

PH 534: Quantum Information and Computing Tutorial

Siddhant Midha

<https://siddhant-midha.github.io/>

January 26, 2023

- States: $|\psi\rangle$
- $|\psi^\pm\rangle = \frac{1}{\sqrt{2}}(|01\rangle \pm |10\rangle)$, $|\phi^\pm\rangle = \frac{1}{\sqrt{2}}(|00\rangle \pm |11\rangle)$,
 $|\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle)$
- Density operators: $\rho, \sigma \dots$
- Hilbert Space: \mathcal{H}
- Set of linear operators from \mathcal{H}_1 to \mathcal{H}_2 : $L(\mathcal{H}_1, \mathcal{H}_2)$
- Set of density matrices on \mathcal{H} : $D(\mathcal{H}) \subset L(\mathcal{H})$
- Set of linear operators from $L(\mathcal{H}_1)$ to $L(\mathcal{H}_2)$: $T(\mathcal{H}_1, \mathcal{H}_2)$
- Set of quantum channels from $L(\mathcal{H}_1)$ to $L(\mathcal{H}_2)$:
 $C(\mathcal{H}_1, \mathcal{H}_2) \subset T(\mathcal{H}_1, \mathcal{H}_2)$

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i ,

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

- The density operator is positive and has unit trace.

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

- The density operator is positive and has unit trace. The converse holds.
- Unitary evolution

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

- The density operator is positive and has unit trace. The converse holds.
- Unitary evolution

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i| \rightarrow \sum_i p_i U |\psi_i\rangle \langle \psi_i| U^* = U \rho U^*$$

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

- The density operator is positive and has unit trace. The converse holds.
- Unitary evolution

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i| \rightarrow \sum_i p_i U |\psi_i\rangle \langle \psi_i| U^* = U \rho U^*$$

- Measurements

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

- The density operator is positive and has unit trace. The converse holds.
- Unitary evolution

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i| \rightarrow \sum_i p_i U |\psi_i\rangle \langle \psi_i| U^* = U \rho U^*$$

- Measurements

$$p(m) = \text{Tr}(M_m^* M_m \rho)$$

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

- The density operator is positive and has unit trace. The converse holds.
- Unitary evolution

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i| \rightarrow \sum_i p_i U |\psi_i\rangle \langle \psi_i| U^* = U \rho U^*$$

- Measurements

$$p(m) = \text{Tr}(M_m^* M_m \rho)$$

$$\rho_m = \frac{M_m \rho M_m^*}{\text{Tr}(M_m^* M_m \rho)}$$

- If a system exists in the states $|\psi_i\rangle$ with probabilities p_i , then its density operator ρ is defined as

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

- The density operator is positive and has unit trace. The converse holds.
- Unitary evolution

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i| \rightarrow \sum_i p_i U |\psi_i\rangle \langle \psi_i| U^* = U \rho U^*$$

- Measurements

$$p(m) = \text{Tr}(M_m^* M_m \rho)$$

$$\rho_m = \frac{M_m \rho M_m^*}{\text{Tr}(M_m^* M_m \rho)}$$

If we have systems A and B , described by the density operator $\rho_{AB} \in D(\mathcal{H}_A \otimes \mathcal{H}_B)$, we define the density operator for the subsystem A as

If we have systems A and B , described by the density operator $\rho_{AB} \in D(\mathcal{H}_A \otimes \mathcal{H}_B)$, we define the density operator for the subsystem A as

$$\rho_A \equiv \text{Tr}_B(\rho_{AB}) \in D(\mathcal{H}_A)$$

If we have systems A and B , described by the density operator $\rho_{AB} \in D(\mathcal{H}_A \otimes \mathcal{H}_B)$, we define the density operator for the subsystem A as

$$\rho_A \equiv \text{Tr}_B(\rho_{AB}) \in D(\mathcal{H}_A)$$

with the partial trace operation Tr_B being defined as

If we have systems A and B , described by the density operator $\rho_{AB} \in D(\mathcal{H}_A \otimes \mathcal{H}_B)$, we define the density operator for the subsystem A as

$$\rho_A \equiv \text{Tr}_B(\rho_{AB}) \in D(\mathcal{H}_A)$$

with the partial trace operation Tr_B being defined as

$$\text{Tr}_B(|a_1\rangle\langle a_2| \otimes |b_1\rangle\langle b_2|) \equiv |a_1\rangle\langle a_2| \text{tr}(|b_1\rangle\langle b_2|)$$

If we have systems A and B , described by the density operator $\rho_{AB} \in D(\mathcal{H}_A \otimes \mathcal{H}_B)$, we define the density operator for the subsystem A as

$$\rho_A \equiv \text{Tr}_B(\rho_{AB}) \in D(\mathcal{H}_A)$$

with the partial trace operation Tr_B being defined as

$$\text{Tr}_B(|a_1\rangle\langle a_2| \otimes |b_1\rangle\langle b_2|) \equiv |a_1\rangle\langle a_2| \text{tr}(|b_1\rangle\langle b_2|)$$

and for general states ρ_{AB} the definition extends by superposition.

If we have systems A and B , described by the density operator $\rho_{AB} \in D(\mathcal{H}_A \otimes \mathcal{H}_B)$, we define the density operator for the subsystem A as

$$\rho_A \equiv \text{Tr}_B(\rho_{AB}) \in D(\mathcal{H}_A)$$

with the partial trace operation Tr_B being defined as

$$\text{Tr}_B(|a_1\rangle\langle a_2| \otimes |b_1\rangle\langle b_2|) \equiv |a_1\rangle\langle a_2| \text{tr}(|b_1\rangle\langle b_2|)$$

and for general states ρ_{AB} the definition extends by superposition. Or,

$$\text{Tr}_B(\rho_{AB}) := \sum_j (I_A \otimes \langle j|_B) \rho_{AB} (I_A \otimes |j\rangle_B)$$

Write the matrix representation, and calculate the reduced single-qubit density matrices ρ_A and ρ_B .

- ① $\rho_{AB} = \frac{1}{3}|00\rangle\langle 00| + \frac{2}{3}|\psi^+\rangle\langle\psi^+|$, where $|\psi^\pm\rangle = \frac{1}{\sqrt{2}}(|01\rangle \pm |10\rangle)$.
- ② $\rho_{AB} = (1-p)\mathbb{I}/4 + p|\psi^-\rangle\langle\psi^-|$, where \mathbb{I} is the 4×4 identity matrix.
- ③ $\rho_{AB} = p|\psi^-\rangle\langle\psi^-| + (1-p)|\psi^+\rangle\langle\psi^+|$.

Theorem 2.7 of QCQI

Let $|\psi\rangle \in \mathcal{H}_{AB}$.

Theorem 2.7 of QCQI

Let $|\psi\rangle \in \mathcal{H}_{AB}$. Then there exist orthonormal states $|i_A\rangle \in \mathcal{H}_A$ and orthonormal states $|i_B\rangle \in \mathcal{H}_B$ such that

Theorem 2.7 of QCQI

Let $|\psi\rangle \in \mathcal{H}_{AB}$. Then there exist orthonormal states $|i_A\rangle \in \mathcal{H}_A$ and orthonormal states $|i_B\rangle \in \mathcal{H}_B$ such that

$$|\psi\rangle = \sum_i \lambda_i |i_A\rangle \otimes |i_B\rangle$$

Theorem 2.7 of QCQI

Let $|\psi\rangle \in \mathcal{H}_{AB}$. Then there exist orthonormal states $|i_A\rangle \in \mathcal{H}_A$ and orthonormal states $|i_B\rangle \in \mathcal{H}_B$ such that

$$|\psi\rangle = \sum_i \lambda_i |i_A\rangle \otimes |i_B\rangle$$

where $\lambda_i \geq 0$ and $\sum_i \lambda_i^2 = 1$.

Question Consider the state, $|\psi\rangle = \frac{1}{\sqrt{3}}(|00\rangle + |01\rangle + |10\rangle)$. Find its Schmidt Decomposition.

- Suppose

$$\rho_A = \sum_i p_i |i_A\rangle \langle i_A|$$

- Introduce R with $\mathcal{H}_R = \mathcal{H}_A$ (can we do better?) and a orthonormal basis $|i_R\rangle$.

- Suppose

$$\rho_A = \sum_i p_i |i_A\rangle \langle i_A|$$

- Introduce R with $\mathcal{H}_R = \mathcal{H}_A$ (can we do better?) and a orthonormal basis $|i_R\rangle$.
- Define

$$|\psi_{AR}\rangle \equiv \sum_i \sqrt{p_i} |i_A\rangle |i_R\rangle$$

- Suppose

$$\rho_A = \sum_i p_i |i_A\rangle \langle i_A|$$

- Introduce R with $\mathcal{H}_R = \mathcal{H}_A$ (can we do better?) and a orthonormal basis $|i_R\rangle$.
- Define

$$|\psi_{AR}\rangle \equiv \sum_i \sqrt{p_i} |i_A\rangle |i_R\rangle$$

- See that

$$\rho_A = \text{Tr}_R(|\psi_{AR}\rangle \langle \psi_{AR}|)$$

- Suppose

$$\rho_A = \sum_i p_i |i_A\rangle \langle i_A|$$

- Introduce R with $\mathcal{H}_R = \mathcal{H}_A$ (can we do better?) and a orthonormal basis $|i_R\rangle$.
- Define

$$|\psi_{AR}\rangle \equiv \sum_i \sqrt{p_i} |i_A\rangle |i_R\rangle$$

- See that

$$\rho_A = \text{Tr}_R(|\psi_{AR}\rangle \langle \psi_{AR}|)$$

Question Compute the purification for the state

$$\rho = \frac{1}{2}(|0\rangle \langle 0| + |+\rangle \langle +|)$$

Further, apply the unitary $U := \mathbb{1} \otimes \sigma_y$ to the purified state, and calculate the reduced state to show whether or not such unitaries create equivalent purifications.

Quantum Operations: Beyond Unitaries

- Natural Extension of unitary operations:

$$\mathcal{E}(\rho) = \text{Tr}_{\text{env}}(U(\rho \otimes \rho_{\text{env}})U^*)$$

- Discard the environment: need some axioms.

Quantum Operations: Beyond Unitaries

- Natural Extension of unitary operations:

$$\mathcal{E}(\rho) = \text{Tr}_{\text{env}}(U(\rho \otimes \rho_{\text{env}})U^*)$$

- Discard the environment: need some axioms.
 - ① $\text{Tr}[\mathcal{E}(\rho)] \in [0, 1] \forall \rho$ ($\text{Tr}[\mathcal{E}(\rho)]$ is the probability that ρ undergoes the transformation \mathcal{E}).

Quantum Operations: Beyond Unitaries

- Natural Extension of unitary operations:

$$\mathcal{E}(\rho) = \text{Tr}_{env}(U(\rho \otimes \rho_{env})U^*)$$

- Discard the environment: need some axioms.
 - ① $\text{Tr}[\mathcal{E}(\rho)] \in [0, 1] \forall \rho$ ($\text{Tr}[\mathcal{E}(\rho)]$ is the probability that ρ undergoes the transformation \mathcal{E}).
 - ② Convex linearity

$$\mathcal{E}\left(\sum_i p_i \rho_i\right) = \sum_i p_i \mathcal{E}(\rho_i)$$

for all density matrices ρ_i and probabilities p_i s.t. $\sum_i p_i = 1$

- ③ \mathcal{E} is completely positive. Not only does \mathcal{E} preserve positivity, $(I \otimes \mathcal{E})$ also preserves positivity for I being the identity on an arbitrarily dimensional system's hilbert space.

Quantum Operations: Beyond Unitaries

- Natural Extension of unitary operations:

$$\mathcal{E}(\rho) = \text{Tr}_{\text{env}}(U(\rho \otimes \rho_{\text{env}})U^*)$$

- Discard the environment: need some axioms.
 - ① $\text{Tr}[\mathcal{E}(\rho)] \in [0, 1] \forall \rho$ ($\text{Tr}[\mathcal{E}(\rho)]$ is the probability that ρ undergoes the transformation \mathcal{E}).
 - ② Convex linearity

$$\mathcal{E}\left(\sum_i p_i \rho_i\right) = \sum_i p_i \mathcal{E}(\rho_i)$$

for all density matrices ρ_i and probabilities p_i s.t. $\sum_i p_i = 1$

- ③ \mathcal{E} is completely positive. Not only does \mathcal{E} preserve positivity, $(I \otimes \mathcal{E})$ also preserves positivity for I being the identity on an arbitrarily dimensional system's hilbert space.

Formally,

$$\mathcal{E} : L(\mathcal{H}_1) \rightarrow L(\mathcal{H}_2) \in \mathcal{C}(\mathcal{H}_1, \mathcal{H}_2)$$

For a single qubit state ρ , a measurement in the computational basis can be described by the operations

$$\mathcal{E}_0(\rho) \equiv |0\rangle\langle 0|\rho|0\rangle\langle 0| \text{ and } \mathcal{E}_1(\rho) \equiv |1\rangle\langle 1|\rho|1\rangle\langle 1|$$

For a single qubit state ρ , a measurement in the computational basis can be described by the operations $\mathcal{E}_0(\rho) \equiv |0\rangle\langle 0|\rho|0\rangle\langle 0|$ and $\mathcal{E}_1(\rho) \equiv |1\rangle\langle 1|\rho|1\rangle\langle 1|$ with probabilities given as $\text{tr}[\mathcal{E}_0(\rho)]$ and $\text{tr}[\mathcal{E}_1(\rho)]$.

For a single qubit state ρ , a measurement in the computational basis can be described by the operations $\mathcal{E}_0(\rho) \equiv |0\rangle\langle 0|\rho|0\rangle\langle 0|$ and $\mathcal{E}_1(\rho) \equiv |1\rangle\langle 1|\rho|1\rangle\langle 1|$ with probabilities given as $\text{tr}[\mathcal{E}_0(\rho)]$ and $\text{tr}[\mathcal{E}_1(\rho)]$. The final state is

$$\frac{\mathcal{E}_i(\rho)}{\text{Tr}([\mathcal{E}_i(\rho)])} \text{ for some } i \in \{0, 1\}$$

For a single qubit state ρ , a measurement in the computational basis can be described by the operations $\mathcal{E}_0(\rho) \equiv |0\rangle\langle 0|\rho|0\rangle\langle 0|$ and $\mathcal{E}_1(\rho) \equiv |1\rangle\langle 1|\rho|1\rangle\langle 1|$ with probabilities given as $\text{tr}[\mathcal{E}_0(\rho)]$ and $\text{tr}[\mathcal{E}_1(\rho)]$. The final state is

$$\frac{\mathcal{E}_i(\rho)}{\text{Tr}([\mathcal{E}_i(\rho)])} \text{ for some } i \in \{0, 1\}$$

That is, if no measurement is happening, the map \mathcal{E} would be a *completely positive trace preserving (CPTP)* map.

Theorem 8.1 of QCQI

The map \mathcal{E} satisfies the axioms for a valid quantum operation iff there exists a set of operators $\{E_i\}$ such that

$$\mathcal{E}(\rho) = \sum_i E_i \rho E_i^*$$

for all valid density matrices ρ and $0 \preceq \sum_i E_i^* E_i \preceq I$

Note: $A \preceq B$ if $B - A$ is PSD.

Theorem 8.1 of QCQI

The map \mathcal{E} satisfies the axioms for a valid quantum operation iff there exists a set of operators $\{E_i\}$ such that

$$\mathcal{E}(\rho) = \sum_i E_i \rho E_i^*$$

for all valid density matrices ρ and $0 \preceq \sum_i E_i^* E_i \preceq I$

Note: $A \preceq B$ if $B - A$ is PSD. So if we are just dealing with CPTP maps, then these E_i satisfy $\sum_i E_i^* E_i = I$, and are called the *kraus operators*.

① **The Vectorization Map:** $\text{vec} : L(\mathcal{H}_2, \mathcal{H}_1) \rightarrow \mathcal{H}_1 \otimes \mathcal{H}_2$,
defined as

$$|\text{vec}(|a_1\rangle \langle b_2|)\rangle := |a_1\rangle \otimes |b_2\rangle$$

① **The Vectorization Map:** $\text{vec} : L(\mathcal{H}_2, \mathcal{H}_1) \rightarrow \mathcal{H}_1 \otimes \mathcal{H}_2$,
defined as

$$|\text{vec}(|a_1\rangle \langle b_2|)\rangle := |a_1\rangle \otimes |b_2\rangle$$

Thus, the vec map takes operators in $L(\mathcal{H})$ to vectors in $\mathcal{H} \otimes \mathcal{H}$.

- ① **The Vectorization Map:** $\text{vec} : L(\mathcal{H}_2, \mathcal{H}_1) \rightarrow \mathcal{H}_1 \otimes \mathcal{H}_2$,
defined as

$$|\text{vec}(|a_1\rangle \langle b_2|)\rangle := |a_1\rangle \otimes |b_2\rangle$$

Thus, the vec map takes operators in $L(\mathcal{H})$ to vectors in $\mathcal{H} \otimes \mathcal{H}$.

- ② **The Choi Representation** $\mathcal{J} : T(\mathcal{H}_1, \mathcal{H}_2) \rightarrow L(\mathcal{H}_2 \otimes \mathcal{H}_1)$

$$\mathcal{J}(\mathcal{E}) := (\mathcal{E} \otimes \mathbb{1}_{L(\mathcal{H}_1)})(|\text{vec}(\mathbb{1}_{\mathcal{H}_1})\rangle \langle \text{vec}(\mathbb{1}_{\mathcal{H}_1})|)$$

- The Vectorization Map:** $\text{vec} : L(\mathcal{H}_2, \mathcal{H}_1) \rightarrow \mathcal{H}_1 \otimes \mathcal{H}_2$, defined as

$$|\text{vec}(|a_1\rangle \langle b_2|)\rangle := |a_1\rangle \otimes |b_2\rangle$$

Thus, the vec map takes operators in $L(\mathcal{H})$ to vectors in $\mathcal{H} \otimes \mathcal{H}$.

- The Choi Representation** $\mathcal{J} : T(\mathcal{H}_1, \mathcal{H}_2) \rightarrow L(\mathcal{H}_2 \otimes \mathcal{H}_1)$

$$\mathcal{J}(\mathcal{E}) := (\mathcal{E} \otimes \mathbb{1}_{L(\mathcal{H}_1)})(|\text{vec}(\mathbb{1}_{\mathcal{H}_1})\rangle \langle \text{vec}(\mathbb{1}_{\mathcal{H}_1})|)$$

Note that $|\text{vec}(\mathbb{1}_{\mathcal{H}_1})\rangle$ is a vector in $\mathcal{H}_1 \otimes \mathcal{H}_1$, and thus $(|\text{vec}(\mathbb{1}_{\mathcal{H}_1})\rangle \langle \text{vec}(\mathbb{1}_{\mathcal{H}_1})|)$ is an operator on $\mathcal{H}_1 \otimes \mathcal{H}_1$. Now, we know that \mathcal{E} takes operators from $L(\mathcal{H}_1)$ to $L(\mathcal{H}_2)$, and $\mathbb{1}_{L(\mathcal{H}_1)}$ takes operators from $L(\mathcal{H}_1)$ to $L(\mathcal{H}_1)$. Thus, $\mathcal{J}(\mathcal{E})$ lies in $L(\mathcal{H}_2 \otimes \mathcal{H}_1)$.

- ① **The Vectorization Map:** $\text{vec} : L(\mathcal{H}_2, \mathcal{H}_1) \rightarrow \mathcal{H}_1 \otimes \mathcal{H}_2$, defined as

$$|\text{vec}(|a_1\rangle \langle b_2|)\rangle := |a_1\rangle \otimes |b_2\rangle$$

Thus, the vec map takes operators in $L(\mathcal{H})$ to vectors in $\mathcal{H} \otimes \mathcal{H}$.

- ② **The Choi Representation** $\mathcal{J} : T(\mathcal{H}_1, \mathcal{H}_2) \rightarrow L(\mathcal{H}_2 \otimes \mathcal{H}_1)$

$$\mathcal{J}(\mathcal{E}) := (\mathcal{E} \otimes \mathbb{1}_{L(\mathcal{H}_1)})(|\text{vec}(\mathbb{1}_{\mathcal{H}_1})\rangle \langle \text{vec}(\mathbb{1}_{\mathcal{H}_1})|)$$

Note that $|\text{vec}(\mathbb{1}_{\mathcal{H}_1})\rangle$ is a vector in $\mathcal{H}_1 \otimes \mathcal{H}_1$, and thus $(|\text{vec}(\mathbb{1}_{\mathcal{H}_1})\rangle \langle \text{vec}(\mathbb{1}_{\mathcal{H}_1})|)$ is an operator on $\mathcal{H}_1 \otimes \mathcal{H}_1$. Now, we know that \mathcal{E} takes operators from $L(\mathcal{H}_1)$ to $L(\mathcal{H}_2)$, and $\mathbb{1}_{L(\mathcal{H}_1)}$ takes operators from $L(\mathcal{H}_1)$ to $L(\mathcal{H}_1)$. Thus, $\mathcal{J}(\mathcal{E})$ lies in $L(\mathcal{H}_2 \otimes \mathcal{H}_1)$.

If $\mathcal{H}_1 = \mathcal{H}_2$,

$$\mathcal{J}(\mathcal{E}) = \sum_{i,j} \mathcal{E}(|i\rangle \langle j|) \otimes |i\rangle \langle j|$$

The Choi Theorem

If $\mathcal{E} : \mathcal{H} \rightarrow \mathcal{H}$ is a quantum channel with kraus operators $\{A_i\}$, then,

$$\mathcal{J}(\mathcal{E}) = \sum_i |\text{vec}(A_i)\rangle \langle \text{vec}(A_i)|$$

The Choi Theorem

If $\mathcal{E} : \mathcal{H} \rightarrow \mathcal{H}$ is a quantum channel with kraus operators $\{A_i\}$, then,

$$\mathcal{J}(\mathcal{E}) = \sum_i |\text{vec}(A_i)\rangle \langle \text{vec}(A_i)|$$

Questions

The Choi Theorem

If $\mathcal{E} : \mathcal{H} \rightarrow \mathcal{H}$ is a quantum channel with kraus operators $\{A_i\}$, then,

$$\mathcal{J}(\mathcal{E}) = \sum_i |\text{vec}(A_i)\rangle \langle \text{vec}(A_i)|$$

Questions

- 1 Compute the Choi operator, $\mathcal{J}(\mathcal{E})$ for the completely depolarizing channel on a qubit,

$$\mathcal{E}_\sigma(\rho) := \sigma \quad \forall \rho \in D(\mathcal{H})$$

The Choi Theorem

If $\mathcal{E} : \mathcal{H} \rightarrow \mathcal{H}$ is a quantum channel with kraus operators $\{A_i\}$, then,

$$\mathcal{J}(\mathcal{E}) = \sum_i |\text{vec}(A_i)\rangle \langle \text{vec}(A_i)|$$

Questions

- 1 Compute the Choi operator, $\mathcal{J}(\mathcal{E})$ for the completely depolarizing channel on a qubit,

$$\mathcal{E}_\sigma(\rho) := \sigma \quad \forall \rho \in D(\mathcal{H})$$

- 2 Suppose,

$$\mathcal{J}(\mathcal{E}) = \frac{2}{3} |\phi^-\rangle \langle \phi^-| + \frac{4}{3} |\psi^+\rangle \langle \psi^+|$$

Calculate the Kraus operators for \mathcal{E} .

An isometry maps a quantum system from a smaller to a larger Hilbert space. For instance, the isometry on a single qubit,

$$U|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |3\rangle)$$

$$U|1\rangle = \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle),$$

maps the qubit from the Hilbert space basis $\{|0\rangle, |1\rangle\}$ to the four-dimensional basis $\{|0\rangle, |1\rangle, |2\rangle, |3\rangle\}$

- 1) Write the isometry U as a 4×2 matrix, and show that it $U^\dagger U = \mathbb{I}$, but $UU^\dagger \neq \mathbb{I}$.
- 2) Find a 4×4 unitary matrix, \tilde{U} , such that it gives U , when the last 2 columns are removed. (Hint: You have to show $\tilde{U}^\dagger \tilde{U} = \tilde{U} \tilde{U}^\dagger = \mathbb{I}$.)
- 3) Alternatively, show that there exists a 4×2 matrix V , such that $U = \tilde{U} \cdot V$. Interpret, the role of V .

Appendix: Interpretations of the OpSum Representation

S+E \rightarrow Kraus

- We have, the system coupled to an environment under unitary evolution as

$$U(\rho \otimes |e_0\rangle \langle e_0|)U^*$$

- Trace out the environment,

$$\sum_k (\mathbb{1} \otimes \langle e_k|) U(\rho \otimes |e_0\rangle \langle e_0|) U^* (\mathbb{1} \otimes |e_k\rangle)$$

- See that

$$(\rho \otimes |e_0\rangle \langle e_0|) = (\mathbb{1} \otimes |e_0\rangle) \rho (\mathbb{1} \otimes \langle e_0|)$$

- Define

$$E_k \equiv (\mathbb{1} \otimes \langle e_k|) U (\mathbb{1} \otimes |e_0\rangle)$$

and see the equivalence.

With the notation in the previous slide, define

$$\rho_k \equiv \frac{E_k \rho E_k^*}{\text{tr}(E_k \rho E_k^*)}$$

Thus, we can consider the act of applying \mathcal{E} as applying U to $\rho \otimes |e_0\rangle \langle e_0|$ and then measuring the environment in the $|e_k\rangle$ basis.

That is equivalent to replacing ρ randomly by ρ_k with the probability $p(k) = \text{tr}(E_k \rho E_k^*)$, thus resulting in

$$\mathcal{E}(\rho) = \sum_k p(k) \rho_k = \sum_k E_k \rho E_k^*$$

as expected.

Let P_m be a projective measurement on the $S+E$. Define

$$\mathcal{E}_m(\rho) = \text{tr}_E(P_m U(\rho \otimes \sigma) U^* P_m)$$

where $\sigma = \sum_j q_j |j\rangle \langle j|$. And let $|e_k\rangle$ be a basis for \mathcal{H}^E . Thus,

$$\mathcal{E}_m(\rho) = \sum_{jk} q_j (\mathbb{1} \otimes \langle e_k |) P_m U(\rho \otimes |j\rangle \langle j|) U^* P_m (\mathbb{1} \otimes |e_k\rangle)$$

Define

$$E_{jk}^m = \sqrt{q_j} (\mathbb{1} \otimes \langle e_k |) P_m U(\mathbb{1} \otimes |j\rangle)$$

and we get

$$\mathcal{E}_m(\rho) = \sum_{jk} E_{jk}^m \rho E_{jk}^{m*}$$

It is easy to see that $\sum_{jk} E_{jk}^{m*} E_{jk}^m \preceq \mathbb{1}$, not $= \mathbb{1}$. The state evolves to $\mathcal{E}_m(\rho)$ with probability $\text{tr}(\mathcal{E}_m(\rho))$.

- We are given the set $\{\mathcal{E}_m\}$. We shall construct a S+E (+measurement) model.
- Let E_k^m be the kraus rep. for the operation \mathcal{E}_m .
- Introduce env E with an orthonormal basis $|m, k\rangle$ with indices in 1-1 correspondence.
- Let $|e_k\rangle$ be a basis for \mathcal{H}^E and define

$$U |\psi\rangle \otimes |e_0\rangle \equiv \sum_{mk} E_{mk} |\psi\rangle |m, k\rangle$$

As done earlier, we know this can be extended to a unitary on the composite system.

- Define

$$P_m \equiv \sum_k |m, k\rangle \langle m, k|$$

Let $\rho = \sum p_i |i\rangle \langle i|$ be a state of our system.

Consider

$$\begin{aligned} U(\rho \otimes |e_0\rangle \langle e_0|)U^* &= \sum_j p_j U(|j\rangle \langle j| \otimes |e_0\rangle \langle e_0|)U^* \\ &= \sum_{mjk} p_j (E_{mk} |j\rangle \langle j| E_{mk}^* \otimes |m, k\rangle \langle m, k|) \end{aligned}$$

Now it is easy (and a bit annoying to write) to see that measuring P_m will result in $\mathcal{E}_m(\rho)$ with probability $\text{tr}(\mathcal{E}_m(\rho))$.

Let $\rho = \sum p_i |i\rangle \langle i|$ be a state of our system.

Consider

$$\begin{aligned} U(\rho \otimes |e_0\rangle \langle e_0|)U^* &= \sum_j p_j U(|j\rangle \langle j| \otimes |e_0\rangle \langle e_0|)U^* \\ &= \sum_{mjk} p_j (E_{mk} |j\rangle \langle j| E_{mk}^* \otimes |m, k\rangle \langle m, k|) \end{aligned}$$

Now it is easy (and a bit annoying to write) to see that measuring P_m will result in $\mathcal{E}_m(\rho)$ with probability $\text{tr}(\mathcal{E}_m(\rho))$.