

## PROBLEM A – PART B

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### **Q1: What do you observe in the histogram after applying $U^\dagger$ ?**

**Observation:** After applying  $U^\dagger$  to Bob's qubit, the histogram shows that Bob measures the state  $|0\rangle$  with approximately **100% probability** ( $\geq 99\%$  in simulation). The outcome  $|1\rangle$  appears with near-zero probability, which is consistent with ideal noiseless simulation.

This is the expected result if teleportation succeeded: Bob's qubit has been transformed into exactly the state  $|\psi\rangle$ , and applying  $U^\dagger$  (the inverse of the original preparation) rotates it back to  $|0\rangle$ .

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### **Q2: Why does applying $U^\dagger$ verify teleportation?**

The original unknown state was prepared as:

$$|\psi\rangle = U|0\rangle \text{ where } U = Rz(\varphi) \cdot Ry(\theta)$$

If teleportation worked correctly, Bob's qubit is in state  $|\psi\rangle$  at the end of the correction step. Applying the inverse  $U^\dagger$  gives:

$$U^\dagger|\psi\rangle = U^\dagger(U|0\rangle) = (U^\dagger U)|0\rangle = I|0\rangle = |0\rangle$$

Since  $U$  is unitary,  $U^\dagger U = I$  (identity). Therefore measuring  $|0\rangle$  with 100% probability confirms that Bob holds exactly  $|\psi\rangle$  — i.e., teleportation was successful. If Bob had a different state, the measurement would yield  $|1\rangle$  with nonzero probability.

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### **Q3: Why were we able to avoid sending classical bits?**

**In the standard protocol**, Alice measures qubits 0 and 1, obtains two classical bits, and sends them to Bob over a classical channel. Bob then applies conditional Pauli corrections (X and/or Z) based on the received bits.

**In this fully quantum version**, we *do not* measure qubits 0 and 1. Instead, we replace the classical feed-forward with coherent controlled quantum gates: - `qc.cx(1, 2)` — replaces the classically-controlled X correction - `qc.cz(0, 2)` — replaces the classically-controlled Z correction

This works because the controlled gates act on the joint state *in superposition*: when qubit 1 is  $|1\rangle$ , CX applies X to qubit 2; when qubit 0 is  $|1\rangle$ , CZ applies Z to qubit 2. The effect is mathematically equivalent to the classical correction for all measurement outcomes simultaneously.

**Key difference:** The real-world protocol requires classical communication and is limited by the speed of light. This fully quantum version is a coherent, unitary circuit that works only in a *closed quantum system* (no actual communication between separated parties). It cannot be used for faster-than-light communication — the classical channel is still required in practice because Alice and Bob are physically separated. The fully quantum version is useful for circuit-based verification and quantum algorithms, not for actual physical teleportation between distant parties.