

I)

$$\left( -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + g|x| \right) \psi(x) = E \psi(x)$$

$$\left[ \frac{\partial^2}{\partial x^2} + \frac{2m}{\hbar^2} (E - g|x|) \right] \psi(x) = 0$$

$$\left[ \frac{mE}{\hbar^2} \right] = L^{-2} \quad , \quad \left[ \frac{mg}{\hbar^2} \right] = L^{-3}$$

$$\left[ \left( \frac{mE}{\hbar^2} \right)^3 \right] = \left[ \left( \frac{mg}{\hbar^2} \right)^2 \right]$$

$$[E] = \left[ \left( \frac{\hbar^2 g^2}{2m} \right)^{1/3} \right]$$

$$\Rightarrow E_n = \left( \frac{\hbar^2 g^2}{m} \right)^{1/3} f(n)$$

$$\psi_{a,c}(x) = c \theta(x+a) \theta(a-x) \left( 1 - \frac{|x|}{a} \right)$$

$$\hookrightarrow 1 \stackrel{!}{=} \int \psi^*(x) \psi(x) dx$$

$$= |c|^2 \int \theta(x+a) \theta(a-x) \left( 1 - \frac{|x|}{a} \right)^2 dx$$

$$= |c|^2 \int_{-a}^a \left( 1 - \frac{|x|}{a} \right)^2 dx = \frac{2a}{3} |c|^2$$

$$\Rightarrow |c|^2 = \frac{3}{2a}$$

$$\bar{H} = \int \psi^* H \psi dx$$

$$= - \underbrace{\frac{\hbar^2}{2m} \int_{-\infty}^{\infty} \psi^*(x) \frac{d^2}{dx^2} \psi(x) dx}_{(I)} + g \underbrace{\int_{-\infty}^{\infty} \psi^*(x) |x| \psi(x) dx}_{(II)}$$

$$(II) = |c|^2 \left( - \int_{-a}^0 x \left(1 - \frac{x}{a}\right)^2 dx + \int_0^a x \left(1 - \frac{x}{a}\right)^2 dx \right)$$

$$= \frac{a^2 |c|^2}{6} = \frac{a}{4}$$

$$\begin{aligned} \frac{d}{dx} \psi(x) &= c \delta(x+a) \theta(a-x) \left(1 - \frac{|x|}{a}\right) \\ &\quad - c \theta(x+a) \delta(x-a) \left(1 - \frac{|x|}{a}\right) \\ &\quad + c \theta(x+a) \theta(x-a) \left(-\frac{|x|}{x} \frac{1}{a}\right) \end{aligned}$$

$$(I) = \int_{-\infty}^{\infty} \psi^*(x) \frac{d^2}{dx^2} \psi(x)$$

$$= \psi^*(x) \frac{d}{dx} \psi(x) \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \left| \frac{d\psi}{dx} \right|^2 dx$$

$$= -|c|^2 \int_{-a}^a \left(-\frac{|x|}{xa}\right)^2 dx = -\frac{2|c|^2}{a} = -\frac{3}{a^2}$$

$$\Rightarrow \bar{H} = \frac{3\hbar^2}{2ma} + \frac{ag}{4}$$

$$\psi_{a,c}(x) = c \theta(x+a) \theta(a-x) \left(1 - \frac{|x|}{a}\right)$$

$$\frac{\delta \bar{H}}{\delta a} = -\frac{3\hbar^2}{ma^3} + \frac{g}{4} \stackrel{!}{=} 0 \Rightarrow a = \left(\frac{12\hbar^2}{gm}\right)^{1/3}$$

$$\bar{H} = \frac{3\hbar^2}{2m} \left(\frac{gm}{12\hbar^2}\right)^{2/3} + \frac{g}{4} \left(\frac{12\hbar^2}{gm}\right)^{1/3} = \frac{3}{4} \left(\frac{3\hbar^2 g^2}{2m}\right)^{1/3}$$

II) Zentralpot.:  $f = 2\hbar$  ,  $s = \hbar$

$$H_{so} = a \hat{L} \hat{S}$$

$$\mathcal{B} = \{H, J^2, J_z, L^2, S^2\}$$

$$\phi_{jmjel_s} \equiv \phi : \quad J^2 \phi_{jmjel_s} = \hbar^2 j(j+1) \phi$$

$$L^2 \phi_{jmjel_s} = \hbar^2 l(l+1) \phi$$

$$S^2 \phi = \hbar^2 s(s+1) \phi$$

$$J_z \phi = \hbar m_j \phi$$

$$J = L + S$$

$$\Rightarrow H_{so} = a L S = \frac{1}{2} a (J^2 - L^2 - S^2)$$

$$\Rightarrow E_{so} = \frac{\hbar^2}{2} a [\bar{j}(\bar{j}+1) - \ell(\ell+1) - s(s+1)]$$

$$= \begin{cases} 2a\hbar^2 & , \bar{j} = 3 \\ -a\hbar^2 & , \bar{j} = 2 \\ -3a\hbar^2 & , \bar{j} = 1 \end{cases}$$

Entartung:

$$d = 2\bar{j} + 1 = \begin{cases} 7 & , \bar{j} = 3 \\ 5 & , \bar{j} = 2 \\ 3 & , \bar{j} = 1 \end{cases}$$