Assignment #1

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1 Solution 1

Take $\mu(x,y,z)=\psi(x,y,z)e^{-ikz}$ into the wave equation $\nabla^2\mu+k^2u=0$,

$$\nabla_t^2(\psi e^{-ikz}) + \frac{\partial^2}{\partial z^2}(\psi e^{-ikz}) + k^2 \psi e^{-ikz} = 0 \tag{1}$$

where

$$\nabla_t^2 = (\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2})$$

Expand $\frac{\partial^2}{\partial z^2}(\psi e^{-ikz})$ we can get,

$$\frac{\partial^2}{\partial z^2} (\psi e^{-ikz}) = (\frac{\partial^2 \psi}{\partial z^2} - 2ik \frac{\partial \psi}{\partial z} - k^2 \psi) e^{-ikz}$$
 (2)

Take equation 2 into equation 1,

$$\nabla_t^2 \psi + \frac{\partial^2 \psi}{\partial z^2} - 2ik \frac{\partial \psi}{\partial z} = 0 \tag{3}$$

In order for paraxial condition to be satisfied, which means $\psi(x,y,z)$ must vary much more slowly along the z-axis than it does along the x- and y-axes. Thus, when substituting equation 2 into the wave equation, $\partial^2 \psi / \partial z^2$ can be ignored compared to the other second derivatives. So we can rewrite equation 1 as

$$\nabla_t^2 \psi - 2ik \frac{\partial \psi}{\partial z} = 0 \tag{4}$$

Take the trial function as,

$$\psi(x,y,z) = \exp\{-i(P(z) + \frac{k}{2q(z)}r^2)\}, r^2 = x^2 + y^2$$
 (5)

and in the cylindrical coordinate $\nabla_t^2 = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}$, substitute equation 5 into equation 4 and organized by the order of r.

$$\frac{k^2}{q(z)^2}(\frac{dq(z)}{dz} - 1)r^2 - 2k(\frac{dP(z)}{dz} + \frac{i}{q(z)}) = 0 \tag{6}$$

Due to that equation 6 is always equal to zero no matter how r varies, which indicates the parameter equals to zero,

$$\frac{dq(z)}{dz} - 1 = 0$$

$$\frac{dP(z)}{dz} + \frac{i}{q(z)} = 0$$

Finally, we can get

$$\frac{dq(z)}{dz} = 1$$

$$\frac{dP(z)}{dz} = -\frac{i}{q(z)}$$
(7)

2 Solution 2

By solving equation 7 we could get, $q(z) = q(z_1) + (z_2 - z_1)$. Write the complex beam parameter in two real parameters, R and ω ,

$$\frac{1}{q} = \frac{1}{R} - i \frac{\lambda}{n\pi\omega^2}.$$
 (8)

Substitute equation 8 into $\mu(x, y, z)$,

$$\mu(x, y, z) = exp\{-i(P(z) + kz + k\frac{r^2}{2R}) - \frac{r^2}{\omega^2}\},\tag{9}$$

where $k=2n\pi/\lambda$ is used. Now consider about the surface of constant phase which is r-dependence, which means the phase parameter in equation 9 is a constant,

$$z + \frac{r^2}{2R} = constant \tag{10}$$

and P(z) has been ignore for that it varies much more slowly than other two terms. R can be interpreted as the radius of wavefront by the equation of calculating the curvature under the paraxial approximation $r \ll R$.

$$\rho = \frac{ds}{d\theta} = \frac{(1 + f'(x))^{3/2}}{f''(x)}.$$
(11)

Substitute

$$z = f(r) = constant - \frac{r^2}{2R},$$

into equation 11,

$$\rho = R(1 + \frac{r}{R})^{3/2} \tag{12}$$

So the curvature of wave front $\rho=R$ under the paraxial condition $r\ll R$. Now consider the minimum size of the spot when the wave front becomes a plan $(R=\infty)$, the corresponding value of q is:

At waist:
$$q \equiv q_0 = i \frac{n\pi\omega^2}{\lambda}$$
 (13)

Setting z as the distance from the waist, ω_0 as the radius of waist, the values of q is,

$$q(z) = q_0 + z = i\frac{n\pi\omega^2}{\lambda} + z,\tag{14}$$

From equation 14 we can derive that,

Distanace to waist: $d = -Re\{q(z)\}$

Rdius of waist:
$$\omega_0 = \sqrt{\frac{\lambda}{n\pi} Im\{q(z)\}}$$
 (15)

By solving equation 8 and 14, we can obtain ω and R as the function of propagation distance z from waist:

$$\omega(z) = \omega_0 [1 + (\frac{\lambda z}{n\pi\omega_0^2})^2]^{1/2}$$

$$R(z) = z [1 + (\frac{n\pi\omega_0^2}{\lambda z})^2].$$
(16)

The figure of $\omega(z)$ and R(z) is drew as in figure 1 and 2,where we have assumed that $\omega_0=1, \lambda=1, n=1$ for simplification.

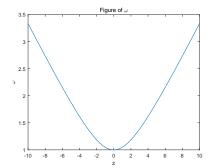


Figure 1: Figure of $\omega(z)$

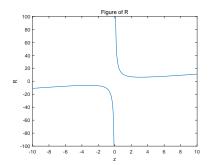


Figure 2: Figure of R(z)