

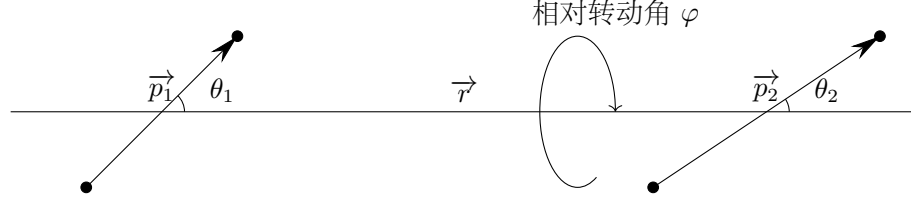
## van der Waals 相互作用

(1) 由偶极矩在外场中的电势能公式

$$U = -\vec{p} \cdot \vec{E} \quad (1)$$

以及电偶极子的电场力公式

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{3(\vec{p} \cdot \hat{r})\hat{r} - \vec{p}}{r^3} \quad (2)$$



可得其电势能大小为

$$\begin{aligned} U(\theta_1, \theta_2, \varphi, r) &= -\vec{p} \cdot \vec{E} \\ &= -\frac{1}{4\pi\epsilon_0} \frac{3(\vec{p}_1 \cdot \hat{r})\hat{r} - \vec{p}_1}{r^3} \cdot \vec{p}_2 \\ &= -\frac{1}{4\pi\epsilon_0} \frac{3(\vec{p}_1 \cdot \hat{r})\hat{r} \cdot \vec{p}_2 - \vec{p}_1 \cdot \vec{p}_2}{r^3} \\ &= -\frac{1}{4\pi\epsilon_0} \frac{3p_1 p_2 \cos \theta_1 \cos \theta_2 - p_1 p_2 (\cos \theta_1 \cos \theta_2 + \sin \theta_1 \sin \theta_2 \cos \varphi)}{r^3} \\ &= -\frac{1}{4\pi\epsilon_0} \frac{2p_1 p_2 \cos \theta_1 \cos \theta_2 - p_1 p_2 \sin \theta_1 \sin \theta_2 \cos \varphi}{r^3} \end{aligned} \quad (3)$$

(2) 按照定义

$$\bar{U} = \frac{\int_1 \int_2 U(\theta_1, \theta_2, \varphi, r) \exp\left(-\frac{U(\theta_1, \theta_2, \varphi, r)}{k_B T}\right) d\Omega_1 d\Omega_2}{\int_1 \int_2 \exp\left(-\frac{U(\theta_1, \theta_2, \varphi, r)}{k_B T}\right) d\Omega_1 d\Omega_2} \quad (4)$$

令

$$\begin{aligned} \beta &= \frac{1}{4\pi\epsilon_0} \frac{p_1 p_2}{r^3} \frac{1}{k_B T} \\ f(\Omega) &= 2 \cos \theta_1 \cos \theta_2 - \sin \theta_1 \sin \theta_2 \cos \varphi \end{aligned} \quad (5)$$

上式化为

$$\begin{aligned} \bar{U} &= \frac{\int_1 \int_2 k_B T \beta f(\Omega) \exp(\beta f(\Omega)) d\Omega_1 d\Omega_2}{\int_1 \int_2 \exp(\beta f(\Omega)) d\Omega_1 d\Omega_2} \\ &= k_B T \beta \frac{-\frac{\partial}{\partial \beta} \int_1 \int_2 \exp(\beta f(\Omega)) d\Omega_1 d\Omega_2}{\int_1 \int_2 \exp(\beta f(\Omega)) d\Omega_1 d\Omega_2} \\ &= -k_B T \beta \frac{\partial}{\partial \beta} \ln \left( \int_1 \int_2 \exp(\beta f(\Omega)) d\Omega_1 d\Omega_2 \right) \end{aligned} \quad (6)$$

易知

$$d\Omega_1 = \sin \theta_1 d\theta_1 d\varphi_1 \quad (7)$$

$$d\Omega_2 = \sin \theta_2 d\theta_2 d\varphi_2 \quad (8)$$

而由于引入了相对转动角，故

$$d\Omega_1 d\Omega_2 = \sin \theta_1 d\theta_1 \sin \theta_2 d\theta_2 d\varphi \quad (9)$$

下求积分

$$I = \int_{\theta_1=0}^{\pi} \int_{\theta_2=0}^{\pi} \int_{\varphi=0}^{2\pi} \exp(\beta f(\Omega)) \sin \theta_1 d\theta_1 \sin \theta_2 d\theta_2 d\varphi \quad (10)$$

代入题给近似

$$U(\theta_1, \theta_2, \varphi, r) \sim \frac{1}{4\pi\epsilon_0} \frac{\vec{p}_1 \cdot \vec{p}_2}{r^3} \ll k_B T \quad (11)$$

与 Tailor 展开公式

$$\exp(x) \approx 1 + x + \frac{1}{2}x^2 \quad x \ll 1 \quad (12)$$

$$\begin{aligned} I &\approx \int_{\theta_1=0}^{\pi} \int_{\theta_2=0}^{\pi} \int_{\varphi=0}^{2\pi} \left( 1 + \beta f(\Omega) + \frac{1}{2} (\beta f(\Omega))^2 \right) \sin \theta_1 d\theta_1 \sin \theta_2 d\theta_2 d\varphi \\ &= I_0 + I_1 + I_2 \end{aligned} \quad (13)$$

其中  $I_0, I_1, I_2$  分别对应其零阶，一阶，二阶近似的积分结果。易知

$$I_0 = 8\pi \quad (14)$$

$$I_1 = 0 \quad (15)$$

$$I_2 = \frac{8\pi}{3} \beta^2 \quad (16)$$

故有

$$\begin{aligned} \bar{U} &\approx -k_B T \beta \frac{\partial}{\partial \beta} \ln \left( 8\pi \left( 1 + \frac{\beta^2}{3} \right) \right) \\ &= -k_B T \beta \frac{1}{1 + \frac{\beta^2}{3}} \cdot \frac{2\beta}{3} \\ &\approx -k_B T \frac{2\beta^2}{3} \\ &= -\frac{2}{3k_B T} \left( \frac{p_1 p_2}{4\pi\epsilon_0} \right)^2 \cdot \frac{1}{r^6} \end{aligned} \quad (17)$$

(3) 在  $p_2$  处，电场强度大小为

$$|\vec{E}| = \frac{p_1}{4\pi\epsilon_0} \frac{\sqrt{1 + 3\cos^2 \theta_1}}{r^3} \quad (18)$$

故其偶极矩大小为

$$|\vec{p}| = \alpha \frac{p_1}{4\pi\epsilon_0} \frac{\sqrt{1 + 3\cos^2 \theta_1}}{r^3} \quad (19)$$

代入 (1) 式

$$U = -\alpha \frac{p_1^2}{(4\pi\epsilon_0)^2} \frac{1 + 3\cos^2 \theta_1}{r^6} \quad (20)$$

(4) 此时，有

$$\overline{U} = \frac{\int_{\theta_1=0}^{\pi} U \exp\left(-\frac{U}{k_B T}\right) \sin \theta_1 d\theta_1}{\int_{\theta_1=0}^{\pi} \exp\left(-\frac{U}{k_B T}\right) \sin \theta_1 d\theta_1} \quad (21)$$

保留至零阶项

$$\overline{U} = -\alpha \frac{2p_1^2}{(4\pi\epsilon_0)^2 r^6} \quad (22)$$