

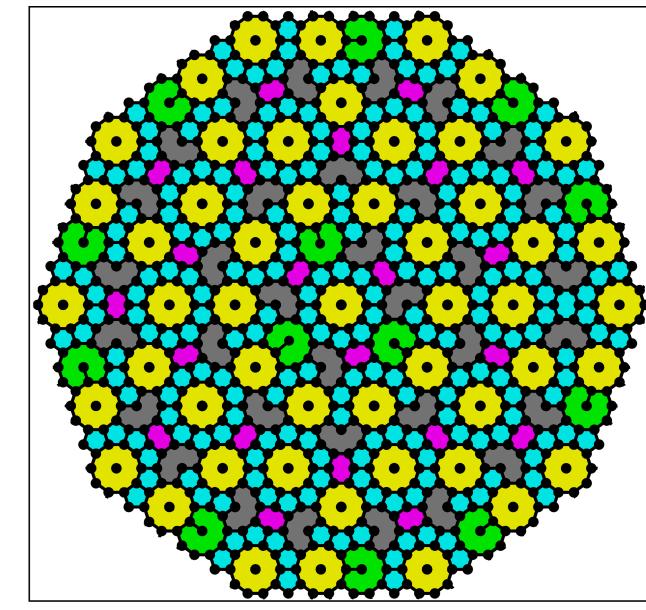
# Thermodynamical Stability Analysis of a Model Quasicrystal

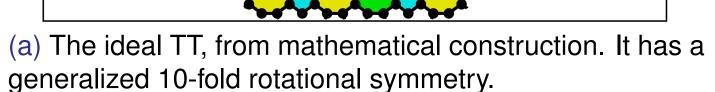
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#### Introduction

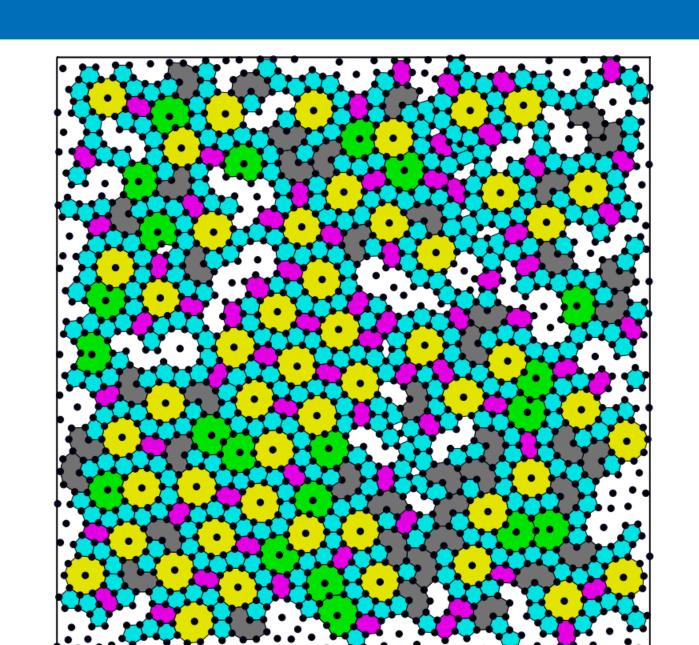
- According to the random tiling hypothesis quasicrystals (QCs) are high temperature phases that stabilize due to large entropy [1].
- We prove the hypothesis by using the decagonal Tübingen tiling as a model QC.
- We apply exclusively geometric methods, namely the polar calculus.

### The Tübingen tiling (TT)

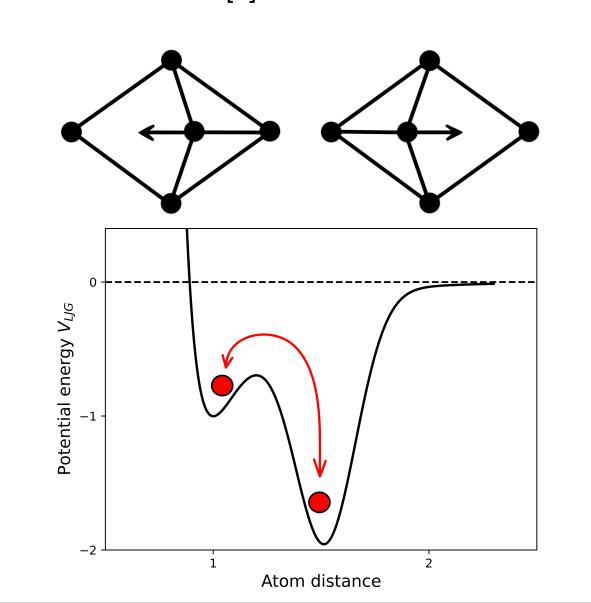




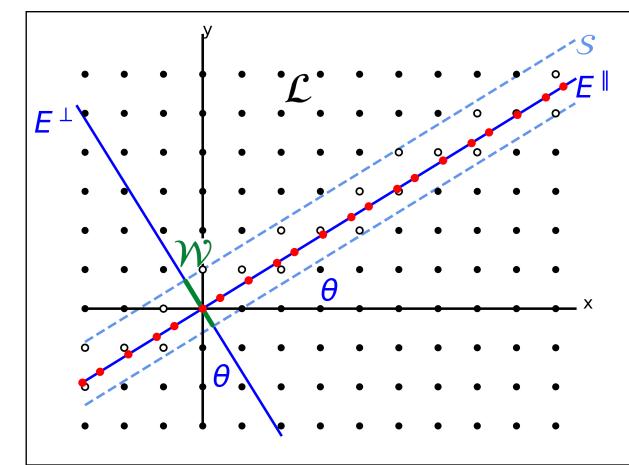
- A random tiling arises from the ideal TT through atom flips.
- Random flips do not change the generalized 10-fold rotational symmetry.
- A randomized TT arises also in a MD-simulation with a Lennard-Jones Gauß-potential.
- Many possible flip configurations lead to a large entropy.



(b) A randomized version of the TT, from a molecular dynamics simulation [2].



### Constructing crystals by cut and project scheme (CPS)

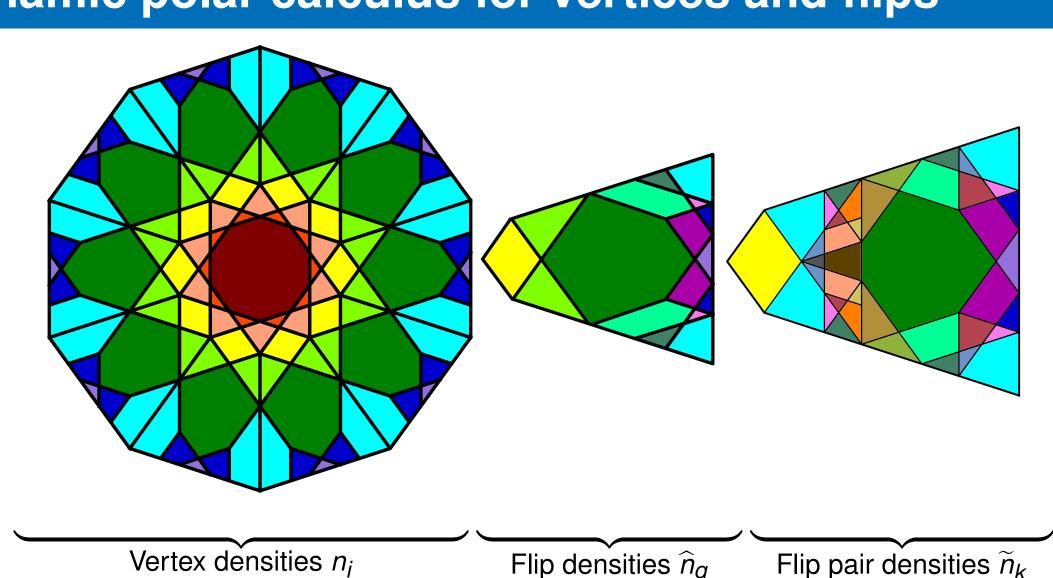


orientation.

(a) Cut space S and projection space  $E^{\parallel}$  have the same (b) Cut space S and projection space  $E^{\parallel}$  have different orientation.

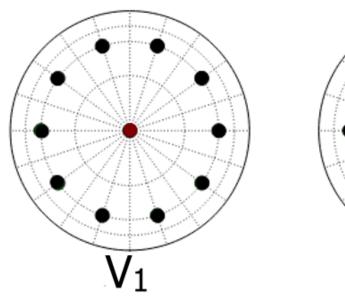
- 1D crystals can be constructed as the  $E^{\parallel}$ -projection of a 2D hyperlattice  $\mathcal{L} \cap \mathcal{S}$ .
- The structure of a projected 1D crystal is determined by phason strain:  $\chi = \tan(\theta \theta')$ .
- The window  $\mathcal{W}$  is the  $E^{\perp}$ -projection of  $\mathcal{L} \cap \mathcal{S}$ .
- Phason strain expresses itself in the tiling by flips.
- The CPS for the TT requires a 4D hyperlattice and yields two independent strain components:  $\chi = (\chi_1, \chi_2)$ .

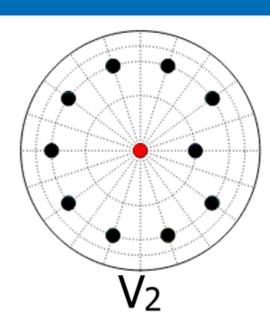
# The dynamic polar calculus for vertices and flips

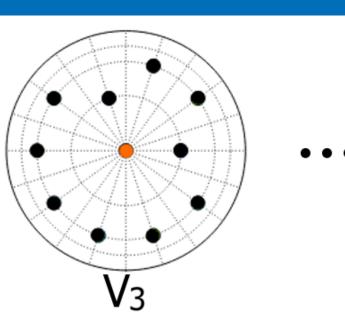


- The TT's window is a regular decagon that will be deformed by applying phason strain.
- ullet Every atom configuration corresponds to a domain in  ${\mathcal W}$ .
- The density of such configurations is given by the domain's area.

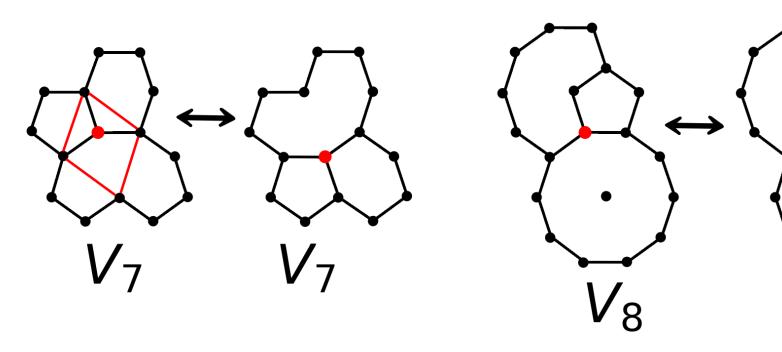
### Atom environments and flip configurations in the TT for a cut-off radius $r_c = 2$



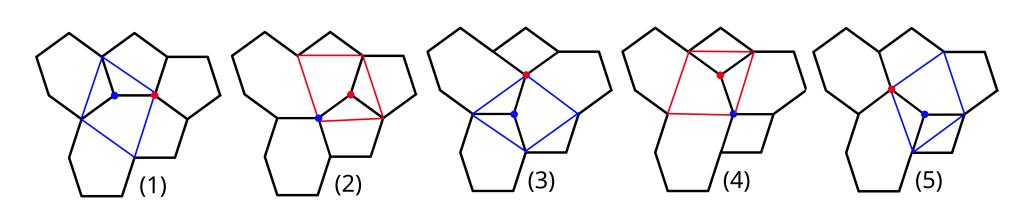




• The energy density of one specific tiling is determined by its atom environments  $V_i$  and their potential energies  $E_i$ .



- ullet A fixed average phason strain  $\chi$  corresponds to certain flip configurations and determines the free energy density  $F(\chi, T)$ .
- Asymmetric flips may increase or decrease the tiling energy depending on the energy difference of their two states,  $\Delta E \leq 0$ .



• The simplest model of correlated flips consists of pairs of coupled flips and has at most 5 non-degenerate states.

### The free energy density

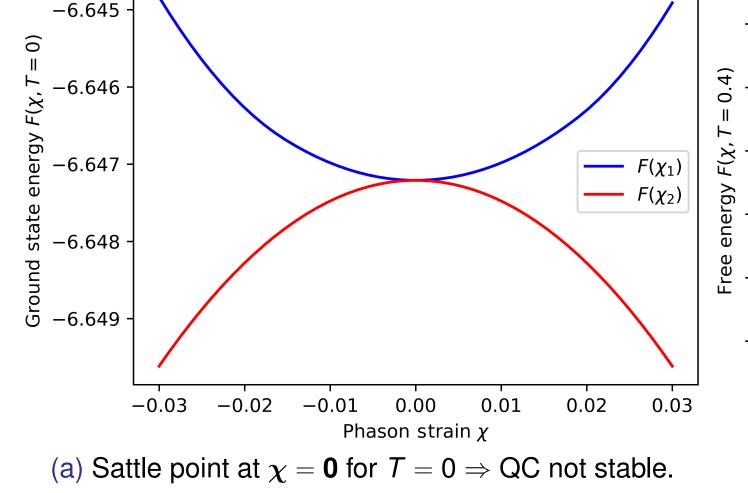
 $\sum n_i(\chi)E_i +$ projected tiling

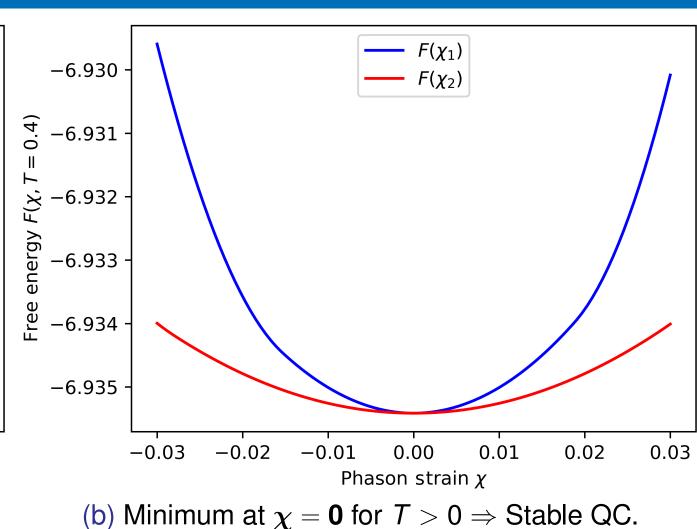
Flips  $q \mid \Delta E_q < 0$ flip relaxation

 $\widehat{n}_q(\chi)\Delta E_q+\sum \widetilde{n}_k(\chi)\ln(\sum \mathrm{e}^{-k_B\Delta E_k^{(i)}/T})$ Flip pairs *k* flip randomization

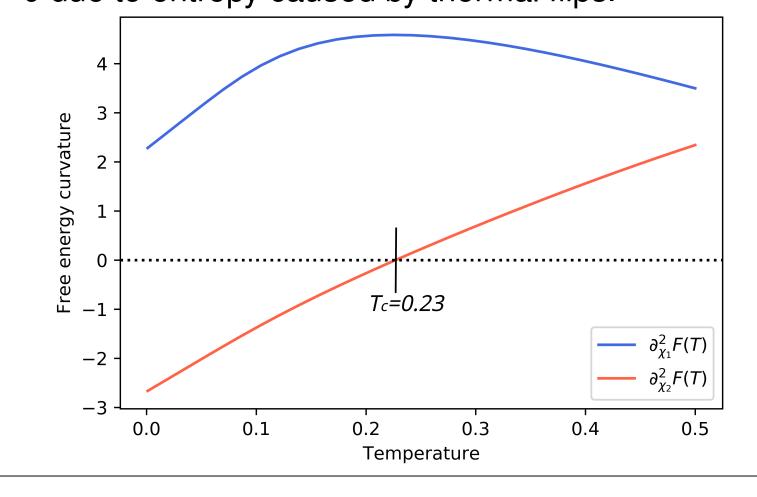
- $n_i$ ,  $\hat{n}_q$  and  $\hat{n}_k$  are the densities of atom environments, flip configurations or flip pair configurations for a given average  $\chi$ .
- These quantities are calculated by a dynamic polar calculus.

# Free energy at low and high temperature





- $F(\chi, T = 0)$  has saddle point at  $\chi = 0$ . •  $F(\chi, T)$  is minimized at  $\chi = \mathbf{0}$  for T > 0 due to entropy caused by thermal flips.
- The phonon free energy is  $\chi$ -independent [3].
- The free energy curvature  $\partial_{\chi_2}^2 F(\chi, T)|_{\chi=0}$  becomes positive above  $T_C = 0.23$ and hence, the QC is stabilized.



#### References:

- [1] Henley C. "Random Tiling Models". In: Quasicrystals The State Of The Art. Singapore: World Scientific Publishing, 1991, pp. 429–524.
- Engel M. "Dynamics and Defects of Complex Crystals and Quasicrystals: Perspectives from Simple Model Systems". PhD thesis. Universität Stuttgart, 2008.
- Kiselev A. "Phasonen in quasikristallinen Strukturen des Lennard-Jones-Gauß-Systems". Diploma thesis. Universität Stuttgart, 2011.

