

AC Quantum Transport: Formalisms and Applications

Siddhant Midha

Department of Electrical Engineering
Indian Institute of Technology Bombay

EE 755 Course Project
May 4, 2023



This talk

This talk

We will talk about:

This talk

We will talk about:

- ➊ Green's what – what?

This talk

We will talk about:

- ① Green's what – what?
- ② NEGF \rightarrow **AC** NEGF.

This talk

We will talk about:

- ① Green's what – what?
- ② NEGF \rightarrow **AC** NEGF.
- ③ **AC** Scattering Matrix. Why?

This talk

We will talk about:

- ① Green's what – what?
- ② NEGF \rightarrow **AC** NEGF.
- ③ **AC** Scattering Matrix. Why?
- ④ Numerics: Simulation results. Cool stuff.

This talk

We will talk about:

- ① Green's what – what?
- ② NEGF \rightarrow **AC** NEGF.
- ③ **AC** Scattering Matrix. Why?
- ④ Numerics: Simulation results. Cool stuff.

Quite technical, so focus on gathering intuition! The paper is there for the details :)

NEGF – Green's Functions; *Quickly*

NEGF – Green's Functions; *Quickly*

- Hamiltonian $\hat{\mathbf{H}} = \sum_{n,m} \mathbf{H}_{nm} c_n^\dagger c_m \rightarrow \mathcal{G}(E) = (E + i\eta - \mathbf{H})^{-1}$

NEGF – Green's Functions; *Quickly*

- Hamiltonian $\hat{\mathbf{H}} = \sum_{n,m} \mathbf{H}_{nm} c_n^\dagger c_m \rightarrow \mathcal{G}(E) = (E + i\eta - \mathbf{H})^{-1}$
- Integrate out electrodes:

NEGF – Green's Functions; *Quickly*

- Hamiltonian $\hat{\mathbf{H}} = \sum_{n,m} \mathbf{H}_{nm} c_n^\dagger c_m \rightarrow \mathcal{G}(E) = (E + i\eta - \mathbf{H})^{-1}$
- Integrate out electrodes: $\mathcal{G}_{\bar{0}\bar{0}}(E) = (E + i\eta - H - \sum_{m=1} \Sigma'(m; E))^{-1}$

NEGF – Green's Functions; *Quickly*

- Hamiltonian $\hat{\mathbf{H}} = \sum_{n,m} \mathbf{H}_{nm} c_n^\dagger c_m \rightarrow \mathcal{G}(E) = (E + i\eta - \mathbf{H})^{-1}$
- Integrate out electrodes: $\mathcal{G}_{\bar{0}\bar{0}}(E) = (E + i\eta - H - \sum_{m=1} \Sigma'(m; E))^{-1}$
- \rightarrow Get observables!

NEGF – Green's Functions; *Quickly*

- Hamiltonian $\hat{\mathbf{H}} = \sum_{n,m} \mathbf{H}_{nm} c_n^\dagger c_m \rightarrow \mathcal{G}(E) = (E + i\eta - \mathbf{H})^{-1}$
- Integrate out electrodes: $\mathcal{G}_{\bar{0}\bar{0}}(E) = (E + i\eta - H - \sum_{m=1} \Sigma'(m; E))^{-1}$
- → Get observables! $I_m = \frac{e}{h} \int dE \sum_{m'=1}^N (f_m - f_{m'}) \text{Tr}[\mathcal{G}_0 \Gamma_{m'} \mathcal{G}_0^\dagger \Gamma_m]$

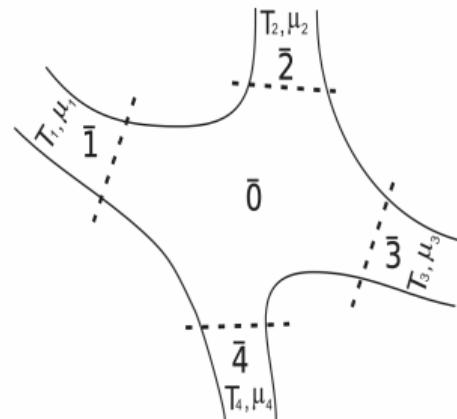


Figure 1: System under study

AC NEGF I [SW13]

- Move into time-domain!

AC NEGF I [SW13]

- Move into time-domain!

$$\mathfrak{G}_{nm}^r(t, t') = -\frac{\iota}{\hbar}\theta(t - t')\langle\{c_n(t), c_m^\dagger(t')\}\rangle, \quad (1)$$

$$\mathfrak{G}_{nm}^<(t, t') = \frac{\iota}{\hbar}\langle c_n^\dagger(t)c_m(t')\rangle \quad (2)$$

AC NEGF I [SW13]

- Move into time-domain!

$$\mathfrak{G}_{nm}^r(t, t') = -\frac{\iota}{\hbar}\theta(t - t')\langle\{c_n(t), c_m^\dagger(t')\}\rangle, \quad (1)$$

$$\mathfrak{G}_{nm}^<(t, t') = \frac{\iota}{\hbar}\langle c_n^\dagger(t)c_m(t')\rangle \quad (2)$$

- AC perturbation: $\mathbf{H} = \mathcal{H} + \mathcal{V}(t)$

AC NEGF I [SW13]

- Move into time-domain!

$$\mathfrak{G}_{nm}^r(t, t') = -\frac{\iota}{\hbar} \theta(t - t') \langle \{c_n(t), c_m^\dagger(t')\} \rangle, \quad (1)$$

$$\mathfrak{G}_{nm}^<(t, t') = \frac{\iota}{\hbar} \langle c_n^\dagger(t) c_m(t') \rangle \quad (2)$$

- AC perturbation: $\mathbf{H} = \mathcal{H} + \mathcal{V}(t)$
- Let \mathfrak{g}^r and $\mathfrak{g}^<$ be the unperturbed Green's functions, then Dyson's equations tell us:

AC NEGF I [SW13]

- Move into **time-domain!**

$$\mathfrak{G}_{nm}^r(t, t') = -\frac{\iota}{\hbar} \theta(t - t') \langle \{c_n(t), c_m^\dagger(t')\} \rangle, \quad (1)$$

$$\mathfrak{G}_{nm}^<(t, t') = \frac{\iota}{\hbar} \langle c_n^\dagger(t) c_m(t') \rangle \quad (2)$$

- AC perturbation: $\mathbf{H} = \mathcal{H} + \mathcal{V}(t)$
- Let \mathfrak{g}^r and $\mathfrak{g}^<$ be the unperturbed Green's functions, then Dyson's equations tell us:

$$\mathfrak{G}^r = \mathfrak{g}^r + \mathfrak{g}^r * \mathcal{V} * \mathfrak{G}^r \quad (3)$$

$$\mathfrak{G}^< = \mathfrak{g}^< + \mathfrak{g}^r * \mathcal{V} * \mathfrak{G}^< + \mathfrak{g}^< * \mathcal{V} * \mathfrak{G}^a \quad (4)$$

AC NEGF I [SW13]

- Move into **time-domain!**

$$\mathfrak{G}_{nm}^r(t, t') = -\frac{\iota}{\hbar}\theta(t - t')\langle\{c_n(t), c_m^\dagger(t')\}\rangle, \quad (1)$$

$$\mathfrak{G}_{nm}^{<}(t, t') = \frac{\iota}{\hbar}\langle c_n^\dagger(t)c_m(t')\rangle \quad (2)$$

- AC perturbation: $\mathbf{H} = \mathcal{H} + \mathcal{V}(t)$
- Let \mathfrak{g}^r and $\mathfrak{g}^{<}$ be the unperturbed Green's functions, then Dyson's equations tell us:

$$\mathfrak{G}^r = \mathfrak{g}^r + \mathfrak{g}^r * \mathcal{V} * \mathfrak{G}^r \quad (3)$$

$$\mathfrak{G}^{<} = \mathfrak{g}^{<} + \mathfrak{g}^r * \mathcal{V} * \mathfrak{G}^{<} + \mathfrak{g}^{<} * \mathcal{V} * \mathfrak{G}^a \quad (4)$$

- Currents! $I_m(t) = e \sum_{i \in \sigma_0, \alpha \in \sigma_m} (\mathcal{V}_{\alpha i}(t)\mathfrak{G}_{i\alpha}^{<}(t, t) - \mathcal{V}_{i\alpha}(t)\mathfrak{G}_{\alpha i}^{<}(t, t))$

AC NEGF I [SW13]

- Two types considered in [SW13]

AC NEGF I [SW13]

- Two types considered in [SW13]
 - ➊ External AC Perturbation: Uniform across the lead lattice points: $\mathbf{W} = eV_{ac}\mathbf{I}_{\bar{m}}$

AC NEGF I [SW13]

- Two types considered in [SW13]
 - ① *External AC Perturbation:* Uniform across the lead lattice points: $\mathbf{W} = eV_{ac}\mathbf{I}_{\bar{m}}$
 - ② *Arbitrary internal AC perturbation:* Arbitrary matrix, single oscillating function, given as,
$$\mathbf{W} = eV_{ac}\mathbf{W}_{\bar{0}\bar{0}}$$

AC NEGF I [SW13]

- Two types considered in [SW13]
 - ① External AC Perturbation: Uniform across the lead lattice points: $\mathbf{W} = eV_{ac}\mathbf{I}_{\bar{m}}$
 - ② Arbitrary internal AC perturbation: Arbitrary matrix, single oscillating function, given as, $\mathbf{W} = eV_{ac}\mathbf{W}_{\bar{0}\bar{0}}$
- Results!

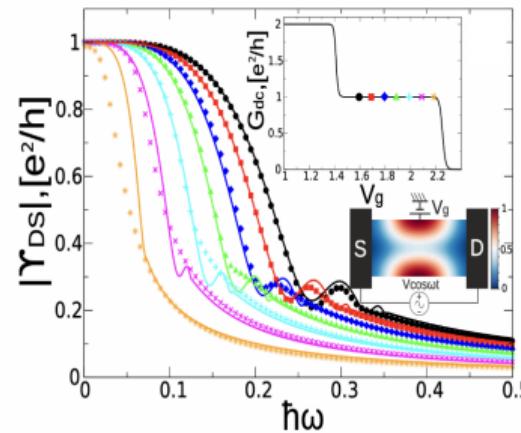


Figure 2: QPC AC conductance

AC NEGF II [PG18]

- AC Bias: $V(t) = V_{AC} \cos \omega t$ applied *to the leads*. First order GFs:

AC NEGF II [PG18]

- AC Bias: $V(t) = V_{AC} \cos \omega t$ applied *to the leads*. First order GFs:

$$G^r(E) = G_0^r(E) = g_\omega^r(E); \quad \Sigma^\gamma(E) = \Sigma_0^\gamma(E) + \sigma_\omega^\gamma(E) \quad (\gamma \in \{r, <\})$$

AC NEGF II [PG18]

- AC Bias: $V(t) = V_{AC} \cos \omega t$ applied *to the leads*. First order GFs:

$$G^r(E) = G_0^r(E) = g_\omega^r(E); \quad \Sigma^\gamma(E) = \Sigma_0^\gamma(E) + \sigma_\omega^\gamma(E) \quad (\gamma \in \{r, <\}) \quad (5)$$

- Correction terms:

AC NEGF II [PG18]

- AC Bias: $V(t) = V_{AC} \cos \omega t$ applied *to the leads*. First order GFs:

$$G^r(E) = G_0^r(E) = g_\omega^r(E); \quad \Sigma^\gamma(E) = \Sigma_0^\gamma(E) + \sigma_\omega^\gamma(E) \quad (\gamma \in \{r, <\}) \quad (5)$$

- Correction terms:

$$\sigma_\omega^\gamma(E) = \frac{eV_{AC}}{\hbar\omega} [\Sigma_0^\gamma(E) - \Sigma_0^\gamma(E_+)] \quad (6)$$

$$g_\omega^r = G_0^r(E_+) [U_\omega + \sigma_\omega^r(E)] G_0^r(E) \quad (7)$$

AC NEGF II [PG18]

- AC Bias: $V(t) = V_{AC} \cos \omega t$ applied *to the leads*. First order GFs:

$$G^r(E) = G_0^r(E) = g_\omega^r(E); \quad \Sigma^\gamma(E) = \Sigma_0^\gamma(E) + \sigma_\omega^\gamma(E) \quad (\gamma \in \{r, <\}) \quad (5)$$

- Correction terms:

$$\sigma_\omega^\gamma(E) = \frac{eV_{AC}}{\hbar\omega} [\Sigma_0^\gamma(E) - \Sigma_0^\gamma(E_+)] \quad (6)$$

$$g_\omega^r = G_0^r(E_+) [U_\omega + \sigma_\omega^r(E)] G_0^r(E) \quad (7)$$

where $E_+ = E + \hbar\omega$.

AC NEGF II [PG18]

- AC Bias: $V(t) = V_{AC} \cos \omega t$ applied *to the leads*. First order GFs:

$$G^r(E) = G_0^r(E) = g_\omega^r(E); \quad \Sigma^\gamma(E) = \Sigma_0^\gamma(E) + \sigma_\omega^\gamma(E) \quad (\gamma \in \{r, <\}) \quad (5)$$

- Correction terms:

$$\sigma_\omega^\gamma(E) = \frac{eV_{AC}}{\hbar\omega} [\Sigma_0^\gamma(E) - \Sigma_0^\gamma(E_+)] \quad (6)$$

$$g_\omega^r = G_0^r(E_+) [U_\omega + \sigma_\omega^r(E)] G_0^r(E) \quad (7)$$

where $E_+ = E + \hbar\omega$.

- (First order) Lesser GF:

AC NEGF II [PG18]

- AC Bias: $V(t) = V_{AC} \cos \omega t$ applied *to the leads*. First order GFs:

$$G^r(E) = G_0^r(E) = g_\omega^r(E); \quad \Sigma^\gamma(E) = \Sigma_0^\gamma(E) + \sigma_\omega^\gamma(E) \quad (\gamma \in \{r, <\}) \quad (5)$$

- Correction terms:

$$\sigma_\omega^\gamma(E) = \frac{eV_{AC}}{\hbar\omega} [\Sigma_0^\gamma(E) - \Sigma_0^\gamma(E_+)] \quad (6)$$

$$g_\omega^r = G_0^r(E_+) [U_\omega + \sigma_\omega^r(E)] G_0^r(E) \quad (7)$$

where $E_+ = E + \hbar\omega$.

- (First order) Lesser GF:

$$g_\omega^<(E) = G_0^r(E_+) \Sigma_0^<(E_+) g_\omega^r(E)^\dagger + G_0^r(E_+) \sigma_\omega^<(E) G_0^r(E)^\dagger + g_\omega^r \Sigma_0^<(E) G_0^r(E)^\dagger \quad (8)$$

AC NEGF II [PG18]

- AC Bias: $V(t) = V_{AC} \cos \omega t$ applied *to the leads*. First order GFs:

$$G^r(E) = G_0^r(E) = g_\omega^r(E); \quad \Sigma^\gamma(E) = \Sigma_0^\gamma(E) + \sigma_\omega^\gamma(E) \quad (\gamma \in \{r, <\}) \quad (5)$$

- Correction terms:

$$\sigma_\omega^\gamma(E) = \frac{eV_{AC}}{\hbar\omega} [\Sigma_0^\gamma(E) - \Sigma_0^\gamma(E_+)] \quad (6)$$

$$g_\omega^r = G_0^r(E_+) [U_\omega + \sigma_\omega^r(E)] G_0^r(E) \quad (7)$$

where $E_+ = E + \hbar\omega$.

- (First order) Lesser GF:

$$g_\omega^<(E) = G_0^r(E_+) \Sigma_0^<(E_+) g_\omega^r(E)^\dagger + G_0^r(E_+) \sigma_\omega^<(E) G_0^r(E)^\dagger + g_\omega^r \Sigma_0^<(E) G_0^r(E)^\dagger \quad (8)$$

- → **Why care?** Simplified equations, inefficient, but doable!

AC NEGF II [PG18]

AC NEGF II [PG18]

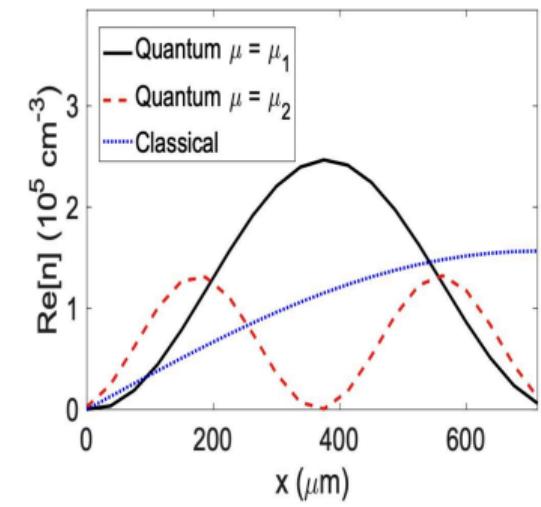
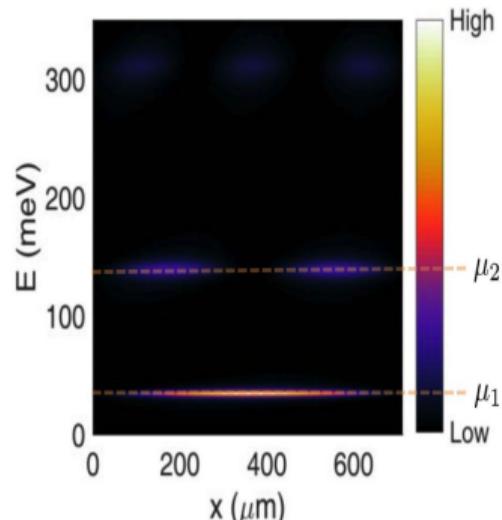
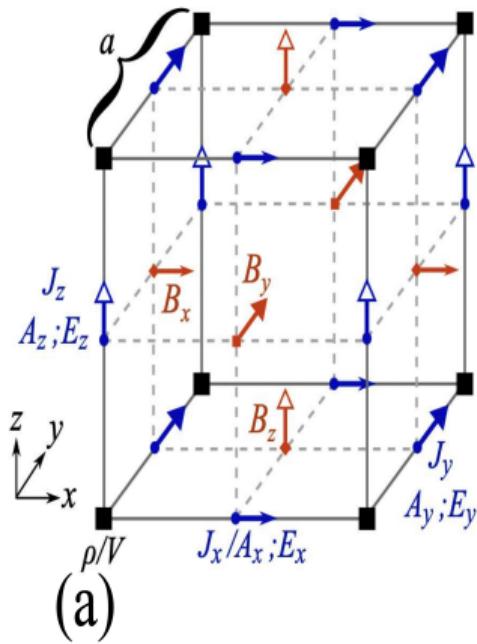
- Development of recursive **AC** GF algorithm.

AC NEGF II [PG18]

- Development of recursive **AC** GF algorithm.
- Couple Maxwell's equations fully! **Self consistently**

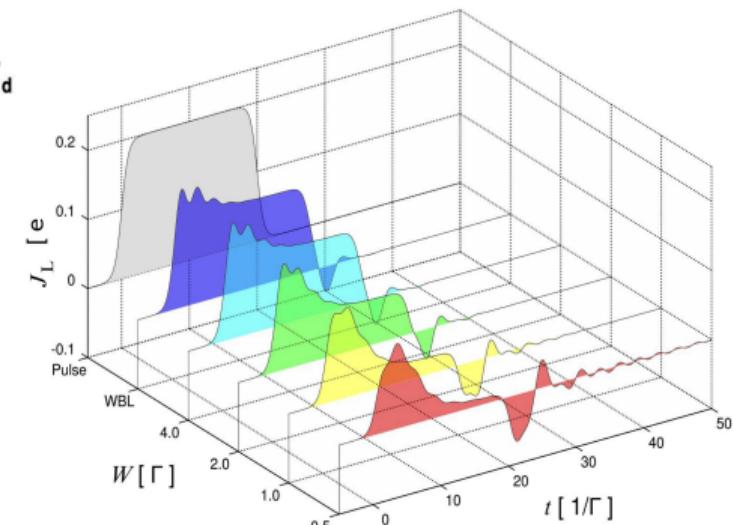
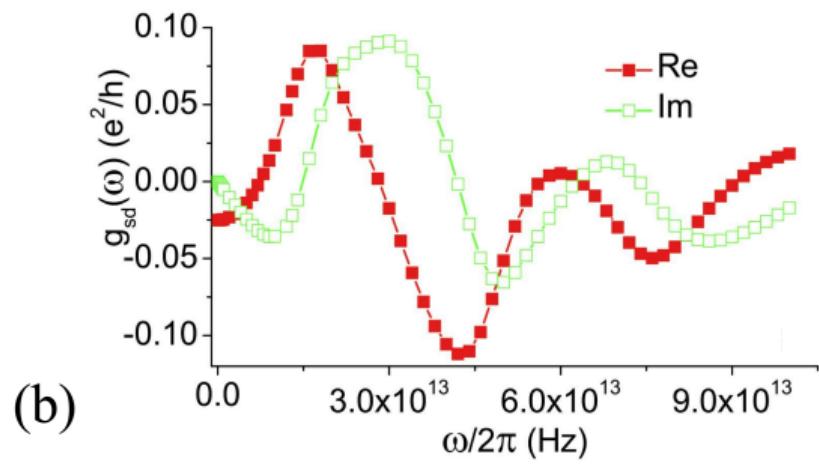
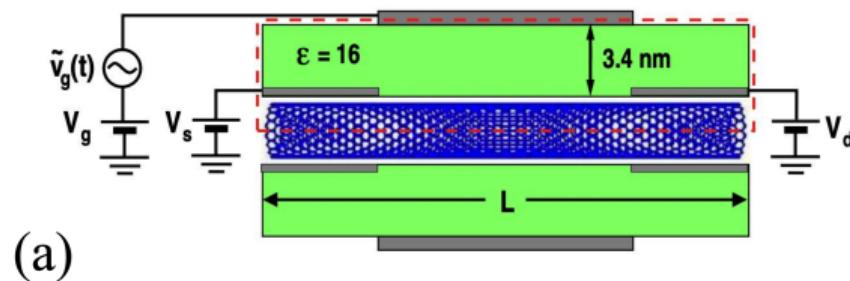
AC NEGF II [PG18]

- Development of recursive **AC** GF algorithm.
- Couple Maxwell's equations fully! **Self consistently**



AC NEGF Others [CS09; PD17; KVL10]

AC NEGF Others [CS09; PD17; KVL10]



AC Scattering Matrix I [Klo+21]

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$
- Scattering wave functions $\psi_{\alpha E}$ at $t < t_0$:

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$
- Scattering wave functions $\psi_{\alpha E}$ at $t < t_0$: labelled by continuous energy E , and discrete α (labels conducting channels at E) s.t. $\mathbf{H}_0 \psi_{\alpha E} = E \psi_{\alpha E}$

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$
- Scattering wave functions $\psi_{\alpha E}$ at $t < t_0$: labelled by continuous energy E , and discrete α (labels conducting channels at E) s.t. $\mathbf{H}_0 \psi_{\alpha E} = E \psi_{\alpha E}$
- For $t > t_0$, we have,

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$
- Scattering wave functions $\psi_{\alpha E}$ at $t < t_0$: labelled by continuous energy E , and discrete α (labels conducting channels at E) s.t. $\mathbf{H}_0 \psi_{\alpha E} = E \psi_{\alpha E}$
- For $t > t_0$, we have,

$$\begin{aligned} i\partial_t \psi_{\alpha E}(t, i) &= \sum_j \mathbf{H}_{ij}(t) \psi_{\alpha E}(t, j) \\ \psi_{\alpha E}(t < t_0, i) &= \psi_{\alpha E}(i) e^{-iEt} \end{aligned}$$

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$
- Scattering wave functions $\psi_{\alpha E}$ at $t < t_0$: labelled by continuous energy E , and discrete α (labels conducting channels at E) s.t. $\mathbf{H}_0 \psi_{\alpha E} = E \psi_{\alpha E}$
- For $t > t_0$, we have,

$$\begin{aligned} i\partial_t \psi_{\alpha E}(t, i) &= \sum_j \mathbf{H}_{ij}(t) \psi_{\alpha E}(t, j) \\ \psi_{\alpha E}(t < t_0, i) &= \psi_{\alpha E}(i) e^{-iEt} \end{aligned} \tag{9}$$

- Some observable A would then be calculated as,

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$
- Scattering wave functions $\psi_{\alpha E}$ at $t < t_0$: labelled by continuous energy E , and discrete α (labels conducting channels at E) s.t. $\mathbf{H}_0 \psi_{\alpha E} = E \psi_{\alpha E}$
- For $t > t_0$, we have,

$$\begin{aligned} i\partial_t \psi_{\alpha E}(t, i) &= \sum_j \mathbf{H}_{ij}(t) \psi_{\alpha E}(t, j) \\ \psi_{\alpha E}(t < t_0, i) &= \psi_{\alpha E}(i) e^{-iEt} \end{aligned} \tag{9}$$

- Some observable A would then be calculated as,

$$\langle \hat{\mathbf{A}} \rangle(t) = \sum_{\alpha ij} \int \frac{dE}{2\pi} f_\alpha(E) \psi_{\alpha E}^*(t, i) \mathbf{A}_{ij} \psi_{\alpha E}(t, j)$$

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$
- Scattering wave functions $\psi_{\alpha E}$ at $t < t_0$: labelled by continuous energy E , and discrete α (labels conducting channels at E) s.t. $\mathbf{H}_0 \psi_{\alpha E} = E \psi_{\alpha E}$
- For $t > t_0$, we have,

$$\begin{aligned} i\partial_t \psi_{\alpha E}(t, i) &= \sum_j \mathbf{H}_{ij}(t) \psi_{\alpha E}(t, j) \\ \psi_{\alpha E}(t < t_0, i) &= \psi_{\alpha E}(i) e^{-iEt} \end{aligned} \tag{9}$$

- Some observable A would then be calculated as,

$$\langle \hat{\mathbf{A}} \rangle(t) = \sum_{\alpha ij} \int \frac{dE}{2\pi} f_\alpha(E) \psi_{\alpha E}^*(t, i) \mathbf{A}_{ij} \psi_{\alpha E}(t, j) \tag{10}$$

- For instance, particle current $I_{ij}(t)$ from site i to site j is

AC Scattering Matrix I [Klo+21]

- Hamiltonian: $\hat{\mathbf{H}} = \hat{\mathbf{H}}_0 + \hat{\mathbf{W}}(t)$, perturbation at $t > t_0$
- Scattering wave functions $\psi_{\alpha E}$ at $t < t_0$: labelled by continuous energy E , and discrete α (labels conducting channels at E) s.t. $\mathbf{H}_0 \psi_{\alpha E} = E \psi_{\alpha E}$
- For $t > t_0$, we have,

$$\begin{aligned} i\partial_t \psi_{\alpha E}(t, i) &= \sum_j \mathbf{H}_{ij}(t) \psi_{\alpha E}(t, j) \\ \psi_{\alpha E}(t < t_0, i) &= \psi_{\alpha E}(i) e^{-iEt} \end{aligned} \tag{9}$$

- Some observable A would then be calculated as,

$$\langle \hat{\mathbf{A}} \rangle(t) = \sum_{\alpha ij} \int \frac{dE}{2\pi} f_\alpha(E) \psi_{\alpha E}^*(t, i) \mathbf{A}_{ij} \psi_{\alpha E}(t, j) \tag{10}$$

- For instance, particle current $I_{ij}(t)$ from site i to site j is

$$I_{ij}(t) = -2 \operatorname{Im} \sum_\alpha \int \frac{dE}{2\pi} f_\alpha(E) \psi_{\alpha E}^*(t, i) \mathbf{H}_{ij} \psi_{\alpha E}(t, j), \tag{11}$$

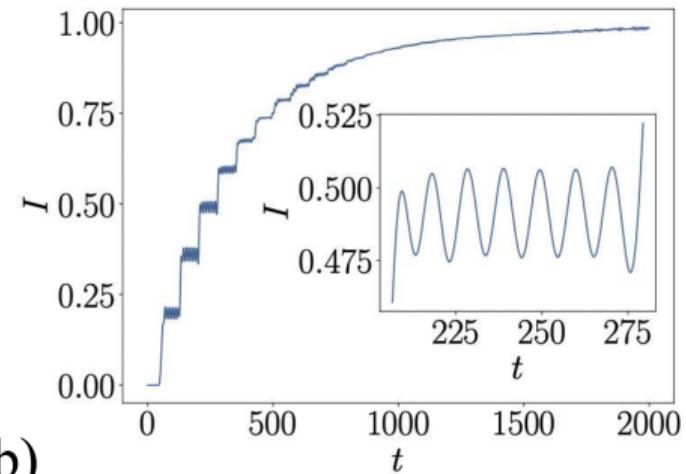
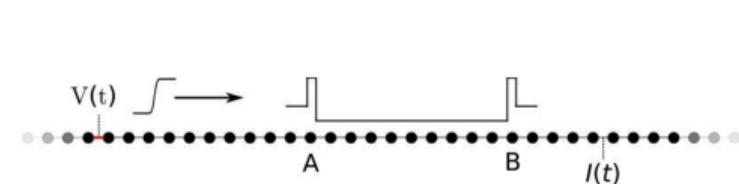
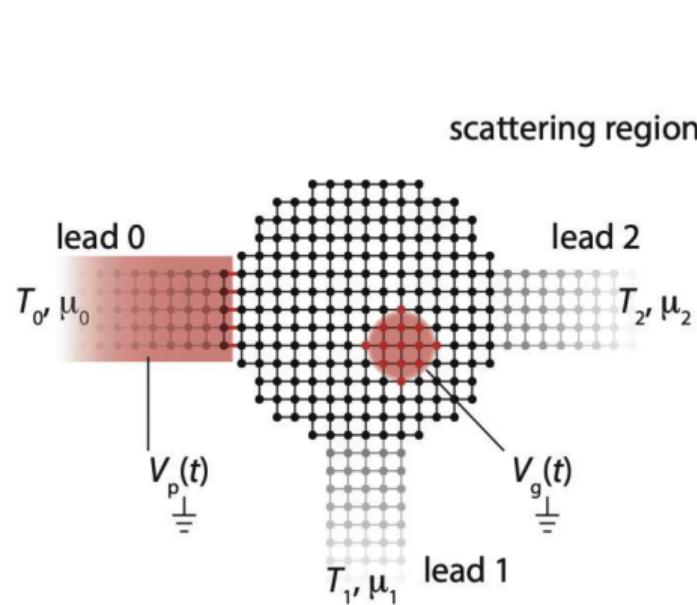
AC Scattering Matrix II [Klo+21]

AC Scattering Matrix II [Klo+21]

- Again, two cases: **uniform** perturbation at a lead (absorbed through a phase factor!), or arbitrary in the channel.

AC Scattering Matrix II [Klo+21]

- Again, two cases: **uniform** perturbation at a lead (absorbed through a phase factor!), or arbitrary in the channel.



(a)

(b)



Numerics: Basic I

Numerics: Basic I

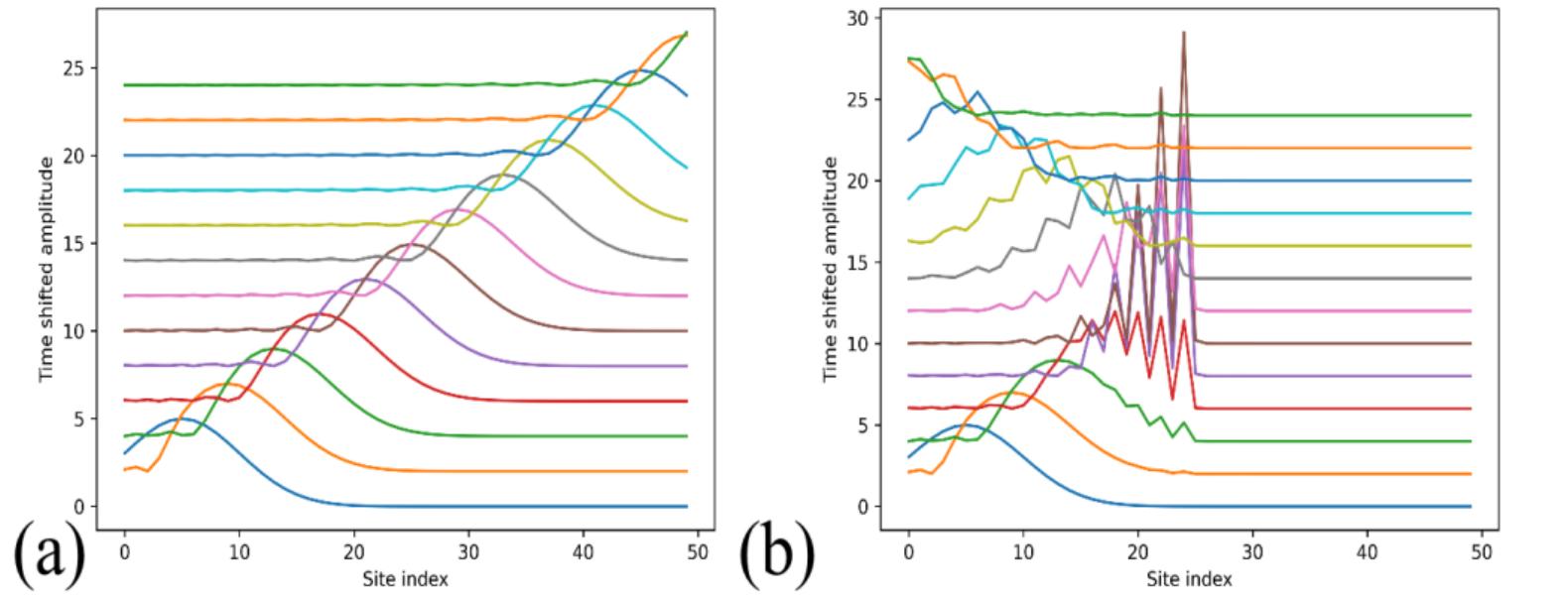


Figure 6: Gaussian wave packet propagation (a) free (b) with barrier

Numerics: Basic II

Numerics: Basic II

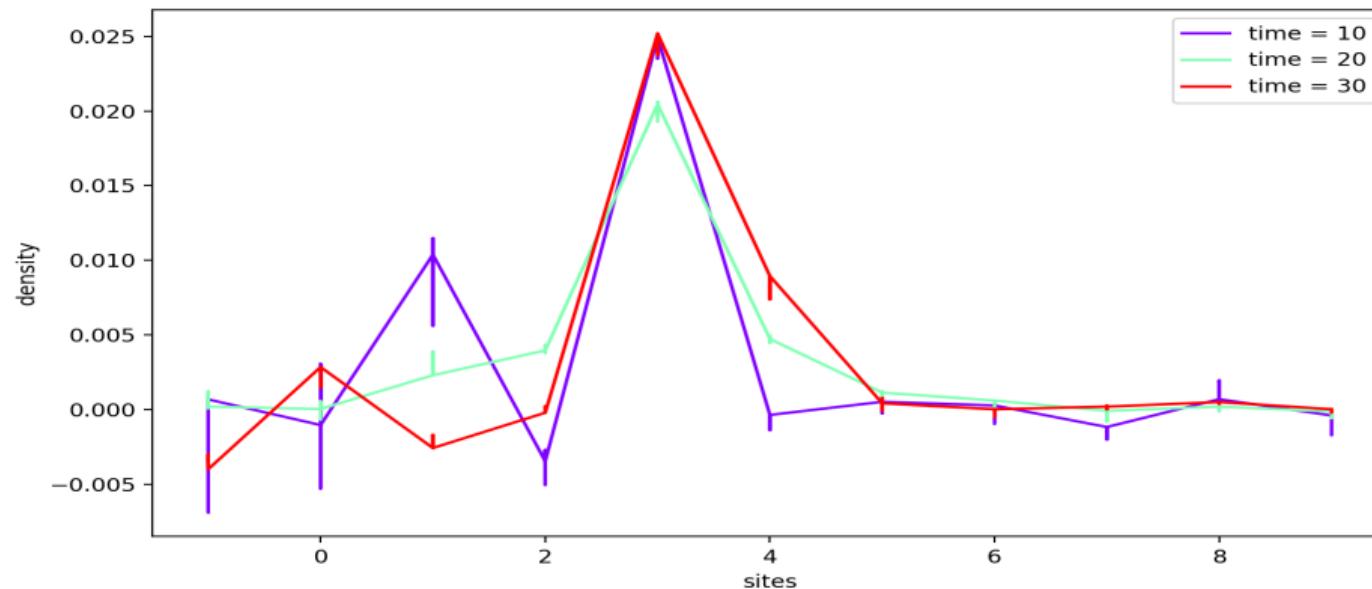


Figure 7: Small square pulse applied across leads of a normal-superconducting junction

Numerics: Basic III

Numerics: Basic III

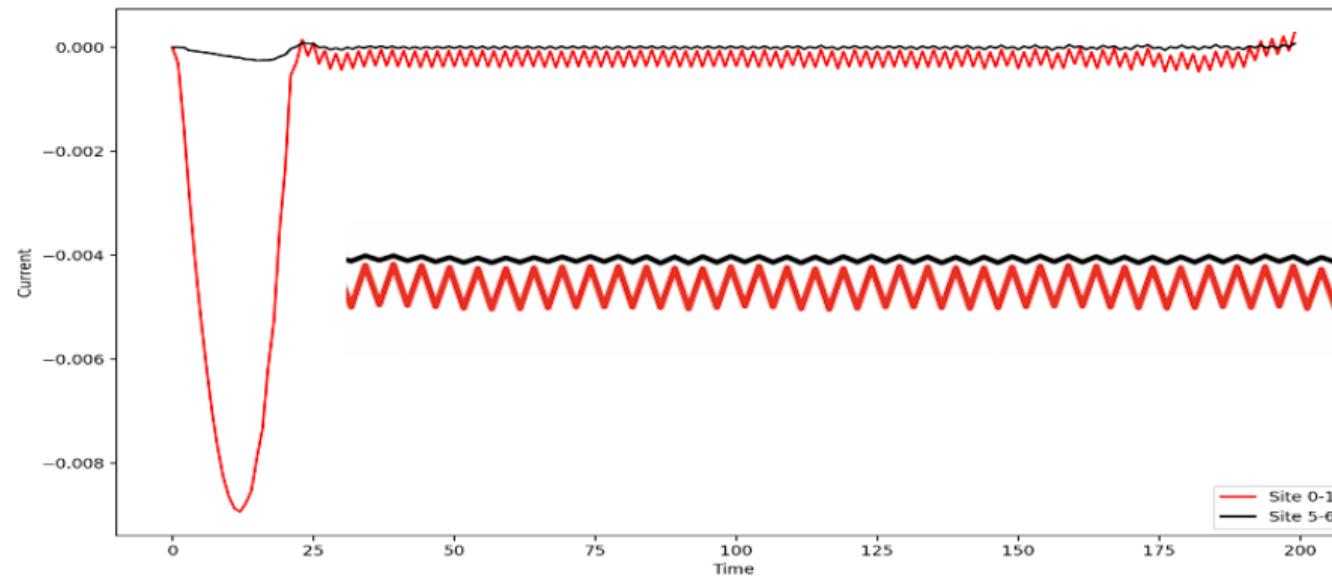


Figure 8: Voltage pulse applied across leads of a topological phase SSH chain

Numerics: Advanced

Numerics: Advanced

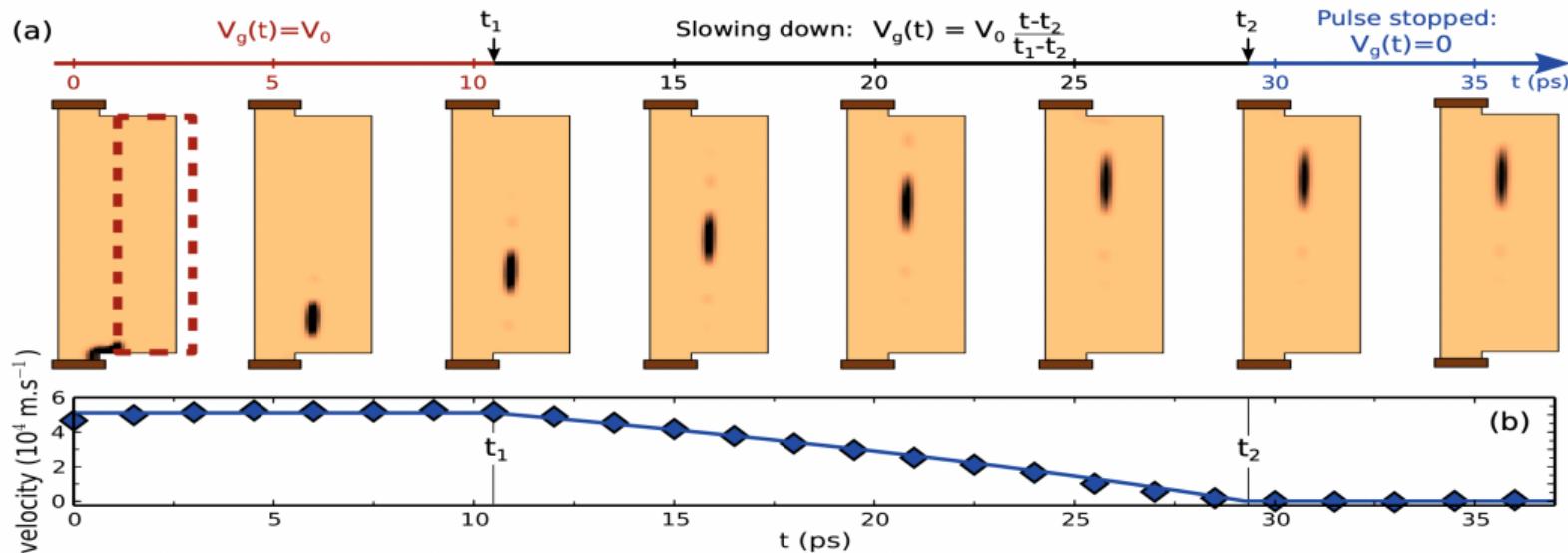


Figure 9: (a) Charge density plotted at various times during the stopping protocol. The gate is polarized for $t < t_1$ (thereby causing the edge mode at the center), and slowly grounded between t_1 and t_2 . At $t = t_2$, the pulse is stopped. (b) Velocity $v(t)$ of the pulse as a function of time. Diamonds correspond to numerical data, the full line to the analytical result. Figure from [\[GWW14\]](#)

That's a wrap!

That's a wrap!

Woop, heavy stuff for 12 minutes, no?

That's a wrap!

Woop, heavy stuff for 12 minutes, no? Thank you :)

References I

- [CS09] Alexander Croy and Ulf Saalmann. "Propagation scheme for nonequilibrium dynamics of electron transport in nanoscale devices". In: *Physical Review B* 80.24 (Dec. 2009). DOI: [10.1103/physrevb.80.245311](https://doi.org/10.1103/physrevb.80.245311). URL: <https://doi.org/10.1103%2Fphysrevb.80.245311>.
- [GWW14] Benoit Gaury, Joseph Weston, and Xavier Waintal. "Stopping electrons with radio-frequency pulses in the quantum Hall regime". In: *Physical Review B* 90.16 (Oct. 2014). DOI: [10.1103/physrevb.90.161305](https://doi.org/10.1103/physrevb.90.161305). URL: <https://doi.org/10.1103%2Fphysrevb.90.161305>.
- [Klo+21] Thomas Kloss et al. "Tkwant: a software package for time-dependent quantum transport". In: *New Journal of Physics* 23.2 (Feb. 2021), p. 023025. DOI: [10.1088/1367-2630/abddf7](https://doi.org/10.1088/1367-2630/abddf7). URL: <https://doi.org/10.1088%2F1367-2630%2Fabddf7>.

References II

- [KVL10] Diego Kienle, Mani Vaidyanathan, and François Léonard. “Self-consistent ac quantum transport using nonequilibrium Green functions”. In: *Physical Review B* 81.11 (Mar. 2010). DOI: [10.1103/physrevb.81.115455](https://doi.org/10.1103/physrevb.81.115455). URL: <https://doi.org/10.1103%2Fphysrevb.81.115455>.
- [PD17] Archak Purkayastha and Yonatan Dubi. “Quantum transport under ac drive from the leads: A Redfield quantum master equation approach”. In: *Phys. Rev. B* 96 (8 Aug. 2017), p. 085425. DOI: [10.1103/PhysRevB.96.085425](https://doi.org/10.1103/PhysRevB.96.085425). URL: <https://link.aps.org/doi/10.1103/PhysRevB.96.085425>.
- [PG18] Timothy M. Philip and Matthew J. Gilbert. “Theory of AC quantum transport with fully electrodynamic coupling”. In: *Journal of Computational Electronics* 17.3 (May 2018), pp. 934–948. DOI: [10.1007/s10825-018-1191-z](https://doi.org/10.1007/s10825-018-1191-z). URL: <https://doi.org/10.1007%2Fs10825-018-1191-z>.

References III

- [SW13] Oleksii Shevtsov and Xavier Waintal. “Numerical toolkit for electronic quantum transport at finite frequency”. In: *Physical Review B* 87.8 (Feb. 2013). DOI: [10.1103/physrevb.87.085304](https://doi.org/10.1103/physrevb.87.085304). URL: <https://doi.org/10.1103%2Fphysrevb.87.085304>.