

Q conceptions

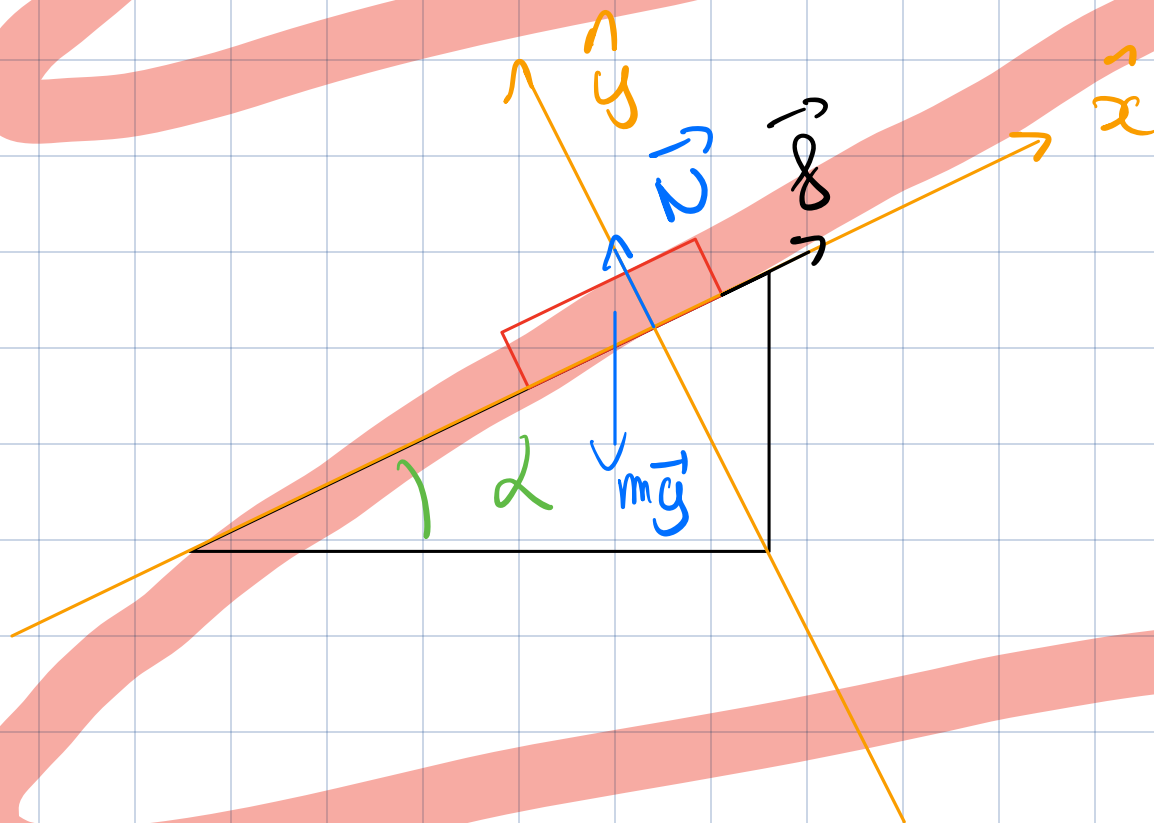
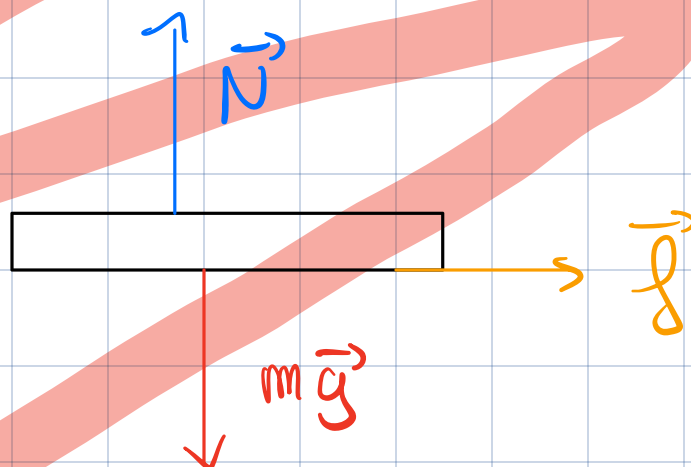
- ① force de frottement, le sol.
- ② force de frottement de l'air dépend de la masse volumique.

Si masse vol égale, alors chute au temps.

- ③ Oui (force normale)

Exercise 1

(a)



$$\sum \vec{F} = m\vec{a}$$

$$= m\vec{g} + \vec{N} + \vec{f}$$

On projette :

$$\begin{cases} m\ddot{y} = N - mg \cos \alpha \\ m\ddot{x} = f - mg \sin \alpha \end{cases}$$

On pose $f = f_{\max} = N \mu_s$.

on a $\ddot{x} = 0$ et $\ddot{y} = 0$.

$$\begin{cases} N = mg \cos \alpha \\ N \mu_s = mg \sin \alpha \end{cases}$$

$$\mu_s = \tan \alpha.$$



Exercise 1

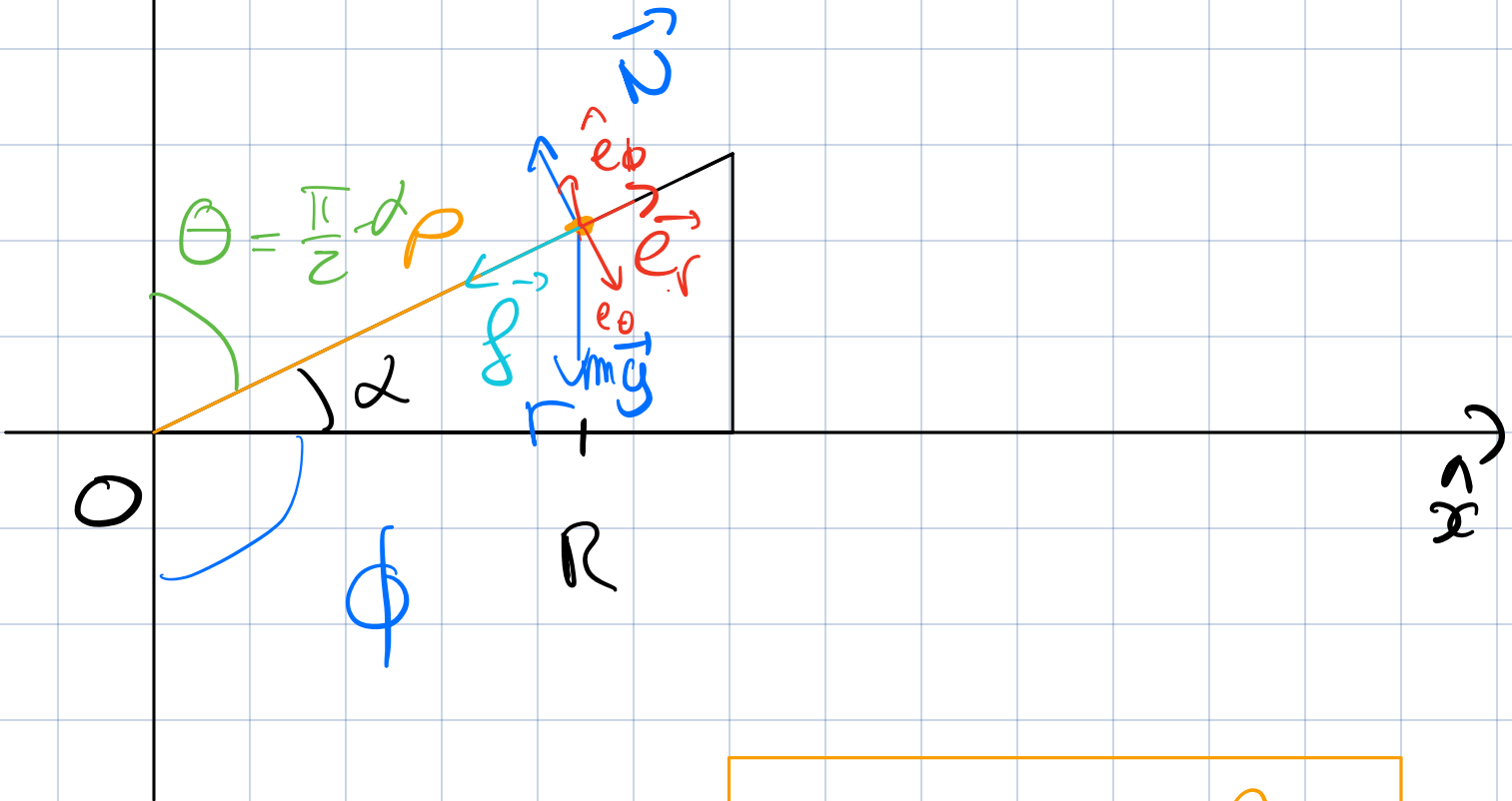
$$a_N = \frac{v^2}{R} \quad \vec{v}_N$$

\vec{y}

$$m\vec{g}_N + \vec{f}_N$$

$$\frac{v^2}{R}$$

$$= mg \sin \alpha$$



$$\sin(\theta) = \frac{R}{\rho}$$

$$\sum \vec{F}_i = m\vec{a}$$

$$= \vec{N} + m\vec{g} + \vec{f}$$

On projette

$$\sin\left(\frac{\pi}{2} - \alpha\right) = \cos(\theta)$$

$$\begin{cases} m_{ap} = f - \sin(\alpha) mg \\ m_{a\theta} = \cos(\alpha) mg + N \\ m_{a\phi} = 0 \end{cases}$$

$\cos\left(\frac{\pi}{2} - \alpha\right) = \sin(\theta)$

- on pose $\rho \text{ constant} = \frac{R}{\sin \theta}$
- $\dot{\phi} \text{ constant} = \frac{v}{R}$
- $\dot{\theta} \text{ constant} = \frac{\pi}{2} - \alpha$
- $\mu_c N = f$

$$\left\{ \begin{array}{l} -m \frac{R}{\cancel{\sin \theta}} \frac{v^2}{R^2} \cancel{\sin^2 \theta} = f - \sin \alpha \, mg \quad (\vec{e}_r) \\ -m \frac{R}{\cancel{\sin \theta}} \frac{v^2}{R^2} \cancel{\sin \theta} \cos \theta \, \vec{e}_\theta = \cos \alpha \, mg - N \quad (\vec{e}_\theta) \\ 0 = 0 \quad (\vec{e}_\phi) \end{array} \right.$$

$$\Rightarrow \left\{ \begin{array}{l} -\frac{mv^2}{R} \cos \alpha = -f - \sin \alpha \, mg \\ -\frac{mv^2}{R} \sin \alpha = \cos \alpha \, mg - N \end{array} \right.$$

$$\Rightarrow \left\{ \begin{array}{l} f = -\sin \alpha \, mg + \frac{mv^2}{R} \cos \alpha \\ N = -\cos \alpha \, mg - \frac{mv^2}{R} \sin \alpha \end{array} \right.$$

$$\Rightarrow \begin{cases} f = mg \cos \alpha \left(\frac{-\sin \alpha \cancel{mg}}{\cancel{mg} \cos \alpha} + \frac{v^2}{Rg} \right) \\ N = \cos \alpha mg \left(1 + \frac{v^2}{Rg} \tan \alpha \right) \end{cases}$$

$$\Rightarrow \begin{cases} f = mg \cos \alpha \left(-\tan \alpha + \frac{v^2}{Rg} \right) \\ N = mg \cos \alpha \left(1 + \frac{v^2}{Rg} \tan \alpha \right) \end{cases}$$

On veut $f \leq \mu_s N$ (la force max qu'on peut exercer avant que la voiture dérape et qu'on passe en frottements cinétiques)

$$f \leq \mu_s N$$

$$\Leftrightarrow (\cancel{mg \cos \alpha}) \left(\frac{v^2}{Rg} - \tan \alpha \right) \leq \mu_s (\cancel{mg \cos \alpha}) \left(1 + \frac{v^2}{Rg} \tan \alpha \right)$$

$$\Leftrightarrow \frac{v^2}{Rg} - \tan \alpha \leq \mu_s \left(1 + \frac{v^2}{Rg} \tan \alpha \right)$$

$$\Leftrightarrow \frac{v^2}{Rg} - \tan \alpha - \mu_s - \frac{v^2}{Rg} \tan \alpha \mu_s \leq 0.$$

$$\Leftrightarrow \frac{v^2}{Rg} (1 - \tan \alpha \mu_s) \leq \mu_s + \tan \alpha.$$

Si $1 - \tan \alpha \mu_s \leq 0$:

pas de dérapage

$$-\tan \alpha \mu_s \leq -1$$

$$\tan \alpha \mu_s > 1$$

$$\Leftrightarrow \boxed{\mu_s \geq \frac{1}{\tan \alpha}}$$

Si $1 - \tan \alpha \mu_s > 0$:

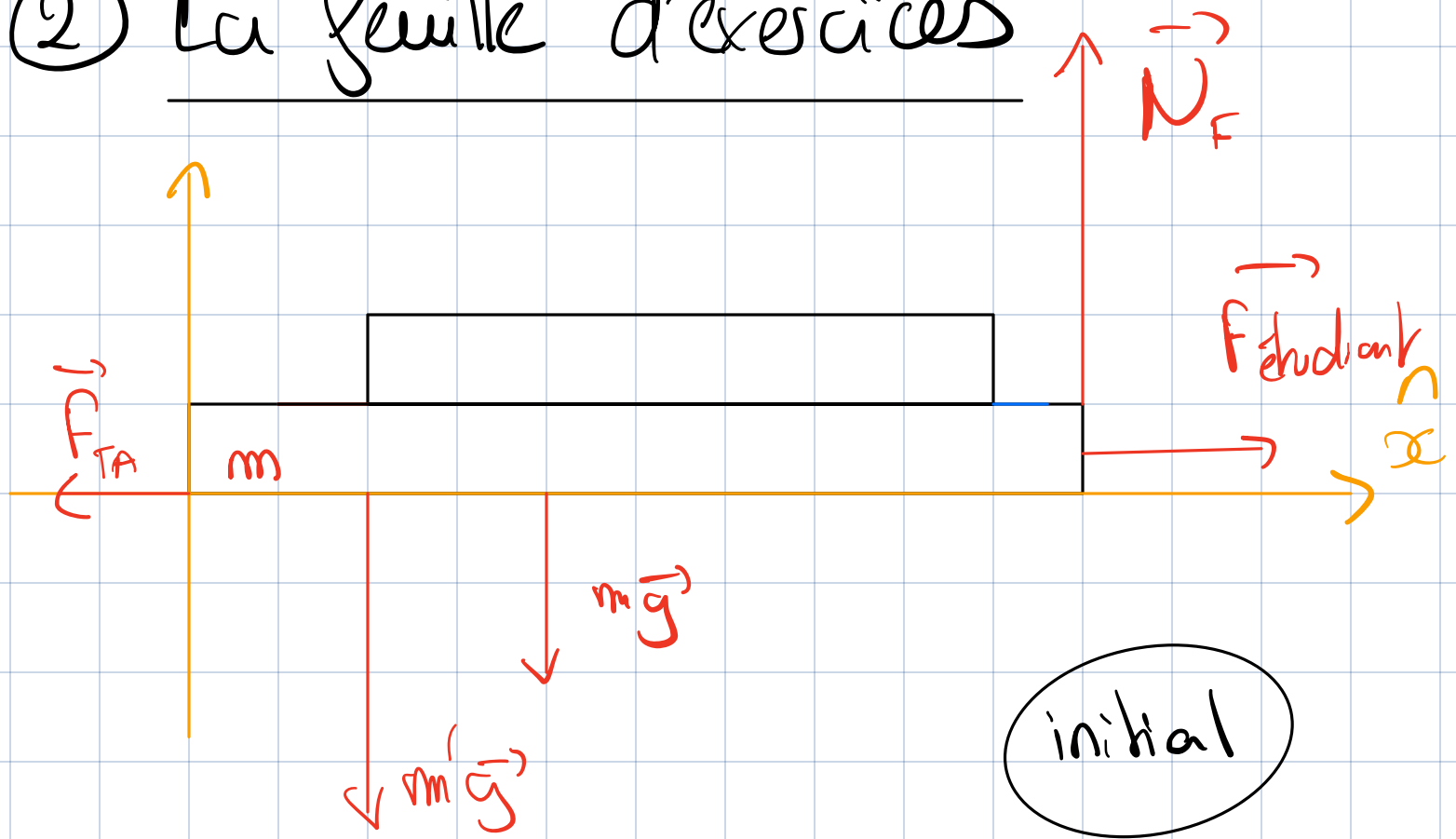
$$\frac{v^2}{R_g} \leq \frac{\mu_s + \tan \alpha}{1 - \tan \alpha \mu_s}$$

$$\Leftrightarrow v \leq \sqrt{R_g \left(\frac{\mu_s + \tan \alpha}{1 - \tan \alpha \mu_s} \right)}$$

Application numérique :

$$v \leq 72 \text{ m} \cdot \text{s}^{-1}.$$

② La feuille d'exercices



Système de la feuille + téléphone

$$\sum \vec{F} = (m + m') \vec{a}$$

$$= m'\vec{g} + m\vec{g} + \vec{N}_F + \vec{F}_e + \vec{F}_{TA}$$

On projette :

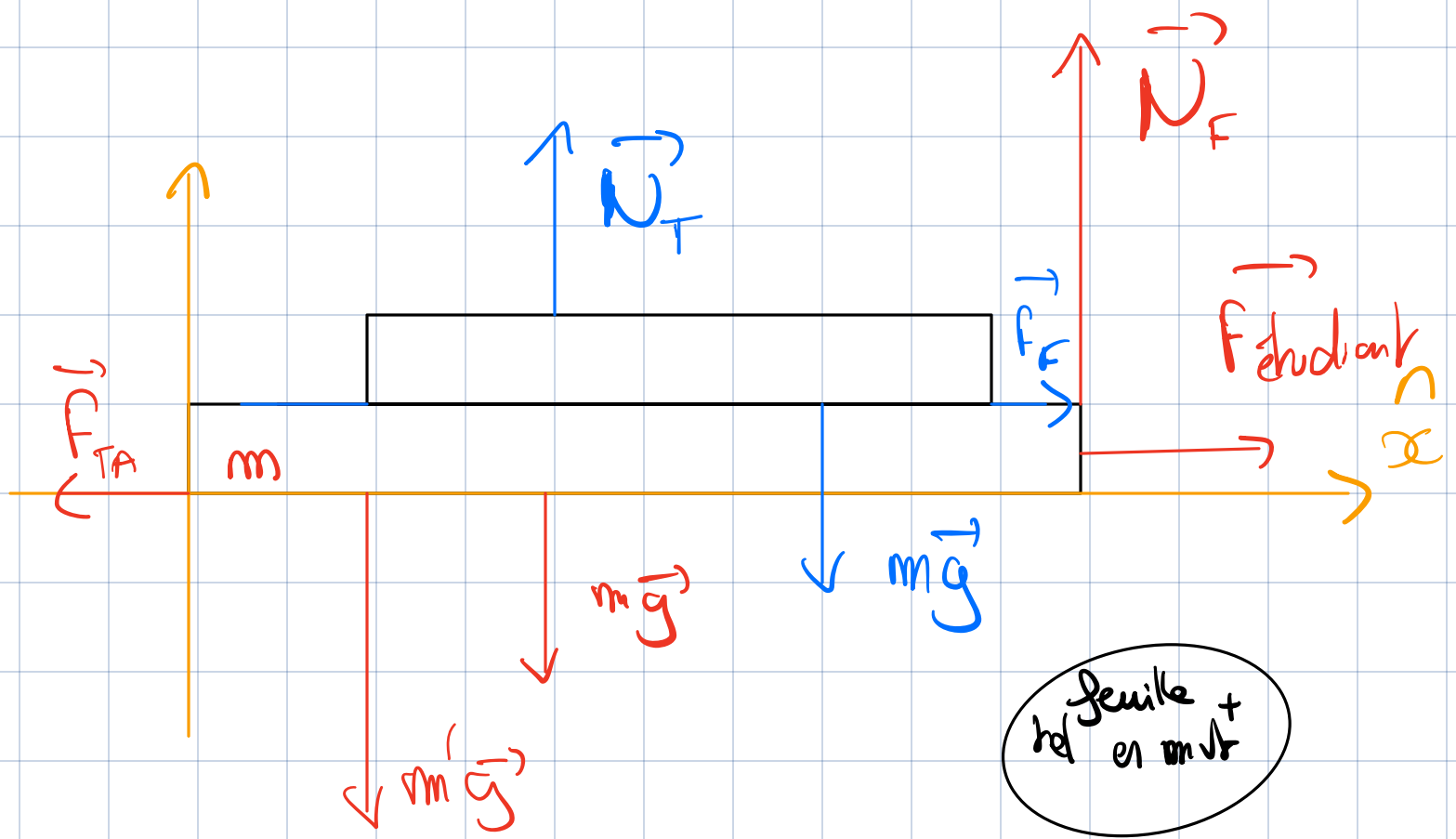
$$\begin{cases} (m+m')\ddot{y} = (m'+m)g - N \\ (m+m')\ddot{x} = F_e - F_{TA} \end{cases}$$

On pose $F_{TA} = \mu_s \cdot N$ et $\ddot{x} = 0$.

$$\begin{cases} N = (m'+m)g \\ 0 = F_e - \mu_s \cdot N - \mu_s' \cdot N \end{cases}$$

$$F_e \geq \mu_s \cdot (m'+m)g$$

Maintenant, on sépare les deux systèmes



Pour le système tel + feuille :

$$\begin{cases} N = (m' + m)g \\ (m + m')\ddot{x} = F_e - \mu_c \cdot (m + m')g. \end{cases}$$

Pour le système téléphone;

$$\Sigma \vec{F}' = m \vec{a}'$$

$$= m \vec{g}' + \vec{N}_T + \vec{F}_F$$

On projette

$$\begin{cases} m \ddot{y}' = mg - N \\ m \ddot{x}' = F_F \end{cases}$$

$$\text{On pose } F_F = \mu_s' \cdot N$$

or, avant que le tel bouge par rapport à la feuille $\ddot{x}' = \ddot{x}$.

$$\Leftrightarrow \text{tant que } F_F < F_F^{\text{max}}.$$

$$\text{on a } \ddot{x} = \frac{F_F}{m},$$

$$\begin{aligned} F_e &= (m + m') \ddot{x} + \mu_c \cdot (m + m') g \\ &= (m + m') \cdot \frac{F_F}{m} + \mu_c (m + m') g \\ &= (m + m') \cdot \frac{\mu_s' \cdot m g}{m} + \mu_c (m + m') g \\ &= (\mu_s' + \mu_c) \cdot (m + m') \cdot g \end{aligned}$$

mais peut-être que $\mu_c > \mu_s$ (la force minimale pour faire bouger la feuille une fois serait trop faible pour la faire

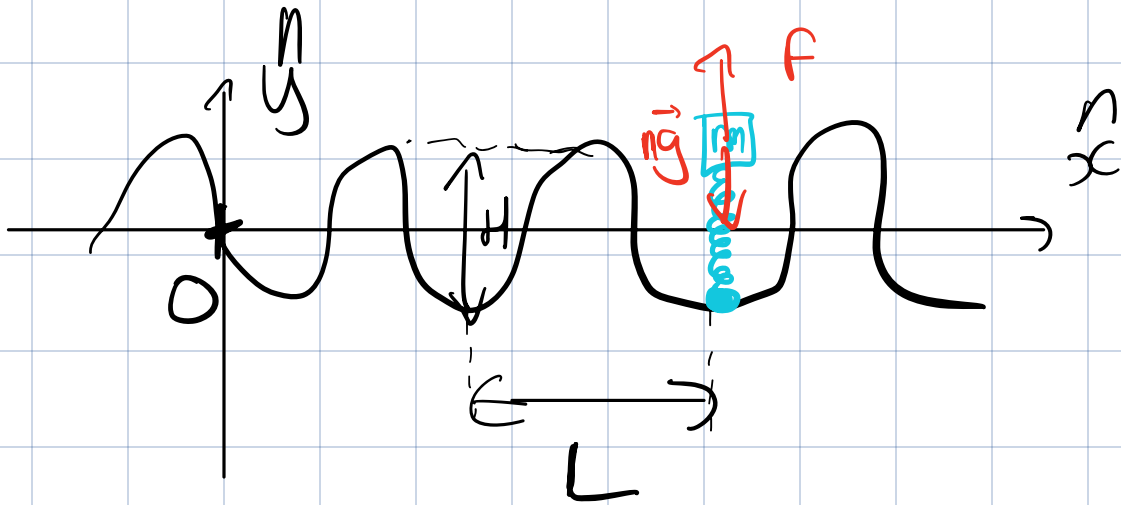
gläser).

$$F_e \geq (\max(\mu_s; (\mu_s' + \mu_c)))$$

$$- (m + m') \cdot g.$$

③ Champ de bosse

$$\dot{x} = v_x \Rightarrow \ddot{x} = 0.$$



$$h(t) = \frac{H}{2} \cos(\omega t)$$

$$T = \frac{2\pi}{\omega} \Leftrightarrow \boxed{\omega = \frac{2\pi}{T}}$$

$$T = \frac{L}{v_x} \quad \text{donc} \quad \omega = \frac{2\pi}{L} \cdot v_x$$

$$h(t) = \frac{K}{2} \cos\left(\frac{2\pi}{L} \cdot v_x t\right)$$

$$\begin{aligned} \sum \vec{F} &= m\vec{a} \\ &= \vec{P} + \vec{F} \end{aligned}$$

$$\begin{cases} m \ddot{y} = mg - k \Delta z \\ m \ddot{x} = 0 \end{cases}$$

$$\Delta y = y(t) - h(t) - l_0$$

$$\begin{cases} m\ddot{y} = mg - k(y - h - l_0) \\ 0 = 0 \end{cases}$$



$$m\ddot{y} = mg - ky + kh + kl_0$$

$$m\ddot{y} = -mg - ky + k\left(\frac{H}{2} \cos\left(\frac{2\pi}{L} \cdot v_{sc} t\right)\right) + kl_0$$

$$\ddot{y} = -g - \frac{k}{m} y + \frac{k}{m} \frac{H}{2} \cos\left(\frac{2\pi}{L} \cdot v_{sc} t\right) + \frac{k}{m} l_0$$

$$\omega_0 = \sqrt{\frac{k}{m}}$$

$$\ddot{y} = -g - \omega_0^2 y + \omega_0^2 \frac{H}{2} \sin\left(\left(\frac{\pi}{2} - \frac{2\pi}{L} \cdot v_{sc}\right)t\right) + \omega_0^2 l_0$$

$$\ddot{y} + g + \omega_0^2 y - \omega_0^2 l_0 = \omega_0^2 \frac{H}{2} (-)$$

$$\Leftrightarrow \ddot{y} + \omega_0^2 \left(\frac{g}{\omega_0^2} + y - l_0\right) = \omega_0^2 \frac{H}{2} (-)$$

$$u = \frac{g}{\omega_0^2} - l_0 + y \quad | \quad \alpha_0 = \omega_0^2 \frac{H}{2}$$

$$\Rightarrow \ddot{v} + \omega_0^2 v = \alpha_0 \sin(\omega t)$$

$$\omega = \frac{\pi}{2} - \frac{2\pi}{c} \cdot v_x$$

$$\textcircled{d} \begin{cases} v(t) = p \sin(\omega t - \phi) & \phi = 0 \\ \dot{v}(t) = p\omega \cos(\omega t) \\ \ddot{v}(t) = -p\omega^2 \sin(\omega t) \end{cases}$$

$$-p\omega^2 \sin(\omega t) + \omega_0^2 (p \sin(\omega t)) = \alpha_0 \sin(\omega t)$$

$$\Leftrightarrow \sin(\omega t) (-p\omega^2 + p\omega_0^2) = \alpha_0 \sin(\omega t)$$

$$\Leftrightarrow -p\omega^2 + p\omega_0^2 = \alpha_0$$

$$\Leftrightarrow p(-\omega^2 + \omega_0^2) = \alpha_0$$

$$\Leftrightarrow p = \frac{\alpha_0}{-\omega^2 + \omega_0^2}$$

$$\rho = \frac{M}{2} \cdot \omega_0^2 \cdot \frac{1}{-\omega^2 + \omega_0^2}$$

$$= \frac{M}{2} \cdot \frac{\cancel{\omega_0^2}}{\cancel{\omega_0^2} \left(1 + \frac{-\omega^2}{\omega_0^2} \right)}$$

$$= \frac{M}{2} \left(\frac{1}{1 - \frac{\omega^2}{\omega_0^2}} \right)$$

$$\omega = \omega_0, \text{ resonance } \quad \times$$

$$-\sqrt{\frac{k}{m}} = \frac{\pi}{2} - \frac{2\pi}{L} V_x$$

$$\frac{2\pi}{L} \cdot k = \frac{\pi}{2}$$

$$-\sqrt{\frac{k}{m}} - \frac{\pi}{2} = -\frac{2\pi}{L} V_x$$

$$k = \frac{L\pi}{4\pi}$$

$$\Rightarrow \frac{-L}{2\pi} \left(-\sqrt{\frac{k}{m}} - \frac{\pi}{2} \right) = V_x$$

$$\Rightarrow V_x = \frac{L}{2\pi} \left(-\sqrt{\frac{k}{m}} + \frac{\pi}{2} \right)$$

$$\left[\begin{array}{l} w \rightarrow +\infty \\ \text{alors } V_x \rightarrow -\infty \\ \text{et } p \rightarrow 0 \end{array} \right.$$

$$\left[\begin{array}{l} w \rightarrow -\infty \\ \text{alors } V_x \rightarrow +\infty \\ p \rightarrow 0 \end{array} \right.$$

$$\left[\begin{array}{l} w \rightarrow 0 \\ \text{alors } V_x \rightarrow \frac{L}{4} \\ \text{et } p \rightarrow 0 \end{array} \right.$$