Distance Measures for Quantum Information

Aaryan Gupta and Siddhant Midha

Last updated on August 27, 2022



We have two different measures of similarity or distance between two probability distributions/states $\{p_x\}$ and $\{q_x\}$ over the same index set x. They are -

1 Trace Distance $= D(p_x, q_x) = \frac{1}{2} \sum_x |p_x - q_x|$

We have two different measures of similarity or distance between two probability distributions/states $\{p_x\}$ and $\{q_x\}$ over the same index set x. They are -

• Trace Distance = $D(p_x, q_x) = \frac{1}{2} \sum_x |p_x - q_x|$ Properties:

- **1** Trace Distance $= D(p_x, q_x) = \frac{1}{2} \sum_x |p_x q_x|$ Properties:
 - It is a metric on probability distributions.

- Trace Distance = $D(p_x, q_x) = \frac{1}{2} \sum_x |p_x q_x|$ Properties:
 - It is a metric on probability distributions.
 - $D(p_x, q_x) = \max_S |\sum_{x \in S} p_x \sum_{x \in S} q_x|$ over all subsets S of the index set $\{x\}$, \to relation to distinguishability.
- 2 Fidelity = $F(p_x, q_x) = \sum_x \sqrt{p_x q_x}$

- **1** Trace Distance $= D(p_x, q_x) = \frac{1}{2} \sum_x |p_x q_x|$ Properties:
 - It is a metric on probability distributions.
 - $D(p_x, q_x) = \max_S |\sum_{x \in S} p_x \sum_{x \in S} q_x|$ over all subsets S of the index set $\{x\}$, \to relation to distinguishability.
- ② Fidelity = $F(p_x, q_x) = \sum_x \sqrt{p_x q_x}$ Properties:

- **1** Trace Distance $= D(p_x, q_x) = \frac{1}{2} \sum_x |p_x q_x|$ Properties:
 - It is a metric on probability distributions.
 - $D(p_x, q_x) = \max_S |\sum_{x \in S} p_x \sum_{x \in S} q_x|$ over all subsets S of the index set $\{x\}$, \to relation to distinguishability.
- ② Fidelity = $F(p_x, q_x) = \sum_x \sqrt{p_x q_x}$ Properties:
 - It is not a metric but $\cos^{-1}(F)$ is a metric

- 1 Trace Distance = $D(p_x, q_x) = \frac{1}{2} \sum_x |p_x q_x|$ Properties:
 - It is a metric on probability distributions.
 - $D(p_x, q_x) = \max_S |\sum_{x \in S} p_x \sum_{x \in S} q_x|$ over all subsets S of the index set $\{x\}$, \to relation to distinguishability.
- ② Fidelity = $F(p_x, q_x) = \sum_x \sqrt{p_x q_x}$ Properties:
 - It is not a metric but $\cos^{-1}(F)$ is a metric
 - It can be geometrically interpreted as the inner product between the unit vectors $\sqrt{p_{\rm x}}$ and $\sqrt{q_{\rm x}}$



- 1 Trace Distance = $D(p_x, q_x) = \frac{1}{2} \sum_x |p_x q_x|$ Properties:
 - It is a metric on probability distributions.
 - $D(p_x, q_x) = \max_S |\sum_{x \in S} p_x \sum_{x \in S} q_x|$ over all subsets S of the index set $\{x\}$, \to relation to distinguishability.
- ② Fidelity = $F(p_x, q_x) = \sum_x \sqrt{p_x q_x}$ Properties:
 - It is not a metric but $\cos^{-1}(F)$ is a metric
 - It can be geometrically interpreted as the inner product between the unit vectors $\sqrt{p_{\rm x}}$ and $\sqrt{q_{\rm x}}$



We have two different measures of similarity or distance between two probability distributions/states $\{p_x\}$ and $\{q_x\}$ over the same index set x. They are -

- **1** Trace Distance $= D(p_x, q_x) = \frac{1}{2} \sum_x |p_x q_x|$ Properties:
 - It is a metric on probability distributions.
 - $D(p_x, q_x) = \max_S |\sum_{x \in S} p_x \sum_{x \in S} q_x|$ over all subsets S of the index set $\{x\}$, \to relation to distinguishability.
- ② Fidelity = $F(p_x, q_x) = \sum_x \sqrt{p_x q_x}$ Properties:
 - It is not a metric but $\cos^{-1}(F)$ is a metric
 - It can be geometrically interpreted as the inner product between the unit vectors $\sqrt{p_{\rm x}}$ and $\sqrt{q_{\rm x}}$

Note that these are *static* measures of distance between probability distributions.



Now let us consider a dynamic measure of distance which encapuslates how well information is preserved by a physical process i.e. noise.

Now let us consider a dynamic measure of distance which encapuslates how well information is preserved by a physical process i.e. noise.



Figure 9.2. Given a Markov process $X \to Y$ we may first make a copy of X, \tilde{X} , before subjecting X to the noise which turns it into Y.

Now let us consider a dynamic measure of distance which encapuslates how well information is preserved by a physical process i.e. noise.



Figure 9.2. Given a Markov process $X \to Y$ we may first make a copy of X, \tilde{X} , before subjecting X to the noise which turns it into Y.

Suppose you have the state X and you subject it to a Markov process to get Y. A natural measure would be $P(X \neq Y)$. Now make a perfectly correlated copy \hat{X} of X. It turns out trace has an intimate relation with this dynamic measure of distance which is

Now let us consider a dynamic measure of distance which encapuslates how well information is preserved by a physical process i.e. noise.



Figure 9.2. Given a Markov process $X \to Y$ we may first make a copy of X, \tilde{X} , before subjecting X to the noise which turns it into Y.

Suppose you have the state X and you subject it to a Markov process to get Y. A natural measure would be $P(X \neq Y)$. Now make a perfectly correlated copy \hat{X} of X. It turns out trace has an intimate relation with this dynamic measure of distance which is

$$D(X,\hat{X})=P(X\neq Y)$$

Define the quantum trace distance between two density operators as

Define the quantum trace distance between two density operators as

$$D(\rho,\sigma) = \frac{1}{2} Tr(|\rho - \sigma|)$$

Define the quantum trace distance between two density operators as

$$D(
ho,\sigma) = \frac{1}{2} Tr(|
ho - \sigma|)$$

where $|A| = \sqrt{A^T A}$ (positive square root) We have,

Define the quantum trace distance between two density operators as

$$D(\rho,\sigma) = \frac{1}{2} Tr(|\rho - \sigma|)$$

where $|A| = \sqrt{A^T A}$ (positive square root) We have,

• The tuple $(Dens(\mathcal{H}), D)$ is a metric space.

Define the quantum trace distance between two density operators as

$$D(
ho,\sigma) = \frac{1}{2} Tr(|
ho - \sigma|)$$

where $|A| = \sqrt{A^T A}$ (positive square root) We have.

- The tuple $(Dens(\mathcal{H}), D)$ is a metric space.
- For commuting states,

$$D(\rho,\sigma)=D(\lambda_i,\mu_i)$$

Define the quantum trace distance between two density operators as

$$D(
ho,\sigma) = \frac{1}{2} Tr(|
ho - \sigma|)$$

where $|A| = \sqrt{A^T A}$ (positive square root) We have.

- The tuple $(Dens(\mathcal{H}), D)$ is a metric space.
- For commuting states,

$$D(\rho,\sigma)=D(\lambda_i,\mu_i)$$

Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$.

Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$. Then,

$$D(\rho,\sigma)=\frac{|\vec{r}-\vec{s}|}{2}$$

Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$. Then,

$$D(\rho,\sigma)=\frac{|\vec{r}-\vec{s}|}{2}$$

Converts to euclidean distance.

Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$. Then,

$$D(\rho,\sigma)=\frac{|\vec{r}-\vec{s}|}{2}$$

- Converts to euclidean distance.
- Hints towards rotation invariance.

Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$. Then,

$$D(\rho,\sigma)=\frac{|\vec{r}-\vec{s}|}{2}$$

- Converts to euclidean distance.
- Hints towards rotation invariance.
- Helpful visualization.

Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$. Then,

$$D(\rho,\sigma)=\frac{|\vec{r}-\vec{s}|}{2}$$

- Converts to euclidean distance.
- Hints towards rotation invariance.
- Helpful visualization.

Relation to distinguishability via measurement.

Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$. Then,

$$D(\rho,\sigma) = \frac{|\vec{r} - \vec{s}|}{2}$$

- Converts to euclidean distance.
- 4 Hints towards rotation invariance.
- Helpful visualization.

Relation to distinguishability via measurement.

$$D(\rho, \sigma) = \max_{P:P \leq I} tr(P(\rho - \sigma))$$

Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$. Then,

$$D(\rho,\sigma) = \frac{|\vec{r} - \vec{s}|}{2}$$

- Converts to euclidean distance.
- Hints towards rotation invariance.
- Helpful visualization.

Relation to distinguishability via measurement.

$$D(\rho,\sigma) = \max_{P:P \leq I} tr(P(\rho - \sigma))$$

Proof.



Bloch sphere. Let $\rho \equiv \vec{r}, \sigma \equiv \vec{s}$. Then,

$$D(\rho,\sigma)=\frac{|\vec{r}-\vec{s}|}{2}$$

- Converts to euclidean distance.
- Hints towards rotation invariance.
- Helpful visualization.

Relation to distinguishability via measurement.

$$D(\rho,\sigma) = \max_{P:P \leq I} tr(P(\rho - \sigma))$$

Proof. Key step, $\rho - \sigma \rightarrow Q - S$.



Quantum Trace distance as an upper bound

Quantum Trace distance as an upper bound

Let $\{E_m\}$ be any POVM, and let $p_m = Tr(\rho E_m)$ and $q_m = Tr(\sigma E_m)$.

Quantum Trace distance as an upper bound

Let $\{E_m\}$ be any POVM, and let $p_m = Tr(\rho E_m)$ and $q_m = Tr(\sigma E_m)$. Then the following holds,

$$D(\rho,\sigma) = max_{E_m}D(p_m,q_m)$$

Quantum Trace distance as an upper bound

Let $\{E_m\}$ be any POVM, and let $p_m = Tr(\rho E_m)$ and $q_m = Tr(\sigma E_m)$. Then the following holds,

$$D(\rho,\sigma) = max_{E_m}D(p_m,q_m)$$

Proof - Show the inequality, and then show existence.

Contractiveness

Quantum Trace distance as an upper bound

Let $\{E_m\}$ be any POVM, and let $p_m = Tr(\rho E_m)$ and $q_m = Tr(\sigma E_m)$. Then the following holds,

$$D(\rho,\sigma) = max_{E_m}D(p_m,q_m)$$

Proof - Show the inequality, and then show existence.

Contractiveness

Suppose \mathcal{E} is a TP map.

Quantum Trace distance as an upper bound

Let $\{E_m\}$ be any POVM, and let $p_m = Tr(\rho E_m)$ and $q_m = Tr(\sigma E_m)$. Then the following holds,

$$D(\rho,\sigma) = max_{E_m}D(p_m,q_m)$$

Proof - Show the inequality, and then show existence.

Contractiveness

Suppose ${\mathcal E}$ is a TP map. Then the following holds

More properties

Quantum Trace distance as an upper bound

Let $\{E_m\}$ be any POVM, and let $p_m = Tr(\rho E_m)$ and $q_m = Tr(\sigma E_m)$. Then the following holds,

$$D(\rho,\sigma) = max_{E_m}D(p_m,q_m)$$

Proof - Show the inequality, and then show existence.

Contractiveness

Suppose ${\mathcal E}$ is a TP map. Then the following holds

$$D(\mathcal{E}(\rho), \mathcal{E}(\sigma)) \leq D(\rho, \sigma) \forall \rho, \sigma$$

More properties

Quantum Trace distance as an upper bound

Let $\{E_m\}$ be any POVM, and let $p_m = Tr(\rho E_m)$ and $q_m = Tr(\sigma E_m)$. Then the following holds,

$$D(\rho,\sigma) = max_{E_m}D(p_m,q_m)$$

Proof - Show the inequality, and then show existence.

Contractiveness

Suppose $\ensuremath{\mathcal{E}}$ is a TP map. Then the following holds

$$D(\mathcal{E}(\rho), \mathcal{E}(\sigma)) \leq D(\rho, \sigma) \forall \rho, \sigma$$

Proof - Along similar lines, use the previous properties.

More properties

Quantum Trace distance as an upper bound

Let $\{E_m\}$ be any POVM, and let $p_m = Tr(\rho E_m)$ and $q_m = Tr(\sigma E_m)$. Then the following holds,

$$D(\rho,\sigma) = max_{E_m}D(p_m,q_m)$$

Proof - Show the inequality, and then show existence.

Contractiveness

Suppose ${\mathcal E}$ is a TP map. Then the following holds

$$D(\mathcal{E}(\rho), \mathcal{E}(\sigma)) \leq D(\rho, \sigma) \forall \rho, \sigma$$

Proof - Along similar lines, use the previous properties. Corollary:

$$D(\rho_A, \sigma_A) \leq D(\rho_{AB}, \sigma_{AB})$$



• Define the quantum fidelity as

$$F(
ho,\sigma) = Tr\sqrt{
ho^{1/2}\sigma
ho^{1/2}}$$

• Define the quantum fidelity as

$$F(\rho,\sigma) = Tr\sqrt{\rho^{1/2}\sigma\rho^{1/2}}$$

Observe,

Define the quantum fidelity as

$$F(\rho,\sigma) = Tr\sqrt{\rho^{1/2}\sigma\rho^{1/2}}$$

- Observe,
 - **1** If ρ and σ commute, then

$$F(\rho,\sigma)=F(\lambda_i,\mu_i)$$

Define the quantum fidelity as

$$F(\rho,\sigma) = Tr\sqrt{\rho^{1/2}\sigma\rho^{1/2}}$$

- Observe,
 - **1** If ρ and σ commute, then

$$F(\rho,\sigma)=F(\lambda_i,\mu_i)$$

$$P(|\psi\rangle, \rho) = \sqrt{\langle \psi | \rho | \psi \rangle}$$

Define the quantum fidelity as

$$F(\rho,\sigma) = Tr\sqrt{\rho^{1/2}\sigma\rho^{1/2}}$$

- Observe,
 - **1** If ρ and σ commute, then

$$F(\rho,\sigma)=F(\lambda_i,\mu_i)$$

- $P(|\psi\rangle, \rho) = \sqrt{\langle \psi | \rho | \psi \rangle}$

Theorem (Uhlmann's Theorem)

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q.

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q. Then,

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q. Then,

$$F(
ho, \sigma) = \max_{|\psi\rangle, |\varphi\rangle} |\langle \psi | \varphi\rangle |$$

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q. Then,

$$F(
ho, \sigma) = \max_{|\psi
angle, |arphi
angle} |\langle\psi| arphi
angle|$$

where $|\psi\rangle$ and $|\varphi\rangle$ are purifications of ρ and σ in $\mathcal{H}_{Q}\otimes\mathcal{H}_{R}$.

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q. Then,

$$F(
ho, \sigma) = \max_{|\psi\rangle, |\varphi\rangle} |\langle \psi | \varphi\rangle |$$

where $|\psi\rangle$ and $|\varphi\rangle$ are purifications of ρ and σ in $\mathcal{H}_Q\otimes\mathcal{H}_R$.

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q. Then,

$$F(
ho, \sigma) = \max_{|\psi\rangle, |\varphi\rangle} |\langle \psi | \varphi\rangle |$$

where $|\psi\rangle$ and $|\varphi\rangle$ are purifications of ρ and σ in $\mathcal{H}_Q\otimes\mathcal{H}_R$.

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q. Then,

$$F(
ho,\sigma) = \max_{|\psi
angle, |arphi
angle} |\langle\psi\midarphi
angle|$$

where $|\psi\rangle$ and $|\varphi\rangle$ are purifications of ρ and σ in $\mathcal{H}_Q\otimes\mathcal{H}_R$.

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q. Then,

$$F(
ho, \sigma) = \max_{|\psi\rangle, |\varphi\rangle} |\langle \psi \mid \varphi\rangle|$$

where $|\psi\rangle$ and $|\varphi\rangle$ are purifications of ρ and σ in $\mathcal{H}_Q\otimes\mathcal{H}_R$.

- $F(\mathcal{E}(\rho), \mathcal{E}(\sigma)) \geq F(\rho, \sigma)$

Theorem (Uhlmann's Theorem)

Suppose $\rho, \sigma \in Dens(\mathcal{H}_Q)$. Let R be a copy of Q. Then,

$$F(
ho, \sigma) = \max_{|\psi\rangle, |\varphi\rangle} |\langle \psi \mid \varphi\rangle|$$

where $|\psi\rangle$ and $|\varphi\rangle$ are purifications of ρ and σ in $\mathcal{H}_Q\otimes\mathcal{H}_R$.

- $F(\mathcal{E}(\rho), \mathcal{E}(\sigma)) \geq F(\rho, \sigma)$
- **3** $A(\rho, \sigma) := \arccos F(\rho, \sigma)$ is a metric on $Dens(\mathcal{H})$



Quantum trace distance and fidelity are qualitatively equivalent measures of distance. In case of pure states, they are exactly equal.

Quantum trace distance and fidelity are qualitatively equivalent measures of distance. In case of pure states, they are exactly equal.

Proof. Let the two pure states be $|a\rangle$ and $|b\rangle$. Let $|a\rangle = |0\rangle$ and $|b\rangle = cos(\theta) |0\rangle + sin(\theta) |1\rangle$. For these, $F(a,b) = |cos(\theta)|$ and $D(a,b) = |sin(\theta)| = \sqrt{1 - F(a,b)^2}$.

Quantum trace distance and fidelity are qualitatively equivalent measures of distance. In case of pure states, they are exactly equal.

Proof. Let the two pure states be $|a\rangle$ and $|b\rangle$. Let $|a\rangle = |0\rangle$ and $|b\rangle = cos(\theta) |0\rangle + sin(\theta) |1\rangle$. For these, $F(a,b) = |cos(\theta)|$ and $D(a,b) = |sin(\theta)| = \sqrt{1 - F(a,b)^2}$.

Theorem

$$1 - F(\rho, \sigma) \le D(\rho, \sigma) \le \sqrt{1 - F(\rho, \sigma)^2}$$
 for any two states ρ, σ .

Quantum trace distance and fidelity are qualitatively equivalent measures of distance. In case of pure states, they are exactly equal.

Proof. Let the two pure states be $|a\rangle$ and $|b\rangle$. Let $|a\rangle = |0\rangle$ and $|b\rangle = cos(\theta) |0\rangle + sin(\theta) |1\rangle$. For these, $F(a,b) = |cos(\theta)|$ and $D(a,b) = |sin(\theta)| = \sqrt{1 - F(a,b)^2}$.

Theorem,

$$1 - F(\rho, \sigma) \le D(\rho, \sigma) \le \sqrt{1 - F(\rho, \sigma)^2}$$
 for any two states ρ, σ .

Proof. To prove the right side of the inequality, consider purifications $|\psi\rangle$, $|\phi\rangle$ chosen such that $F(\rho, \sigma) = \langle \psi | \phi \rangle = F(\psi, \phi)$.

Quantum trace distance and fidelity are qualitatively equivalent measures of distance. In case of pure states, they are exactly equal.

Proof. Let the two pure states be $|a\rangle$ and $|b\rangle$. Let $|a\rangle = |0\rangle$ and $|b\rangle = cos(\theta) |0\rangle + sin(\theta) |1\rangle$. For these, $F(a,b) = |cos(\theta)|$ and $D(a,b) = |sin(\theta)| = \sqrt{1 - F(a,b)^2}$.

Theorem

$$1 - F(\rho, \sigma) \le D(\rho, \sigma) \le \sqrt{1 - F(\rho, \sigma)^2}$$
 for any two states ρ, σ .

Proof. To prove the right side of the inequality, consider purifications $|\psi\rangle\,, |\phi\rangle$ chosen such that $F(\rho,\sigma)=\langle\psi|\phi\rangle=F(\psi,\phi)$. Since $|\psi\rangle\,, |\phi\rangle$ are pure states and tace distance is non- increasing under partial trace

$$D(\rho, \sigma) \le D(\psi, \phi) \le \sqrt{1 - F(\rho, \sigma)^2}$$



Quantum trace distance and fidelity are qualitatively equivalent measures of distance. In case of pure states, they are exactly equal.

Proof. Let the two pure states be $|a\rangle$ and $|b\rangle$. Let $|a\rangle = |0\rangle$ and $|b\rangle = cos(\theta) |0\rangle + sin(\theta) |1\rangle$. For these, $F(a,b) = |cos(\theta)|$ and $D(a,b) = |sin(\theta)| = \sqrt{1 - F(a,b)^2}$.

Theorem

$$1 - F(\rho, \sigma) \le D(\rho, \sigma) \le \sqrt{1 - F(\rho, \sigma)^2}$$
 for any two states ρ, σ .

Proof. To prove the right side of the inequality, consider purifications $|\psi\rangle\,, |\phi\rangle$ chosen such that $F(\rho,\sigma)=\langle\psi|\phi\rangle=F(\psi,\phi)$. Since $|\psi\rangle\,, |\phi\rangle$ are pure states and tace distance is non- increasing under partial trace

$$D(\rho, \sigma) \le D(\psi, \phi) \le \sqrt{1 - F(\rho, \sigma)^2}$$

The left side of the inequality can be proved using POVMs and simple mathematical manipulation.

Consider the system to be in $|\psi\rangle$ initially.

¹pure states are enough, because $F_{min} \leq F(\rho, \mathcal{E}(\rho)) \forall \rho$ can be shown

Consider the system to be in $|\psi\rangle$ initially. Suppose it undergoes evolution under \mathcal{E} .

¹pure states are enough, because $F_{min} \leq F(\rho, \mathcal{E}(\rho)) \forall \rho$ can be shown $\bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet$

Consider the system to be in $|\psi\rangle$ initially. Suppose it undergoes evolution under $\mathcal{E}.$ The information preserved can be estimated by

$$F(|\psi\rangle, \mathcal{E}(|\psi\rangle\langle\psi|)$$

¹pure states are enough, because $F_{min} \leq F(\rho, \mathcal{E}(\rho)) \forall \rho$ can be shown $\bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet$

Consider the system to be in $|\psi\rangle$ initially. Suppose it undergoes evolution under $\mathcal{E}.$ The information preserved can be estimated by

$$F(\ket{\psi}, \mathcal{E}(\ket{\psi} \bra{\psi}))$$

If we talk about the channel alone, we can minimize over all states ¹

$$F_{min} = min_{|\psi\rangle}F(\ket{\psi},\mathcal{E}(\ket{\psi}ra{\psi})$$

Consider the system to be in $|\psi\rangle$ initially. Suppose it undergoes evolution under \mathcal{E} . The information preserved can be estimated by

$$F(\ket{\psi}, \mathcal{E}(\ket{\psi} \bra{\psi}))$$

If we talk about the channel alone, we can minimize over all states ¹

$$F_{ extit{min}} = extit{min}_{|\psi
angle} F(\ket{\psi}, \mathcal{E}(\ket{\psi}ra{\psi})$$

Illustrations

- Depolarizing Channel
- Phase Damping Channel
- Gate Fidelities

¹pure states are enough, because $F_{min} \leq F(\rho, \mathcal{E}(\rho)) \forall \rho$ can be shown $\bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet$

Ensemble Oriented

A QIS is an entity which produces states ρ_i with probability p_i for all $i \in I$.

Ensemble Oriented

A QIS is an entity which produces states ρ_i with probability p_i for all $i \in I$.

With this and the previous slide, we can talk about how well a source is preserved under a channel with the following quantity

$$ar{F} := \sum_j p_j F(
ho_j, \mathcal{E}(
ho_j))^2$$

Ensemble Oriented

A QIS is an entity which produces states ρ_i with probability p_i for all $i \in I$.

With this and the previous slide, we can talk about how well a source is preserved under a channel with the following quantity

$$ar{F} := \sum_j p_j F(
ho_j, \mathcal{E}(
ho_j))^2$$

This is called the *ensemble average fidelity*. Provided $\bar{F}\approx 1$ one can be confident that the source is being preserved by the channel.

This notion is inspired from the idea of converting dynamic distance to correlation between a system and its copy. The correlation here translates to entanglement – A channel which preserves information well is one which preserves entanglement well.

This notion is inspired from the idea of converting dynamic distance to correlation between a system and its copy. The correlation here translates to entanglement – A channel which preserves information well is one which preserves entanglement well.

Entanglement Oriented

A QIS (Q,R) is a system Q in some state which is entangled to some environment R. WLOG it is in the state $\rho = Tr_R(|RQ\rangle \langle RQ|)$.

This notion is inspired from the idea of converting dynamic distance to correlation between a system and its copy. The correlation here translates to entanglement – A channel which preserves information well is one which preserves entanglement well.

Entanglement Oriented

A QIS (Q,R) is a system Q in some state which is entangled to some environment R. WLOG it is in the state $\rho = Tr_R(|RQ\rangle \langle RQ|)$.

We define entanglement fidelity,

This notion is inspired from the idea of converting dynamic distance to correlation between a system and its copy. The correlation here translates to entanglement – A channel which preserves information well is one which preserves entanglement well.

Entanglement Oriented

A QIS (Q,R) is a system Q in some state which is entangled to some environment R. WLOG it is in the state $\rho = Tr_R(|RQ\rangle \langle RQ|)$.

We define entanglement fidelity,

$$F_{e}(\rho, \mathcal{E}) := F(|RQ\rangle, |RQ'\rangle)$$

= $\langle RQ| [(I_{R} \otimes \mathcal{E})(|RQ\rangle \langle RQ|)] |RQ\rangle$

where $\rho = Tr_R(|RQ\rangle \langle RQ|)$. It can be shown that only the choice of ρ and \mathcal{E} affect the EF, not the choice of purification.

• There exists a nice formula to compute the F_e .

① There exists a nice formula to compute the F_e . Let E_i be kraus elements of $\mathcal{E} \otimes I_R$

• There exists a nice formula to compute the F_e . Let E_i be kraus elements of $\mathcal{E} \otimes I_R$

$$F_{e}(\mathcal{E}, \rho) = \langle RQ | \rho_{RQ'} | RQ \rangle$$

$$= \sum_{i} |\langle RQ | E_{i} | RQ \rangle|^{2}$$

$$= \sum_{i} |tr(\rho E_{i})|^{2}$$

① There exists a nice formula to compute the F_e . Let E_i be kraus elements of $\mathcal{E} \otimes I_R$

$$F_{e}(\mathcal{E}, \rho) = \langle RQ | \rho_{RQ'} | RQ \rangle$$

$$= \sum_{i} |\langle RQ | E_{i} | RQ \rangle|^{2}$$

$$= \sum_{i} |tr(\rho E_{i})|^{2}$$

② $F_e(\rho,\mathcal{E}) \leq (F(\rho,\mathcal{E}(\rho))^2$. Intuitively, attempt (2) is stronger than attempt (1). It is tougher to preserve the entanglement and the state than just the state.

① There exists a nice formula to compute the F_e . Let E_i be kraus elements of $\mathcal{E} \otimes I_R$

$$F_{e}(\mathcal{E}, \rho) = \langle RQ | \rho_{RQ'} | RQ \rangle$$

$$= \sum_{i} |\langle RQ | E_{i} | RQ \rangle|^{2}$$

$$= \sum_{i} |tr(\rho E_{i})|^{2}$$

- ② $F_e(\rho, \mathcal{E}) \leq (F(\rho, \mathcal{E}(\rho))^2$. Intuitively, attempt (2) is stronger than attempt (1). It is tougher to preserve the entanglement and the state than just the state.
- F_e is convex. Now, we have,

① There exists a nice formula to compute the F_e . Let E_i be kraus elements of $\mathcal{E} \otimes I_R$

$$F_{e}(\mathcal{E}, \rho) = \langle RQ | \rho_{RQ'} | RQ \rangle$$

$$= \sum_{i} |\langle RQ | E_{i} | RQ \rangle|^{2}$$

$$= \sum_{i} |tr(\rho E_{i})|^{2}$$

- ② $F_e(\rho, \mathcal{E}) \leq (F(\rho, \mathcal{E}(\rho))^2$. Intuitively, attempt (2) is stronger than attempt (1). It is tougher to preserve the entanglement and the state than just the state.
- F_e is convex. Now, we have,

$$F_{e}(\sum_{j} p_{j}\rho_{j}, \mathcal{E}) \leq \sum_{j} p_{j}F_{e}(\rho_{j}, \mathcal{E}) \leq \sum_{j} p_{j}F(\rho_{j}, \mathcal{E}(\rho_{j}))^{2}$$



① There exists a nice formula to compute the F_e . Let E_i be kraus elements of $\mathcal{E} \otimes I_R$

$$F_{e}(\mathcal{E}, \rho) = \langle RQ | \rho_{RQ'} | RQ \rangle$$

$$= \sum_{i} |\langle RQ | E_{i} | RQ \rangle|^{2}$$

$$= \sum_{i} |tr(\rho E_{i})|^{2}$$

- ② $F_e(\rho, \mathcal{E}) \leq (F(\rho, \mathcal{E}(\rho))^2$. Intuitively, attempt (2) is stronger than attempt (1). It is tougher to preserve the entanglement and the state than just the state.
- F_e is convex. Now, we have,

$$F_{e}(\sum_{j} p_{j}\rho_{j}, \mathcal{E}) \leq \sum_{j} p_{j}F_{e}(\rho_{j}, \mathcal{E}) \leq \sum_{j} p_{j}F(\rho_{j}, \mathcal{E}(\rho_{j}))^{2}$$

Thus, $F_e < \bar{F}!$



Concluding Remarks

Thus, any quantum channel $\mathcal E$ which does a good job of preserving the entanglement between a source described by a density operator and a reference system will automatically do a good job of preserving an ensemble source described by probabilities p_j and states ρ_j such that $\rho = \sum_j p_j \rho_j$. In this sense the notion of a quantum source based on entanglement fidelity is a more stringent notion than the ensemble definition.