

Integral Equation Model for Electromagnetic Scattering from a Rough Surface

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Overview

1 Integral Equation Model

2 Modified IEM

Forward Model

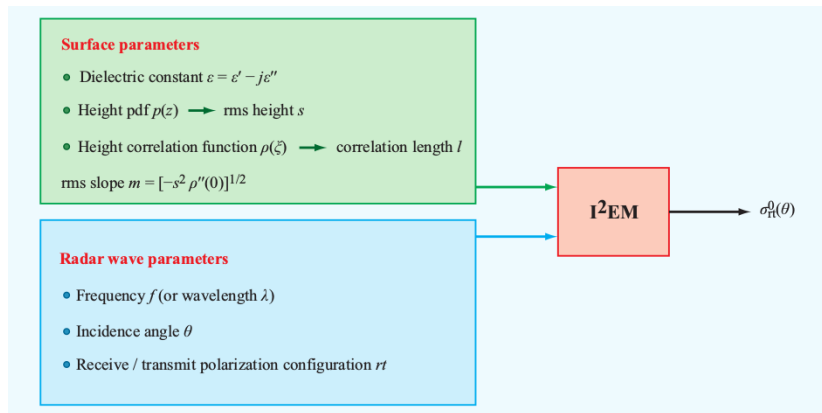


Figure: Improved Integral Equation Model (I^2EM)

Scattering from a rough surface

Incident Field

$$\vec{E}^i = \hat{a}E_0 \exp(-jk_1 \hat{n}_i \cdot r) \quad (1)$$

²F. T. Ulaby, R. K. Moore and A. K. Fung, Microwave Remote Sensing, vol. 2, chapter 12. Reading, MA Addison-Wesley, 1982.

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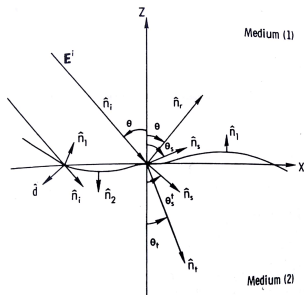


Figure: Geometry of surface scattering problem²

²F. T. Ulaby, R. K. Moore and A. K. Fung, Microwave Remote Sensing, vol. 2, chapter 12. Reading, MA Addison-Wesley, 1982.

Standard Integral Equations

Vector Wave Equation for Electric field

$$\nabla \times \nabla \times \vec{E} - k^2 \vec{E} = -j\omega\mu \vec{J} - \nabla \times \vec{K} \quad (2)$$

Standard Integral Equations

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Vector Greens Theorem

$$\int_V (\vec{Q} \cdot \nabla \times \nabla \times \vec{P} - \vec{P} \cdot \nabla \times \nabla \times \vec{Q}) dv = \int_\Sigma (\vec{P} \times \nabla \times \vec{Q} - \vec{Q} \times \nabla \times \vec{P}) \cdot d\vec{s} \quad (3)$$

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Substituting $\vec{P} = \vec{E}$ and $\vec{Q} = \phi \hat{p}$

$$\int_V j\omega\mu \vec{J} \phi + \vec{K} \times \nabla \phi - (\rho/\epsilon) \nabla \phi dv = \int_\Sigma j\omega\mu (\hat{n}_s \times \vec{H}) \phi - (\hat{n}_s \times \vec{E}) \times \nabla \phi - (\hat{n}_s \cdot \vec{E}) \nabla \phi ds \quad (4)$$

Standard Integral Equations

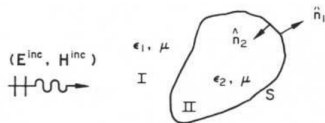


Figure: Scattering from a dielectric body

Standard Integral Equations

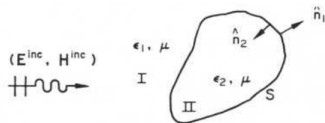


Figure: Scattering from a dielectric body

When source recedes to infinity and the volume currents are zero

$$\vec{E}(x) = \vec{E}^i - \frac{1}{4\pi} \oint_S j\omega\mu(\hat{n}' \times \vec{H})\phi - (\hat{n}' \times \vec{E}) \times \nabla'\phi - (\hat{n}' \cdot \vec{E})\nabla'\phi ds' \quad (5)$$

Standard Integral Equations³

Surface Tangential Fields

$$\hat{n} \times \vec{E} = 2\hat{n} \times \vec{E}^i - \frac{2}{4\pi} \hat{n} \times \int \vec{\mathcal{E}} ds' \quad (6)$$

$$\hat{n} \times \vec{H} = 2\hat{n} \times \vec{H}^i + \frac{2}{4\pi} \hat{n} \times \int \vec{\mathcal{H}} ds' \quad (7)$$

³Poggio, A. J., and E. K. Miller, "Integral equation solution of three dimensional scattering problems," Computer Techniques for Electromagnetics, Chapter 4, Pergmon, New York, 1973.

Standard Integral Equations³

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where

$$\vec{\mathcal{E}} = jk_1 \eta_1 (\hat{n}' \times \vec{H}) G_1 - (\hat{n}' \times \vec{E}_1) \times \nabla' G - (\hat{n}' \cdot \vec{E}) \nabla' G \quad (8)$$

$$\vec{\mathcal{H}} = jk/\eta_1 (\hat{n}' \times \vec{E}) G_1 + (\hat{n}' \times \vec{H}) \times \nabla' G + (\hat{n}' \cdot \vec{H}) \nabla' G \quad (9)$$

³Poggio, A. J., and E. K. Miller, "Integral equation solution of three dimensional scattering problems," Computer Techniques for Electromagnetics, Chapter 4, Pergmon, New York, 1973.

Surface Tangential fields

A pair of integral equations

$$\hat{n} \times \vec{E} = (\hat{n} \times \vec{E})_k + (\hat{n} \times \vec{E})_c \quad (10)$$

$$\hat{n} \times \vec{H} = (\hat{n} \times \vec{H})_k + (\hat{n} \times \vec{H})_c \quad (11)$$

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Fields under Kirchoff Approximation

$$(\hat{n}_1 \times \vec{E})_k = (\hat{n} \times \vec{E}^i)_k + (\hat{n} \times \vec{E}^r)_k \quad (12)$$

$$(\hat{n} \times \vec{H})_k = (\hat{n} \times \vec{H}^i)_k + (\hat{n} \times \vec{H}^r)_k \quad (13)$$

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Complementary surface fields are defined as

$$(\hat{n} \times \vec{E})_c = \hat{n} \times (\vec{E}^i - \vec{E}^r) - \frac{2}{4\pi} \hat{n}_1 \times \int \vec{\mathcal{E}} ds' \quad (14)$$

$$(\hat{n} \times \vec{H})_c = \hat{n} \times (\vec{H}^i - \vec{H}^r) + \frac{2}{4\pi} \hat{n} \times \int \vec{\mathcal{H}} ds' \quad (15)$$

Estimation of Tangential Surface Fields

Vertical Polarization

$$(\hat{n} \times \vec{E})_{kv} \approx (1 - R_{\parallel}) \hat{n} \times \hat{v} E^i \quad (16)$$

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Estimation of Tangential Surface Fields

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Horizontal Polarization

$$(\hat{n} \times \vec{E})_{kh} \approx (1 + R_{\perp}) \hat{n} \times \hat{h} E^i \quad (17)$$

Cross Polarization

$$(\hat{n} \times \vec{E})_{kh} \approx (1 - (R_{\perp} + R_{\parallel})/2) (\hat{n} \times \hat{p}) E^i \quad (18)$$

Far zone backscattered field

Stratton-Chu Integral

$$E_{qp}^s = -\frac{jk}{r\pi R} \exp(-jkR) \int \hat{q} \cdot [-\hat{k}_i \times (\hat{n} \times \vec{E}_p) + \eta(\hat{n} \times \vec{H}_p)] \quad (19)$$

⁴ where

$$\vec{k}_i = k\vec{k}_i = k(\hat{x}\sin\theta - \hat{z}\cos\theta) = \hat{x}k_x - \hat{z}k_z \quad (20)$$

⁴Stratton, J.A., Electromagnetic Theory, McGraw-Hill, NY, 1941

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Polarisation vectors \hat{q} and \hat{p} may be equal to \hat{h} or \hat{v}

$$\hat{h} = \hat{y} \quad (21)$$

$$\hat{v} = \hat{x}\cos\theta + \hat{z}\sin\theta \quad (22)$$

⁴Stratton, J.A., Electromagnetic Theory, McGraw-Hill, NY, 1941

Far zone backscattered field

Kirchoff and Complementary Components of scattered field

$$E_{qp}^k = CE_0 \int f_{qp} \exp(-2j\vec{k}_i \cdot \vec{r}) dx dy \quad (23)$$

$$E_{qp}^c = \frac{CE_0}{8\pi^2} \int F_{qp} \exp[ju(x - x') + jv(y - y') - j\hat{k}_i \cdot \vec{r} - j\vec{k}_i \cdot \vec{r}'] dx dy dx' dy' du dv \quad (24)$$

Far zone backscattered field

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Field Coefficients

$$f_{qp} = \left[\hat{k}_i \times \hat{q} \cdot (\hat{n} \times \vec{E})_{kp} + \eta \hat{q} \cdot (\hat{n} \times \vec{H})_{kp} \right] \cdot (D_1/E^i) \quad (25)$$

$$F_{vv} = 8\pi^2 \left[\hat{k}_i \times \hat{q} \cdot (\hat{n} \times \vec{E})_{cp} + \eta \hat{q} \cdot (\hat{n} \times \vec{H})_{cp} \right] \cdot (D_1/E^i) \quad (26)$$

where $D_1 = \sqrt{Z_x^2 + Z_y^2 + 1}$

Power Calculation

Incoherent power = total power - subtract mean squared power

$$\begin{aligned}
 \langle E_{qp}^s E_{qp}^{s*} \rangle - \langle E_{qp}^s \rangle \langle E_{qp}^s \rangle^* &= \langle E_{qp}^k E_{qp}^{k*} \rangle - \langle E_{qp}^k \rangle \langle E_{qp}^k \rangle^* \\
 &\quad + 2\text{Re} \left[\langle E_{qp}^c E_{qp}^{k*} \rangle - \langle E_{qp}^c \rangle \langle E_{qp}^k \rangle^* \right] \\
 &\quad + \langle E_{qp}^c E_{qp}^{c*} \rangle - \langle E_{qp}^c \rangle \langle E_{qp}^c \rangle^* \quad (27)
 \end{aligned}$$

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Kirchhoff Power term (assuming Gaussain height distribution)

$$P_{qp}^k = |CE_0 f_{qp}|^2 A_0 \exp(-4k_z^2 \sigma^2) \cdot \int [\exp[4k_z^2 \sigma^2 \rho(\xi, \zeta)] - 1] e^{-2jk_x \xi} d\xi d\zeta \quad (28)$$

where $\zeta = x - x'$, $\xi = y - y'$ and A_0 is the illuminated area, σ^2 is the variance of surface and $\rho(\xi, \zeta)$ is the surface correlation.

Backscattering Coefficient

$$\sigma_{qp}^0 = \sigma_{qp}^k + \sigma_{qp}^{kc} + \sigma_{qp}^c = (4\pi R^2 P_{qp}) / (E_0^2 A_0) \quad (29)$$

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Kirchhoff Backscatter term

$$\sigma_{qp}^k = \frac{k^2}{4\pi} [|f_{qp}|^2 \exp(-4k_z^2 \sigma^2)] \cdot \int [\exp[4k_z^2 \sigma^2 \rho(\xi, \zeta)] - 1] e^{-2jk_x \xi} d\xi d\zeta \quad (30)$$

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Spectral Representation of Greens Function

$$G = \left(\frac{1}{2\pi} \right) \int \frac{j}{q} \exp[(-jq)|z - z'| + ju(x - x') + jv(y - y')] du dv \quad (31)$$

Soil Profile

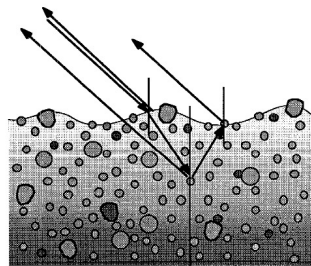
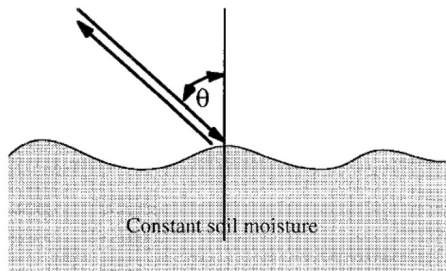
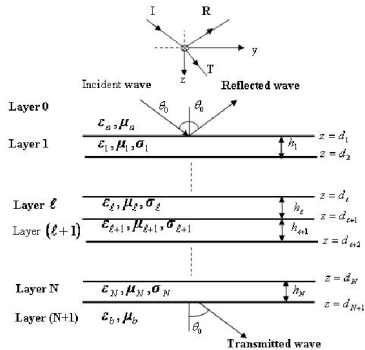


Figure: Homogeneous and Heterogeneous Soil Profile⁵

⁵A. K. Fung et al., "A modified IEM model for: scattering from soil surfaces with application to soil moisture sensing," *Geoscience and Remote Sensing Symposium*, 1996.

Figure: Multilayer dielectric structure⁶

$$R_{ij} = \frac{n_i \cos \theta_i - n_j \cos \theta_j}{n_i \cos \theta_i + n_j \cos \theta_j} \quad (32)$$

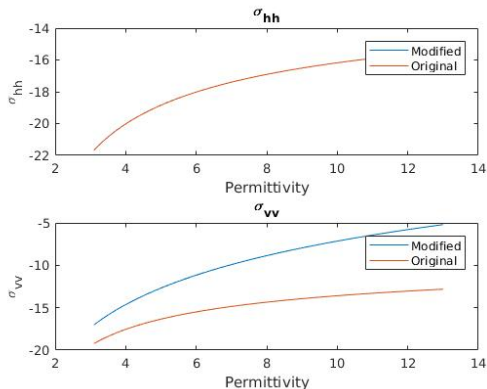
$$T_{ij} = \frac{2n_i \cos \theta_i}{n_i \cos \theta_i + n_j \cos \theta_j} \quad (33)$$

⁶Analysis of planar dielectric multilayers as fss by Transmission Line Transfer Matrix Method (TLTMM)

Two layer IEM

Fresnel Reflection Coefficient

$$R_{\perp} = T_{12}R_{23} \left(\frac{T_{21}e^{-2jk_2d}}{1 - R_{21}R_{23}e^{-2jk_2d}} \right) \quad (34)$$



Modified IEM

Transitional dielectric layer

$$\epsilon_r(z) = 1 + \epsilon_{r\infty} \frac{\exp(mz)}{1 + \exp(mz)} \quad (35)$$

Modified IEM

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Fresnel Coefficient for Perpendicular Polarization

$$R_{\perp} = \frac{\Gamma(jS \cos \theta)}{\Gamma(-jS \cos \theta)} \frac{\Gamma \left[(-jS/2)(\cos \theta + \sqrt{\cos^2 \theta + \epsilon_{r\infty} - 1}) \right]}{\Gamma \left[(jS/2)(\cos \theta - \sqrt{\cos^2 \theta + \epsilon_{r\infty} - 1}) \right]} \\ \times \frac{\Gamma \left[1 - (jS/2)(\cos \theta + \sqrt{\cos^2 \theta + \epsilon_{r\infty} - 1}) \right]}{\Gamma \left[1 + (jS/2)(\cos \theta - \sqrt{\cos^2 \theta + \epsilon_{r\infty} - 1}) \right]} \quad (36)$$

where $S = 2k/m$ and m is the transition rate factor

Modified IEM

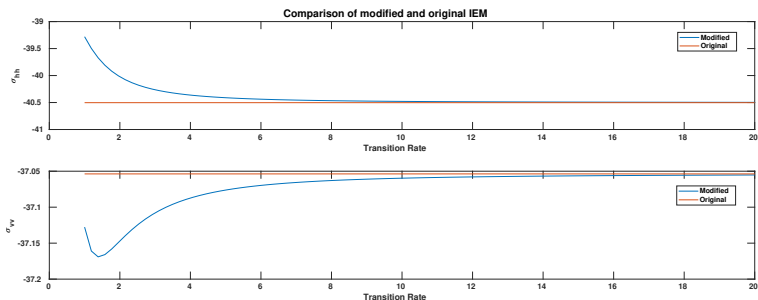
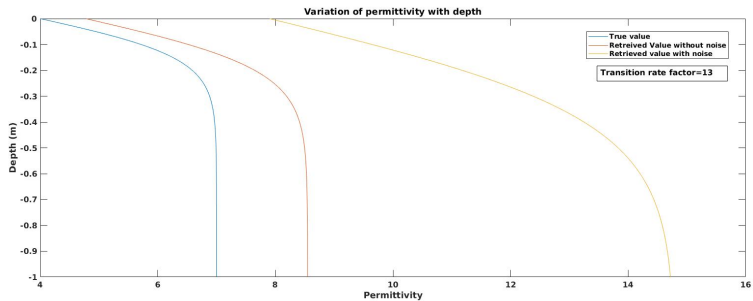
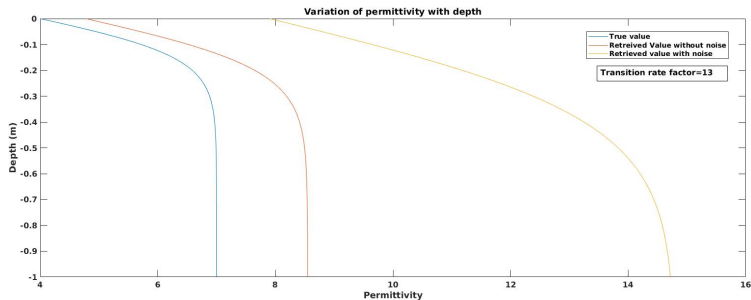


Figure: Backscattering Coefficients calculated from standard and modified IEM

Retrieved Permittivity



Retrieved Permittivity



RMS Errors:

Type	h	l	ϵ	m
Validation	0.068	5.23	1.79	1.79
Test data	0.1	5.73	2.92	3.27

Future work

- 1 Make a modified forward model in which multilayer formulation for moisture can be applied
- 2 Current model is only for bare soil. Suitable model has to be built that takes into account vegetation such as tree, canopy, etc

References

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-  K.-L. Chen, K.-S. Chen, Z.-L. Li, and Y. Liu, "Extension and Validation of an Advanced Integral Equation Model for Bistatic Scattering from Rough Surfaces," Progress In Electromagnetics Research, Vol. 152, 59-76, 2015.
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Thank You