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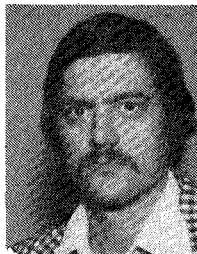
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# Determination of Loaded, Unloaded, and External Quality Factors of a Dielectric Resonator Coupled to a Microstrip Line

APS KHANNA AND Y. GARAUULT

**Abstract**—In the case of a dielectric resonator coupled to a microstrip line, the relations for the determination of unloaded, loaded, and external quality factors, in terms of directly measurable reflection or transmission coefficient at the resonant frequency, have been derived and represented on the corresponding vectorial and scalar planes. Construction of a linear frequency scale and a graphical method to accurately determine the unloaded  $Q$  from the loaded  $Q$  measurement presented.

## I. INTRODUCTION

**T**HE ADVENT OF temperature stable, high- $Q$  and low-loss ceramic materials has created a considerable interest in the application of dielectric resonators in the

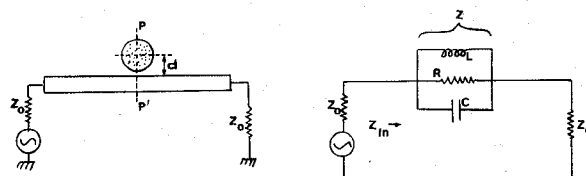


Fig. 1. Dielectric resonator coupled to a microstrip line and its equivalent circuit in symmetry plane  $PP'$ .

microwave integrated circuits [1]. The dielectric resonator coupled to a microstrip line has been used for realizing a number of stable oscillators [2]–[4] and filters [5]. Complete characterization of a microstrip coupled dielectric resonator (Fig. 1) is necessary for the analysis and synthesis of these integrated circuits. Ginzton [6] introduced the plotting of loci of various quality factors and the construction of a linear frequency scale in the case of a single-port

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reflection mode resonant metallic cavity. Podcameni [7] recently suggested a relation for the determination of unloaded  $Q$  of the dielectric resonator, using an indirect method. In this paper, we present the necessary relations, in terms of directly measurable quantities, to draw loci of unloaded, loaded, and external quality factors, on the impedance as well as on the transmittance plane, in the case of a dielectric resonator coupled to a microstrip line (Fig. 1). The construction of a linear frequency scale is presented and a graphical method to accurately determine the unloaded  $Q$  from the loaded  $Q$  measurement is given.

## II. COUPLING FACTOR OF A DIELECTRIC RESONATOR COUPLED TO A MICROSTRIP LINE

Equivalent circuit in the symmetry plane  $PP'$  of a  $TE_{01p}$ -mode cylindrical resonator, magnetically coupled to a 50- $\Omega$  microstrip line, under necessary shielding conditions, is shown in Fig. 1 [5]. The coupling factor  $\beta$  is defined as the ratio of the resonator coupled resistance  $R$  at the resonant frequency to the resistance external to the resonator or

$$\beta = \frac{R}{R_{\text{ext}}} = \frac{R}{2Z_0} = \frac{S_{11_0}}{1 - S_{11_0}} = \frac{1 - S_{21_0}}{S_{21_0}} = \frac{S_{11_0}}{S_{21_0}} \quad (1)$$

where  $S_{11_0}$  and  $S_{21_0}$  are the real quantities representing the reflection and transmission coefficients respectively, at the resonant frequency, in the same symmetry plane  $PP'$  (Fig. 1).

The critical coupling ( $\beta = 1$ ) occurs when the power dissipated in the resonator  $P_d$  is equal to the power dissipated in the external circuit, which is equally divided into the power reflected to the generator ( $P_r = S_{11_0}^2$ ) and the power transmitted to the load ( $P_t = S_{21_0}^2$ ), i.e.,  $P_t = P_r = P_d/2$ .  $P_d$ ,  $P_r$ , and  $P_t$  represent here the RF powers normalized with respect to the incident power.

In the shielded resonator configuration, from the conservation of energy, power dissipated in the resonator being given by

$$P_d = 1 - |S_{11_0}|^2 - |S_{21_0}|^2 \quad (2)$$

the critical coupling corresponds to  $S_{11_0} = S_{21_0} = 0.5$ .

The coupling factor  $\beta$  is a function of the distance between the dielectric resonator and the microstrip line under fixed shielding conditions.  $\beta$  also relates the various quality factors by the well-known relation

$$Q_u = Q_L(1 + \beta) = \beta Q_{\text{ex}} \quad (3)$$

where  $Q_u$ ,  $Q_L$ , and  $Q_{\text{ex}}$  represent unloaded, loaded, and external quality factor, respectively.

## III. LOCI OF UNLOADED, LOADED, AND EXTERNAL QUALITY FACTORS

The aim is to draw the loci of the points, on the impedance ( $S_{11}$ ) and transmittance ( $S_{21}$ ) planes, for the frequency deviations corresponding to  $Q_u$ ,  $Q_L$ , and  $Q_{\text{ex}}$ . In Fig. 1,  $Z_{\text{in}}$  in the plane  $PP'$  of the dielectric resonator is

given by [5]

$$Z_{\text{in}} = Z_0 + \frac{R}{1 + j2Q_u\delta} \quad (4)$$

where  $\delta = (f - f_0)/f_0$  is the normalized frequency deviation.

Using (1) and (3) the normalized input impedance  $z_{\text{in}} = Z_{\text{in}}/Z_0$  can be written as

$$\begin{aligned} z_{\text{in}} &= 1 + \frac{2\beta}{1 + j2Q_u\delta} = 1 + \frac{2\beta}{1 + j2Q_L(1 + \beta)\delta} \\ &= 1 + \frac{2\beta}{1 + j2Q_{\text{ex}}\beta\delta}. \end{aligned} \quad (5)$$

The normalized frequency deviations corresponding to various quality factors are given by

$$\delta_u = \pm \frac{1}{2Q_u}, \quad \delta_L = \pm \frac{1}{2Q_L}, \quad \text{and } \delta_{\text{ex}} = \pm \frac{1}{2Q_{\text{ex}}}. \quad (6)$$

The impedance locus of  $Q_u$ , for example, can be determined by using (6) in (5), and is given by

$$(z_{\text{in}})_u = 1 + \frac{2\beta}{1 \pm j} \quad (7)$$

or using (1)

$$(z_{\text{in}})_u = \frac{1}{1 - S_{11_0}} \pm j \frac{S_{11_0}}{1 - S_{11_0}}. \quad (8)$$

The corresponding reflection coefficient  $S_{11_u}$  can be determined from

$$S_{11_u} = \frac{(z_{\text{in}})_u - 1}{(z_{\text{in}})_u + 1}$$

to be

$$S_{11_u} = \frac{S_{11_0}}{\sqrt{S_{11_0}^2 - 2S_{11_0} + 2}} e^{\pm j \tan^{-1}(1 - S_{11_0})}. \quad (9)$$

The coupled dielectric resonator being a series impedance in the equivalent circuit, the reflection and transmission coefficients are related by [8]

$$S_{11} + S_{21} = 1. \quad (10)$$

The relation for the  $Q_u$  locus, in the transmittance ( $S_{21}$ ) plane, can now be obtained, using the above relation.

In general, Table I represents the relations for the loci of various quality factors, for the reflection coefficient  $S_{11}$  and transmission coefficient  $S_{21}$  planes. Fig. 2(a) and (b) represent these loci in addition to some impedance ( $S_{11}$ ) and transmittance ( $S_{21}$ ) curves for various values of the coupling factor  $\beta$  between the resonator and the microstrip line. As an example,  $Q_u$  can be determined from

$$Q_u = f_0/(f_1 - f_2).$$

Care should be taken while making measurements using the scalar network measurement system giving only the modulus of reflection or transmission coefficients. It may be noted, for example, that in the reflection coefficient plane the  $Q_L$  points are always 3 dB below their value at

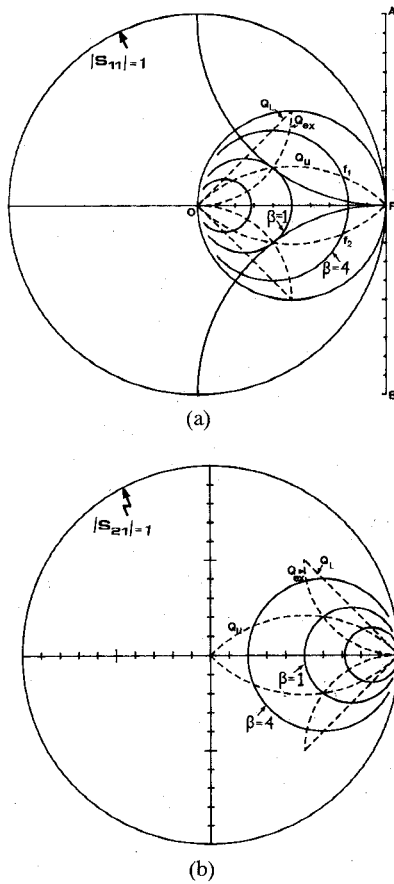


Fig. 2. Loci of  $Q_u$ ,  $Q_L$ , and  $Q_{ex}$  on (a) reflection coefficient plane and on (b) transmission coefficient plane. Reflection and transmission coefficient curves for  $\beta = 0.4, 1$ , and  $4$ .

TABLE I

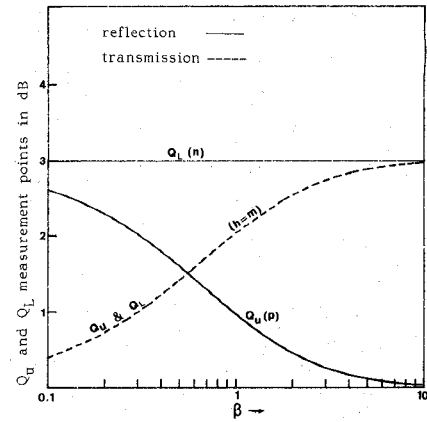
	REFLECTION COEFFICIENT $S_{11}$		TRANSMISSION COEFFICIENT $S_{21}$	
	$ S_{11} $	$\angle S_{11}$	$ S_{21} $	$\angle S_{21}$
$Q_u$	$\frac{S_{11_0}}{\sqrt{S_{11_0}^2 - 2S_{11_0} + 2}}$	$\pm \tan^{-1}(1 - S_{11_0})$	$S_{21} \sqrt{\frac{2}{1 + S_{21_0}^2}}$	$\pm \tan^{-1}\left(\frac{1 - S_{21_0}}{1 + S_{21_0}}\right)$
$Q_L$	$\frac{S_{11_0}}{\sqrt{2}}$	$\pm \frac{\pi}{4}$	$\sqrt{\frac{1 + S_{21_0}^2}{2}}$	$\pm \tan^{-1}\left(\frac{1 - S_{21_0}}{1 + S_{21_0}}\right)$
$Q_{ex}$	$\frac{S_{11_0}}{\sqrt{1 + S_{11_0}^2}}$	$\pm \tan^{-1} S_{11_0}$	$\frac{\sqrt{2(S_{21_0}^2 + 1) - 6S_{21_0}^2(S_{21_0}^2 + 1) + 9S_{21_0}^4}}{S_{21_0}^2 - 2S_{21_0} + 2}$	$\pm \tan^{-1}\left(\frac{S_{21_0}^2 - 2S_{21_0} + 1}{S_{21_0}^2 - S_{21_0} + 1}\right)$

the resonance frequency, while in the transmission coefficient plane  $Q_L$  points are functions of the coupling factor (or  $S_{21_0}$ ). Fig. 3(a) shows the  $Q_u$  and  $Q_L$  measurement points in the reflection and transmission coefficient magnitude planes as a function of coupling factor  $\beta$  and Fig. 3(b) represents various terms used in Fig. 3(a).

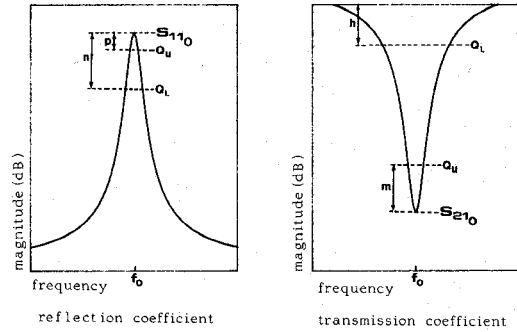
#### IV. CONSTRUCTION OF A LINEAR FREQUENCY SCALE

For a given coupling factor  $\beta$ , a linear frequency scale is obtained along the straight line  $AB$  drawn perpendicular to the  $X$ -axis and passing through point  $F$  (Fig. 4).

The frequency deviation at any point  $R$  on the imped-



(a)



(b)

Fig. 3. (a)  $Q_u$  and  $Q_L$  measurement points in the reflection and transmission coefficient magnitude planes as a function of coupling factor  $\beta$  and (b) definition of various terms used.

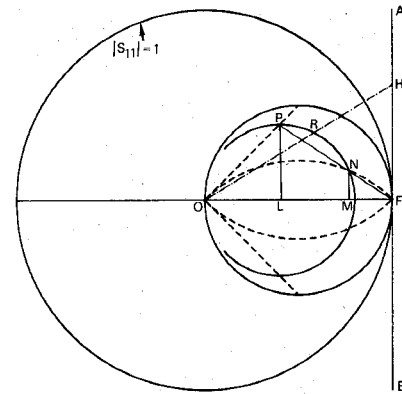


Fig. 4. Construction of a linear frequency scale and determination of unloaded  $Q$  from the loaded  $Q$  measurement.

ance locus can be shown to be proportional to the ordinate of the intersection  $H$  of  $OR$  with  $AB$ . The scale is established from the measurement of frequencies at any two points on the impedance locus.

*Proof:*  $OR$  being the reflection coefficient, from (5)

$$\begin{aligned} S_{11R} &= |S_{11R}| e^{jS_{11R}} = \frac{z_{in} - 1}{z_{in} + 1} \\ &= \frac{\beta(\beta + 1) - j2\beta Q_u \delta}{(\beta + 1)^2 + 4Q_u^2 \delta^2} \end{aligned} \quad (11)$$

$$\text{angle } ROF = \angle S_{11R} = \tan^{-1} \left( \frac{2Q_u \delta}{\beta + 1} \right). \quad (12)$$

The distance FH, being proportional to  $\tan \angle ROF$ , is hence directly proportional to the frequency  $\delta$ .

#### V. GRAPHICAL METHOD TO DETERMINE UNLOADED $Q$ FROM THE LOADED $Q$ MEASUREMENT

Unloaded quality factor of a shielded dielectric resonator being high, its manual direct measurement can cause considerable errors [7].  $Q_u$  can, however, be accurately determined from the measurement of  $Q_L$  only, with the help of a simple graphical method using a Smith chart. This is possible because it can be shown that the  $Q_L$  point  $P$  (Fig. 4) when jointed to  $F$ , intersects the impedance locus at the point  $N$  which corresponds to  $Q_u$ . To prove this, it is sufficient to show that the angle  $PFO$  is the same as angle  $NFO$  (Fig. 4). Using the values given in Table I we have

In triangle  $OPF$ :

$$\text{angle } POF = 45^\circ$$

$$OP = S_{11_0}/\sqrt{2}.$$

Drawing  $PL$  perpendicular to  $OF$

$$OL = PL = S_{11_0}/2$$

$$\text{and } LF = 1 - S_{11_0}/2.$$

$$\text{Hence, angle } PFO = \tan^{-1} \frac{PL}{LF} = \tan^{-1} \frac{S_{11_0}}{2 - S_{11_0}}.$$

In triangle  $ONF$ :

From Table I

$$\text{angle } NOF = \tan^{-1} (1 - S_{11_0}) = \theta_u$$

$$\text{and } ON = \frac{S_{11_0}}{\sqrt{S_{11_0}^2 - 2S_{11_0} + 2}}.$$

Drawing  $NM$  perpendicular to  $OF$

$$NM = ON \sin \theta_u$$

$$OM = ON \cos \theta_u$$

$$\text{and } MF = 1 - ON \cos \theta_u$$

$$\begin{aligned} \text{angle } NFO &= \tan^{-1} \frac{NM}{MF} = \tan^{-1} \frac{ON \sin \theta_u}{1 - ON \cos \theta_u} \\ &= \tan^{-1} \frac{S_{11_0}}{2 - S_{11_0}}. \end{aligned}$$

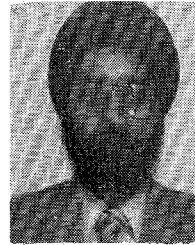
Hence, angle  $\angle NFO$  is the same as angle  $\angle PFO$ .

Once the point  $N$  is determined, the corresponding frequency, and hence  $Q_u$ , can be evaluated accurately. The  $Q_u$  determined experimentally using this method did not show more than  $\pm 1$  percent variation for various values of coupling factor  $\beta$ .

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