Continuum Mechanics - A (2022 Fall)

(Reference Solutions Manual)

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评分总则:满分 200,共 8 大题。每大题 25 分。按照预期每次作业占总评 100 分的 2 分,也就是作业的每 100 分相当于总评的 1 分。接受误判及分数错误求和这两种情况的 argument。

Section 7 pp.57

Exercise 1. Given one-parameter family of deformations with $|\nabla u_{\epsilon}| = \epsilon$ small, show that:

1.
$$E_{\epsilon} = U_{\epsilon} - I + o(\epsilon) = V_{\epsilon} - I + o(\epsilon)$$

2.
$$\det F_{\epsilon} - 1 = div(u_{\epsilon}) + o(\epsilon)$$

Give a physical interpretation of det $F_{\epsilon} - 1$.

Solution.

For the sake of briefness, we omit the subscript ϵ .

For 1, recall that
$$U = \sqrt{C} = \sqrt{F^T F}$$
, $C = I + \nabla u + \nabla u^T + \nabla u^T \nabla u$, thus

$$C = I + 2E + \nabla u^T \nabla u$$

Let $H := \nabla u$, by definition of Frechet derivative, along with the smiling inequality $|AB| \le C(n)|A||B|$, we have

$$C(H) = C(0) + DC(0)[H] + o(H)$$

$$= C(0) + 2DU(0)[H] + o(H)$$

$$= C(0) + 2E(H) + o(H)$$

Therefore,

$$U(H) = U(0) + DU(0)[H] + o(\epsilon)$$
$$= U(0) + E + o(\epsilon)$$
$$= I + E + o(\epsilon)$$

Thus $E = U - I + o(\epsilon)$. $E = V - I + o(\epsilon)$ is obtain in an analogous way.

For 2,

$$F = \begin{pmatrix} u_{1,1} + 1 & u_{1,2} & u_{1,3} \\ u_{2,1} & u_{2,2} + 1 & u_{2,3} \\ u_{3,1} & u_{3,2} & u_{3,3} + 1 \end{pmatrix}$$

Then,

$$\det F = (u_{1,1} + 1)(u_{2,2} + 1)(u_{3,3} + 1) + o(\epsilon)$$
$$= u_{1,1} + u_{2,2} + u_{3,3} + o(\epsilon)$$
$$= div(u) + o(\epsilon)$$

Physical interpretation is given in solution to exercise 4 in hw3: that is, the volume change per unit volume under the deformation f_{ϵ} .

评分细则:

Exercise 4. Let $W = \frac{1}{2}(\nabla u - \nabla u^T)$. Show that

1.
$$|E|^2 + |W|^2 = |\nabla u|^2$$

$$2. |E|^2 - |W|^2 = \nabla u \cdot \nabla u^T$$

Solution.

$$|E|^{2} = \frac{1}{2}(\nabla u + \nabla u^{T}) \cdot \frac{1}{2}(\nabla u + \nabla u^{T})$$

$$= \frac{1}{4}(\nabla u \cdot \nabla u + \nabla u \cdot \nabla u^{T} + \nabla u^{T} \cdot \nabla u + \nabla u^{T} \cdot \nabla u^{T})$$

$$|W|^{2} = \frac{1}{2}(\nabla u - \nabla u^{T}) \cdot \frac{1}{2}(\nabla u - \nabla u^{T})$$

$$= \frac{1}{4}(\nabla u \cdot \nabla u - \nabla u \cdot \nabla u^{T} - \nabla u^{T} \cdot \nabla u + \nabla u^{T} \cdot \nabla u^{T})$$

Therefore,

$$|E|^{2} + |W|^{2} = \frac{1}{2}\nabla u \cdot \nabla u + \frac{1}{2}\nabla u^{T} \cdot \nabla u^{T}$$

$$= \frac{1}{2}|\nabla u|^{2} + \frac{1}{2}|\nabla u|^{2}$$

$$= |\nabla u|^{2}$$

$$|E|^{2} - |W|^{2} = \frac{1}{2}\nabla u \cdot \nabla u^{T} + \frac{1}{2}\nabla u^{T} \cdot \nabla u$$

$$= \nabla u \cdot \nabla u^{T}$$

评分细则:

Exercise 5. (Korn's inequality) Let $u \in C^2(\bar{\mathscr{R}})$ and suppose that u = 0 on $\partial \mathscr{R}$. Show that

$$\int_{\mathscr{R}} |\nabla u|^2 \, dx \le 2 \int_{\mathscr{R}} |E|^2 \, dx.$$

Solution. By exercise 4,

$$2\int_{\mathscr{R}} |E|^2 dx = \int_{\mathscr{R}} (|E|^2 + |W|^2) + (|E|^2 - |W|^2) dx$$
$$= \int_{\mathscr{R}} |\nabla u|^2 dx + \int_{\mathscr{R}} \nabla u \cdot \nabla u^T dx$$

It suffices to show that the second term is nonnegative. Use exercise 9 (b) in section 4, along with divergence theorem,

$$\int_{\mathcal{R}} \nabla u \cdot \nabla u^T dx = \int_{\mathcal{R}} \operatorname{div} ((\nabla u)u - (\operatorname{div} u)u) + (\operatorname{div} u)^2 dx
= \int_{\mathcal{R}} (\operatorname{div} u)^2 dx + \int_{\partial \mathcal{R}} ((\nabla u)u - (\operatorname{div} u)u) \cdot \overrightarrow{\nu} ds
= \frac{u|_{\partial \mathcal{R}} = 0}{D} \int_{\mathcal{R}} (\operatorname{div} u)^2 dx
\ge 0.$$

评分细则:

注: 广义 Korn's 不等式的意义在于定义索伯列夫空间 $H_0^1(\mathcal{R})$ 上 u 的等价范数,进而方便处理弹性力学中的 PDE 边界值问题。这里是具有二阶正则性时的 strong version。

Exercise 6. Pure torsion is given as

$$x_1 = r\cos(\theta), \quad p_1 = R\cos(\Theta),$$

 $x_2 = r\sin(\theta), \quad p_2 = R\sin(\Theta),$
 $x_3 = z, \qquad p_3 = Z,$

where

$$r = R$$
, $\theta = \Theta + \alpha Z$, $z = Z$

Find the displacement components in terms of p_i , i = 1, 2, 3 and $\beta = \alpha p_3$. Furthermore, show that

1.

$$\nabla u \to 0$$
 and $u \to 0$ as $\alpha \to 0$;

2.

$$u_1(p,\alpha) = -\alpha p_2 p_3 + o(\epsilon),$$

$$u_2(p,\alpha) = \alpha p_1 p_3 + o(\epsilon).$$

Solution.

$$u_1(p) = R\cos(\Theta + \alpha Z) - R\cos(\Theta)$$

$$= R\left[\cos(\Theta)\cos(\beta) - \sin(\Theta)\sin(\beta)\right] - R\cos(\Theta)$$

$$= p_1\cos(\beta) - p_2\sin(\beta) - p_1$$

$$= p_1(\cos(\beta) - 1) - p_2\sin(\beta)$$

In the same fashion, it can be easily concluded that

$$u_2(p) = p_2(\cos(\beta) - 1) + p_1\sin(\beta)$$

Finally,

$$u_3(p) = z - Z = 0$$

Since $\sin(\gamma) \sim \gamma$, $\cos(\gamma) \sim 1$ for small γ , the asymptotic expressions for u_1, u_2 as $\alpha \to 0$ are

$$u_1(p,\alpha) = p_1(\cos(\beta) - 1) - p_2\sin(\beta)$$

$$= -p_2\beta + o(\alpha)$$

$$= -\alpha p_2 p_3 + o(\alpha), \text{ as } \alpha \to 0$$

$$u_2(p,\alpha) = p_2(\cos(\beta) - 1) + p_1\sin(\beta)$$

$$= p_1\beta + o(\alpha)$$

$$= -\alpha p_1 p_3 + o(\alpha), \text{ as } \alpha \to 0$$

Compute the values at $\alpha = 0$:

$$u(p,0) = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$$

$$\nabla_p u(p,0) = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

By the sequential continuity of u and ∇u , we obtain

$$u \xrightarrow{\alpha \to 0} 0, \quad \nabla u \xrightarrow{\alpha \to 0} 0$$

评分细则:没有过程扣5分。

Section 8 pp.66

Exercise 2. Show that

$$\dot{C} = 2F^T D_m F$$

Solution.

$$\dot{C}(p,t) = F^T \dot{F} + (\dot{F}^T) F$$

$$= F^T \dot{F} + (\dot{F})^T F$$

$$2F^T D_m F = F^T (\operatorname{grad} v)_m F + F^T \left[(\operatorname{grad} v)^T \right]_m F$$

$$= F^T (\operatorname{grad} v)_m F + F^T \left[(\operatorname{grad} v)_m \right]^T F$$
(1)

To prove $\dot{C} = 2F^T D_m F$, it suffices to show that

$$\dot{F}(p,t) = (\operatorname{grad} v)_m F(p,t)$$

because $\dot{F}^T = F^T \left[(\operatorname{grad} v)_m \right]^T$. Equation (1) is already given in eqn. (8)₁ in section 8. 评分细则:

Exercise 4. Given a motion \mathscr{E}

$$x_1 = p_1 e^t,$$

$$x_2 = p_2 + t,$$

$$x_3 = p_3.$$

in some Cartesian frame. Compute the spatial velocity field v and determine the streamlines.

Solution. The reference mapping is

$$\begin{cases} p_1 = x_1 e^{-t} \\ p_2 = x_2 - t \\ p_3 = x_3 \end{cases}$$

The velocity field is by definition,

$$\dot{x}(p,t) = \left(\begin{array}{c} p_1 e^t \\ 1 \\ 0 \end{array}\right)$$

Applying the reference mapping p to obtain the spatial velocity field:

$$v(x,t) = \begin{pmatrix} x_1 \\ 1 \\ 0 \end{pmatrix}$$

The steamlines ODEs are

$$\begin{cases} \dot{s}_1 = s_1 \\ \dot{s}_2 = 1 \\ \dot{s}_3 = 0 \end{cases}$$

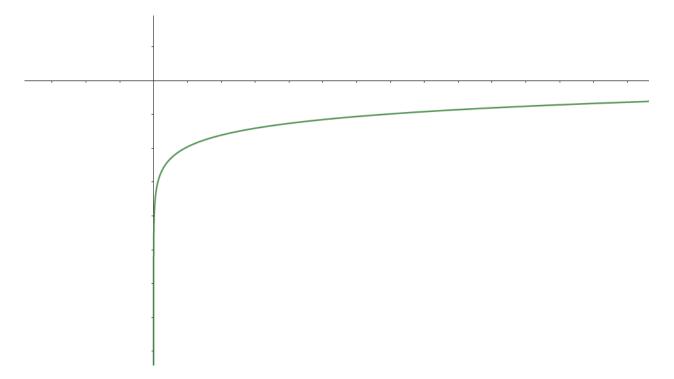
Integration shows that

$$\int \frac{ds_1}{s_1} = \int d\lambda$$
$$\ln^{|s_1|} = \lambda + \tilde{C}$$
$$s_1(\lambda) = Ce^{\lambda}$$

where \tilde{C} is integration constant. $s_1(0) = x_1(p,0) = y_1$, thus, $C = y_1$. After performing the same process for s_2, s_3 , we finally obtain the streamlines equations at time t passing through (y_1, y_2, y_3) at $\lambda = 0$:

$$\begin{cases} s_1(\lambda) = y_1 e^{\lambda} \\ s_2(\lambda) = \lambda + y_2 \\ s_3(\lambda) = y_3 \end{cases}$$

The sketch is



Exercise 5. Consider the motion ξ defined by

$$x(p,t) = p_0 + U(t)[p - p_0],$$

where

$$U(t) = \sum_{i=1}^{3} \alpha_i(t)e_i \otimes e_i$$

 $\alpha_i > 0$ are smooth and $\{e_i\}$ is an orthonormal basis. Compute p, v, L and determine the streamlines.

Solution.

We make the convention that superscript (i) always stands for the i-th Cartesian

component. Since

$$\begin{cases} x^{(1)} = p_0^{(1)} + \alpha_1(t)[p^{(1)} - p_0^{(1)}] \\ x^{(2)} = p_0^{(2)} + \alpha_2(t)[p^{(2)} - p_0^{(2)}] \\ x^{(3)} = p_0^{(3)} + \alpha_3(t)[p^{(3)} - p_0^{(3)}] \end{cases}$$

We have the reference mapping p:

$$\begin{cases} p^{(1)} = p_0^{(1)} + \frac{x^{(1)} - p_0^{(1)}}{\alpha_1(t)} \\ p^{(2)} = p_0^{(2)} + \frac{x^{(2)} - p_0^{(2)}}{\alpha_2(t)} \\ p^{(3)} = p_0^{(3)} + \frac{x^{(3)} - p_0^{(3)}}{\alpha_3(t)} \end{cases}$$

The velocity field:

$$\dot{x}(p,t) = \begin{pmatrix} \dot{\alpha}_1(t)[p^{(1)} - p_0^{(1)}] \\ \dot{\alpha}_2(t)[p^{(2)} - p_0^{(2)}] \\ \dot{\alpha}_3(t)[p^{(3)} - p_0^{(3)}] \end{pmatrix}$$

Apply reference mapping p to obtain the spatial velocity field:

$$v(x,t) = \begin{pmatrix} \dot{\alpha}_1(t) \frac{x^{(1)} - p_0^{(1)}}{\alpha_1(t)} \\ \dot{\alpha}_2(t) \frac{x^{(2)} - p_0^{(2)}}{\alpha_2(t)} \\ \dot{\alpha}_3(t) \frac{x^{(3)} - p_0^{(3)}}{\alpha_3(t)} \end{pmatrix}$$

The velocity gradient L:

$$L = \operatorname{grad} v = \begin{pmatrix} \frac{\dot{\alpha}_1(t)}{\alpha_1(t)} & 0 & 0\\ 0 & \frac{\dot{\alpha}_2(t)}{\alpha_2(t)} & 0\\ 0 & 0 & \frac{\dot{\alpha}_3(t)}{\alpha_3(t)} \end{pmatrix}$$

Streamlines equations are given by

$$\dot{s}(\lambda) = \begin{pmatrix} \dot{\alpha}_1(\tau) \frac{s^{(1)}(\lambda) - p_0^{(1)}}{\alpha_1(t)} \\ \dot{\alpha}_2(\tau) \frac{s^{(2)}(\lambda) - p_0^{(2)}}{\alpha_2(t)} \\ \dot{\alpha}_3(\tau) \frac{s^{(3)}(\lambda) - p_0^{(3)}}{\alpha_3(t)} \end{pmatrix}$$

By change of variables $w^{(i)} = s^{(i)} - p_0^{(i)}$, it becomes ODEs of $w^{(i)}$'s. It follows that

$$w^{(i)}(\lambda) = Ce^{\frac{\dot{\alpha}^{(i)}(\tau)}{\alpha^{(i)}(\tau)}\lambda}$$

Integration constant C is determined by the initial conditions. In conclusion, the streamlines are:

$$s^{(i)}(\lambda) = p_0^{(i)} + (y^{(i)} - p_0^{(i)}) e^{\frac{\dot{\alpha}^{(i)}(\tau)}{\alpha^{(i)}(\tau)}\lambda}$$

评分细则:

Exercise 7. Consider a surface $\mathscr{P} = \{p \in \mathscr{D} \subset \mathscr{R} | \varphi(p) = 0\}$, where \mathscr{D} is an open subset of \mathscr{R} . φ is a smooth function defined on \mathscr{D} , whose gradient doesn't vanish on \mathscr{P} . Consider the trajectory \mathscr{P}_t associated with motion x:

$$\mathscr{P}_t = \{ x \in \mathscr{D}_t | \psi(x, t) = 0 \},$$

where $\mathcal{D}_t = x(\mathcal{D}, t), \ \psi(x, t) = \varphi(p(x, t)).$

Show that

- (a) $\nabla \varphi(\mathbf{p}) \ (p \in \mathscr{P})$ is normal to \mathscr{P} ;
- (b) grad $\psi(x,t)$ $(x \in \mathscr{P}_t)$ is normal to \mathscr{P}_t ;
- (c) $\nabla \varphi = F^T(\operatorname{grad} \psi)_m$, and hence $\operatorname{grad} \psi(x,t)$ never vanishes on \mathscr{P}_t ;
- (d) $|\nabla \varphi|^2 = (\operatorname{grad} \psi)_m \cdot B(\operatorname{grad} \psi)_m, B = FF^T;$
- (e) $\psi' = -v \cdot \operatorname{grad} \psi$.

Solution.

Firstly, note that 0 is a regular value of φ , the surface is thus a regular surface, for which tangent plane and the normal is well defined.

By observation, (b) follows immediately from (a). Indeed, it suffices to formally replace counterparts for (b) in (a); (d) follows immediately from (c). Indeed, by associative law,

$$|\nabla \varphi|^2 = \nabla \varphi \cdot \nabla \varphi = (F^T(\operatorname{grad} \psi)_m)^T (F^T(\operatorname{grad} \psi)_m) = (\operatorname{grad} \psi)_m \cdot FF^T(\operatorname{grad} \psi)_m.$$

Besides, (e) follows immediately from eqn. (4) in section 8, since $\dot{\psi} = (\dot{\varphi}_{\sigma}) = (\dot{\varphi})_{\sigma} = 0$.

Therefore, it suffices to validate (a) and (b).

For (a): Given any but fixed $p \in \mathscr{P}$, we arbitrarily take a regular curve restricted on surface $\mathscr{P} \alpha : (-\epsilon, \epsilon) \mapsto \mathbb{R}^3$ such that $\alpha(0) = p$.

$$\varphi \circ \alpha(t) = \varphi(\alpha_1(t), \alpha_2(t), \alpha_3(t)) = 0$$

Taking derivative w.r.t. t at t=0 leads to

$$\varphi_{p_1} \cdot \dot{\alpha}_1 + \varphi_{p_2} \cdot \dot{\alpha}_2 + \varphi_{p_3} \cdot \dot{\alpha}_3 = \nabla \varphi(p) \cdot \dot{\alpha}(t) = 0$$

Thus, $\nabla \varphi(p)$ is normal to any tangents at p.

For (c):

$$\psi_m\left(x(p,t),t\right) = \psi(\cdot,t) \circ x(p,t)$$

$$\nabla \psi_m = \nabla_p \psi_m = \begin{pmatrix}
\frac{\partial x_1}{\partial p_1} & \frac{\partial x_2}{\partial p_1} & \frac{\partial x_3}{\partial p_1} \\
\frac{\partial x_1}{\partial p_2} & \frac{\partial x_2}{\partial p_2} & \frac{\partial x_3}{\partial p_2} \\
\frac{\partial x_1}{\partial p_3} & \frac{\partial x_2}{\partial p_3} & \frac{\partial x_3}{\partial p_3}
\end{pmatrix}
\begin{pmatrix}
\frac{\partial \psi}{\partial x_1}(x(p,t),t) \\
\frac{\partial \psi}{\partial x_2}(x(p,t),t) \\
\frac{\partial \psi}{\partial x_3}(x(p,t),t)
\end{pmatrix} = F^T(\operatorname{grad}\psi)_m$$

One the other hand, recall that $\varphi(p(x,t)) = \psi(x,t)$, hence

$$\psi_m(p,t) = [\varphi(p(x,t))]_m$$
$$= \varphi(p(x(p,t),t))$$
$$= \varphi(p)$$

and $\nabla \psi_m = \nabla \varphi$. We've done.