

Quantum Transport Simulations with Tight-Binding Models

Hamiltonian : A mathematical desc. of the total energy of a system
↳ KE + PE

very important - allows us to solve the Schrödinger eqn.

Schrödinger eqn. : calculates how the system will evolve or change in time.

↑ basically tells us everything useful we need to know

Now:

What is the Tight-Binding Model (TBM) ?

- It's an approximation
- Used in solid state physics,
- To describe how electrons move in a material

Assumptions of the TBM :

↓ (like atoms in a lattice)

- 1.) Electrons are tightly bound to atomic sites
- 2.) Electrons can hop to neighbouring states w/ a probability determined by a param., "hopping amplitude" (t).
- 3.) Each site has an on-site potential that represents the energy of an electron at that site

1D tight-binding Hamiltonian :

$$H = \sum_i \underbrace{\epsilon_i c_i^+ c_i}_\text{Number Operator} + \sum_{i,j} \underbrace{t_{ij} c_i^+ c_j}_\text{hopping term}$$

$\langle i, j \rangle$ represents the on-site energy of an e^- at site i

represents the hopping term, which allows an e^- to move between neighboring sites i and j

Where :

→ c_i^+ and c_i are **creation** and **annihilation** operators for electrons at site i .

→ ϵ_i is the on-site energy (potential E @ site i)

→ t_{ij} is the hopping amplitude \leftrightarrow sites i and j

An Aside on Creation and Annihilation Operators:

→ These operators describe how particles (like e^- s) are added or removed from a quantum system

Creation Operator $[c_i^+]$ → Adds a particle (e^-) to site i

Annihilation Operator $[c_i]$ → Removes a particle (e^-) from site i

Mathematically, they act on quantum states in the following way:

$$c_i^+ |0\rangle = |1\rangle$$

$$c_i |1\rangle = |0\rangle$$

where: $|0\rangle$ = No e^- at site i

$|1\rangle$ = An electron is present at site i

These ops. obey the anti-commutation relations

(due to Pauli Exclusion Principle) :

$$\{c_i, c_i^\dagger\} = c_i c_i^\dagger + c_j^\dagger c_i = \delta_{ij}$$

$$\{c_j, c_j^\dagger\} = \{c_i^\dagger, c_j^\dagger\} = 0$$

which enforce Fermionic behavior (e^- are fermions)

and obey the Pauli Exclusion Principle, meaning

two e^- s cannot occupy the same quantum state.

Breaking Down the Hamiltonian:

$\rightarrow c_i^\dagger c_i$ (Number Operator)

- Counts whether there is an e^- at site i
- If there is an e^- at site i , the system has energy E_i

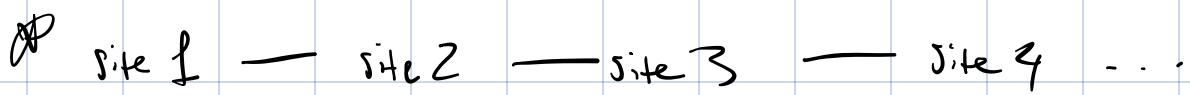
$\rightarrow c_i^\dagger c_j$ (Hopping Term)

- Moves an e^- from site j to site i
- The hopping amplitude $|t_{ij}|$ determines how strongly e^- s move between sites.



Physical Meaning of the TBM

→ Imagine an array of atoms where e^- s can hop between neighboring sites.

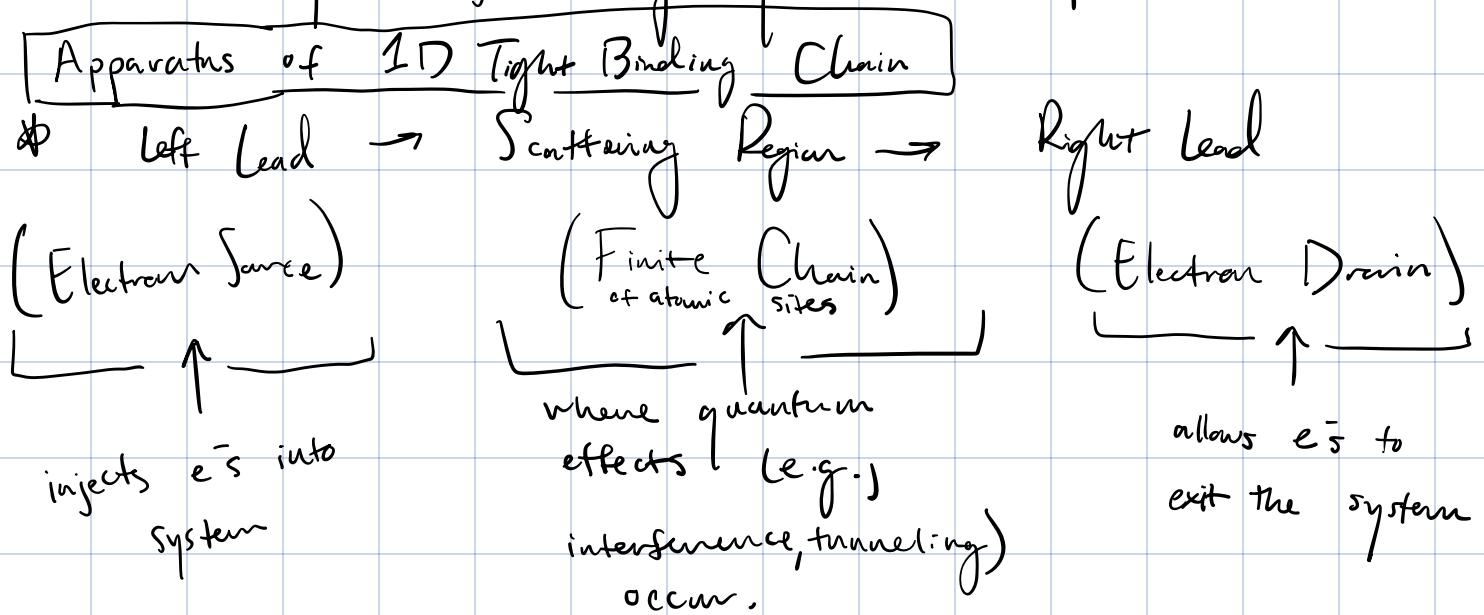


→ if $t_{ij} = 0$, electrons are localized at their sites

→ if $t_{ij} \neq 0$, e_i 's can spread out and visit both sites

- The energy spectrum of this system depends on the hopping strength and boundary conditions

→ In our 1D TBM, the leads allows e^- s to enter and exit the system enabling quantum transport



leads are modeled as infinite conductors, providing a well defined boundary condition.

Overview: How does Quantum Transport Work in this System?

- 1.) An e^- wave enters the system from left lead
- 2.) This e^- propagates through finite chain
- 3.) Some part of the wave transmits through to the right lead, and some part reflects back.

→ The conductance (electron flow) is computed using the Landauer formula:

$$G = \frac{2e^2}{h} T(E)$$

where $T(E)$ is the transmission probability at energy E

Now, let's add some disorder to our system.



Anderson Localization

→ What is it??

- In a perfect system, e^- 's propagate freely through the lattice
- If we introduce random disorder (on-site potential variations), electron waves interfere destructively, leading to this localization effect.
- In 1D systems, even weak disorder can cause

localization, meaning e^- 's stop propagating and conductance stops

(remember, this is just on-site potential variations)

- Expected results of adding disorder
 - w/o disorder \rightarrow conductance remains at $1.0e^2/h$
 - w disorder \rightarrow sharp dips in conductance appears due to localized states
 - Stronger disorder \rightarrow conductance may drop to 0 in some energy ranges

What happens if we add a magnetic field? (lattice)

\rightarrow Quantum Hall Effect (QHE)

\hookleftarrow What is it?

(B)

- when you apply a magnetic field \perp to a 2D material, e^- 's move in cyclotron orbits instead of straight paths
- In a TBM, we introduce a magnetic phase factor (Peierls phases) to simulate the effect of a magnetic field.

Edge states
emerge,
carrying
current
w/o
backscattering

- This breaks time-reversal symmetry and leads to the formation of Floquet levels.

definitions aside

"Peierls Phase" \rightarrow refers to a phase factor introduced in condensed matter physics to account for the effect of a magnetic field on the hopping of electrons in a lattice

↪ essentially represents the additional phase an e^- acquires when moving through a magnetic field

"Landau levels" are just the quantized energy levels of e^-s in a magnetic field.

Going back to QHE ... how do we implement a magnetic field in Kwant?

\rightarrow Modify hopping term to introduce a magnetic phase using Peierls substitution:

$$t_{ij} \rightarrow t_{ij} e^{i 2\pi \phi_{ij}}$$

* where ϕ_{ij} is related to the magnetic flux per plaquette

- So ϕ represents the magnetic flux per unit cell

- $e^{i 2\pi \phi}$ modifies the hopping term to introduce a phase shift

What to expect when adding the magnetic field

$$(\phi = 0)$$

→ w/o a magnetic field → conductance remains unchanged

→ w/ a magnetic field ($\phi > 0$) :

- Shifts in conductance plateaus (energy levels due to cyclotron orbits)
- Gaps where conductance drops (formation of Landau levels)
- Stronger fields lead to well-separated plateaus

conductance will show plateaus at integer multiples of e^2/h

Now... what happens if we add disorder to this QHE system?

→ Two competing effects:

A. Landau levels and Quantum Hall Plateaus

→ w/o disorder, e^- form Landau levels, leading to conductance plateaus

→ As disorder increases, localized states appear b/w Landau levels.

B. Anderson Localization

→ Weak disorder → quantum hall plateaus remain intact

→ Strong disorder → States b/w Landau levels become localized, meaning e^- 's become trapped instead of conducting

→ very strong disorder \rightarrow localization dominates, and the QHE breaks down

So: by gradually increasing disorder strength, we can observe the transition from quantum hall transport to localized insulating behavior.

Expected Results:

Disorder Strength = 0.0

→ sharp conductance plateaus at integer multiples of e^2/h

Weak Disorder ~ 1.0

→ slight rounding of plateaus, but QHE remains intact

Moderate Disorder ~ 2.0

→ some conductance dips appear due to localized states

Strong disorder > 3.0

→ conductance breaks down, as all states become localized

→ quantum hall plateaus disappear, and the system becomes more like an insulator

lets Compute the hamiltonians for the 1D and 2D TBMs

1D:

w/o magnetic phase factor *

$$H = \sum_i \varepsilon_i c_i^\dagger c_i + \sum_{(i,j)} t_{ij} c_i^\dagger c_j$$

- where ε_i is on-site energy at site i
- c_i^\dagger, c_i are creation / annihilation operators for e^- s at site i
- t_{ij} is the hopping amplitude btw neighboring sites

- So system consists of single row sites
- Electrons only hop left or right (nearest-neighbor hopping)
- Dispersion relation :

$$E(k) = \varepsilon - 2t \cos(k)$$

- The bandwidth is $4t$ (since energy varies from $-2t$ to $2t$)



2D : Allow hopping in both x and y directions

$$H = \sum_{x,y} \epsilon_{x,y} C_{x,y}^\dagger C_{x,y} + \sum_{(x',y')} t_{x',y'} e^{i\phi_{x',y'}} C_{x,y}^\dagger C_{x',y'}$$

Where :

- $\epsilon_{x,y}$ is the on-site energy at each site (x,y)
- $C_{x,y}^\dagger, C_{x,y}$ are creation / annihilation
- hopping now occurs in x and y directions
- A magnetic phase factor $e^{i\phi}$ is introduced in the y-direction to simulate the field.

So :

- e's now hop in both x and y directions
- Peierls phase $e^{i\phi}$ modifies hopping in the y direction (introducing the cyclotron motion in a magnetic field)
- Dispersion relation is more complex bc of Landau Levels :

$$E(k_x, k_y) = \epsilon - 2t(\cos k_x + \cos k_y)$$

- The magnetic field quantizes the energy into Landau levels which are missing in 1D.

Next Step in Experimentation:

Compute the energy spectrum to observe
Landau quantization

Extend the system to a honeycomb lattice
→ Graphene QHE