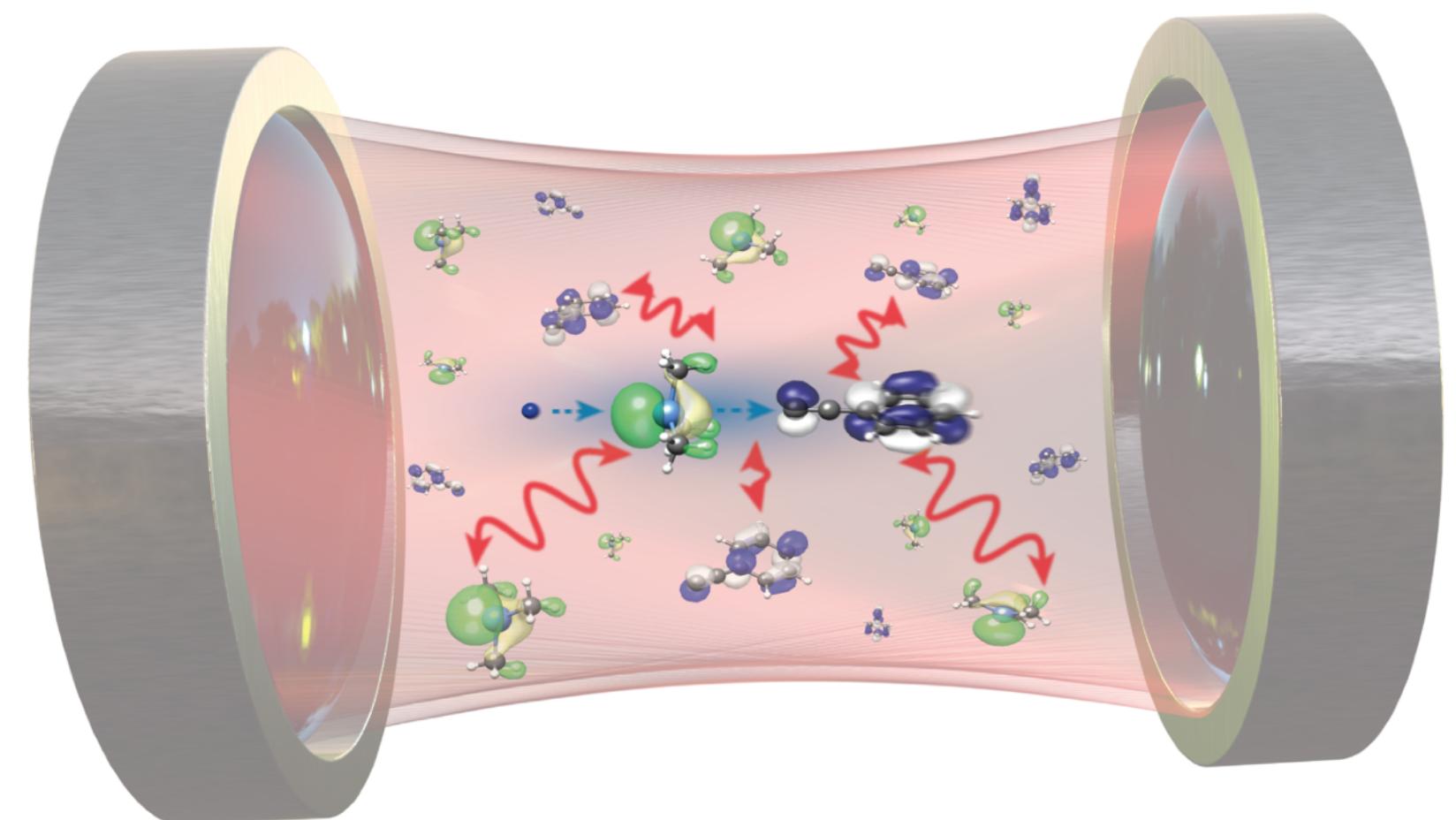
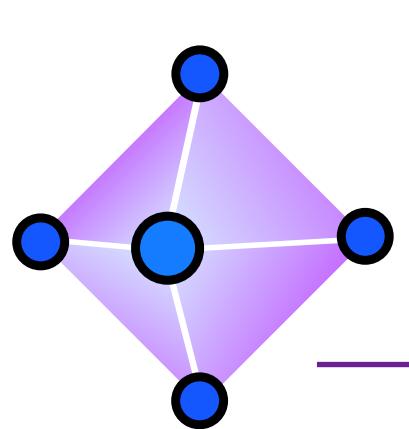


Dominik Sidler, January 2025

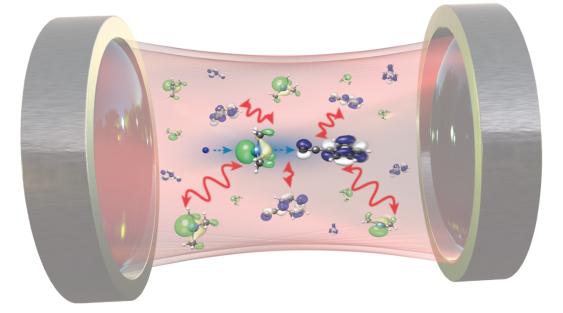
# Polaritonic / QED Chemistry

## Lecture 2: Quantum Optics

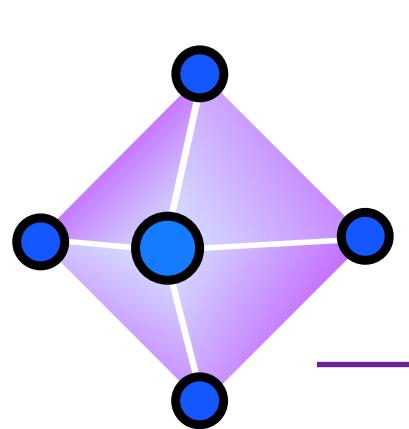




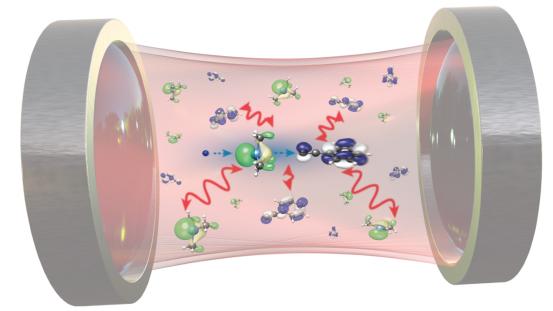
# Lecture Outline



1. Repeat quantum mechanics: Notation and important concepts with toy examples.
2. Introduce Jaynes-Cummings model that provides simplest, analytically solvable model for strong (quantized) light-matter interaction.
3. Discuss model extensions numerically (Rabi model) and analytically (Tavis-Cummings model).
4. Understand concept of „collective“ strong coupling.



# Quantum Mechanics Repetition



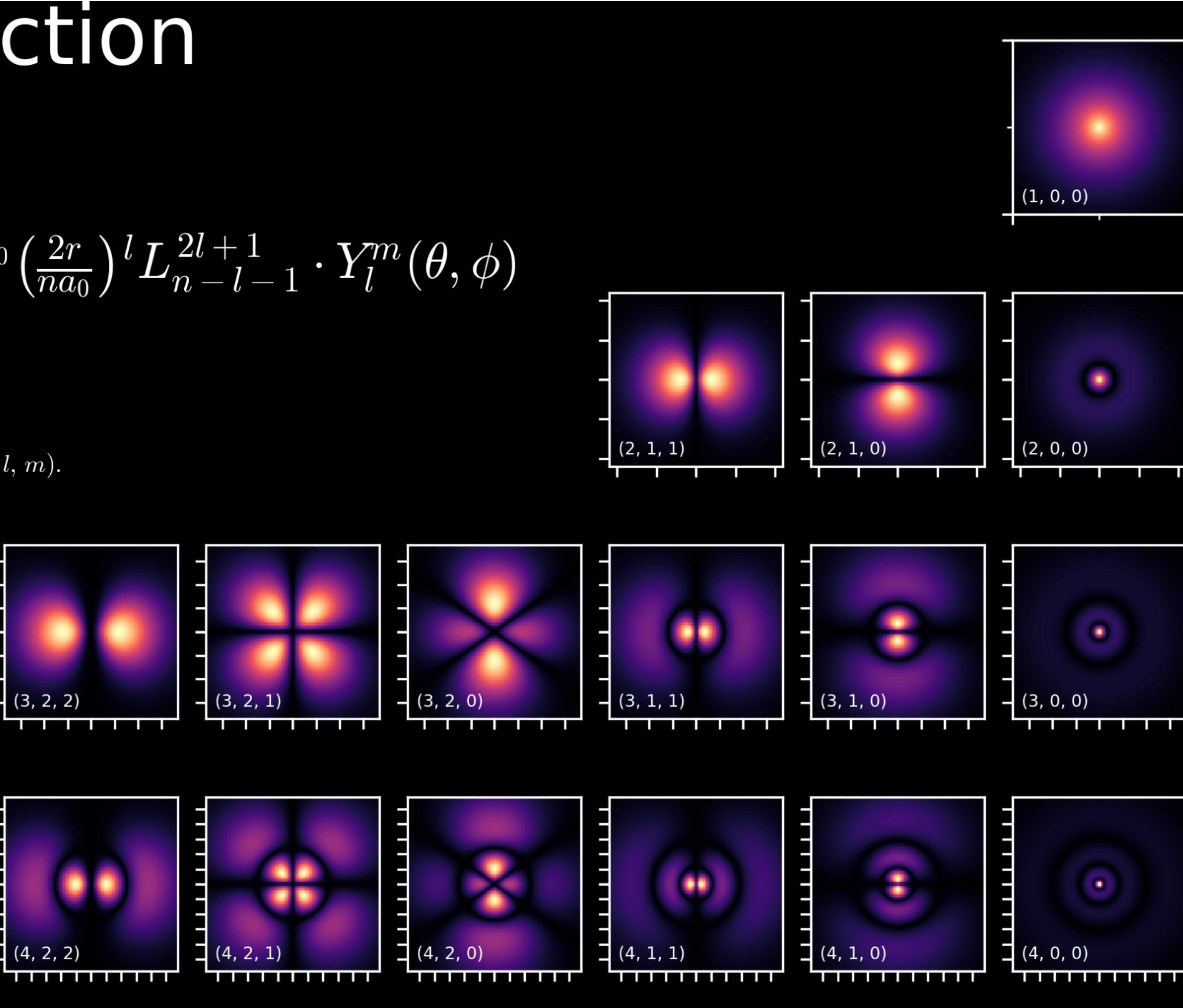
## Blackboard

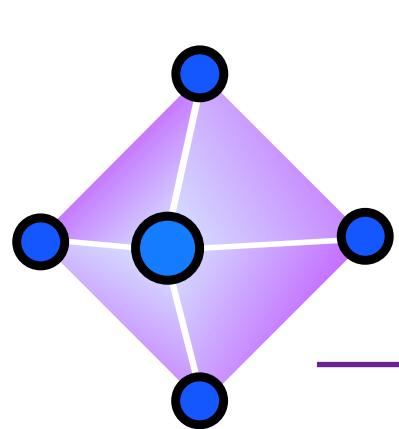
### Hydrogen Wave Function Probability Density

$$\psi_{nlm}(r, \theta, \phi) = \sqrt{\left(\frac{2}{na_0}\right)^3 \frac{(n-l-1)!}{2n[(n+l)!]}} e^{-r/na_0} \left(\frac{2r}{na_0}\right)^l L_{n-l-1}^{2l+1} \cdot Y_l^m(\theta, \phi)$$

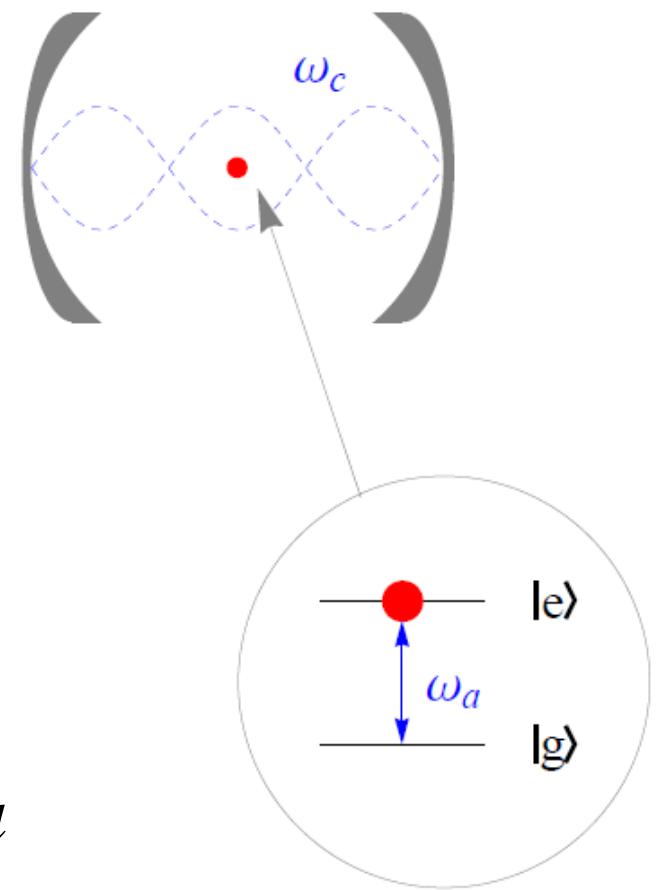
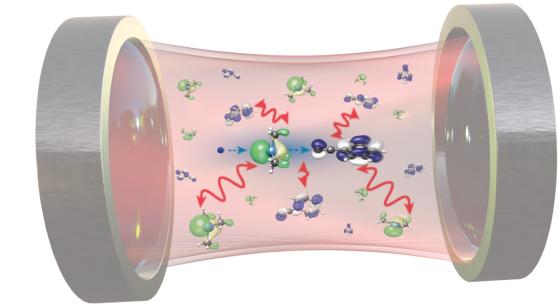
NOTE:

1. The distance between two major ticks in the plot is  $5a_0$ .
2. The numbers in the brace are the three quantum numbers  $(n, l, m)$ .





# Jaynes-Cummings (JC) model



## Strong Coupling - Vacuum Rabi Oscillations

$$\hat{H} = \hat{H}_{\text{field}} + \hat{H}_{\text{atom}} + \hat{H}_{\text{int}}$$

$$\hat{H}_{\text{field}} = \hbar\omega_c \hat{a}^\dagger \hat{a}$$

Angular (photon) mode frequency  $\omega_c$

$$\hat{H}_{\text{atom}} = \hbar\omega_a \frac{\hat{\sigma}_z}{2}$$

$\hat{\sigma}_z = |e\rangle\langle e| - |g\rangle\langle g|$  matter excitation frequency  $\omega_a$

$$\hat{H}_{\text{int}} = \frac{\hbar\Omega}{2} \hat{E} \hat{S}.$$

$\hat{E} = E_{\text{ZPF}} (\hat{a} + \hat{a}^\dagger)$  raising / lowering operator of **2-level matter representation**

Polarization operator  $\hat{S} = \hat{\sigma}_+ + \hat{\sigma}_-$ ,  $\hat{\sigma}_+ = |e\rangle\langle g|$ ,  $\hat{\sigma}_- = |g\rangle\langle e|$ ,

Schrödinger picture:  $H_S = H_{0,S} + H_{1,S}$ . Interaction picture:  $H_{0,I}(t) = e^{iH_{0,S}t/\hbar} H_{0,S} e^{-iH_{0,S}t/\hbar} = H_{0,S}$

$$H_{1,I}(t) = e^{iH_{0,S}t/\hbar} H_{1,S} e^{-iH_{0,S}t/\hbar}$$

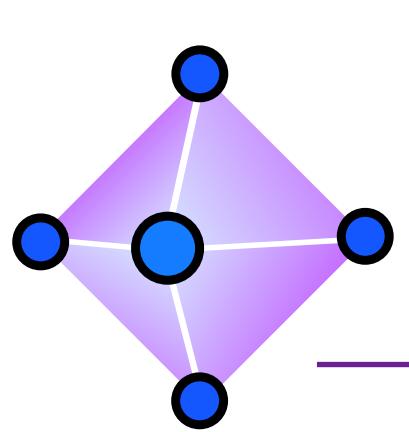
$$\hat{H}_0 = \hat{H}_{\text{field}} + \hat{H}_{\text{atom}}$$

$$\hat{H}_{\text{int}}(t) = \frac{\hbar\Omega}{2} \left( \hat{a}\hat{\sigma}_- e^{-i(\omega_c+\omega_a)t} + \hat{a}^\dagger\hat{\sigma}_+ e^{i(\omega_c+\omega_a)t} + \hat{a}\hat{\sigma}_+ e^{i(-\omega_c+\omega_a)t} + \hat{a}^\dagger\hat{\sigma}_- e^{-i(-\omega_c+\omega_a)t} \right)$$

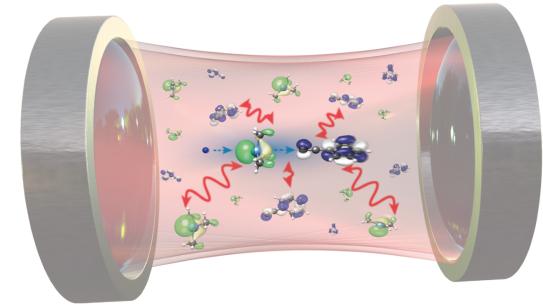
**Rotating wave approximation (RWA):** Discard quickly oscillating states that couple energetically far separated state (mix only little for moderate **light-matter coupling strength  $\Omega$** )

Schrödinger picture + RWA (**Jaynes-Cummings model**)

$$\boxed{\hat{H}_{\text{JC}} = \hbar\omega_c \hat{a}^\dagger \hat{a} + \hbar\omega_a \frac{\hat{\sigma}_z}{2} + \frac{\hbar\Omega}{2} (\hat{a}\hat{\sigma}_+ + \hat{a}^\dagger\hat{\sigma}_-)}$$

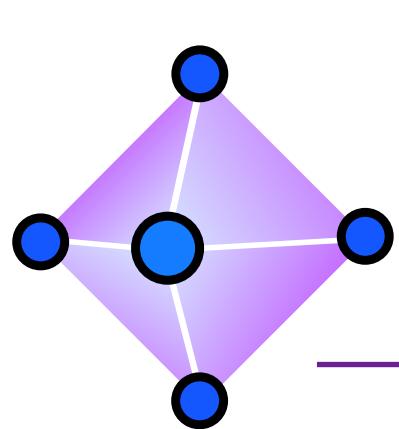


# Solution of JC model

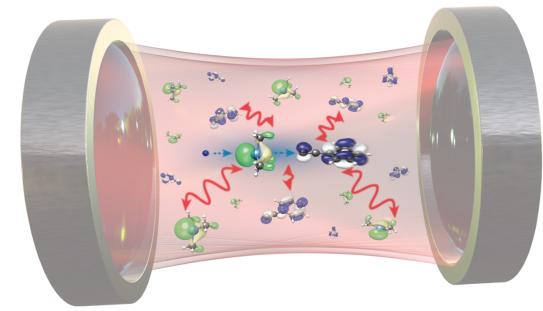


## Blackboard

$$\hat{H}_{\text{JC}} = \hbar\omega_c \hat{a}^\dagger \hat{a} + \hbar\omega_a \frac{\hat{\sigma}_z}{2} + \frac{\hbar\Omega}{2} (\hat{a}\hat{\sigma}_+ + \hat{a}^\dagger \hat{\sigma}_-)$$



# JC Summary / Properties



$$\hat{H}_{\text{JC}} = \hbar\omega_c \hat{a}^\dagger \hat{a} + \hbar\omega_a \frac{\hat{\sigma}_z}{2} + \frac{\hbar\Omega}{2} (\hat{a}\hat{\sigma}_+ + \hat{a}^\dagger\hat{\sigma}_-)$$

Detuning  $\delta = \omega_a - \omega_c$

**Vacuum Rabi frequency (splitting):**  $\Omega_n(\delta) = \sqrt{\delta^2 + \Omega^2(n+1)}$

**Excited state energy eigenvalues (Jaynes-Cummings ladder)**

$$E_{\pm}(n) = \hbar\omega_c \left( n + \frac{1}{2} \right) \pm \frac{1}{2} \hbar\Omega_n(\delta),$$

**Lower polariton** (eigenstates)

$$|n, -\rangle = \sin\left(\frac{\alpha_n}{2}\right) |\psi_{1n}\rangle - \cos\left(\frac{\alpha_n}{2}\right) |\psi_{2n}\rangle$$

**Upper polariton** (eigenstates)

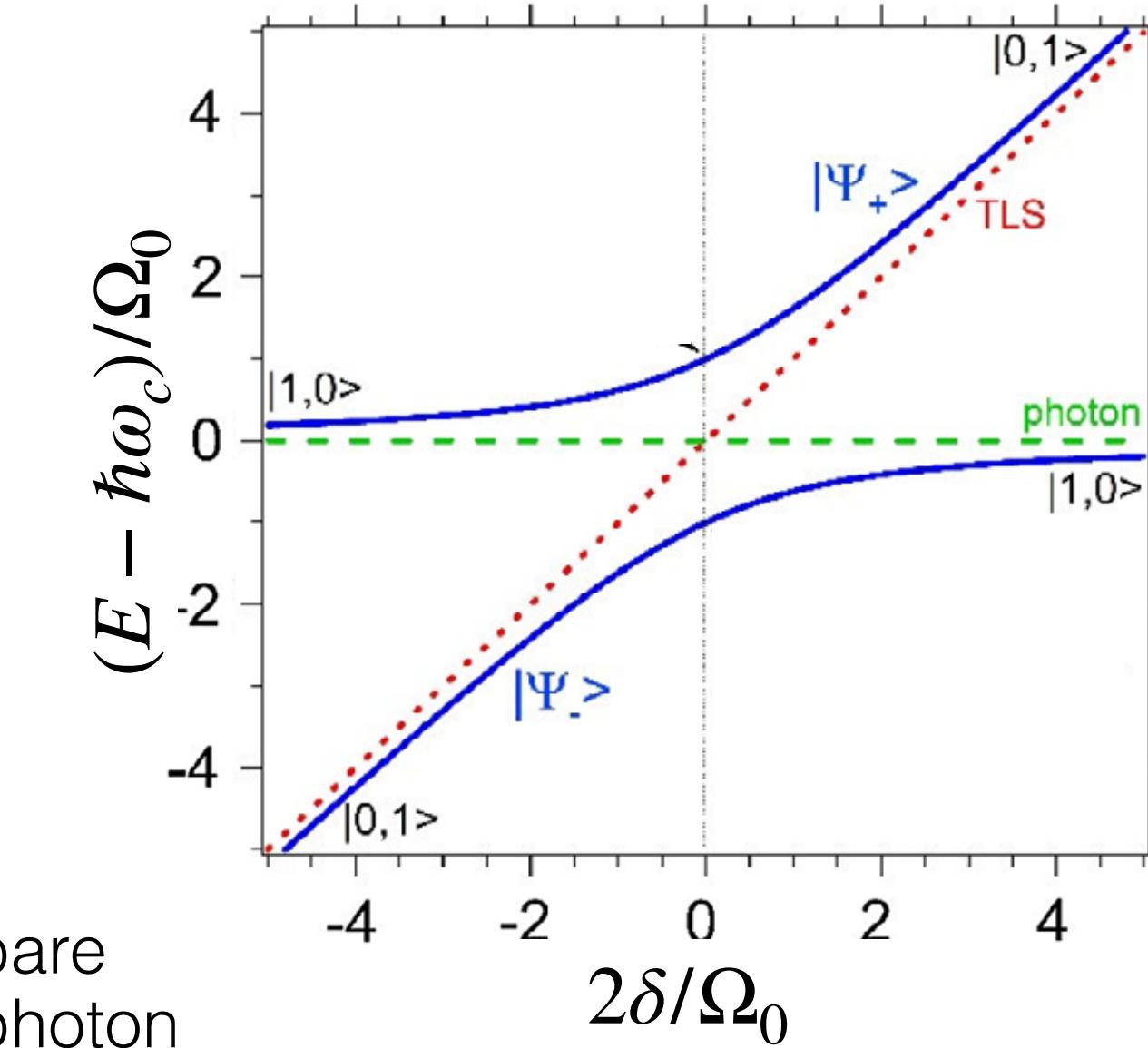
$$|n, +\rangle = \cos\left(\frac{\alpha_n}{2}\right) |\psi_{1n}\rangle + \sin\left(\frac{\alpha_n}{2}\right) |\psi_{2n}\rangle$$

**Time-evolution:**

Excited initial state (represented in eigenstates):  $|\psi_{\text{tot}}(0)\rangle = \sum_n C_n |n, e\rangle = \sum_n C_n \left[ \cos\left(\frac{\alpha_n}{2}\right) |n, +\rangle + \sin\left(\frac{\alpha_n}{2}\right) |n, -\rangle \right]$

**Vacuum Rabi Oscillations:**  
(back to Schrödinger picture)

$$|\psi_{\text{tot}}(t)\rangle = e^{-i\hat{H}_{\text{JC}}t/\hbar} |\psi_{\text{tot}}(0)\rangle = \sum_n C_n \left[ \cos\left(\frac{\alpha_n}{2}\right) |n, +\rangle e^{-iE_+(n)t/\hbar} + \sin\left(\frac{\alpha_n}{2}\right) |n, -\rangle e^{-iE_-(n)t/\hbar} \right]$$

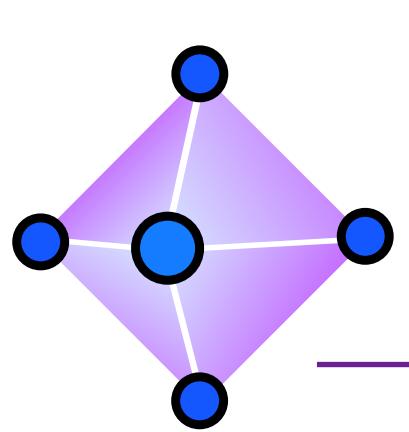


bare  
photon  
excitation

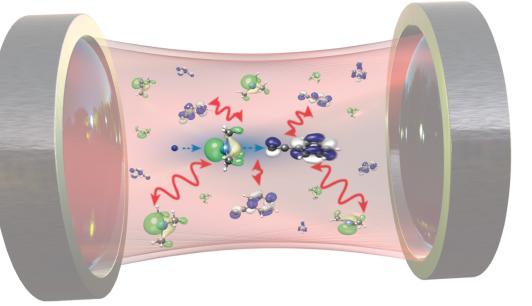
$$|\psi_{1n}\rangle := |n, e\rangle$$

bare matter

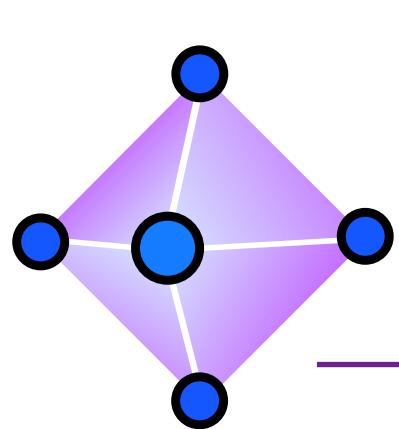
$$\alpha_n := \tan^{-1} \left( \frac{\Omega\sqrt{n+1}}{\delta} \right).$$



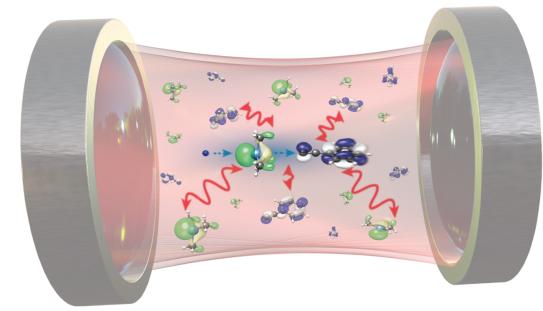
# Python Exercise 2



## Polaritons / Rabi Model (beyond RWA)



# Quantum Rabi Model



## Difference to JC model

Including counter-rotating terms destroys block-diagonal nature of JC-Hamiltonian in terms of **bare** matter and **bare** photon basis set.

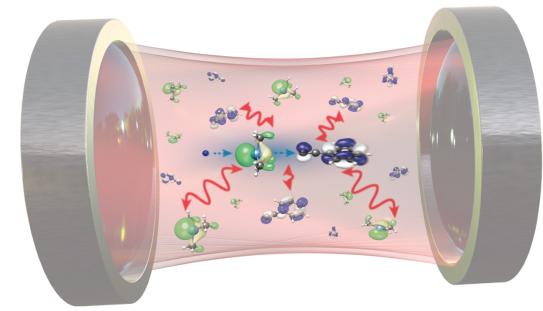
=> Eigenstates of Rabi model require numerical solution, i.e., diagonalization of entire Hamiltonian matrix.

In practice, numerical convergence for relatively small bare matter basis set sizes.

**Remark:** Discarding counter-rotating terms, i.e., the  $(A^2)$ -operator or „dipole self-energy“, removes many physically / chemically relevant effects that could be decisive in real chemical systems! (see upcoming lessons / active area of research)



# Collective Light-Matter Coupling



## Tavis-Cummings model (many emitters)

$$\hat{H}_{TC} = \hbar\omega_c a^\dagger a + \frac{\hbar\omega_a}{2} S_z + \frac{\hbar\Omega}{2} (a^\dagger S_- + a S_+)$$

$$S_\pm = \sum_{i=1}^N \sigma_\pm^{(i)} \quad S_z = \sum_{i=1}^N \sigma_z^{(i)}$$

Primakoff transformation (spin to bosonic Fock operator representation) + assume large ensembles of molecules, e.g.  
 $N \lesssim N_A = 6.022 \times 10^{23}$

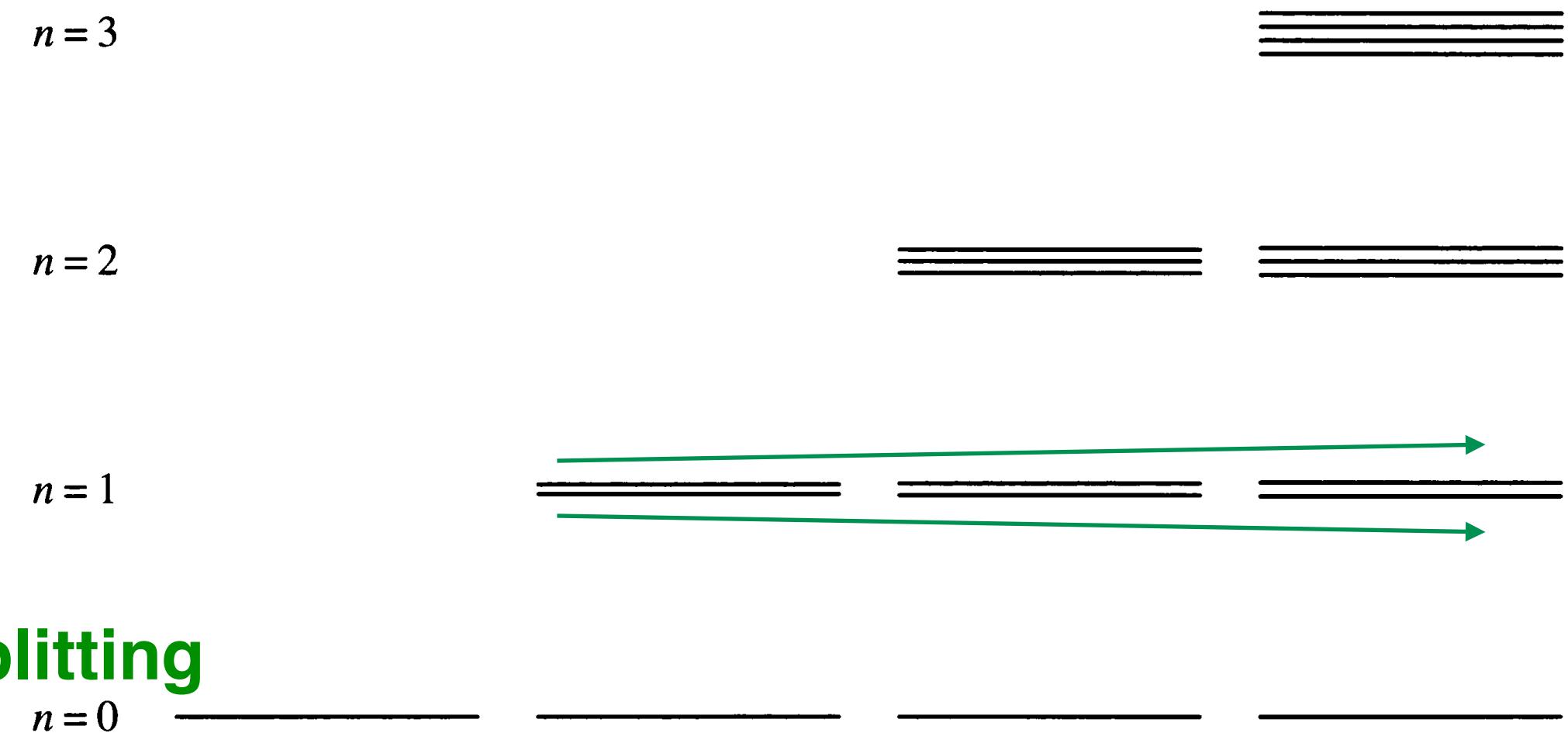
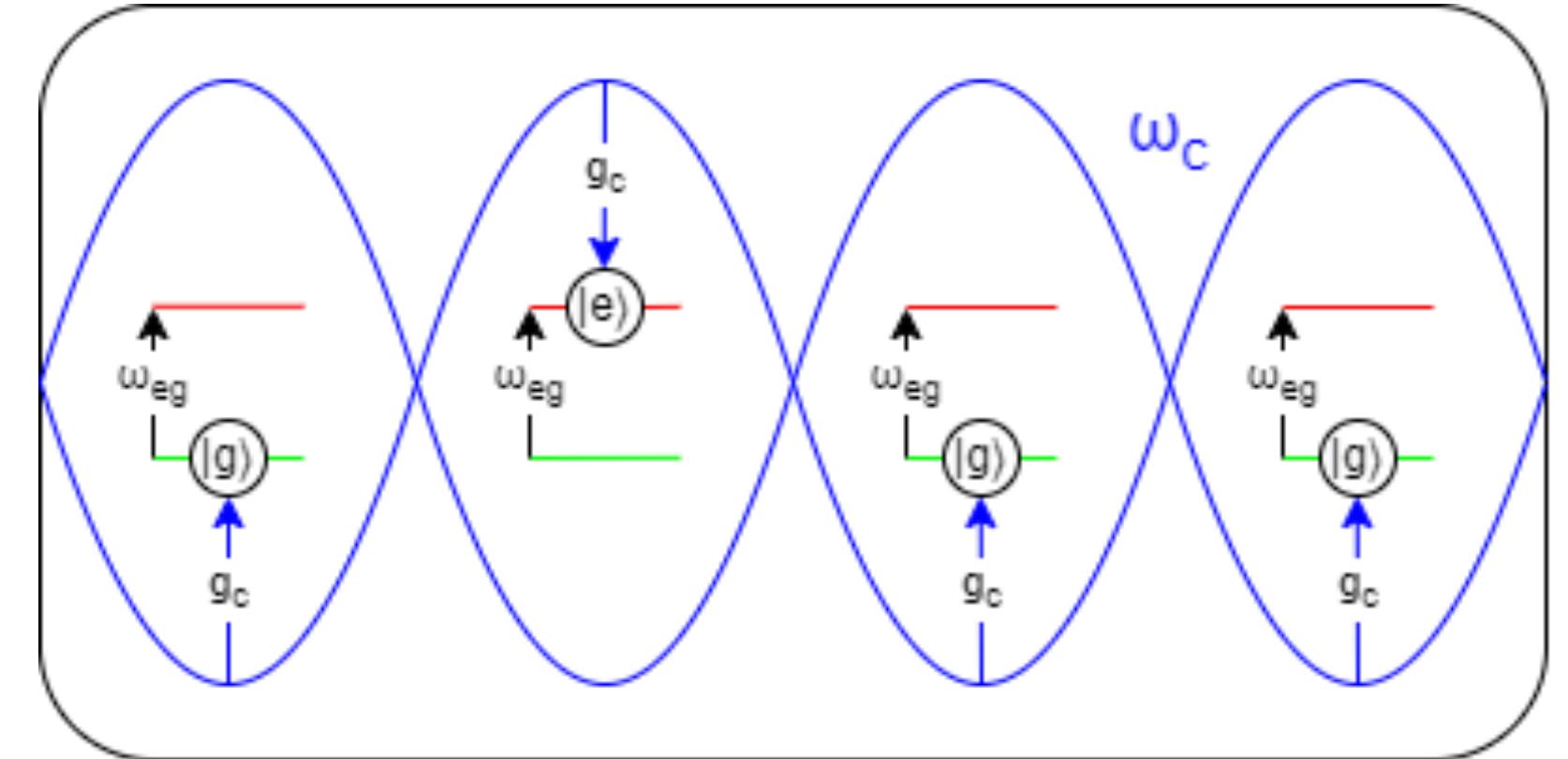
$$S_+ = b^\dagger (N - b^\dagger b)^{1/2} \stackrel{N \gg 1}{\approx} b^\dagger \sqrt{N}, \quad S_- = (N - b^\dagger b)^{1/2} b \stackrel{N \gg 1}{\approx} b \sqrt{N}, \quad S_z = b^\dagger b - \frac{N}{2}$$

TC resembles approximately two coupled harmonic oscillators:

$$\hat{H}_{TC} \sim \hbar\omega_c a^\dagger a + \frac{\hbar\omega_a}{2} \left( -\frac{N}{2} + b^\dagger b \right) + \frac{\hbar\Omega}{2} \boxed{\sqrt{N}(a^\dagger b + ab^\dagger)}$$

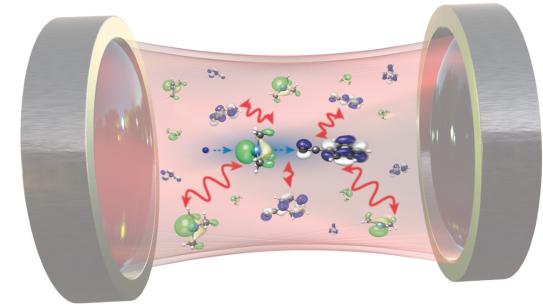
**Collective coupling strength**  $\Omega \mapsto \sqrt{N}\Omega$  introduces **collective Rabi-splitting**

$$E_{TC}(n, j) \sim \left( n - \frac{N}{2} \right) \omega_c - \frac{n}{2} \delta + \boxed{j\sqrt{\hbar^2\Omega^2 N + \delta^2}}, \quad -n/2 \leq j \leq n/2, \quad 0 \leq n < N, \quad \Delta j = \in \mathbb{Z}$$



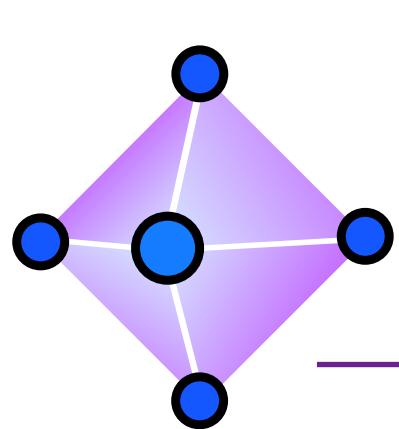


# Common (Misleading) Conclusions from JC / TC Models

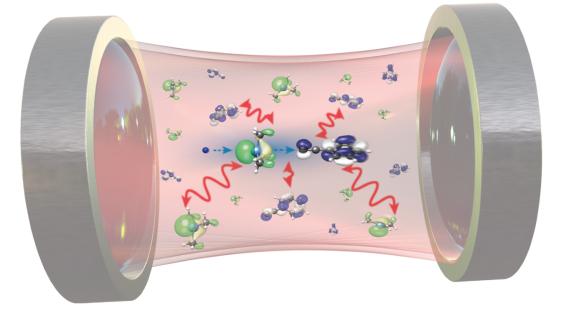


1. Cavity-modifications of the groundstate are minuscule, and would only appear, when going beyond RWA approximation. => Ground-state (chemistry) virtually not affected by optical cavities.
2. In typical experimental setups, a macroscopic amount of molecules is collectively coupled  $N \ggg 1$ . While this suggests that the collective coupling strength / Rabi splitting  $\Omega\sqrt{N}$  may be significant, the coupling between a single molecule and the cavity field must be vanishingly small  $\Omega \approx 0$ . This suggests that chemistry does not change locally under collective strong coupling.

**However: Experiments proof the contrary! What is going on?**

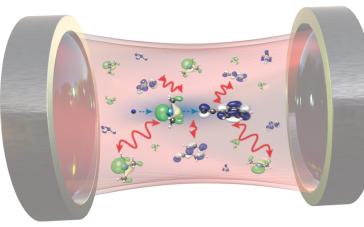


# Overview of Standard 2-level Models



Model	Counter-rotating terms?	Number of two-level systems
Jaynes-Cummings	no	1
Tavis-Cummings	no	N
Rabi	yes	1
Dicke	yes	N

**Remark:** Collective models (Dicke and TC) are far more complex than what we discussed for our approximate solution. More accurate / extensive treatments reveal for example **dark states** that do not couple to the cavity field. Or the models entail phase transitions (e.g. super-radiant phase) and quantum entanglement etc. In particular, the relevance / existence of a super-radiant phase in real systems is heavily disputed and an active field of research.



# Summary and Conclusion

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1. JC model reveals vacuum Rabi splitting (and oscillations) for strongly coupled light and matter. Rabi splitting is a typical experimental observable, which can be fitted to the JC/TC model for interpretation.
2. Counter-rotating light-matter interaction are often discarded in models (and also some ab-initio descriptions). Warning: This is believed to remove many physically/chemically relevant effects (active area of research)
3. Increasing the number of 2-level emitters, increases the Rabi-splitting. Light-matter coupling strength can collectively be increased.