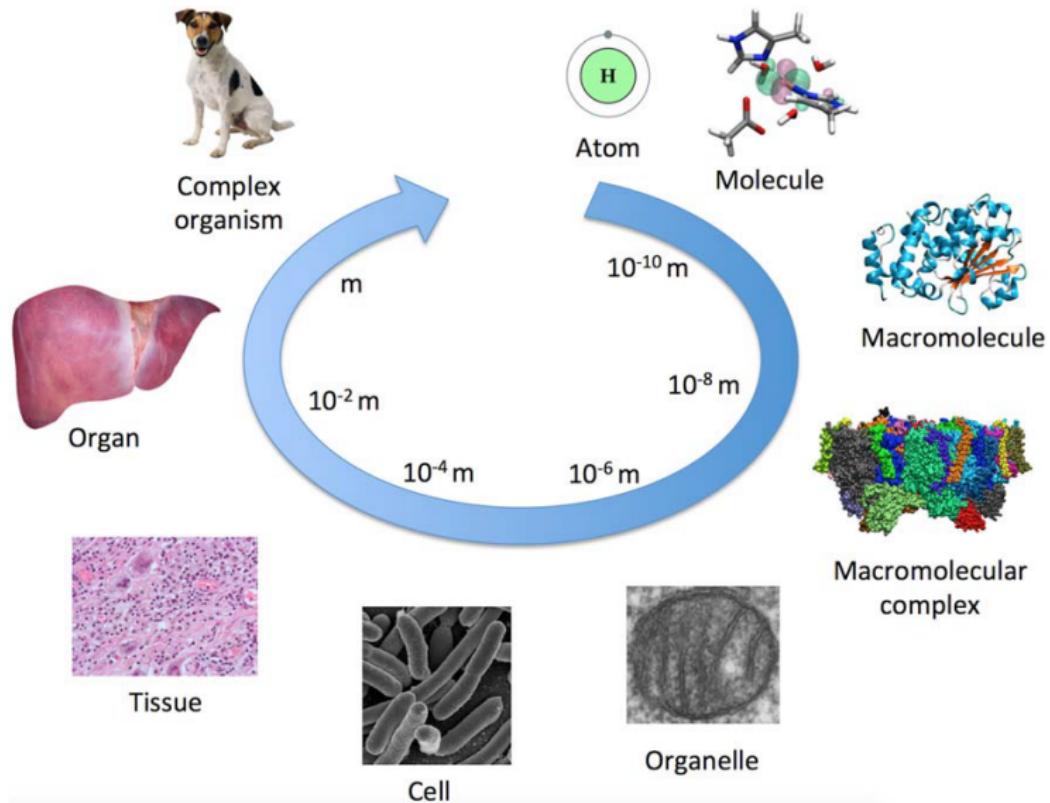


Hybrid particle-field molecular dynamics for biological systems

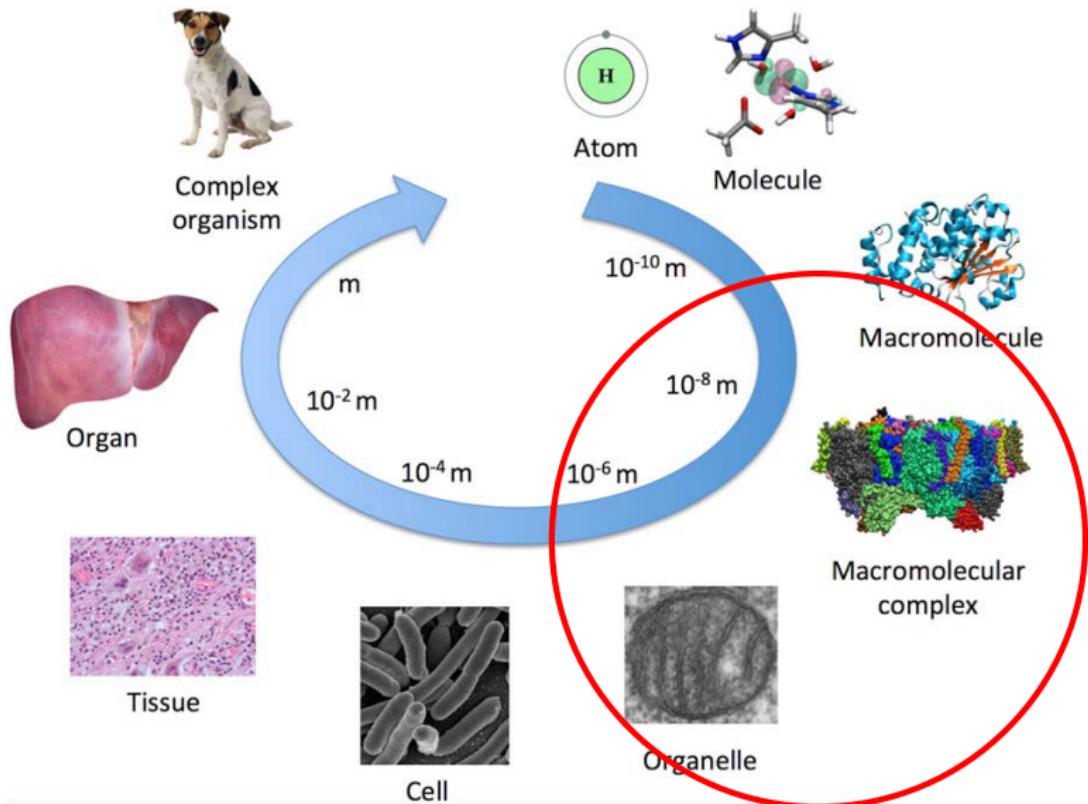
Sigbjørn Løland Bore
University of Oslo, Norway

Monday, September 30, 2019

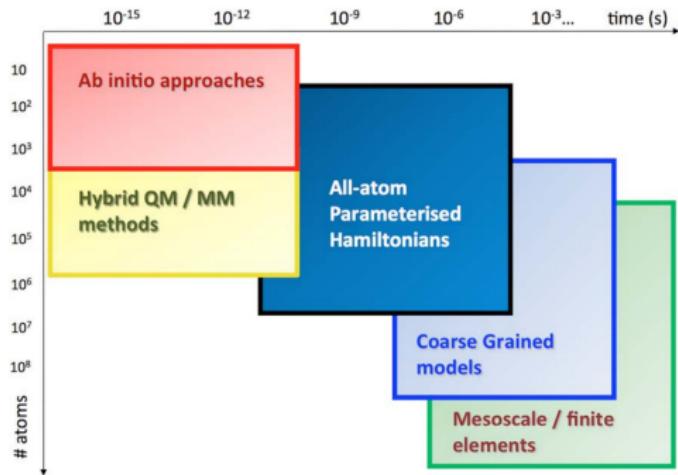
Biological scales



Biological scales

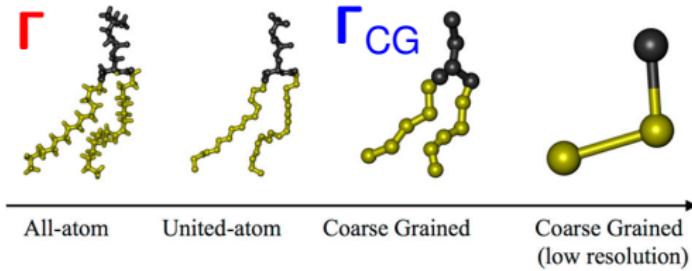
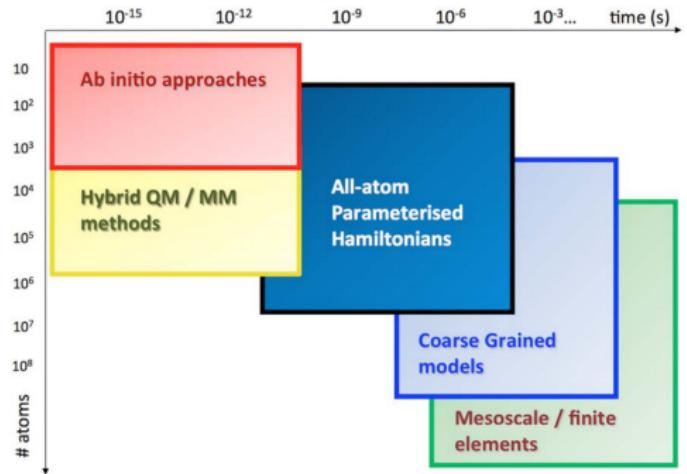


Coarse-graining

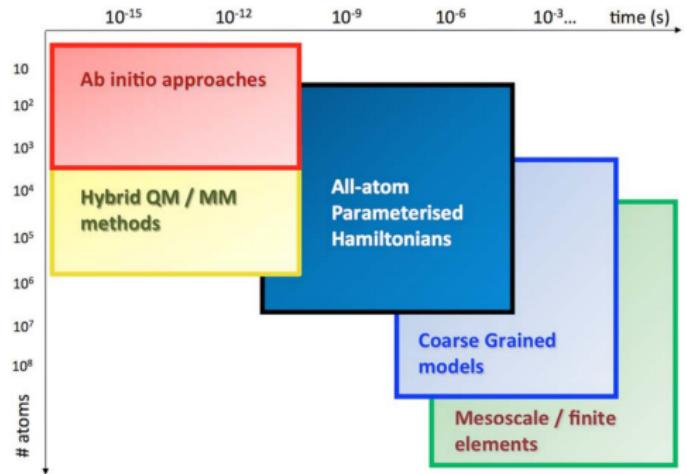


*M. Cascella and S. Vanni, Chem. Modell., (2016), 12
T. A. Soares, et al., J. Phys. Chem. Lett., (2017) 8(15)*

Coarse-graining



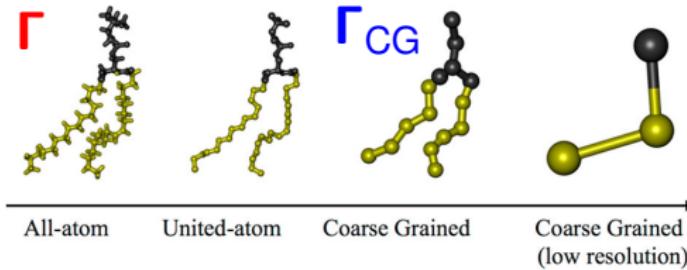
Coarse-graining



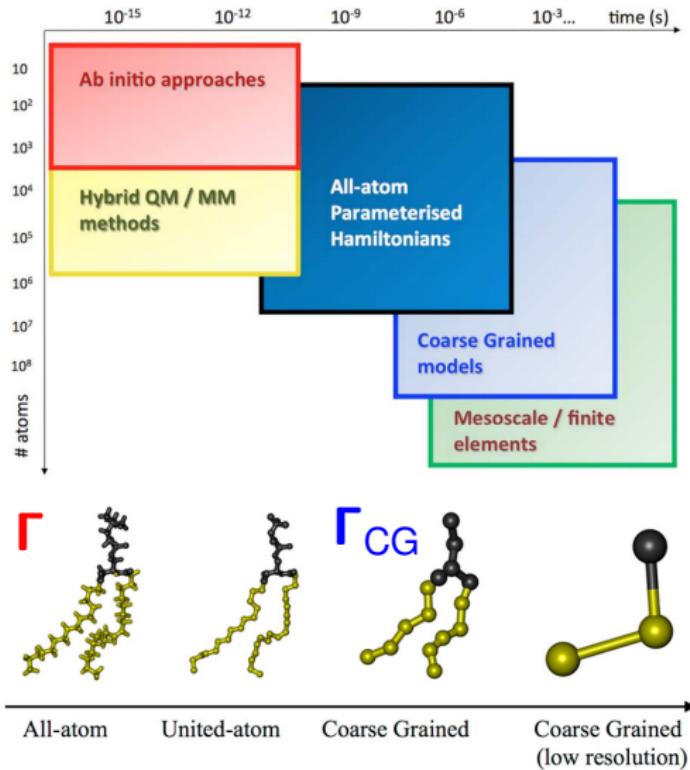
$$Z = \int d\Gamma e^{-\beta H(\Gamma)}$$

↓

$$Z \simeq \int d\Gamma_{CG} e^{-\beta H(\Gamma_{CG})}$$



Coarse-graining



$$Z = \int d\Gamma e^{-\beta H(\Gamma)}$$

↓

$$Z \simeq \int d\Gamma_{CG} e^{-\beta H(\Gamma_{CG})}$$

- ▶ Fewer degrees of freedom
- ▶ Speedup of dynamics

Hybrid Particle-field method

$$H(\{\mathbf{r}\}) = \sum_{m=1}^{N_{\text{mol}}} \underbrace{H_0(\{\mathbf{r}_m\})}_{\text{Intra-molecular}} + \underbrace{W[\{\phi(\mathbf{r})\}]}_{\text{Inter-molecular}}$$

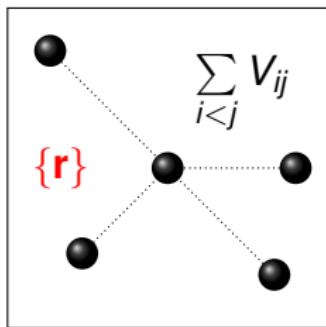
$\{\phi\} \equiv \{\phi_1, \dots, \phi_M\}$, particle-type number densities.

Hybrid Particle-field method

$$H(\{\mathbf{r}\}) = \sum_{m=1}^{N_{\text{mol}}} \underbrace{H_0(\{\mathbf{r}_m\})}_{\text{Intra-molecular}} + \underbrace{W[\{\phi(\mathbf{r})\}]}_{\text{Inter-molecular}}$$

Intermolecular interactions

Particle-particle



$\{\mathbf{r}\} \equiv \{\mathbf{r}_1, \dots, \mathbf{r}_N\}$, *particle positions*.

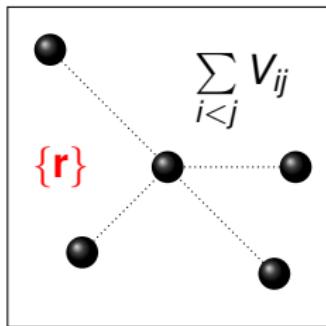
$\{\phi\} \equiv \{\phi_1, \dots, \phi_M\}$, *particle-type number densities*.

Hybrid Particle-field method

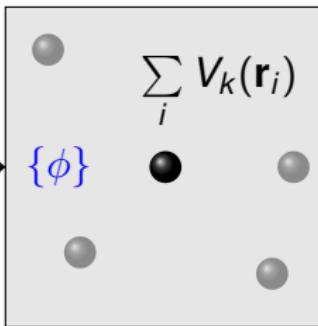
$$H(\{\mathbf{r}\}) = \sum_{m=1}^{N_{\text{mol}}} \underbrace{H_0(\{\mathbf{r}_m\})}_{\text{Intra-molecular}} + \underbrace{W[\{\phi(\mathbf{r})\}]}_{\text{Inter-molecular}}$$

Intermolecular interactions

Particle-particle



Particle-field



$\{\mathbf{r}\} \equiv \{\mathbf{r}_1, \dots, \mathbf{r}_N\}$, *particle positions*.

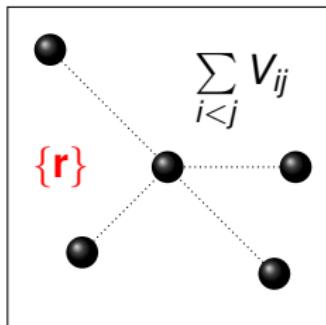
$\{\phi\} \equiv \{\phi_1, \dots, \phi_M\}$, *particle-type number densities*.

Hybrid Particle-field method

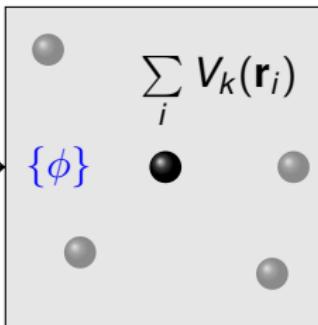
$$H(\{\mathbf{r}\}) = \sum_{m=1}^{N_{\text{mol}}} \underbrace{H_0(\{\mathbf{r}_m\})}_{\text{Intra-molecular}} + \underbrace{W[\{\phi(\mathbf{r})\}]}_{\text{Inter-molecular}}$$

Intermolecular interactions

Particle-particle



Particle-field



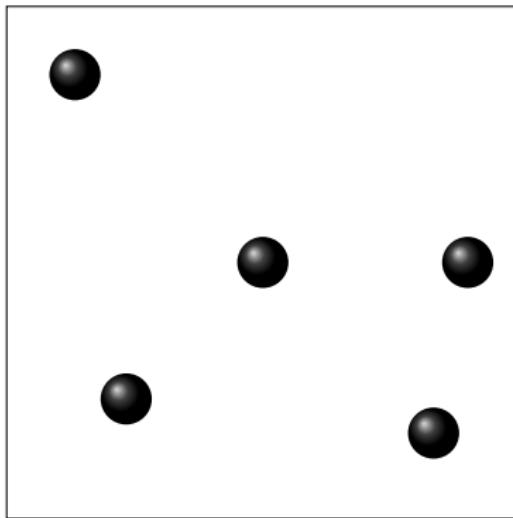
$$V_k(\mathbf{r}) = \frac{\delta W[\{\phi\}]}{\delta \phi_k(\mathbf{r})}$$

$$\mathbf{F}_i = -\nabla_i V_k(\mathbf{r}_i)$$

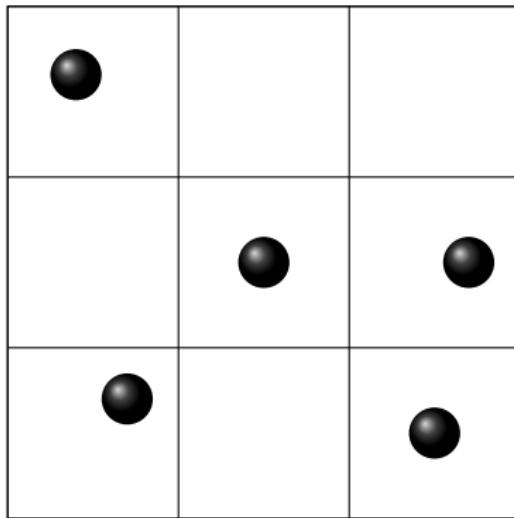
$\{\mathbf{r}\} \equiv \{\mathbf{r}_1, \dots, \mathbf{r}_N\}$, *particle positions*.

$\{\phi\} \equiv \{\phi_1, \dots, \phi_M\}$, *particle-type number densities*.

Force computation: Particle-mesh

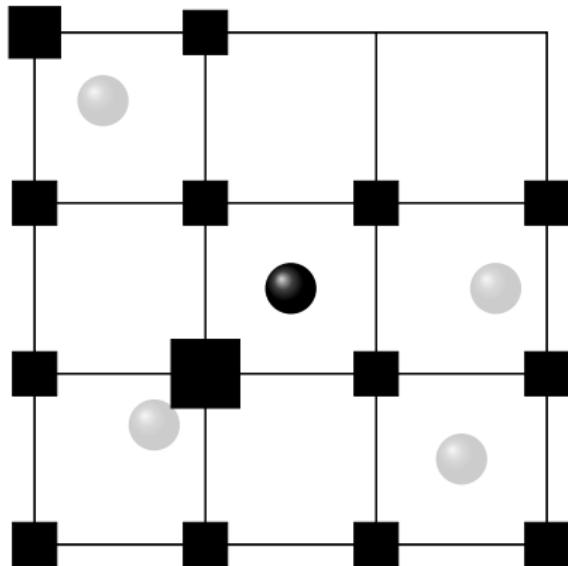


Force computation: Particle-mesh



1) Linear interpolation:
 $\{\mathbf{r}\} \rightarrow \{\phi_{nml}\}$

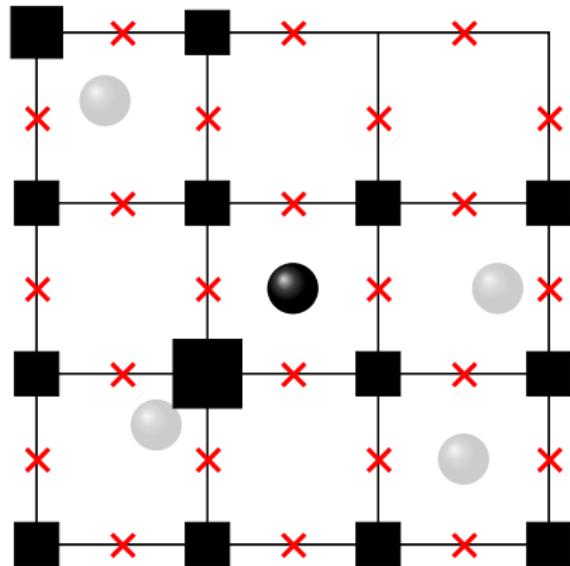
Force computation: Particle-mesh



1) Linear interpolation:
 $\{\mathbf{r}\} \rightarrow \{\phi_{nml}\}$

■ : ϕ_{nml}

Force computation: Particle-mesh



1) Linear interpolation:

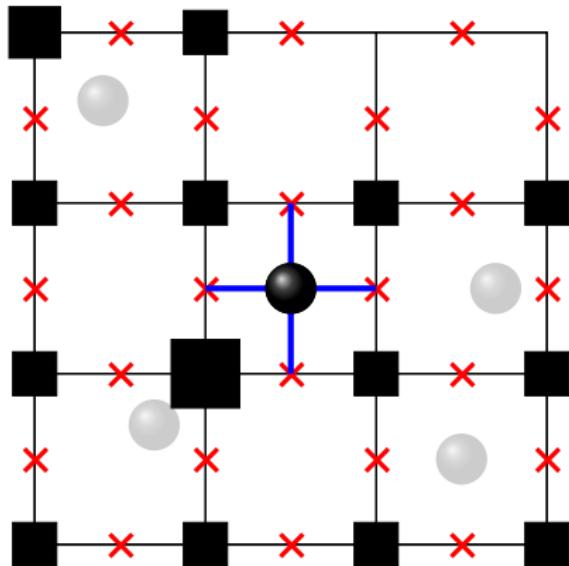
$$\{\mathbf{r}\} \rightarrow \{\phi_{nml}\}$$

2) Finite-differences:

$$\{\phi_{nml}\} \rightarrow \{\nabla \phi_{nml}\} \rightarrow \{\nabla V_{nml}\}$$

■ : ϕ_{nml} ✕ : ∇V_{nml}

Force computation: Particle-mesh



1) Linear interpolation:

$$\{\mathbf{r}\} \rightarrow \{\phi_{nml}\}$$

2) Finite-differences:

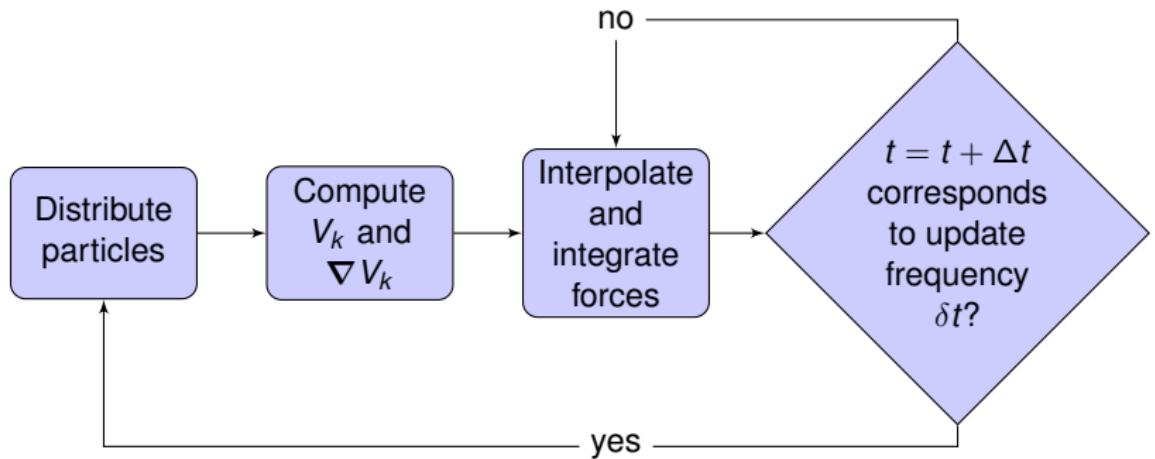
$$\{\phi_{nml}\} \rightarrow \{\nabla \phi_{nml}\} \rightarrow \{\nabla V_{nml}\}$$

3) Force interpolation:

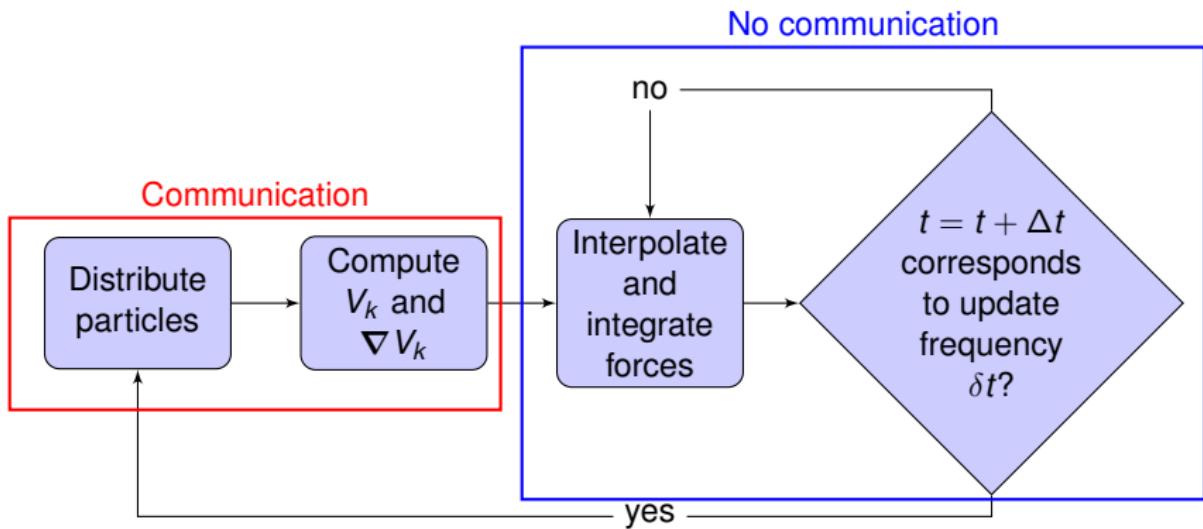
$$\{\nabla V_{nml}\} \rightarrow \mathbf{F}_i$$

■ : ϕ_{nml} ✗ : ∇V_{nml}

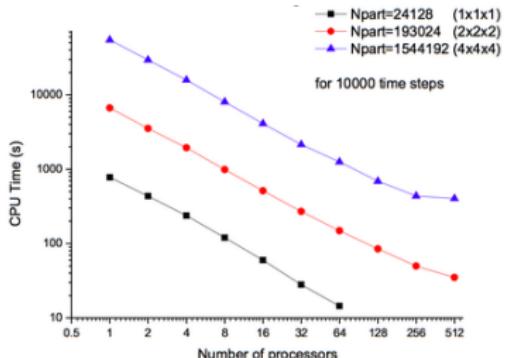
Implementation and parallelization



Implementation and parallelization



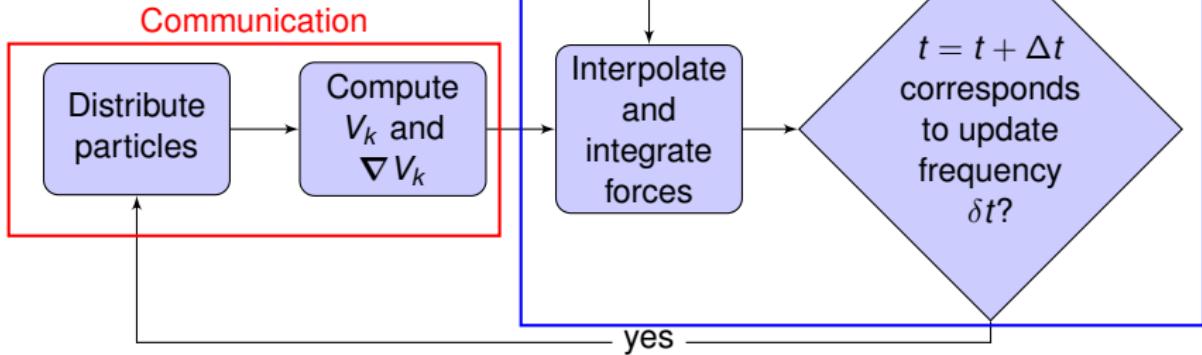
Implementation and parallelization



Excellent scaling for small and large systems!



No communication



Interaction energy: Polymer-theory

$$W[\phi] = \int d\mathbf{r} \frac{1}{\rho_0} \left(\sum_{k\ell} \frac{\tilde{\chi}_{k\ell}}{2} \phi_k(\mathbf{r}) \phi_\ell(\mathbf{r}) + \frac{1}{\kappa} \left(\sum_\ell \phi_\ell(\mathbf{r}) - \rho_0 \right)^2 \right)$$

ρ_0 : density-parameter related to the volume per bead.

Interaction energy: Polymer-theory

$$W[\phi] = \int d\mathbf{r} \frac{1}{\rho_0} \left(\underbrace{\sum_{k\ell} \frac{\tilde{\chi}_{k\ell}}{2} \phi_k(\mathbf{r}) \phi_\ell(\mathbf{r})}_{\text{Mixing}} + \frac{1}{\kappa} \left(\sum_\ell \phi_\ell(\mathbf{r}) - \rho_0 \right)^2 \right)$$

$\tilde{\chi}_{k\ell} > 0 \rightarrow$ Likes not to mix

$\tilde{\chi}_{k\ell} \leq 0 \rightarrow$ Likes to mix

ρ_0 : density-parameter related to the volume per bead.

Interaction energy: Polymer-theory

$$W[\phi] = \int d\mathbf{r} \frac{1}{\rho_0} \left(\underbrace{\sum_{k\ell} \frac{\tilde{\chi}_{k\ell}}{2} \phi_k(\mathbf{r}) \phi_\ell(\mathbf{r})}_{\text{Mixing}} + \underbrace{\frac{1}{\kappa} \left(\sum_\ell \phi_\ell(\mathbf{r}) - \rho_0 \right)^2}_{\text{Compressibility}} \right)$$

$\tilde{\chi}_{k\ell} > 0 \rightarrow$ Likes not to mix

$\tilde{\chi}_{k\ell} \leq 0 \rightarrow$ Likes to mix

$\kappa \sim 0 \rightarrow$ incompressible

$\kappa \gg 0 \rightarrow$ very compressible

ρ_0 : density-parameter related to the volume per bead.

Interaction energy: Polymer-theory

$$W[\phi] = \int d\mathbf{r} \frac{1}{\rho_0} \left(\underbrace{\sum_{k\ell} \frac{\tilde{\chi}_{k\ell}}{2} \phi_k(\mathbf{r}) \phi_\ell(\mathbf{r})}_{\text{Mixing}} + \underbrace{\frac{1}{\kappa} \left(\sum_\ell \phi_\ell(\mathbf{r}) - \rho_0 \right)^2}_{\text{Compressibility}} \right)$$

$\tilde{\chi}_{k\ell} > 0 \rightarrow$ Likes not to mix

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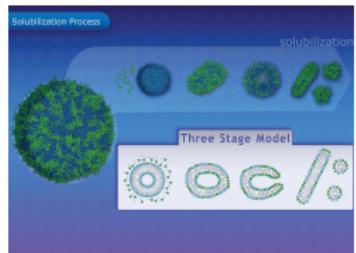
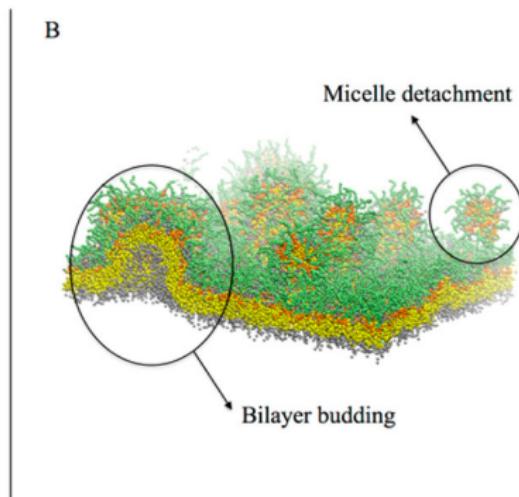
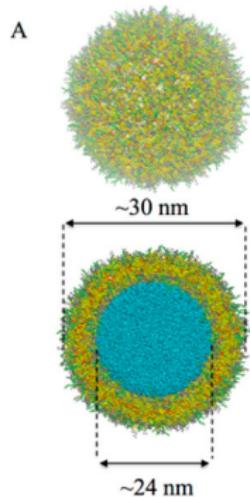
$\kappa \gg 0 \rightarrow$ very compressible

$$\xrightarrow{\text{Net effect}} V_k(\mathbf{r}) = \frac{1}{\rho_0} \left(\sum_\ell \tilde{\chi}_{k\ell} \phi_\ell(\mathbf{r}) + \frac{1}{\kappa} \left(\sum_\ell \phi_\ell(\mathbf{r}) - \rho_0 \right) \right)$$

ρ_0 : density-parameter related to the volume per bead.

Prototypic application

Large lipid vesicles interacting with surfactants



Measuring research from the group of Dr Giuseppe Minervini,
University of Salerno, Italy, and Dr S. Zanchetta,
University of Salerno, Salerno, Italy
Biomimetic solubilization mechanism by Tiberio S. 2020
Accepted manuscript online before print

The solubilization of a lipid vesicle by Tiberio S. 2020 is investigated at molecular resolution using molecular dynamics simulations. The solubilization mechanism is explained by the hybrid particle-field approach. It is possible to monitor the whole process, and to correlate the three stage model for the solubilization mechanism.



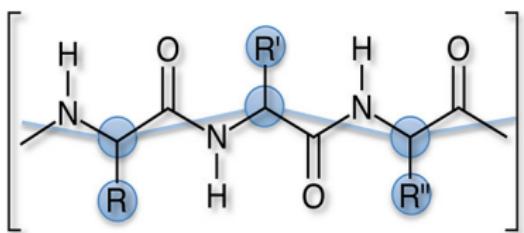
T. A. Soares *et al.*, *JPCL* 2017

A. Pizzirusso *et al.*, *PCCP* 2017

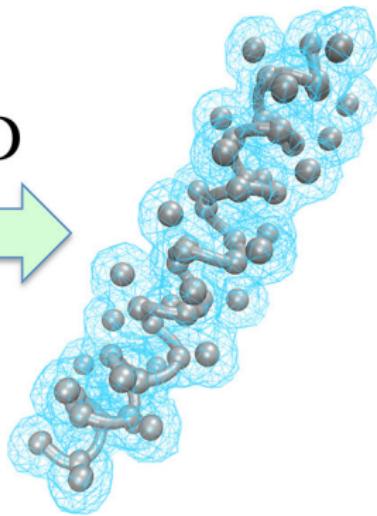


rsc.li/pccp

Hybrid Particle-Field Model for Conformational Dynamics of Peptide Chains



hPF-MD



$$H = H_0(\{\mathbf{r}\}) + W[\rho(\mathbf{r})]$$

S.L Bore et al., JCTC, 2018

Underlying model



Michele Cascella



JCTC 2008

JCTC 2010

JCTC 2013

Underlying model



Michele Cascella

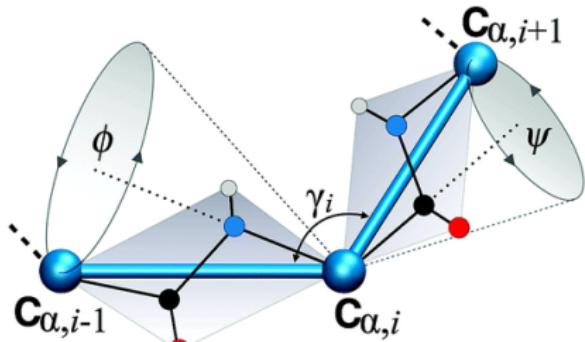


JCTC 2008

JCTC 2010

JCTC 2013

C_{α} -representation



Underlying model



Michele Cascella

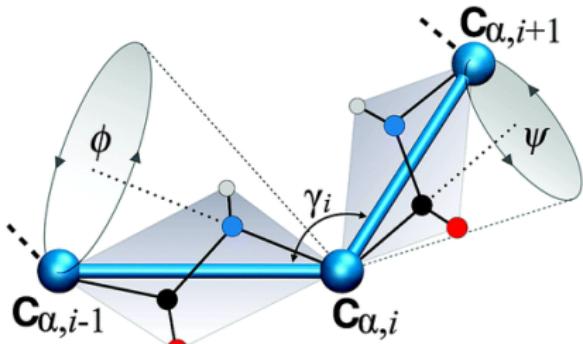


JCTC 2008

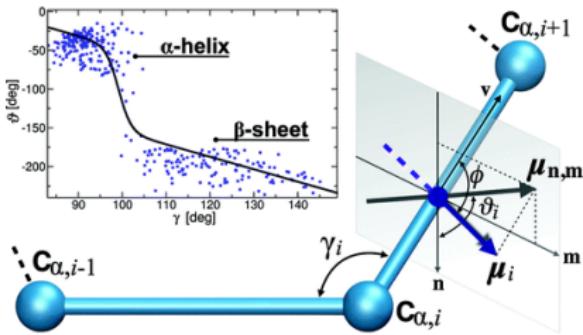
JCTC 2010

JCTC 2013

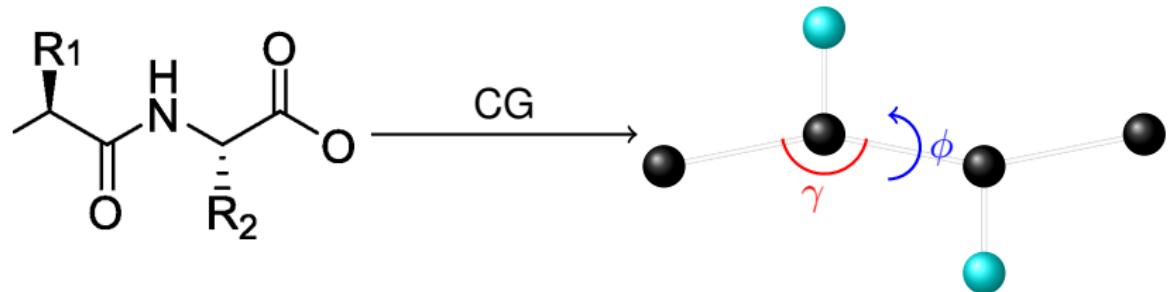
$C\alpha$ -representation



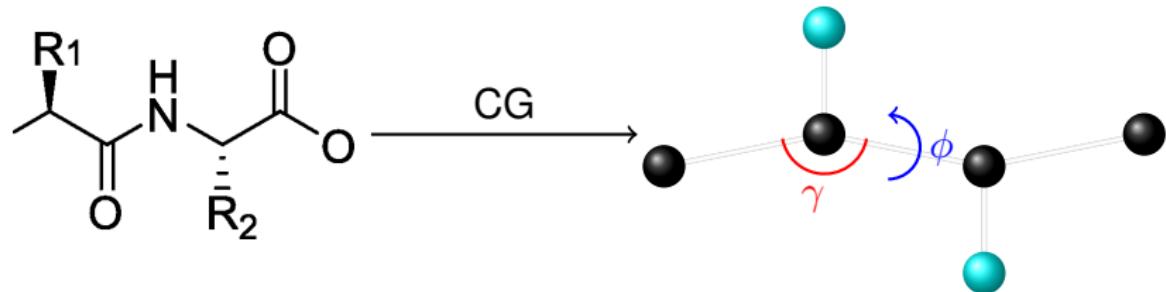
Reconstruction of dipole



Extension of the model: Two-bead representation

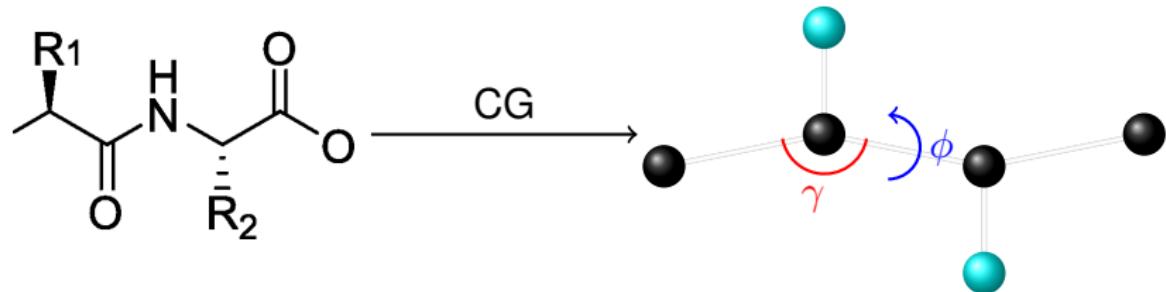


Extension of the model: Two-bead representation



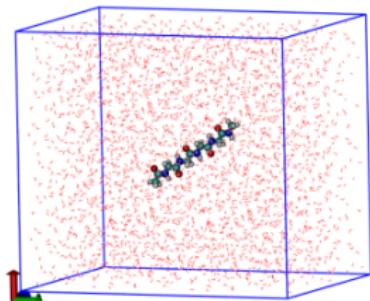
$$V(\gamma, \phi) = V_1(\gamma) + V_2(\phi)?$$

Extension of the model: Two-bead representation

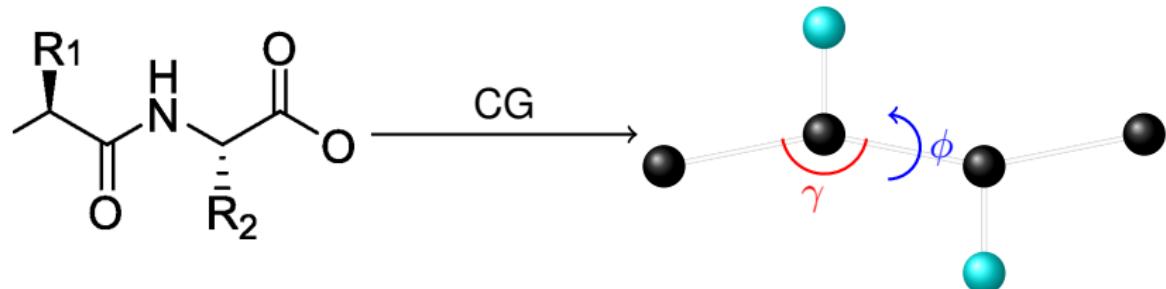


$$V(\gamma, \phi) = V_1(\gamma) + V_2(\phi)?$$

4-Alanine

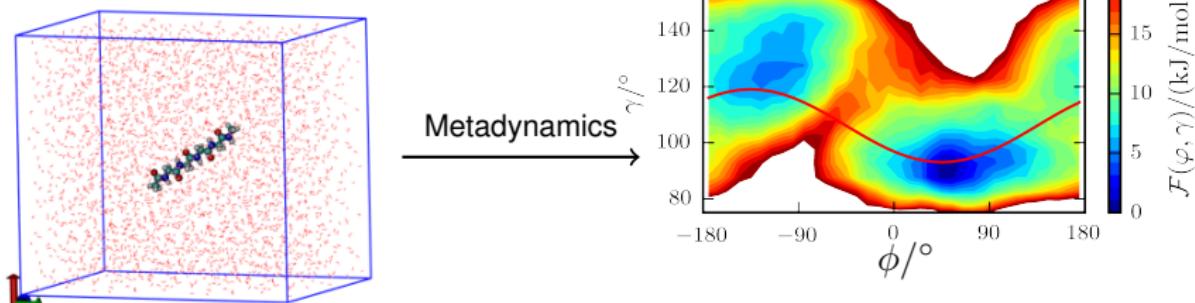


Extension of the model: Two-bead representation

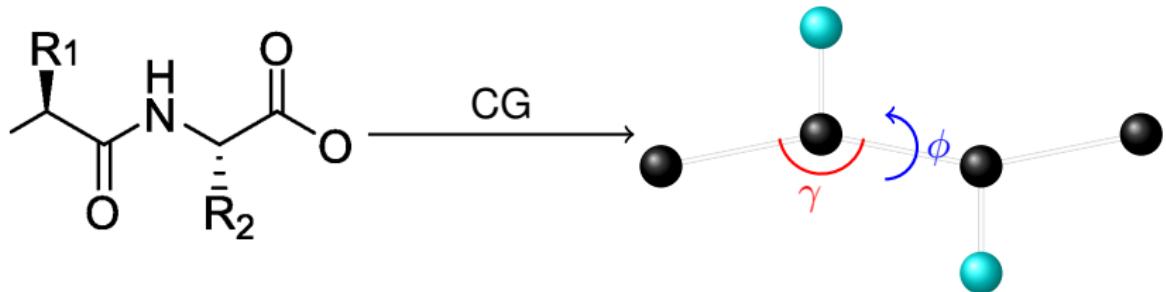


$$V(\gamma, \phi) = V_1(\gamma) + V_2(\phi) \quad \text{X}$$

4-Alanine

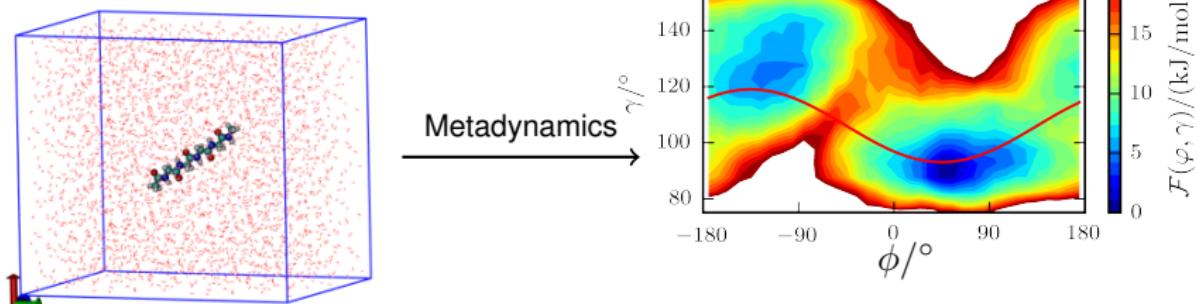


Extension of the model: Two-bead representation



$$V(\gamma, \phi) = \frac{1}{2}k(\phi)(\gamma - \gamma_0(\phi))^2 + V_{\text{prop}}(\phi, \lambda) \quad \checkmark$$

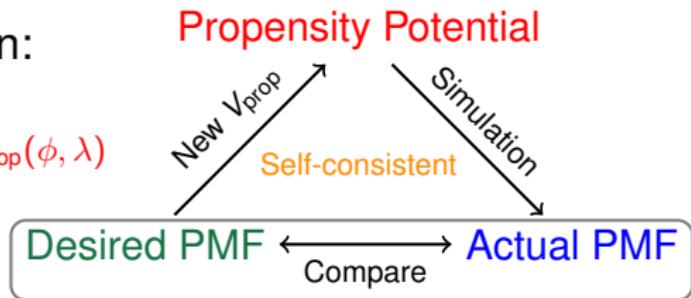
4-Alanine



Extension of the model: Propensity potential

Iterative Boltzmann-inversion:

$$V(\gamma, \phi) = \frac{1}{2}k(\phi)(\gamma - \gamma_0(\phi))^2 + V_{\text{prop}}(\phi, \lambda)$$

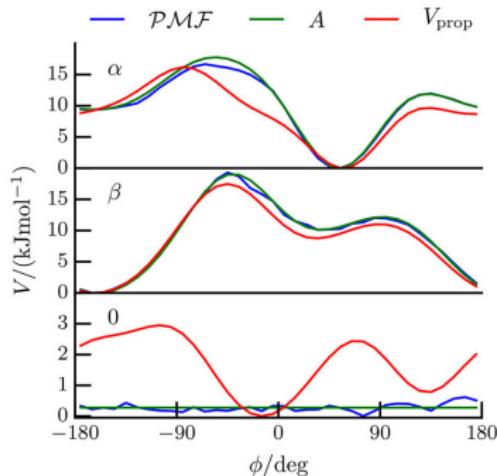
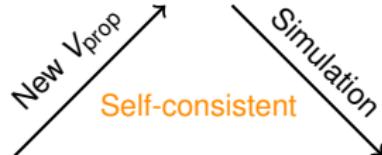


Extension of the model: Propensity potential

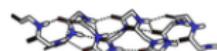
Iterative Boltzmann-inversion:

$$V(\gamma, \phi) = \frac{1}{2}k(\phi)(\gamma - \gamma_0(\phi))^2 + V_{\text{prop}}(\phi, \lambda)$$

Propensity Potential



$\lambda = -1 \implies \text{Helical}$



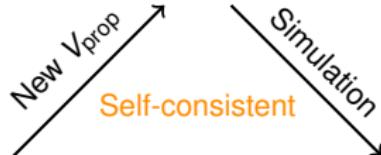
Propensity λ : Linear interpolation between $V_\alpha - V_0 - V_\beta$

Extension of the model: Propensity potential

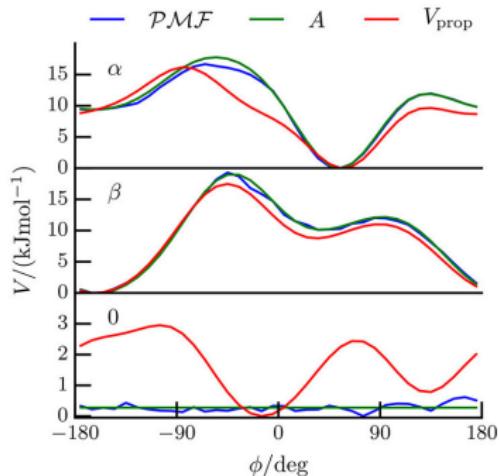
Iterative Boltzmann-inversion:

$$V(\gamma, \phi) = \frac{1}{2}k(\phi)(\gamma - \gamma_0(\phi))^2 + V_{\text{prop}}(\phi, \lambda)$$

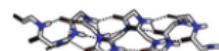
Propensity Potential



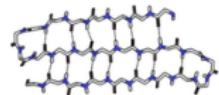
Desired PMF \longleftrightarrow Actual PMF
Compare



$\lambda = -1 \implies \text{Helical}$



$\lambda = 1 \implies \text{Extended}$

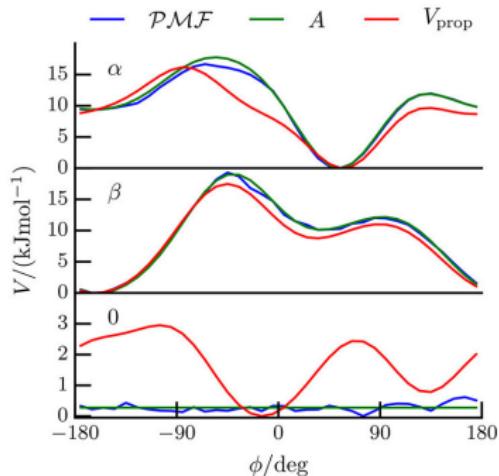
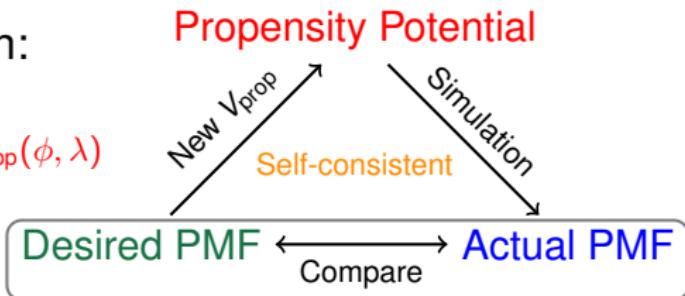


Propensity λ : Linear interpolation between $V_\alpha - V_0 - V_\beta$

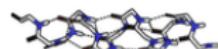
Extension of the model: Propensity potential

Iterative Boltzmann-inversion:

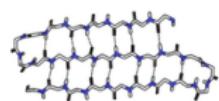
$$V(\gamma, \phi) = \frac{1}{2}k(\phi)(\gamma - \gamma_0(\phi))^2 + V_{\text{prop}}(\phi, \lambda)$$



$\lambda = -1 \Rightarrow \text{Helical}$



$\lambda = 1 \Rightarrow \text{Extended}$



$\lambda = 0 \Rightarrow \text{Random}$



Propensity λ : Linear interpolation between $V_\alpha - V_0 - V_\beta$

Extension of model: χ_{kl} -interactions

Hydrophobic



$4 \cdot \text{H}_2\text{O}$

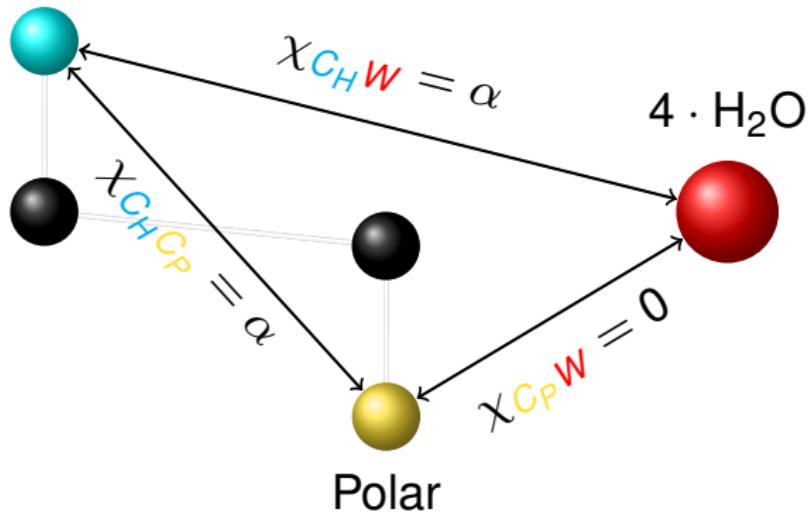


Polar

$$V_k(\mathbf{r}) = \frac{1}{\rho_0} \left(\sum_{\ell} \tilde{\chi}_{k\ell} \phi_{\ell}(\mathbf{r}) + \frac{1}{\kappa} \left(\sum_{\ell} \phi_{\ell}(\mathbf{r}) - \rho_0 \right) \right)$$

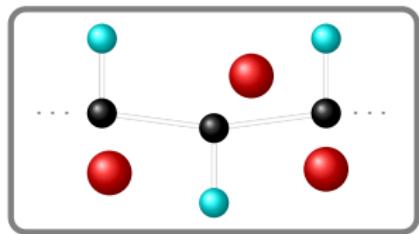
Extension of model: χ_{kl} -interactions

Hydrophobic



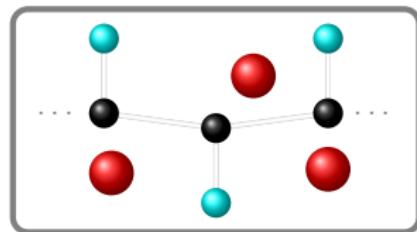
$$V_k(\mathbf{r}) = \frac{1}{\rho_0} \left(\sum_{\ell} \tilde{\chi}_{k\ell} \phi_{\ell}(\mathbf{r}) + \frac{1}{\kappa} \left(\sum_{\ell} \phi_{\ell}(\mathbf{r}) - \rho_0 \right) \right)$$

Phase-diagram: Solvated homo-poly-peptide



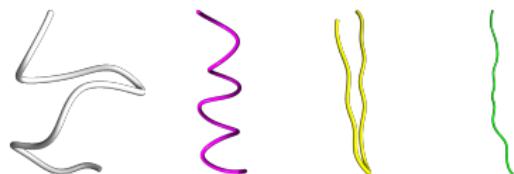
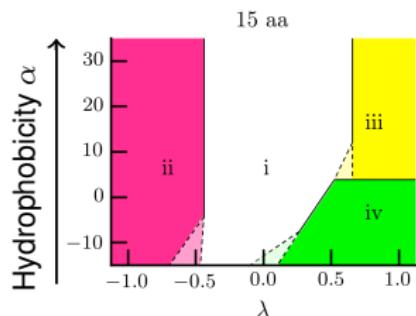
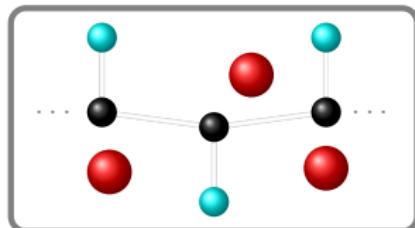
Phase-diagram: Solvated homo-poly-peptide

$\tilde{\chi}_{k\ell}$	CB	$4 \cdot H_2O$
CB	0	α
$4 \cdot H_2O$	α	0



Phase-diagram: Solvated homo-poly-peptide

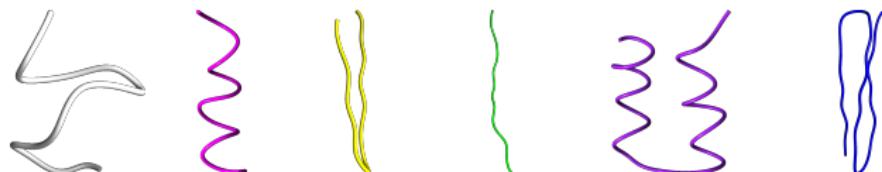
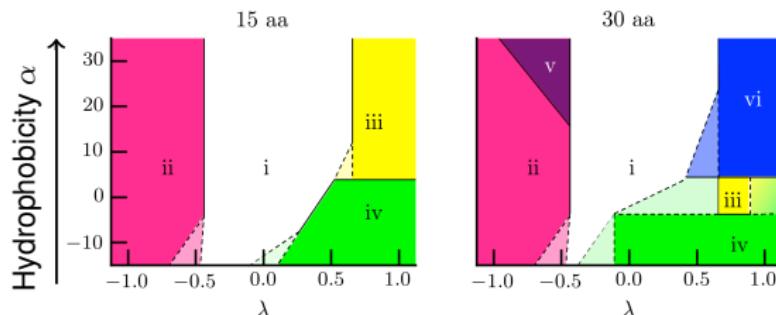
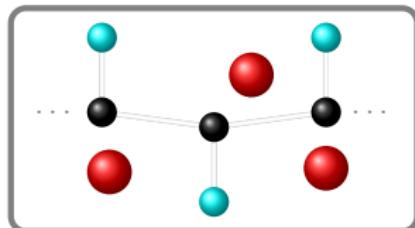
$\tilde{\chi}_{k\ell}$	CB	$4 \cdot H_2O$
CB	0	α
$4 \cdot H_2O$	α	0



i: Random coil ii: α -helix iii: β -hairpin iv: Extended

Phase-diagram: Solvated homo-poly-peptide

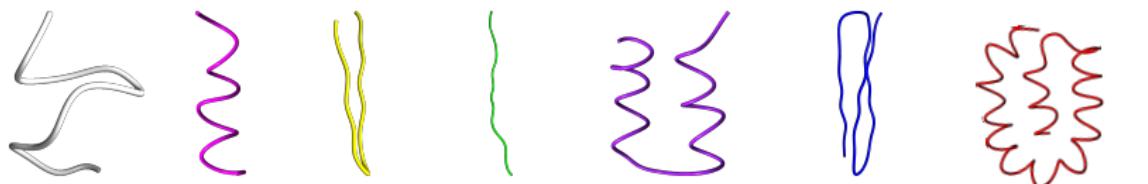
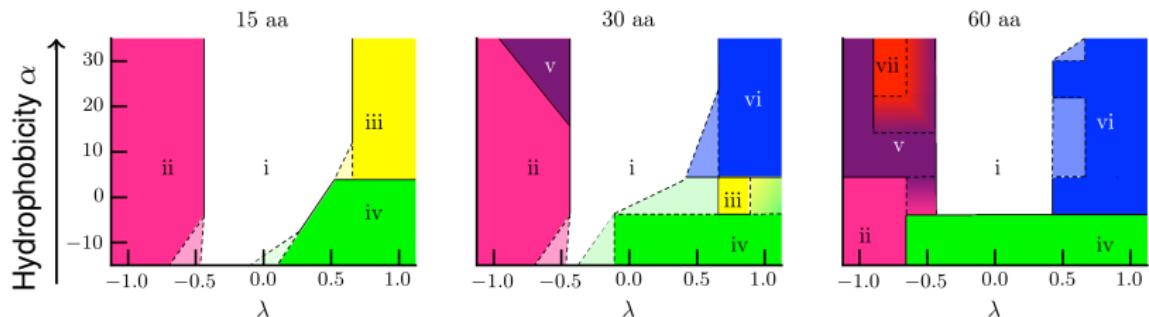
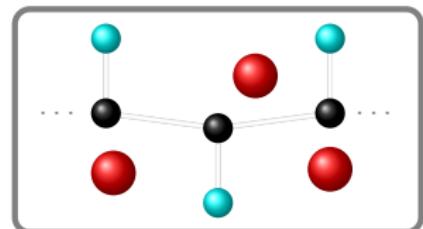
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i: Random coil ii: α -helix iii: β -hairpin iv: Extended v: Helix-coil-helix vi: β -floor/helix

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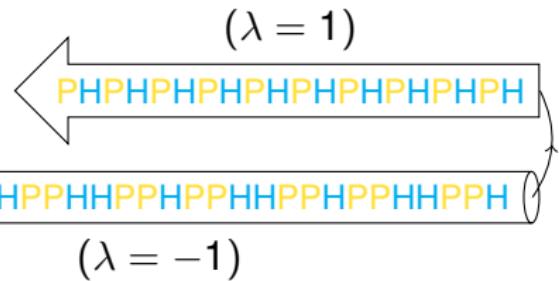
i: Random coil ii: α -helix iii: β -hairpin iv: Extended v: Helix-coil-helix vi: β -floor/helix vii: Helical bundle

Amphilic polypeptides: hydrophobic-polar

$\tilde{\chi}_{k\ell}$	CB _P	CB _H	4·H ₂ O
CB _P	0	α	0
CB _H	α	0	α
4·H ₂ O	0	α	0

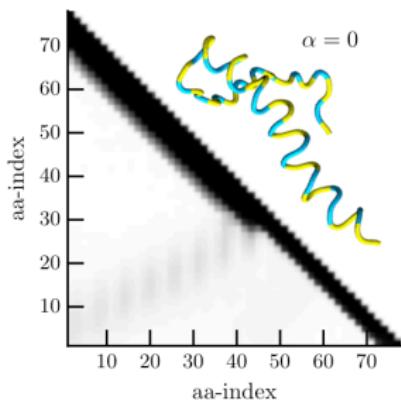
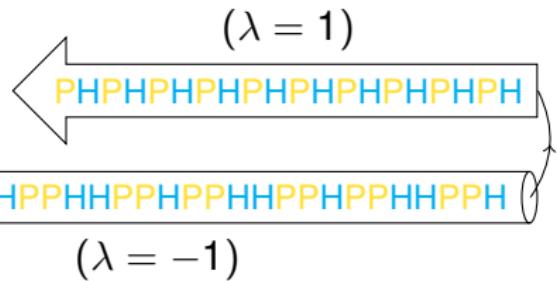
Amphiphilic polypeptides: hydrophobic-polar

$\tilde{\chi}_{k\ell}$	CB_P	CB_H	$4 \cdot H_2O$
CB_P	0	α	0
CB_H	α	0	α
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Amphiphilic polypeptides: hydrophobic-polar

$\tilde{\chi}_{k\ell}$	CB _P	CB _H	4·H ₂ O
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Amphiphilic polypeptides: hydrophobic-polar

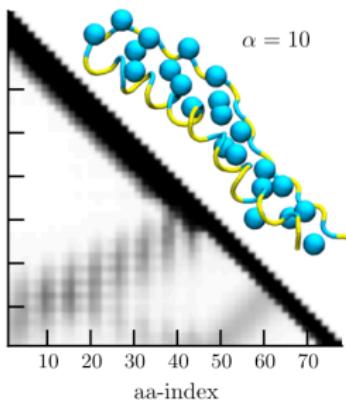
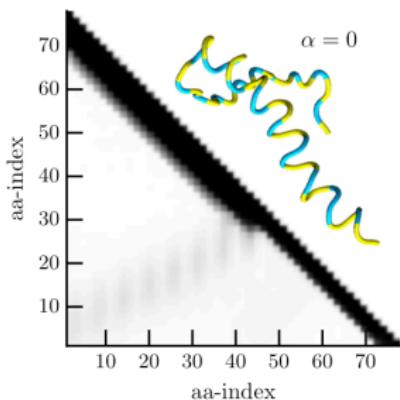
$\tilde{\chi}_{k\ell}$	CB _P	CB _H	4·H ₂ O
CB _P	0	α	0
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4·H ₂ O	0	α	0

$$(\lambda = 1)$$

PHPHPHPHPHPHPHPHPHPHPHPH

(PHPPHHPPHPPHHPPHPPHHPPHPPH) \ominus

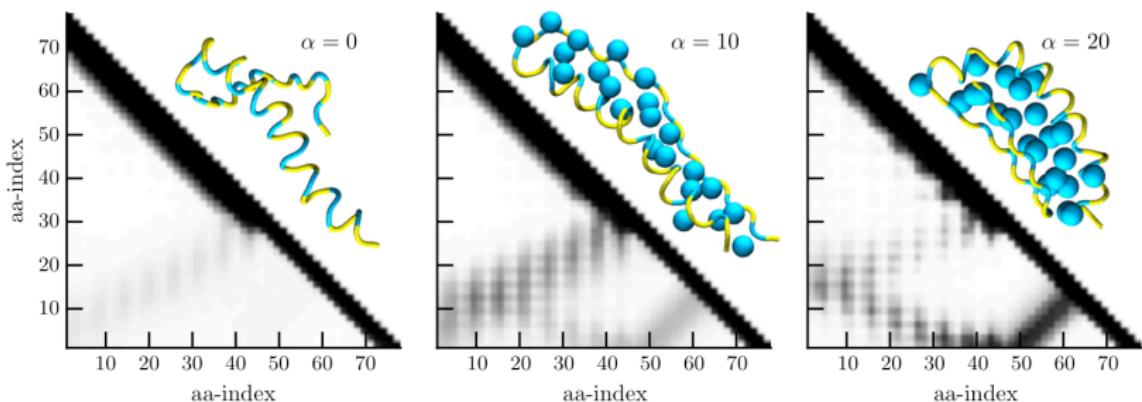
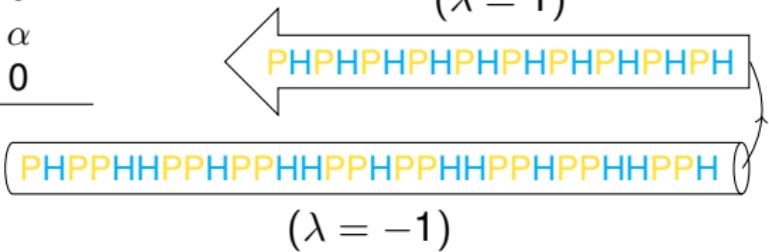
$$(\lambda = -1)$$



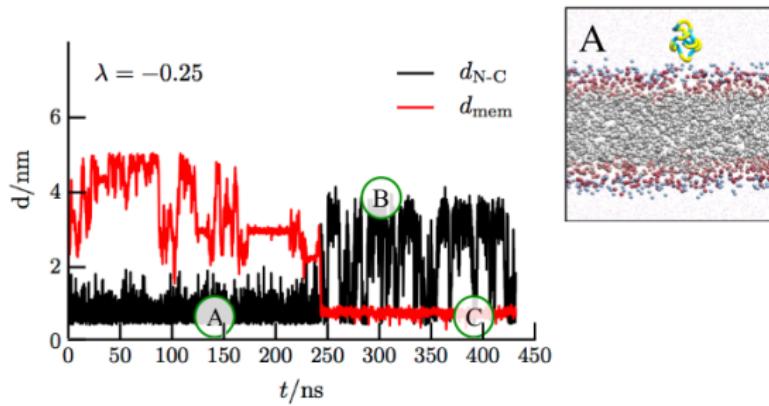
Amphiphilic polypeptides: hydrophobic-polar

$\tilde{\chi}_{k\ell}$	CB _P	CB _H	4·H ₂ O
CB _P	0	α	0
CB _H	α	0	α
4·H ₂ O	0	α	0

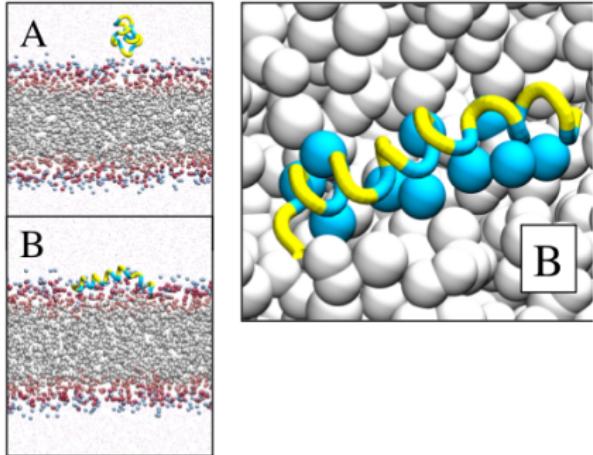
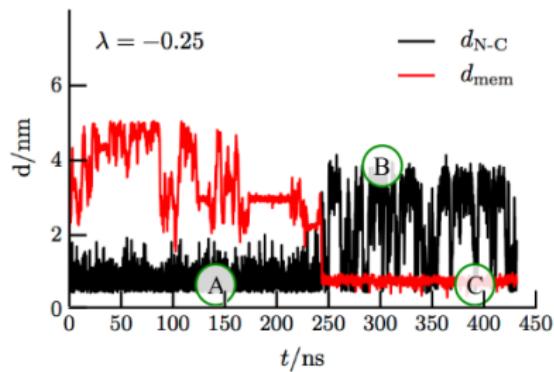
$$(\lambda = 1)$$



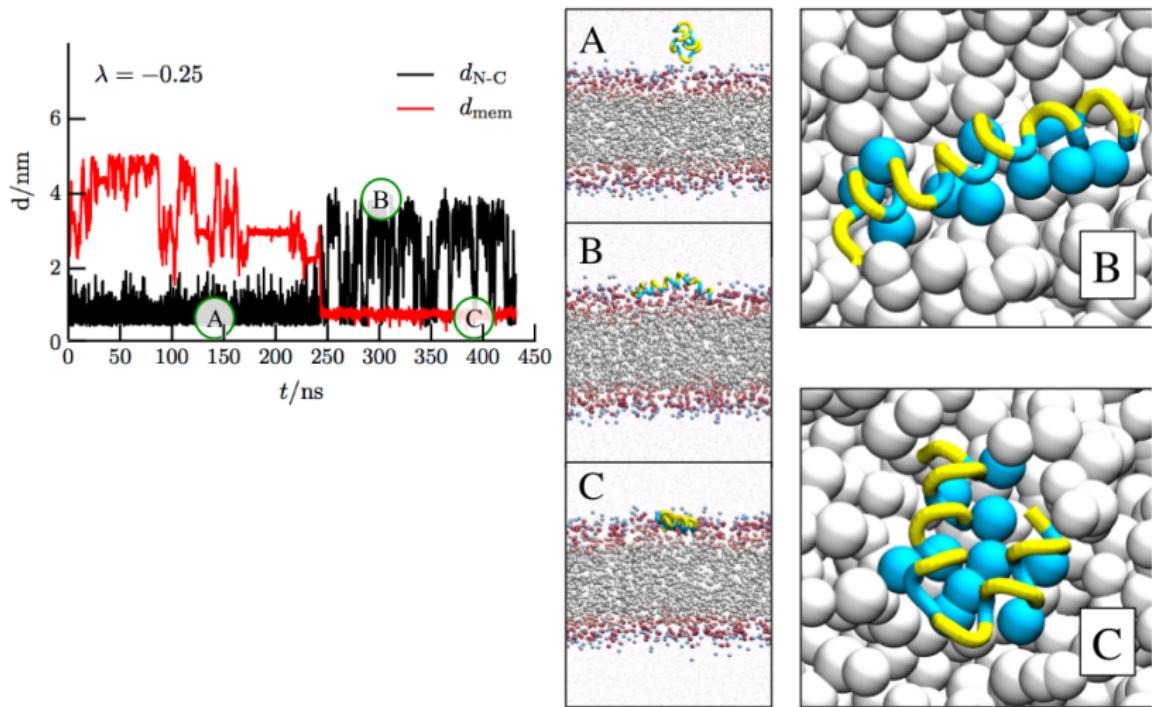
HP-polymer interacting with membrane



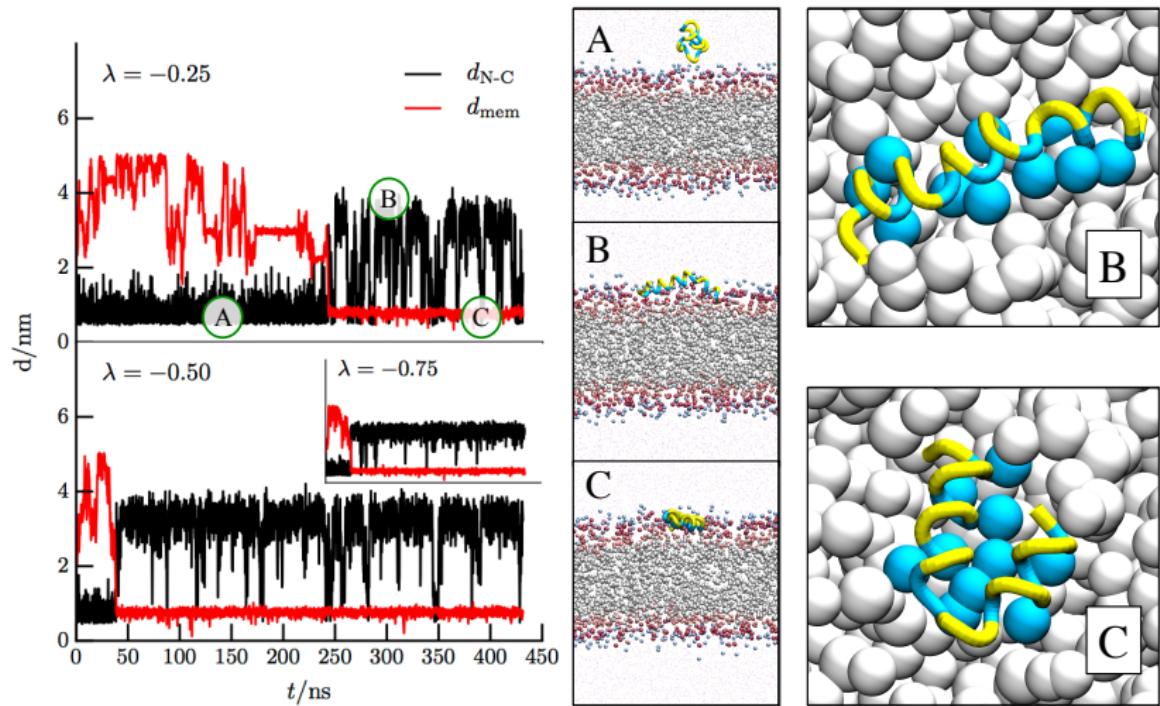
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Outlook

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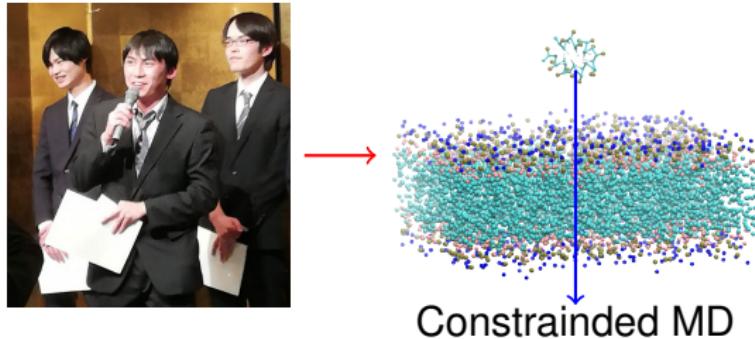
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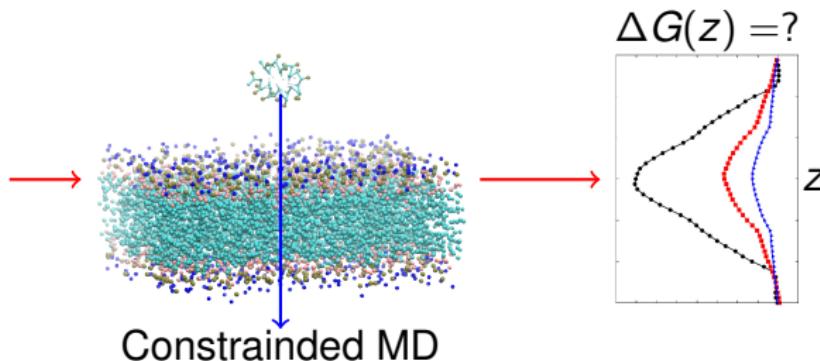
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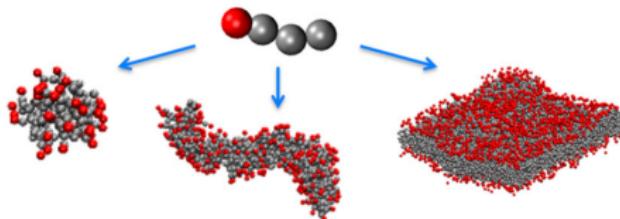
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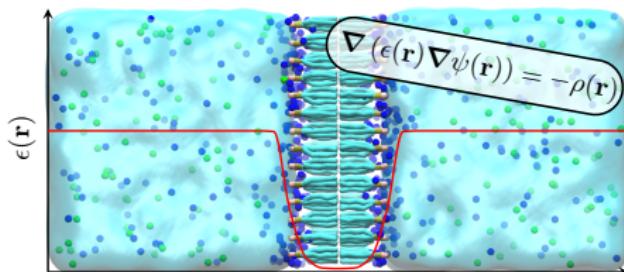


Electrostatics in Hybrid particle-field

hPF-MD + electrostatics



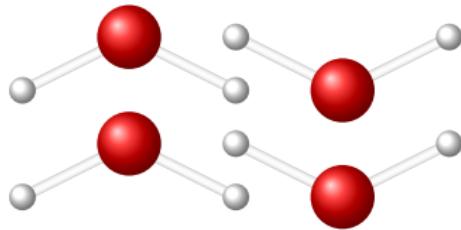
H. B. Kolli et al. JCTC, 2018



S. L. Bore et al., JCTC, 2018

Electrostatic screening: Atomistic vs coarse-grained

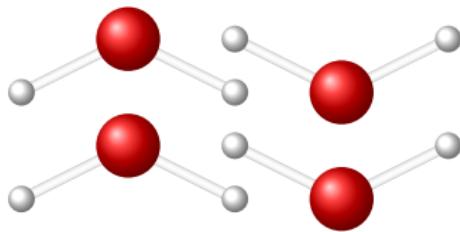
Atomistic molecular dynamics:



Electrostatic screening: Atomistic vs coarse-grained

Atomistic molecular dynamics:

- ▶ *Charges are resolved*
- ▶ *Screening is modeled directly*



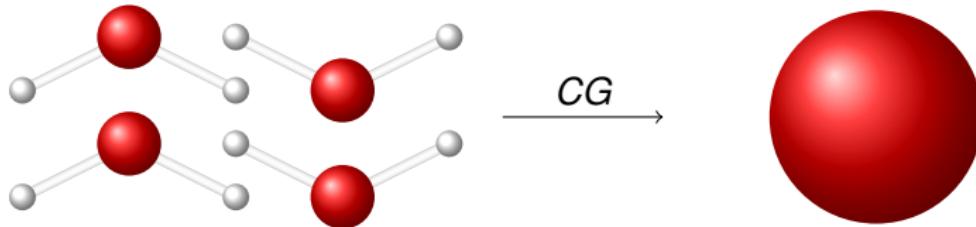
Electrostatic screening: Atomistic vs coarse-grained

Atomistic molecular dynamics:

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Coarse-grained molecular dynamics:

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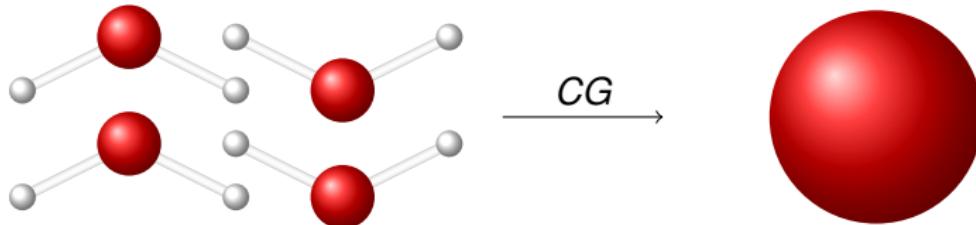
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Coarse-grained molecular dynamics:

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Idea : $\nabla \cdot (\epsilon(\mathbf{r}) \nabla \psi(\mathbf{r})) = -\rho(\mathbf{r})$ (Generalized Poisson equation)

External potential in a density dependent dielectric

Electrostatic interaction energy:

$$W_{\text{elec}}[\{\phi(\mathbf{r})\}] = \frac{1}{2} \int d\mathbf{r} \frac{\mathbf{D}(\mathbf{r}) \cdot \mathbf{D}(\mathbf{r})}{\epsilon(\mathbf{r})},$$

$\{\phi\}$: number densities. \mathbf{D} : displacement field. ϵ : permittivity.

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Potential felt by particles of type K :

$$V_{\text{ext},K}(\mathbf{r}) = \frac{\delta W_{\text{elec}}}{\delta \phi_K(\mathbf{r})}$$

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ψ : electrostatic potential. \mathbf{E} : electrostatic field ($\mathbf{E} = -\nabla \psi = \epsilon \mathbf{D}$).

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ψ : electrostatic potential. \mathbf{E} : electrostatic field ($\mathbf{E} = -\nabla \psi = \epsilon \mathbf{D}$).

Modelling: density dependence of the dielectric

Modelling: density dependence of the dielectric

Density weighted average:

$$\epsilon(\{\phi(\mathbf{r})\}) = \frac{\sum_K^M \epsilon_K \phi_K(\mathbf{r})}{\phi_0(\mathbf{r})},$$

ϵ_K : dielectric of particle type K . ϕ_0 : local total particle density.

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Potential felt by particle of type K :

$$V_{\text{ext},K}(\mathbf{r}) = q_K \psi(\mathbf{r}) - \frac{1}{2} \frac{\epsilon_K - \epsilon(\mathbf{r})}{\phi_0(\mathbf{r})} |\mathbf{E}(\mathbf{r})|^2,$$

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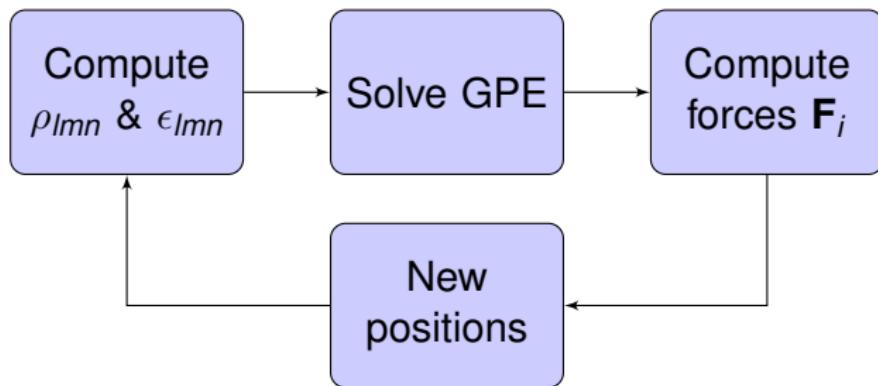
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Forces on particle of type K :

$$\mathbf{F}_K = -\nabla V_{\text{ext},K}(\mathbf{r}) = q_K \mathbf{E}(\mathbf{r}) + \frac{1}{2} \nabla \left(\frac{\epsilon_K - \epsilon(\mathbf{r})}{\phi_0(\mathbf{r})} |\mathbf{E}(\mathbf{r})|^2 \right)$$

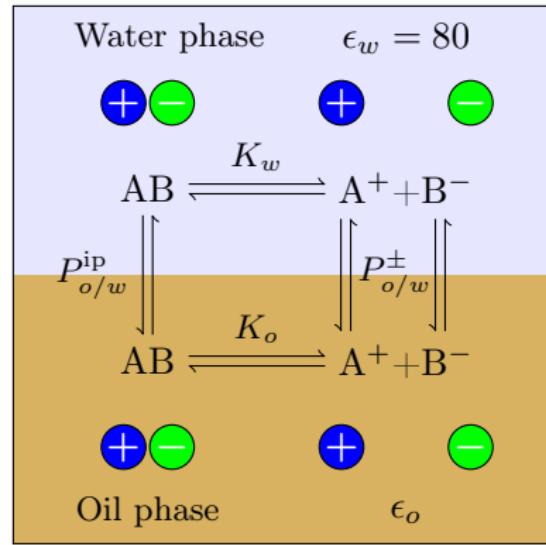
Force computation and molecular dynamics

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Partitioning of ions (1)

Ions in a phase separated oil/water mixture of ϵ_o and ϵ_w .
($RT \times \chi_{ow} = 30 \text{ kJ mol}^{-1}$)



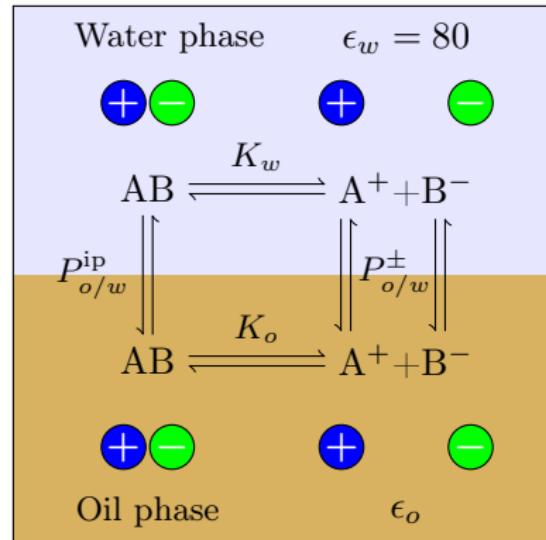
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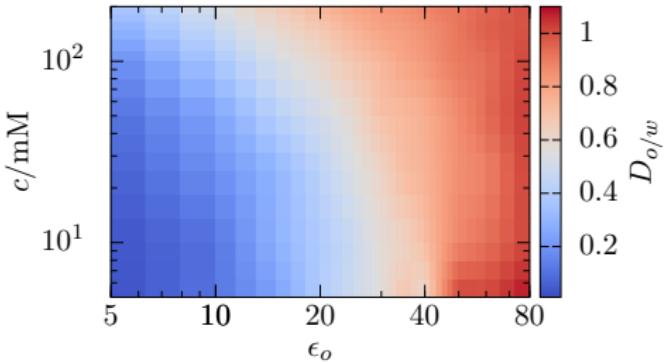
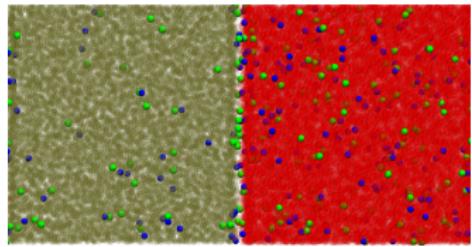
Distribution coefficient:

$$D_{o/w} = \frac{c_o}{c_w}$$

c_o and c_w : concentration of ions within each phase.



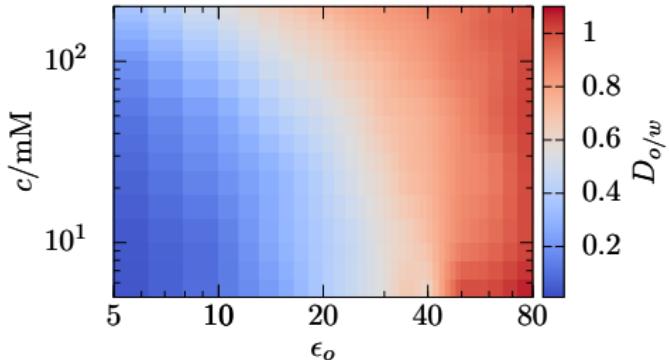
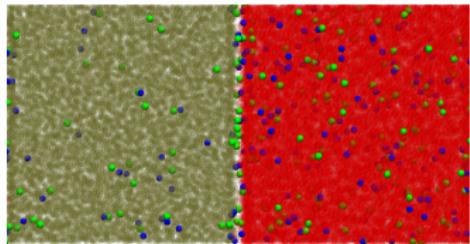
Partitioning of ions (2)



$$D_{o/w} = f(c, P_{o,w}^\pm, P_{o,w}^{\text{ip}}, K_w)$$

c: concentration of ions.

Partitioning of ions (2)

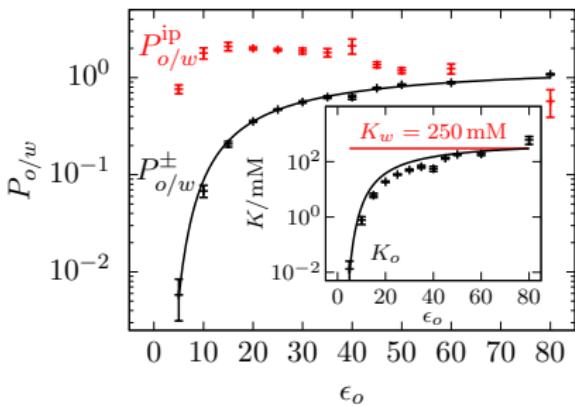


$$D_{o/w} = f(c, P_{o,w}^\pm, P_{o,w}^{\text{ip}}, K_w)$$

c: concentration of ions.

Born theory of ions:

$$\log P_{o/w}^\pm = \gamma \left(\frac{1}{\epsilon_w} - \frac{1}{\epsilon_o} \right)$$



Outlook

Applications

- ▶ Antonio De Nicola: Model for charged lipids
- ▶ Ken Schäfer: Molecular packing of charged surfactants
- ▶ Victoria Ariel Bjørnestad: Simulation and experiments on charged surfactants

Outlook

Applications

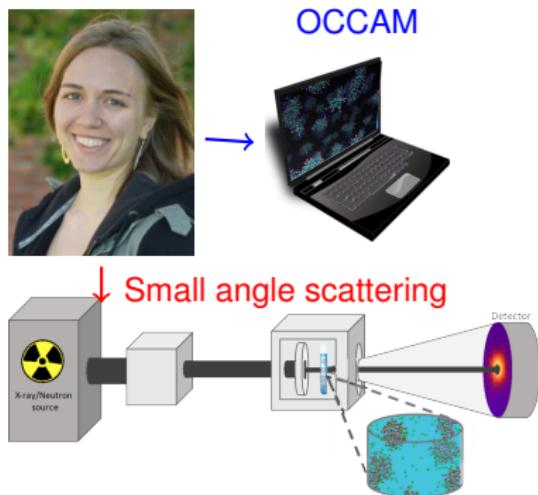
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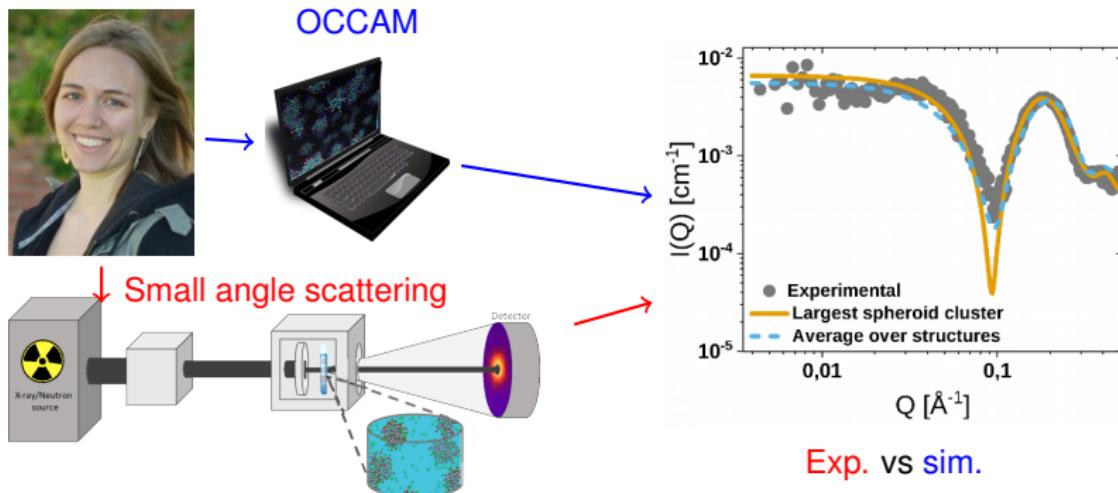
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Acknowledgements

University of Oslo, Norway:

Michele Cascella

Hima Bindu Kolli

Morten Ledum

Victoria Ariel Bjørnestad

Reidar Lund



OCCAM
Molecular Dynamics

The logo for Hylleraas consists of two overlapping circles, one light blue and one light red, followed by the word 'Hylleraas' in a bold, black, sans-serif font.

Yamagata University, Japan:

Giuseppe Milano

Antonio De Nicola

Tsudo Yamanaka

Sendai University, Japan:

Toshihiro Kawakatsu



UiO The logo for the University of Oslo (UiO) consists of the letters 'UiO' in a large, black, sans-serif font, with a red circle to the right.