlation

As in the last section, define

$$q_S = \frac{S^{1-\beta}}{1-\beta}, \quad q_K = \frac{K^{1-\beta}}{1-\beta}, \quad \bar{q} = q_S q_K, \quad b = \frac{q_S^2 + q_K^2}{2q_S q_K}, \quad (42)$$

the call option price is given by

$$\begin{aligned}
C(T, S, K) &= (S - K)^+ \\
&+ \frac{\sqrt{SK}}{\pi} \left(\int_0^\pi \frac{\sin \phi \sin(\nu \phi)}{b - \cos \phi} \frac{G(\gamma^2 T, s(\phi))}{\cosh(s(\phi))} d\phi \right. \\
&\quad \left. + \sin(\nu \pi) \int_0^{+\infty} \frac{e^{-\nu \psi} \sinh \psi}{b + \cosh \psi} \frac{G(\gamma^2 T, s(\psi))}{\cosh(s(\psi))} d\psi \right), \quad (43)
\end{aligned}$$

where

$$\begin{aligned}
\sinh(s(\phi)) &= \frac{\gamma}{\sigma_0} \sqrt{2\bar{q}(b - \cos \phi)}, \\
\sinh(s(\psi)) &= \frac{\gamma}{\sigma_0} \sqrt{2\bar{q}(b + \cosh \psi)}.
\end{aligned} \quad (44)$$

¹Simply record the known result here.

Here,

$$G(t, s) = \frac{2e^{-t/8}}{t\sqrt{\pi t}} \int_s^{+\infty} u e^{-u^2/2t} \sqrt{\cosh u - \cosh s} du, \quad (45)$$

is an integration kernel, which has a highly accurate series expansion approximation,

$$G(t, s) = \sqrt{\frac{\sinh s}{s}} e^{-\frac{s^2}{2t} - \frac{t}{8}} (R(t, s) + \delta R(t, s)), \quad (46)$$

where

$$R(t, s) = 1 + \frac{3tg(s)}{8s^2} - \frac{5t^2(-8s^2 + 3g^2(s) + 24g(s))}{128s^4} + \frac{35t^3(-40s^2 + 3g^3(s) + 24g^2(s) + 120g(s))}{1024s^6}, \quad (47)$$

with $g(s) = s \coth(s) - 1$, and

$$\delta R(t, s) = e^{\frac{t}{8}} - \frac{3072 + 384t + 24t^2 + t^3}{3072}. \quad (48)$$

In the limit of $s \rightarrow 0$, $G(t, s)$ has the following MacLaurin expansion,

$$G(t, s) = e^{-\frac{s^2}{2t}} \left(1 + \left(\frac{1}{12} - \frac{2688 + 80t^2 + 21t^3}{322560} e^{-\frac{t}{8}} \right) s^2 \right). \quad (49)$$

5.2 Non-zero correlation

For non-zero correlation model, we need to map it to a zero correlation model. Denote the mapped parameters with a tilde, then the mapped β and γ can be chosen as

$$\tilde{\beta} = \beta, \quad \tilde{\gamma}^2 = \gamma^2 - \frac{3}{2} \left(\gamma^2 \rho^2 + \sigma_0 \gamma \rho (1 - \beta) S^{\beta-1} \right). \quad (50)$$

Then, the mapped initial volatility can be expressed as an expansion in maturity

$$\tilde{\sigma}_0 = \tilde{\sigma}_0^{(0)} + T \tilde{\sigma}_0^{(1)}. \quad (51)$$

The zeroth order term is given by

$$\tilde{\sigma}_0^{(0)} = \frac{2\Phi \delta \tilde{q} \tilde{\gamma}}{\Phi^2 - 1}, \quad (52)$$

where

$$\Phi = \left(\frac{\sigma_{\min} + \rho \sigma_0 + \gamma \delta q}{(1 + \rho) \sigma_0} \right)^{\frac{\tilde{\gamma}}{\gamma}}, \quad (53)$$

and

$$\delta q = \frac{K^{1-\beta} - S^{1-\beta}}{1 - \beta}, \quad \delta \tilde{q} = \frac{K^{1-\tilde{\beta}} - S^{1-\tilde{\beta}}}{1 - \tilde{\beta}} = \delta q, \quad (54)$$

since $\tilde{\beta} = \beta$. Also,

$$\sigma_{\min}^2 = \gamma^2 \delta q^2 + 2\rho \gamma \delta q \sigma_0 + \sigma_0^2. \quad (55)$$

The first order term is given by

$$\frac{\tilde{\sigma}_0^{(1)}}{\tilde{\sigma}_0^{(0)}} = \tilde{\gamma}^2 \sqrt{1 + \tilde{R}^2} \frac{\frac{1}{2} \log \left(\frac{\sigma_0 \sigma_{\min}}{\tilde{\sigma}_0^{(0)} \tilde{\sigma}_{\min}} \right) - \mathcal{B}_{\min}}{\tilde{R} \log \left(\tilde{R} + \sqrt{\tilde{R}^2 + 1} \right)}, \quad (56)$$

where

$$\tilde{R} = \frac{\delta q \tilde{\gamma}}{\tilde{\sigma}_0^{(0)}}, \quad \tilde{\sigma}_{\min}^2 = \tilde{\gamma}^2 \delta q^2 + \left(\tilde{\sigma}_0^{(0)} \right)^2, \quad (57)$$

and \mathcal{B}_{\min} is given by

$$\mathcal{B}_{\min} = -\frac{1}{2} \frac{\beta}{1 - \beta} \frac{\rho}{\sqrt{1 - \rho^2}} (\pi - \varphi_0 - \arccos \rho - I). \quad (58)$$

Here,

$$\varphi_0 = \arccos \left(-\frac{\gamma \delta q + \rho \sigma_0}{\sigma_{\min}} \right), \quad L = \frac{\sigma_{\min}}{\gamma q_K \sqrt{1 - \rho^2}}, \quad (59)$$

and

$$I = \frac{2}{\sqrt{1 - L^2}} \left(\arctan \left(\frac{u_0 + L}{\sqrt{1 - L^2}} \right) - \arctan \left(\frac{L}{\sqrt{1 - L^2}} \right) \right) \quad L < 1 \quad (60)$$

$$I = \frac{1}{\sqrt{L^2 - 1}} \log \left(\frac{1 + u_0 (L + \sqrt{L^2 - 1})}{1 + u_0 (L - \sqrt{L^2 - 1})} \right) \quad L > 1 \quad (61)$$

with the limit $I(L = 1) = 2u_0/(1 + u_0)$, where

$$u_0 = \frac{\delta q \gamma \rho + \sigma_0 - \sigma_{\min}}{\delta q \gamma \sqrt{1 - \rho^2}}. \quad (62)$$

Finally, the call option price of the non-zero correlation model can be obtained by replace the parameters in Eq. (43) with the mapped parameters.

5.2.1 Limiting cases

In the following limiting cases, the mapped γ parameter will be simplified.

i) At-the-money strike. For $K = S$, we have

$$\tilde{\sigma}_0^{(0)} = \sigma_0, \quad (63)$$

and

$$\frac{\tilde{\sigma}_0^{(1)}}{\tilde{\sigma}_0^{(0)}} = \frac{1}{12} \left(1 - \frac{\tilde{\gamma}^2}{\gamma^2} - \frac{3}{2} \rho^2 \right) \gamma^2 + \frac{1}{4} \beta \rho \gamma \sigma_0 S^{\beta-1}. \quad (64)$$

ii) *Negative $\tilde{\gamma}^2$* . If the mapping of $\tilde{\gamma}$ leads to a negative $\tilde{\gamma}^2$, we can set $\tilde{\gamma} = 0$, and the mapped $\tilde{\sigma}_0$ will take the following limiting form,

$$\tilde{\sigma}_0^{(0)} = \frac{\gamma \delta \tilde{q}}{\log \left(\frac{\sigma_{\min} + \rho \sigma_0 + \gamma \delta q}{(1 + \rho) \sigma_0} \right)}, \quad (65)$$

and

$$\frac{\tilde{\sigma}_0^{(1)}}{\tilde{\sigma}_0^{(0)}} = \frac{\gamma^2}{\log^2 \left(\frac{\sigma_{\min} + \rho \sigma_0 + \gamma \delta q}{(1 + \rho) \sigma_0} \right)} \left(\frac{1}{2} \log \left(\frac{\sigma_0 \sigma_{\min}}{\tilde{\sigma}_0^{(0)} \tilde{\sigma}_{\min}} \right) - \mathcal{B}_{\min} \right). \quad (66)$$

Since $\tilde{\gamma} = 0$, the mapped model will degenerate to the CEV model, and the call option price can be obtained by replacing the parameters in Eq. (39) with the mapped parameters.

6 Free boundary CEV

The CEV process (1) can be modified to allow negative forward in the following way,

$$dS_t = \sigma |S_t|^\beta dW_t, \quad (67)$$

with $0 < \beta < 1/2$. The corresponding Fokker-Planck equation is then given by

$$-\frac{\partial p(t, s)}{\partial t} + \frac{1}{2} \sigma^2 \frac{\partial^2}{\partial s^2} \left(|s|^{2\beta} p(t, s) \right) = 0, \quad (68)$$

with initial condition (3). The solution to the Fokker-Planck equation (68) can be constructed as

$$p(t, s) = \frac{1}{2} \left(p_R(t, |s|) + \text{sgn}(s) p_A(t, |s|) \right), \quad (69)$$

where $p_A(t, s)$ and $p_R(t, s)$ are the solution to the Fokker-Planck equation (2) under the absorbing and reflecting boundary conditions, respectively. Following the time value of the option (18), it can be shown that the call option price in the free boundary CEV model is

$$\begin{aligned} C(T, S, K) &= (S - K)^+ \\ &+ \frac{\sqrt{|SK|}}{\pi} \left(\mathbf{1}_{K \geq 0} \int_0^\pi \frac{\sin \phi \sin(\nu \phi)}{b - \cos \phi} \exp \left(-\frac{\bar{q}(b - \cos \phi)}{\sigma^2 T} \right) d\phi \right. \\ &\quad \left. + \sin(\nu \pi) \int_0^{+\infty} \frac{d\psi}{b + \cosh \psi} \exp \left(-\frac{\bar{q}(b + \cosh \psi)}{\sigma^2 T} \right) \right. \\ &\quad \left. \times \left(\mathbf{1}_{K \geq 0} \cosh(\nu \psi) + \mathbf{1}_{K < 0} \sinh(\nu \psi) \right) \sinh \psi \right), \quad (70) \end{aligned}$$

where

$$\bar{q} = q_K q_S, \quad b = \frac{q_K^2 + q_S^2}{2q_K q_S}, \quad q_K = \frac{|K|^{1-\beta}}{1-\beta}, \quad q_S = \frac{|S|^{1-\beta}}{1-\beta}. \quad (71)$$

7 Free boundary SABR

Similar to the free boundary CEV process, the forward process of the SABR model can also be modified to allow negative forward in the following way,

$$dS_t = \sigma |S_t|^\beta dW_t, \quad (72)$$

with $0 < \beta < 1/2$, which becomes the free boundary SABR model. When the correlation is zero, $\rho = 0$, the call option value has the similar form to the SABR model,

$$\begin{aligned} C(T, S, K) = & (S - K)^+ \\ & + \frac{\sqrt{|SK|}}{\pi} \left(\mathbf{1}_{K \geq 0} \int_0^\pi \frac{\sin \phi \sin(\nu \phi)}{b - \cos \phi} \frac{G(\gamma^2 T, s(\phi))}{\cosh(s(\phi))} d\phi \right. \\ & \quad \left. + \sin(\nu \pi) \int_0^{+\infty} \frac{d\psi}{b + \cosh \psi} \frac{G(\gamma^2 T, s(\psi))}{\cosh(s(\psi))} \right. \\ & \quad \left. \times \left(\mathbf{1}_{K \geq 0} \cosh(\nu \psi) + \mathbf{1}_{K < 0} \sinh(\nu \psi) \right) \sinh \psi \right). \quad (73) \end{aligned}$$

For non-zero correlation, the parameters can be mapped to those of a zero correlation model as in the SABR case. The only difference is that

$$\delta q = \frac{|k|^{1-\beta} - |S|^{1-\beta}}{1 - \beta}, \quad \delta \tilde{q} = \frac{|k|^{1-\tilde{\beta}} - |S|^{1-\tilde{\beta}}}{1 - \tilde{\beta}} = \delta q, \quad (74)$$

where

$$k = \max(K, 0.1S) \quad (75)$$

and

$$L = \frac{\sigma_{\min}(1 - \beta)}{\gamma k^{1-\beta} \sqrt{1 - \rho^2}}, \quad (76)$$

If the mapped $\tilde{\gamma}^2$ is negative, we can similarly set $\tilde{\gamma}$ to 0, and the free boundary SABR model will degenerate to the corresponding free boundary CEV model.

References

- [1] See <http://www.lesniewski.us/papers/working/NotesOnCEV.pdf>.
- [2] See <http://mathworld.wolfram.com/BesselDifferentialEquation.html>.
- [3] See <http://dlmf.nist.gov/10.22.E67>.
- [4] See <http://dlmf.nist.gov/10.32.E4>.