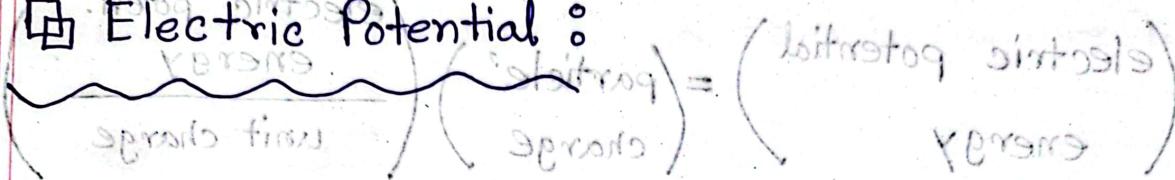
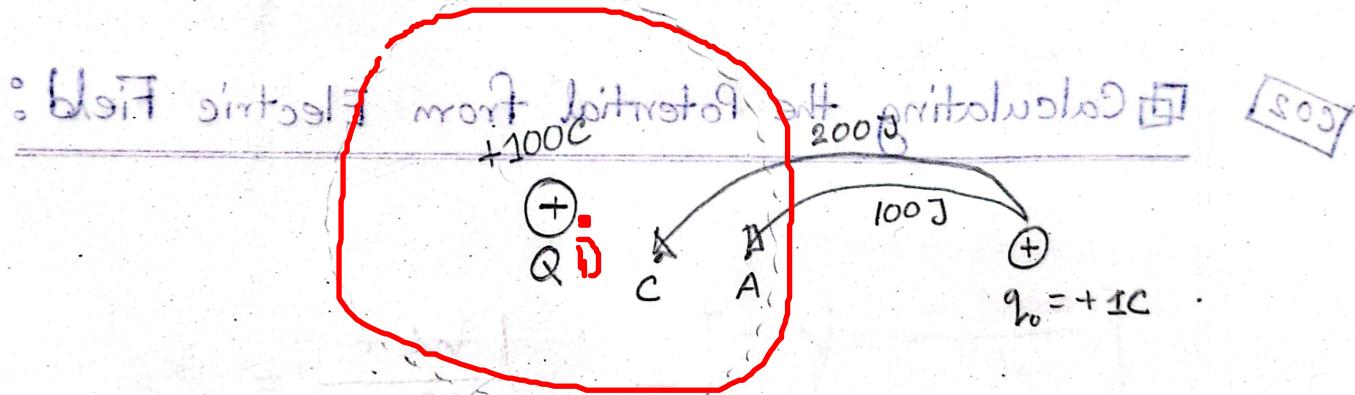


## Potential

Electric Potential :



Electric potential is the amount of electric potential energy per unit charge at a specific point in an electric field.



Potential,

$$V = \frac{-W_{\infty}}{q_0} = \frac{U}{q_0}$$

$$[U = -W]$$

$$\text{if } q_0 = +1C, \quad V = -W_{\infty} = U$$

$$\bar{s}b \leftarrow \boxed{\bar{s}b}$$

$$\bar{s}b \cdot \bar{s} = wb \therefore$$

Electric Potential

$$* \left( \frac{\text{electric potential}}{\text{energy}} \right) = \left( \frac{\text{particle}}{\text{charge}} \right) \left( \frac{\text{electric potential energy}}{\text{unit charge}} \right)$$

$U = qV$

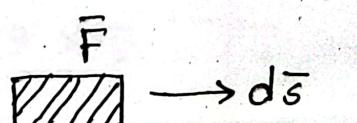
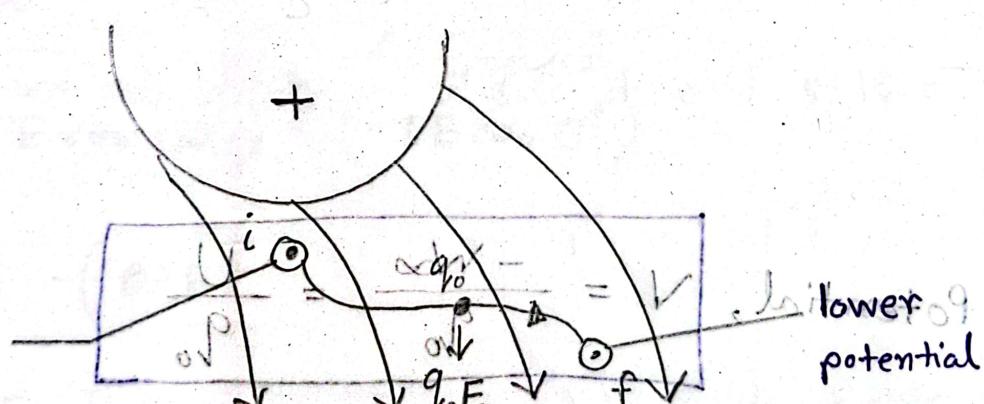
$$V = \frac{U}{q_0}$$

**Q2** Calculating the Potential from Electric Field:

$W = U$   
higher potential.

(চান্দে বেঁচি  
কাহে, তাই

বিকল্প বেঁচি:  $U = -W = V$



$$\therefore dw = \bar{F} \cdot \bar{dS}$$

Now, state to a Clouded visitor.

$$\underline{dW} = \vec{F} \cdot d\vec{s}$$

$$= \underbrace{\left( q_0 \vec{E} \right)}_{\text{Electric field}} \cdot d\vec{s} \quad \left[ \because \vec{F} = q_0 \vec{E} \right]$$

$$\therefore W = \int dW$$

$$= \int_{\text{boundary}} q_0 \vec{E} \cdot d\vec{s}$$

$$= q_0 \int_i^f \vec{E} \cdot d\vec{s}$$

$$\text{Now, } \Delta V = \frac{-W}{q_0} \quad \left[ \therefore V = \frac{-W_{\infty}}{q_0} \right]$$

$$\Rightarrow \frac{V_f - V_i}{d} = \frac{1}{q_0} \cdot -\vec{E} \cdot d\vec{s}$$

$$\Rightarrow V_f - V_i = - \int \vec{E} \cdot d\vec{s}$$

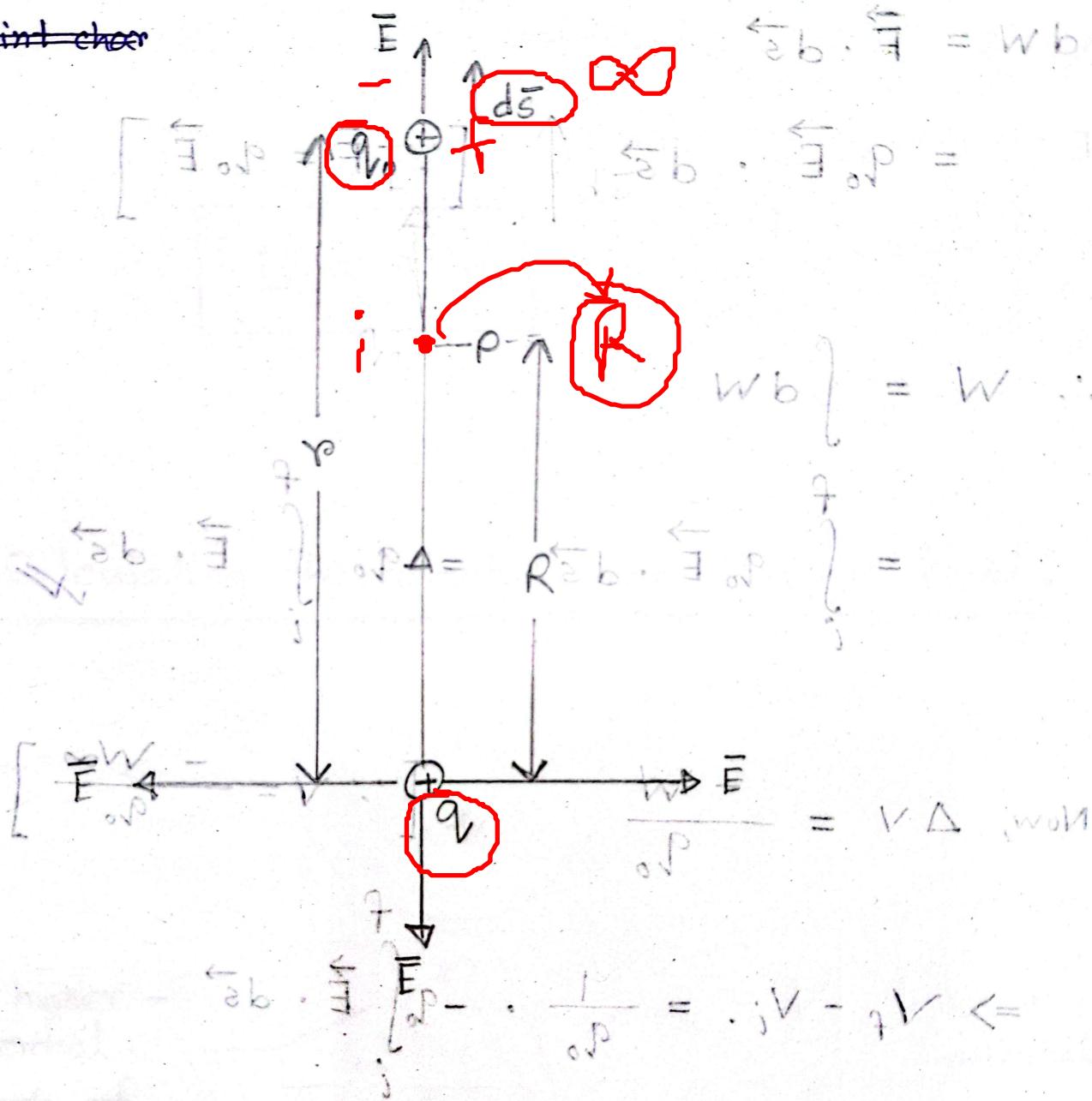
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(sisterly soft moist)

Spring 23

## Box Potential due to a Charged Particle

(Point charge)



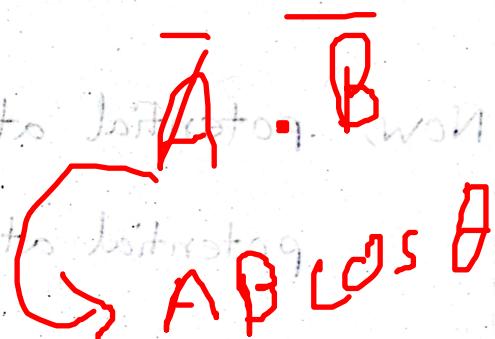
The particle with positive charge  $q$

produces an electric field  $\vec{E}$  and an electric potential  $V$  at point  $P$  (which is at distance  $R$  from the particle)

Now, to find the potential, we move a test charge  $q_0$  from P to infinity, with a displacement of  $d\vec{s}$ .

Now, we know,

$$V_f - V_i = - \int_i^f \vec{E} \cdot d\vec{s}$$



The differential displacement  $d\vec{s}$  of the test particle along its path has the same direction as  $\vec{E}$ .  
 $\therefore \theta = 0^\circ$

Now, because the path is radial,

let us write  $ds$  as  $dr$ . Then, substituting the limits R and  $\infty$  we rewrite eqn  $\textcircled{i}$ :

$$V_f - V_i = \int_R^\infty E dr$$

(i)



Now, potential at infinity  $V_f = 0$

potential at distance  $R$ ,  $V_i = V$

$$\therefore 0 - V = - \int_R^\infty E dr$$

$$\Rightarrow V = \int_R^\infty \left( \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r^2} \right) dr$$

test charge का  
रखना infinity  
ते निये शायद  
अमर्य कोना एक  
अमर्य Electric  
field. घटान है।

$$\Rightarrow V = \frac{q}{4\pi\epsilon_0} \int_R^\infty \frac{1}{r^2} dr$$

$$\Rightarrow V = \frac{q}{4\pi\epsilon_0} \int_R^\infty r^{-2} dr$$

$$= \frac{q}{4\pi\epsilon_0} \left[ \frac{r^{-1}}{-1} \right]_R^\infty$$

$$= - \frac{q}{4\pi\epsilon_0} \left[ r^{-1} \right]_R^\infty$$

$$= - \frac{q}{4\pi\epsilon_0} \left[ \frac{1}{\infty} - \frac{1}{R} \right]$$

$$\Rightarrow V = - \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{R}$$

$$F = C \cdot \frac{q_1 q_2}{r^2}$$

$$E = C \cdot \frac{Q}{r^2}$$

$$V = C \cdot \frac{q}{r}$$

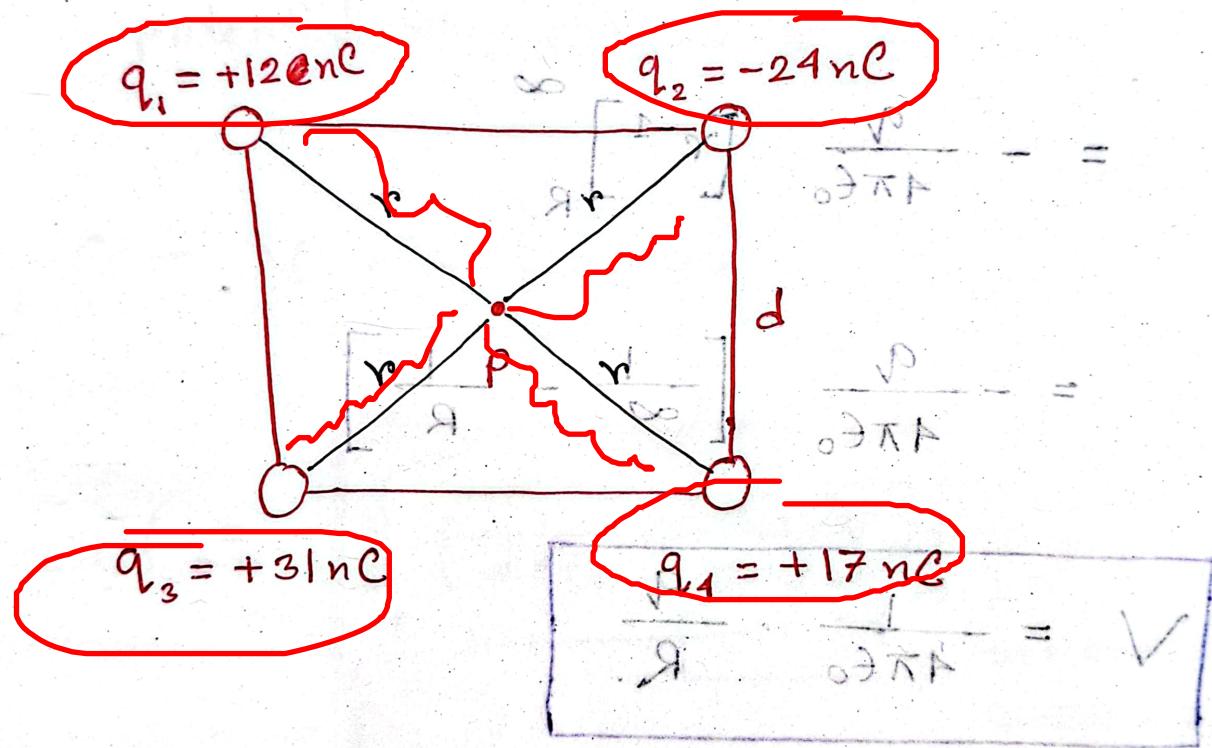
Solving for  $V$  and switching  $R$  to  $r$ ,

Potential due to a charged particle  $q$  at any radial distance  $r$  from the particle:

$$V = \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r}$$

C

Q What is the electric potential at point P, located at the center of the square of charged particles shown below? The distance d is 1.3 m.



Sol:

or of R primitive basis V not envlo

Distance between P and each charge,

$$r = \frac{d}{\sqrt{2}} = \frac{1.3}{\sqrt{2}} = 0.919 \text{ m}$$

$$C \cdot \frac{Q}{r}$$

$$E/F$$

Now, net potential,  $V_{\text{net}}$  is of sub atomic potential.

$$V_{\text{net}} = V_1 + V_2 + V_3 + V_4$$

$$= \frac{1}{4\pi\epsilon_0} \cdot \frac{1}{r} (q_1 + q_2 + q_3 + q_4)$$

$$= \frac{1}{4\pi\epsilon_0} \cdot \frac{1}{r} \left\{ (12 - 24 + 31 + 17) \times 10^{-9} \right\}$$

$$= 9 \times 10^9 \cdot \frac{1}{0.919} (36 \times 10^{-9})$$

$$= 352.56 \text{ V}$$

(Ans.)

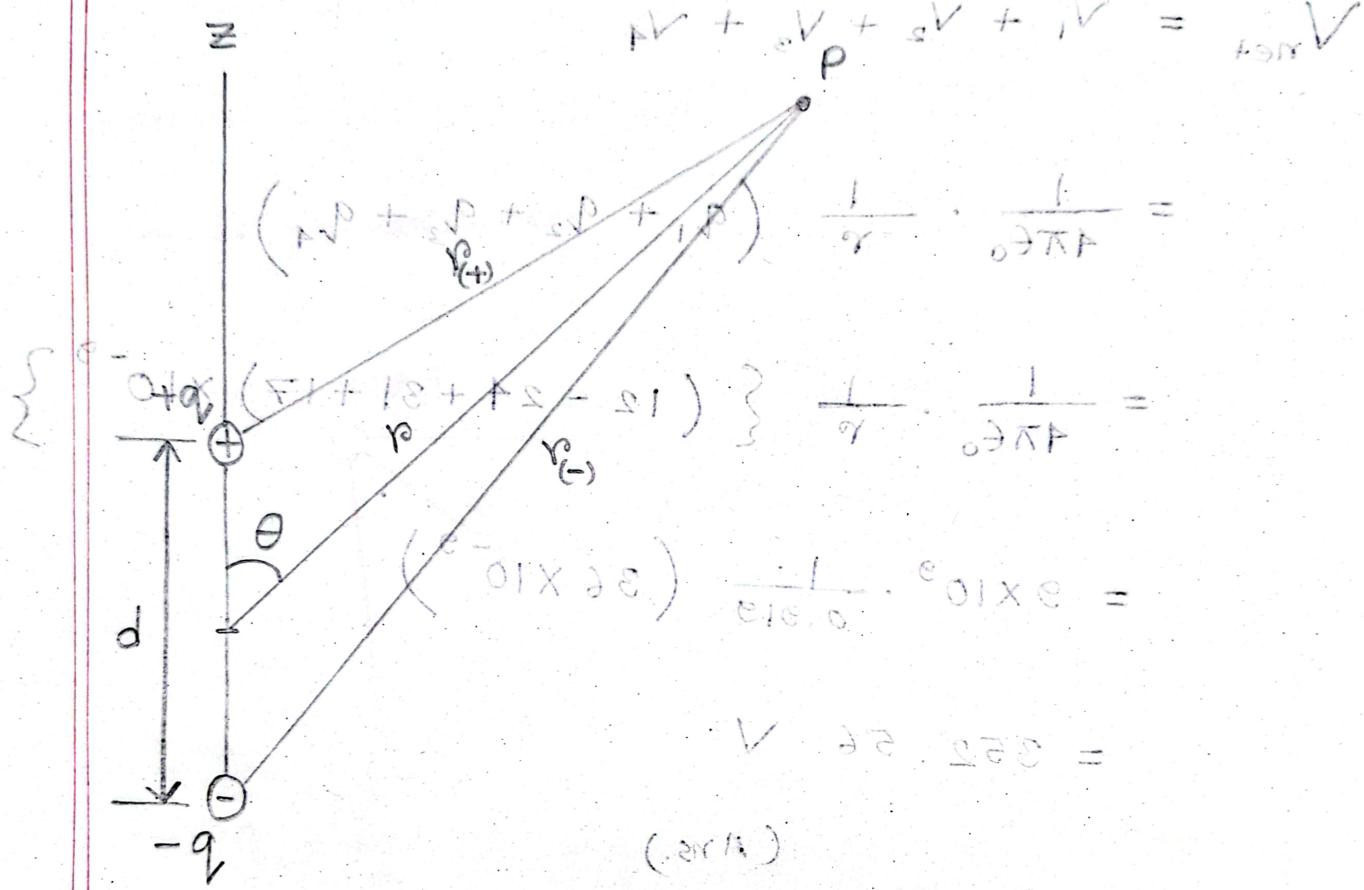
\* Check slide 7 and CW for more questions to practice.

at slab interface is of opposite sign

trapping potential is to attract and bind

Spring 23

## Potential due to an Electric Dipole



We know, potential  $V = \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r}$

Using this eqn to an dip electric dipole to find the potential at an arbitrary point P :

$$V = V_{(+)} + V_{(-)}$$

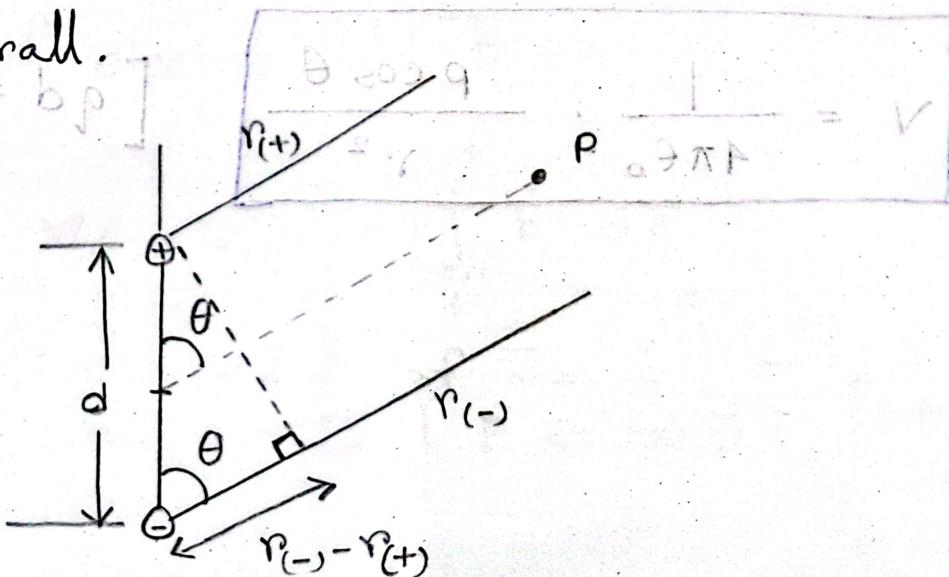
$$= \frac{1}{4\pi\epsilon_0} \left( \frac{+q}{r_{(+)}} + \frac{-q}{r_{(-)}} \right)$$

$$= \frac{q}{4\pi\epsilon_0} \left( \frac{1}{r_{(+)}} - \frac{1}{r_{(-)}} \right)$$

$$= \frac{q}{4\pi\epsilon_0} \left\{ \frac{r_{(-)} - r_{(+)}}{r_{(+)} r_{(-)}} \right\}$$

Now, the distance between the charges of the dipole are quite small. So,  $r \gg d$

So, we can approximate the two lines to P as being parallel.



$\therefore$  Now,

$$\cos \theta = \frac{r_{(-)} - r_{(+)}}{d}$$

$$\left( \frac{p-}{(-)r} + \frac{p+}{(+r)} \right) \cdot \frac{1}{\epsilon_0 \pi A} =$$

$$\Rightarrow r_{(-)} - r_{(+)} = d \cos \theta$$

$$\left( \frac{1}{(-)r} - \frac{1}{(+r)} \right) \cdot \frac{p}{\epsilon_0 \pi A} =$$

and,

$$r_{(+)} \approx r$$

$$r_{(-)} \approx r$$

$$\therefore r_{(+)} r_{(-)} = r \cdot r = r^2$$

After expressing salt resulted capacitance with  $r$

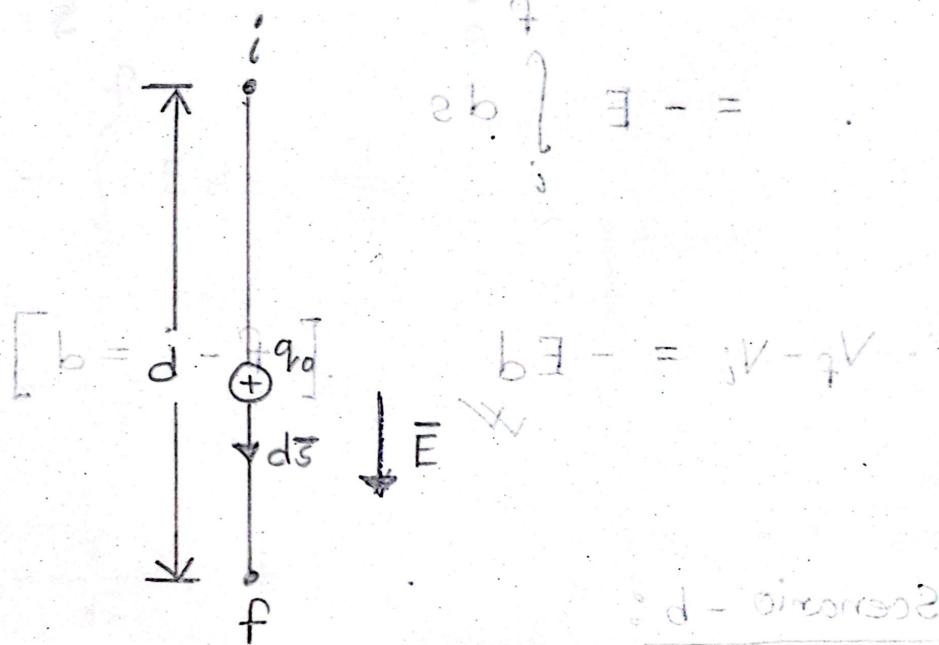
$$\therefore V = \frac{q}{4\pi\epsilon_0} \cdot \frac{d \cos \theta}{r^2}$$

$$\Rightarrow V = \frac{1}{4\pi\epsilon_0} \cdot \frac{p \cos \theta}{r^2} \quad [qd = p]$$

Done

Fall 23

## Relation between potential difference and path:



Scenario - a

In this scenario, a test charge  $q_0$  moves in a straight line from point  $i$  to point  $f$ , along the direction of electric field  $\vec{E}$ .

$$\text{We know, } V_f - V_i = - \int_i^f \vec{E} \cdot d\vec{s}$$

$$= - \int_i^f E ds \cos 0^\circ \quad [\text{both is same dir. } d\vec{s} \text{ and } \vec{E}]$$

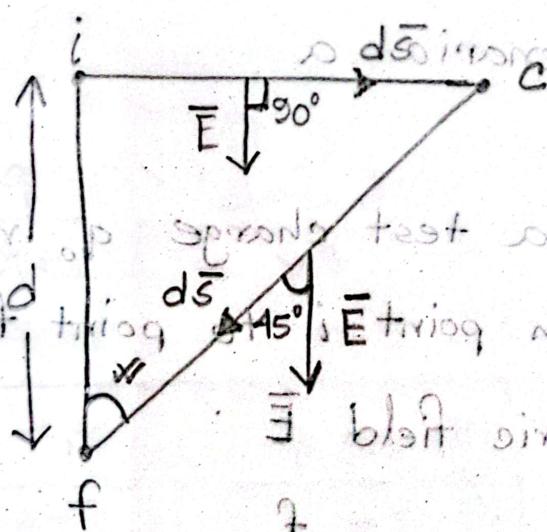
Rechts

$$V_f - V_i = - \int_i^f E ds$$

$$= - E \int_i^f ds$$

$\therefore V_f - V_i = - Ed$   $[f - i = d]$

Scenario - b:



$$\bar{E} \cdot \bar{E} = \bar{E} \cdot \bar{E}$$

$$V_c - V_i = - \int_i^c \bar{E} \cdot d\bar{s} = - \int_i^c E ds \cos 90^\circ = 0$$

Then,  $\nabla_f - \nabla_c = - \int_c^f \bar{E} \cdot d\vec{s} = - \int_c^f E ds \cos 45^\circ$

Electric field along the length of the bar  $= \text{constant} \int_c^f E ds \cdot \frac{1}{\sqrt{2}}$

Given electric field along the length of the bar  $= - \frac{E}{\sqrt{2}} \int_c^f ds$

Here,

$$\sin 45^\circ = \frac{d}{cf}$$

Electric field along the length of the bar  $= - \frac{E}{\sqrt{2}} (cf)$

$$= - \frac{E}{\sqrt{2}} (d \sqrt{2})$$

$$= - Ed$$

$$\sin 45^\circ = \frac{d}{\sqrt{2}}$$

$$= d \sqrt{2}$$

Finally,

$$\nabla_f - \nabla_i = (\nabla_c - \nabla_i) + (\nabla_f - \nabla_c)$$

$$= 0 + (-Ed)$$

$$= -Ed$$

$\therefore$  Potential difference in ~~two~~ both scenario is the same.

Moral: When you want to find the potential difference between two points by moving a test charge between them, you can save time and work by choosing a direct path between those points.

$$(\bar{v}b) \frac{\bar{E}}{b} =$$

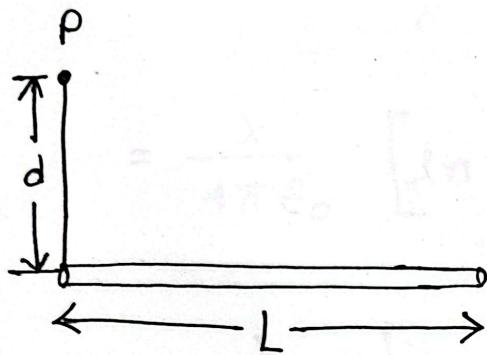
$$b\bar{E} =$$

$$(v - jV) + (jV - v) = jV - jV$$

$$(b\bar{E}) + 0 =$$

$$b\bar{E} =$$

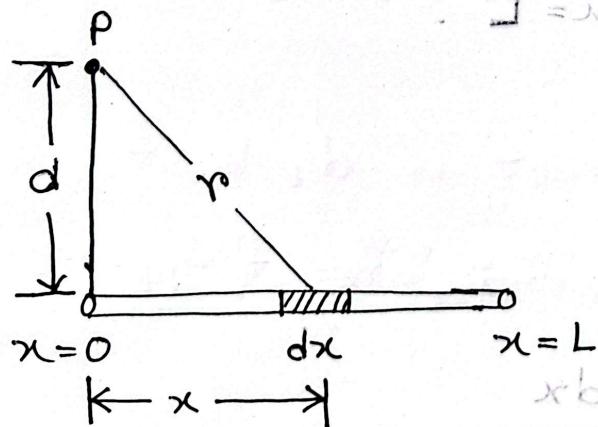
## Potential due to a Continuous Charge Distribution:



Let, charge per unit length =  $\lambda$

Let, a thin nonconducting rod of length  $L$ .

Let us determine the electric potential  $V$  due to the rod at point  $P$ , a perpendicular distance  $d$  from the left end of the rod.



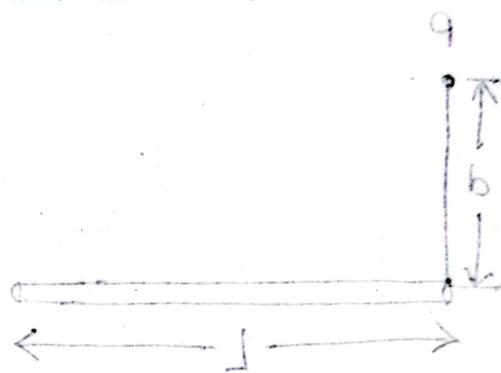
$\therefore$  Charge in differential element  $dx$ ,  $dq = \lambda \cdot dx$

Treating the elements as point charges, we can write the potential  $dV$  as

~~reg spreads to~~

$$dV = \frac{\text{charge}}{4\pi\epsilon_0} \cdot \frac{dq}{r}$$

$$= \frac{1}{4\pi\epsilon_0} \cdot \frac{\lambda dx}{\sqrt{x^2 + d^2}}$$



Integrate to get potential with respect to  $x$

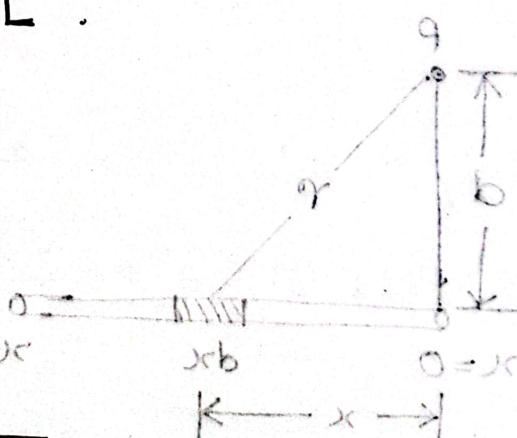
Now, we'll find the total potential  $dV$  produced by the rod at point P by integrating

produced by the rod at point P by integrating along the length of the rod, from  $x=0$  to  $x=L$ .

reg spreads to

$$xb \cdot \lambda \cdot V = \int_{0}^{L} \lambda dx$$

$$= \int_{0}^{L} \frac{1}{4\pi\epsilon_0} \cdot \frac{\lambda dx}{\sqrt{x^2 + d^2}}$$



Final Gauss Law

$$= \frac{\lambda}{4\pi\epsilon_0} \int_0^L \frac{1}{\sqrt{x^2 + d^2}} dx$$

$$= \frac{\lambda}{4\pi\epsilon_0} \left[ \ln(x + \sqrt{x^2 + d^2}) \right]_0^L$$

$$= \frac{\lambda}{4\pi\epsilon_0} \left[ \ln(L + \sqrt{L^2 + d^2}) - \ln d \right]$$

$$\boxed{V_{no} = \frac{\lambda}{4\pi\epsilon_0} \cdot \ln \left( \frac{L + \sqrt{L^2 + d^2}}{d} \right)}$$

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expressions diff, resulted since diff no thing is

blown off by air current when

If slogan is to turn around like expression diff

exact opposite direction, then slogan diff  
turn around like expression diff

• slogan loops to

Equipotential surfaces:  $(\phi_1 + \phi_2) = \frac{q}{4\pi\epsilon_0 r}$

An equipotential is a surface where every point on it has the same electric potential.

This means that, no work is required to move a charge along the surface because the electric potential difference between any two points on the surface is zero.

