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2 ENERGY CORRELATORS AS A PROBE OF THE HARD PROCESS IN RUN 24 $p - p$
3 COLLISIONS AT $\sqrt{s} = 200 \text{ GeV}$ WITH THE sPHENIX DETECTOR AT RHIC

4 by

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6 A dissertation submitted to the Graduate Faculty in Physics in partial fulfillment of the
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Abstract

ENERGY CORRELATORS AS A PROBE OF THE HARD PROCESS IN RUN 24 $p - p$

COLLISIONS AT $\sqrt{s} = 200$ GeV WITH THE SPHENIX DETECTOR AT RHIC

by

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²³ Adviser: Professor Stefan Bathe

²⁴ This dissertation consists of three parts in addition to a literature review establishing the
²⁵ state of the art of jet physics and the Energy-Energy correlator at high energies...

26 Part 1: Hardware This part discusses the physical hardware used to make the measure-
27 ments of the energy in the jets created by the proton-proton collisions. This part contains a
28 discussion of the sPHENIX detector and an in-depth look at the relevant susbsystems for the
29 meausements in this analysis. In addition, the Monte Carlo methods and models to discern
30 the expected respnse and calibrate the detector are discussed in detail, with both subparts
31 coming together in a disucssion of backgrounds and error calculation in general and for the
32 observables at hand.

³³ **Part 2: Technicalities** This part builds from first principles up to the jet measurement
³⁴ in a theoretical light, and then discusses the additional computational techniques on display
³⁵ that will provide the bridge between the theory of these observables and the practical ap-
³⁶ plication to sPHENIX. This part contains a chapter that is a pared down discussion of a
³⁷ midstream paper proposed as part of the thesis process that establishes the safety of this
³⁸ observable against jet finding choices.

³⁹ **Part 3: Experimental Output** This part is the meat and potatoes of the dissertation,
⁴⁰ discussing results from the experiment and the analysis in light of the previous parts and
⁴¹ comparing to the world results from related systems.

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⁹⁰ Chapter 1

⁹¹ Literature Review

⁹² 1.1 Jets Definitions

⁹³ The study of Quantum Chromodynamics in high energy collisions, such as those at the
⁹⁴ Relativistic Heavy Ion Collider or the Large Hadron Collider, is often carried out through
⁹⁵ investigation of the kinematics of jets. Jets, as objects, sit in the boundary between experiment
⁹⁶ and theory, being an experimental signature corresponding to final states of quarks and
⁹⁷ gluons produced in collisions. Jets are a cluster of final state particles that result from the
⁹⁸ showering and hadronization of the initial parton, that are identified in experiment through
⁹⁹ use of one of the multiple jet reconstruction algorithms.

¹⁰⁰ 1.1.1 Jet Reconstruction Algorithms

¹⁰¹ In Ryan Atkin's paper *Review of jet reconstruction algorithms* [1], Atkin provides an overview
¹⁰² of the standard algorithms with a focus on practical implementation for experimental usage.
¹⁰³ Atkin discusses a number of algorithms that fall into two categories: Cone based and Se-
¹⁰⁴ quential Clustering

¹⁰⁵ Cone Methods

¹⁰⁶ In Atkin's paper, Atkin discusses two iterative cone procedures—Iterative Cone with Pro-
¹⁰⁷ gressive Removal and with Split Merge—and one more modern cone method, Seedless Infrared
¹⁰⁸ Safe Cone. These iterative cone procedures are described as simple to implement in an ex-
¹⁰⁹ perimental context, however, they both suffer deficiencies from a theoretical perspective as
¹¹⁰ they fail to be IRC safe [1] [2].

¹¹¹ 1.1.2 Quark and Gluon Jets

¹¹² In their paper *An Operational Definition of Quark and Gluon Jets* [Kimiske2018] Kimiske,
¹¹³ Metodiev and Thaler present a theoretical definition of jet flavor with an eye towards appli-

114 cation at colliders. They aim to present a precise and practical definition of jet flavor at the
115 hadron level, thus allowing for direct measurement from collider experiments rather than
116 a measurement by proxy. While the gluon or quark jet is well defined at the level of the
117 hard scattering, corresponding to an intiating parton of the primary vertex [**Jones1989**]
118 [**Fodor1990**]. Experimentally, there is signifigant intrest in being able to discriminate
119 quark and gluon jets in searches for BSM physics, searchs for Dark Matter, in jet sub-
120 structure analysis, in QGP physics and in understanding the processes of hadronization
121 [**Gallicchio2011**][**Lima2017**][**Bhattecherjee2017**]]. The deterimation of relative popula-
122 tions of quark and gluon jets in a dijet sample is one of the primary tasks of this thesis and
123 will be discussed at length in chapters ?? and ??.

124 Kimiske et al. begin their

125 1.2 Jet Substucure Measurements

126 1.3 N-Point Energy Correlator

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Part I

128

Hardware: The Detector and Simulations

129

¹³⁰

Chapter 2

¹³¹

The sPHENIX detector

¹³² The sPHENIX detector is part of the Relativistic Heavy Ion Collider (RHIC) complex at
¹³³ Brookhaven National Laboratories (BNL) in Upton, New York. RHIC is one of two major
¹³⁴ circular particle colliders in the world, the other of which is the Large Hadron Collider at
¹³⁵ CERN in Geneva, Switzerland. Over the course of the sPHENIX experiment which will
¹³⁶ run from run 23 until RHIC ceases operations after run 25, RHIC will be colliding gold
¹³⁷ nuclei and protons with run 23 having been dedicated to commissioning Au+Au running,
¹³⁸ run 24 dedicated to p+p with three weeks of Au+Au for further tracking commissioning
¹³⁹ and run 25 dedicated to physics running in Au+Au collisions. RHIC runs with a center of
¹⁴⁰ mass energy of $\sqrt{s} = 200\text{GeV}$ for both protons and gold, and additionally is the sole col-
¹⁴¹ linder that collides polarized protons, which allows for more in depth studies of spin physics.
¹⁴² At RHIC, sPHENIX's forerunner, the Pioneering High Energy Nuclear Interaction eXper-
¹⁴³ iment (PHENIX), performed the first measurement of the Quark Gluon Plasma droplets
¹⁴⁴ [QGP droplets] in nuclei-nuclei collisions.

¹⁴⁵ sPHENIX is designed to offer significant upgrades to PHENIX, specifically providing in-
¹⁴⁶ creased coverage with hadronic and electromagnetic calorimeters, increasing effectiveness in
¹⁴⁷ studies of jets and heavy flavor at mid-to-high Bjorken x [sPHENIX TDR][sPHENIX whitepaper][sPHENIX paper]

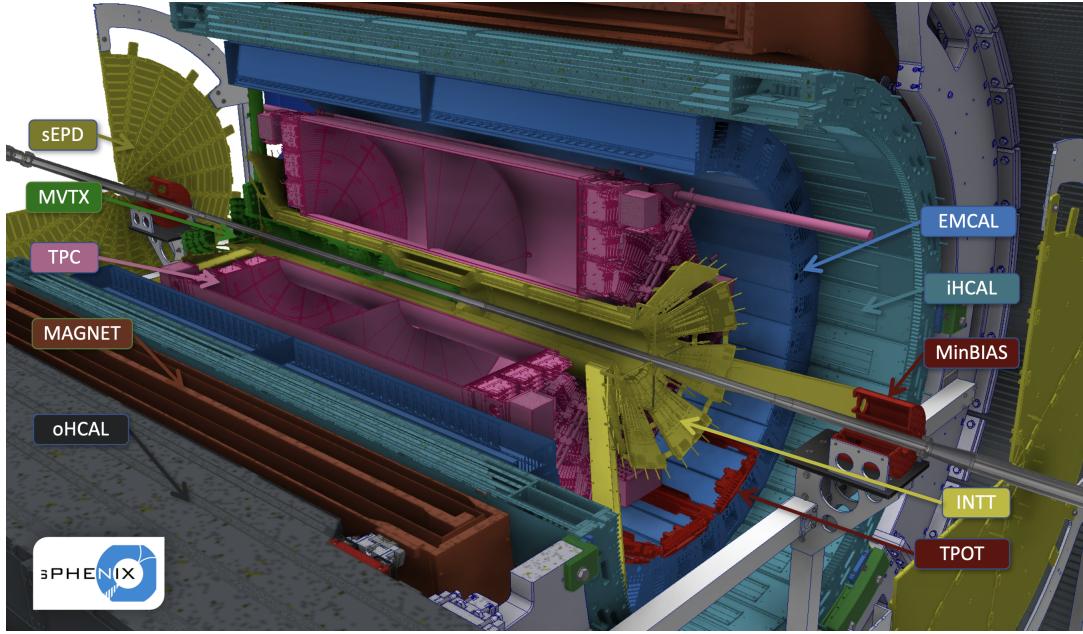


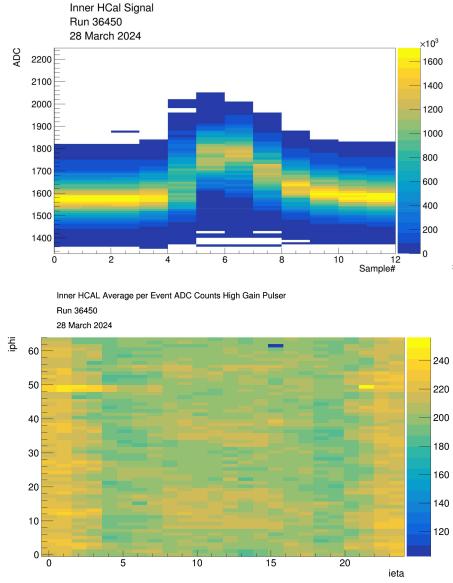
Figure 2.1: The sPHENIX detector. The Zero Degree Calorimeter is not pictured but is located further along the beam pipe.

148 Broadly, sPHENIX can be broken into three categories of subsystem: Calorimeters, Tracking
149 and Event characterization. sPHENIX is constructed around a 1.5 Tesla superconducting
150 magnet, cooled by liquid helium, which was previously used on the BABAR experiment at
151 SLAC.

152 2.1 Calorimetry

153 2.1.1 Hadronic Calorimeters

154 sPHENIX's power in jet physics comes in large part from its calorimeter systems, including
155 two separate hadronic calorimeters. The outer and inner hadronic calorimeter (OHCAL and
156 IHCal respectively) are both divided into 24 bins in η with coverage of $|\eta| \leq 1.1$ and full
157 φ coverage with 64 bins in φ , grouped in φ into 32 sectors in each calorimeter. Each
158 hadronic calorimeter therefore has tower size of $\delta\eta = 0.092$ and $\delta\varphi = 0.098$. The outer



(a) Charge Injection Pulse. This system allows for direct testing of electronics without detector response

Figure 2.2: Top: Persistence Waveform. Bottom: Energy distribution averaged across 10^5 events

(b) LED pulse. This allows for testing of the full readback system and light response

hadronic calorimeter is constructed of steel and sits outside of the magnet, with inner radius of $r = 1.82$ m and outer radius of $r = 2.69$ m. The inner hadronic calorimeter is constructed of aluminum and sits just inside the magnet, with inner radius of $r = 1/16$ m and outer radius of $r = 1.37$ m. Each tower corresponds to a readout board that sums readout from four scintillator paddles embeded into the calorimeter as show in fig. ???. These interface boards form a single readout channel, and can be individually teste via a pulser system that injects charge into the interface board testing readback independent of scintillator response, and LED system that injects a fixed pulse of light into the scintilators to test behavior of the Silicon Photomultipliers (SiPMs). Output of each of these systems is shown in fig 2.2

¹⁶⁸ **Calibration**

¹⁶⁹ The hadronic calorimeters were intially calibrated through use of cosmic ray muons, matching
¹⁷⁰ the spectra to Monte Carlo generated through EcoMug and GEANT4 [**HCal'Calib**]. This
¹⁷¹ calibration is updated through out the run via continued cosmic ray studies in between
¹⁷² physics data taking in addition to ongoing work on in-situ calibrations using tower-slope
¹⁷³ methods [**tower'slope'hcal**]. These calibrations have yeilded an average conversion factor
¹⁷⁴ for the OHCal of and the IHCal of

¹⁷⁵ **2.1.2 Electromagnetic Calorimeter**

¹⁷⁶ Moving in one layer from the hadronic calorimeter, the Electromagnetic Calorimeter (EM-
¹⁷⁷ Cal) has the same coverage in $\eta - \varphi$ space, but has 16 times as many towers with 96 bins in
¹⁷⁸ η and 256 bins in φ . The EMCAL is constructed of blocks of tungsten with embedded scin-
¹⁷⁹ tilationg fibers, and had inner radius $r = 0.9$ m and outer radius of $r = 1.16$ m. Similarly
¹⁸⁰ to the HCals, the EMCAL is also equipped with a pulser and LED, although the increased
¹⁸¹ number of towers and higher varaiblility in response requires that different pulse widths be
¹⁸² used for seperate sets of towers to prevent saturation and clipping on the LED pulse.

¹⁸³ **Calibration**

¹⁸⁴ The EMCAL was calibrated to the π^0 mass through the $\pi^0 \rightarrow \gamma\gamma$ decay channel with additional
¹⁸⁵ corrections applied via the tower slope. The mean m_{π^0} and width are used as quality
¹⁸⁶ assurance plots to monitor radiation damage to the EMCAL on an ongoing basis.

¹⁸⁷ **2.1.3 Zero Degree Calorimeter**

¹⁸⁸ The Zero Degree Calorimeter (ZDC) is a small transverse energy hadron detector that po-
¹⁸⁹ sitioned outside of the experiment hall, in the tunnels along the beam pipe, 18 m from he

190 interaction point on both side. The ZDC is used to detect spectator neutrons to allow for
191 the calculation of multiplicity and event geometry, makind it an overlapping detector be-
192 tween the calorimeters and the event characterization [3]. The sPHENIX ZDC is one of two
193 detectors that was preserved from PHENIX, with minor repairs and upgrades to improve
194 running. In pp collisions, the ZDC and the Minimum Bias detector

195 **2.2 Tracking**

196 The sPHENIX tracking system consits of the Silicon vertexing detectors and the gas based
197 trackers.

198 **2.2.1 INTT**

199 **2.2.2 Micromegas Vertex Detector**

200 The Micromegas Vertex detector (MVTX) is a

201 **2.2.3 Time Projection Chamber**

202 The Time Projection Chamber (TPC)

203 **2.2.4 TPC Outer Tracker**

204 The TPC Outer Tracker (TPOT) was designed as a minimal add-on to improve tracking
205 effeciency and resolution by providing additional track points.

206 2.3 Event Characterization**207 2.3.1 sPhenix Event Plane Detector****208 2.3.2 Minimum Bias Detector**

209 The sPHENIX Minimum Bias Detector (MBD) is the other detector preserved from PHENIX,
210 when it was called the Beam-Beam Counter (BBC). This detector is situated

211 2.4 Triggers

212 The sPHENIX detector has a number of triggers at its disposal to allow for online event type
213 selections for the triggered readout system: the EMCal, the HCals, the ZDC, the MBD, the
214 sEPD and—for run 2024—the TPC. Starting in run 2024, available triggers are Minimum Bias

²¹⁵ **Chapter 3**

²¹⁶ **Monte Carlo Simulations**

²¹⁷ We have some of them, there is some data from the 2 and three point energy correlator, and
²¹⁸ its important for training of our discriminator

²¹⁹ Chapter 4

²²⁰ Determining Backgrounds and Errors

²²¹

Part II

²²²

Technicalities: The finer points of
theory and computing

²²³

²²⁴ **Chapter 5**

²²⁵ **The Energy Correlator and the
primary vertex**

²²⁷ So the energy correlator allows a look back, I guess that's the relevant bit? Point being, if
²²⁸ the jet points back to the primary vertex, the share of the energy between the jets and the
²²⁹ composition of those jets is directly determined by the $2 \rightarrow 2$ scattering and is thus sensitive
²³⁰ to the parton distribution functions of the beam particles. Starting from the definition of
²³¹ the energy correlator given in [4]:

$$\mathcal{E}_N = \int \frac{d\Omega}{2} \int d\vec{n}_{12} \frac{\langle \varepsilon(\vec{n}_1) \dots \varepsilon(\vec{n}_2) \rangle}{Q^2} \delta(\vec{n}_1 \cdot \vec{n}_2 - 2|\vec{n}_1| |\vec{n}_2| \cos \theta) \quad (5.1)$$

232 **Chapter 6**

233 **Jets in Vacuum and the PDF**

234 **6.1 Jet Identification Algorithms**

235 As discussed in chapter 1, there are a variety of jet finding algorithms that prioritize different
236 theoretical aspects of the underlying physics while being experimentally realizable [5] [1].

237 In general, a jet identification algorithm needs to be IRC safe. That is, the jet object
238 needs to display invariance in the Infrared (IR) and Collinear regimes, managing real-virtual
239 cancellation and keeping results meaningful for emission and splitting respectively.

²⁴⁰ Chapter 7

²⁴¹ So exactly how intelligent is AI

²⁴² Chapter 8

²⁴³ Proof Solving and Validation as
²⁴⁴ Quality Control Mechanisms

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Part III

246

Experimental Output: The Main Event

247

248 Chapter 9

249 So is sPHENIX actually working?

250 Chapter 10

251 Measuring the Energy Correlator

²⁵² Chapter 11

²⁵³ The Power of the ENC: α_s at the few
²⁵⁴ GeV scale

255 Chapter 12

256 Event-by-Event distinguishing

²⁵⁷ Chapter 13

²⁵⁸ Comparison is the Theft of Joy: What
²⁵⁹ does the LHC say, and how about
²⁶⁰ STAR?

²⁶¹ Chapter 14

²⁶² Entanglement and other Lofty Goals

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Part IV

264

Wrapping it all up

²⁶⁵ Chapter 15

²⁶⁶ Observable prospects for Run 25 and
²⁶⁷ the EIC

²⁶⁸ Chapter 16

²⁶⁹ Remaining Questions

²⁷⁰ Chapter 17

²⁷¹ Implementation and application for
²⁷² the remaining sPHENIX data

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