

The Jones matrix of $\frac{\lambda}{2}$ wave plate is.

$$A_{\frac{\lambda}{2}} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

of which we choose horizontal as fast axis.

If we rotate $A_{\frac{\lambda}{2}}$ by angle $-\theta$ and $A_{\frac{\lambda}{4}}$ by angle $-\alpha$ we can get.

$$A_1 = S(+\theta) A_{\frac{\lambda}{2}} S(-\theta)$$

$$= \begin{bmatrix} \cos 2\theta & -\sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{bmatrix}$$

$$A_2 = S(\alpha) A_{\frac{\lambda}{4}} S(\alpha)$$

$$= \begin{bmatrix} \cos^2 \alpha + i \sin^2 \alpha & (i-1) \frac{\sin 2\alpha}{2} \\ (i-1) \frac{\sin 2\alpha}{2} & -\sin^2 \alpha + i \cos^2 \alpha \end{bmatrix}$$

~~If we plan the $A_{\frac{\lambda}{2}}$ before~~

If the light travel through $A_{\frac{\lambda}{2}}$ first then $A_{\frac{\lambda}{4}}$.

Suppose the input beam is.

$$E_0 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Therefore we can get the transmission light as.

$$E_t = A_2 A_1 E.$$

$$= A_2 \begin{bmatrix} \cos 2\theta & -\sin 2\theta \\ -\sin 2\theta & -\cos 2\theta \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

$$= \begin{bmatrix} \cos^2 \alpha + i \sin^2 \alpha & (i-1) \frac{\sin 2\alpha}{2} \\ (i-1) \frac{\sin 2\alpha}{2} & -\sin^2 \alpha + i \cos^2 \alpha \end{bmatrix} \begin{bmatrix} \cos 2\theta \\ -\sin 2\theta \end{bmatrix}$$

$$= \begin{bmatrix} \cos \alpha \cos(2\theta - \alpha) + i \sin \alpha \sin(\alpha - 2\theta) \\ -\sin \alpha \cos(2\theta - \alpha) + i \cos \alpha \sin(\alpha - 2\theta) \end{bmatrix}$$

By carefully choose α, θ we can get the ~~out~~ transmission beam as.

$$1) \sin(\alpha - 2\theta) = 0 \Rightarrow \alpha - 2\theta = n\pi \quad (n \in \mathbb{Z})$$

$$E_t = \begin{bmatrix} \cos \alpha \\ -\sin \alpha \end{bmatrix}$$

Therefore E_t can represent any linear polarization.

$$2) \cos(2\theta - \alpha) = 0 \Rightarrow 2\theta - \alpha = (n + \frac{1}{2})\pi$$

$$E_t = \begin{bmatrix} i \sin \alpha \\ \cos \alpha \end{bmatrix}$$

Therefore E_t can represent any ~~the~~ circular polarization light.

2. problem 4.1 in textbook.

4.1

The Jones of a beam splitter has two parts, the reflection parts and transmission part.

Assume that there is no loss in the beam splitter.

$$A_T = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}, \quad A_R = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$$

And the rotation angle is θ , therefore

$$S(\theta) = \begin{bmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{bmatrix}, \quad S(\theta)^{-1} = \begin{bmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{bmatrix}$$

Therefore, after rotation

$$A_T' = S(\theta) A_T S^{-1}(\theta) = \begin{bmatrix} \cos^2\theta & -\sin\theta\cos\theta \\ -\sin\theta\cos\theta & \sin^2\theta \end{bmatrix}$$

$$A_R' = S(\theta) A_R S^{-1}(\theta) = \begin{bmatrix} \sin^2\theta & \sin\theta\cos\theta \\ \sin\theta\cos\theta & \cos^2\theta \end{bmatrix}$$

Problem 4.2

From the description we know that the Jones matrix of this element is

$$A = S\left(\frac{\pi}{4}\right) A_H S\left(\frac{\pi}{4}\right) S\frac{\pi}{4}$$

$$= \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$

$$= \frac{1}{2} \begin{bmatrix} 1 & i \\ 1 & i \end{bmatrix}$$

Of which we choose vertical as fast axis.

If the input beam is a left circular beam for example.

$$E_i = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix}$$

$$E_t = A E_i = \frac{1}{2\sqrt{2}} \begin{bmatrix} 1 & i \\ 1 & i \end{bmatrix} \begin{bmatrix} 1 \\ i \end{bmatrix} = \frac{1}{2\sqrt{2}} \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Therefore the incident light will be rejected.

If it is a left circular beam

$$E_i = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix}$$

$$E_t = A E_i = \frac{1}{2\sqrt{2}} \begin{bmatrix} 2 \\ 2 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Which indicates that it will become a linear polarization light. If we rotate the device by $-\theta$. Then we can get

$$A_1 = S(\theta) A S(-\theta)$$

$$= \frac{1}{2} \begin{pmatrix} \cos^2\theta + \sin\theta\cos\theta + i(\sin^2\theta + \sin\theta\cos\theta) & -\sin\theta\cos\theta - \sin^2\theta + i \\ (\sin\theta\cos\theta + \cos^2\theta) & \\ -\sin\theta\cos\theta - i\sin^2\theta + \cos^2\theta + i\sin\theta\cos\theta & \sin^2\theta + i\sin\theta\cos\theta - \\ \sin\theta\cos\theta + i\cos^2\theta & \end{pmatrix}$$

Therefore we A_1 acts on different circular wave. We can get that.

$$A_1 \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix} = \begin{pmatrix} \cos^2\theta + \sin\theta\cos\theta + i(\sin^2\theta + \sin\theta\cos\theta) - i \dots \\ -\sin\theta\cos\theta - i\sin^2\theta + \cos^2\theta + i\sin\theta\cos\theta + \dots \end{pmatrix}$$

$$= \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$A_1 \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix} = \begin{bmatrix} (\sin\theta + \cos\theta)(\cos\theta + i\sin\theta) \\ (\cos\theta - \sin\theta)(\cos\theta + i\sin\theta) \end{bmatrix}$$

$$= (\cos\theta + i\sin\theta) \begin{bmatrix} \sin\theta + \cos\theta \\ \cos\theta - \sin\theta \end{bmatrix}$$

Which is still a linear polarization beam.

Problem 4.

The zero-order waveplate can be made by align the fast axis of the first plate to the slow axis of the second waveplate. Suppose that the refractive index of these two waveplate are n_e, n_o, n_e', n_o' . For a $k\lambda$ ($k = \frac{1}{2}$ or $\frac{1}{4} \dots$) waveplate, there should be.

$$d_1(n_e - n_o) = (m + k_1)\lambda$$

$$d_2(n_o' - n_e') = (n + k_2)\lambda$$

If we want to make a zero-order ~~pta~~ waveplate. Therefore we can get that.

$$d_1(n_e - n_o) - d_2(n_o' - n_e') = (n - m + k_1 - k_2)\lambda = k\lambda$$

That is the rule how we choose d_1 and d_2 .

If $n_e = n_e', n_o = n_o'$, we can simplify it as.

$$(d_1 - d_2)(n_e - n_o) = k\lambda \quad (k = \frac{1}{2}, \frac{1}{4})$$

Suppose that n_e and n_o ^{are} functions of temperature and wavelength. Let $k\lambda = \text{OPD}(\lambda)$ (optical path difference)

$$\frac{\partial \text{OPD}(\lambda)}{\partial T} = (d_1 - d_2) \left(\frac{\partial n_e}{\partial T} - \frac{\partial n_o}{\partial T} \right)$$

Compare with that of multi-order waveplate

$$\frac{\partial \text{OPD}_1(\lambda)}{\partial T} = d_1 \left(\frac{\partial n_e}{\partial T} - \frac{\partial n_o}{\partial T} \right)$$

$$\frac{\partial \text{OPD}_2(\lambda)}{\partial T} = d_2 \left(\frac{\partial n_o'}{\partial T} - \frac{\partial n_e'}{\partial T} \right)$$

$$\text{Therefore } \frac{\partial \text{OPD}}{\partial T} = \frac{\partial \text{OPD}_1}{\partial T} - \frac{\partial \text{OPD}_2}{\partial T}.$$

~~Hence it has small~~

$$\text{Due to that } \begin{cases} d_1 - d_2 < d_1 \\ d_1 - d_2 < d_2 \end{cases}$$

Therefore zero-order wave plate has smaller dependence on temperature.
That's the same for wavelength.