



**Advanced Statistical Physics  
Wintersemester 2022/23**

Problem sheet 2, handed out: Friday 4<sup>th</sup> November, 2022  
due: **Sunday 13<sup>th</sup> November, 2022 (via email/ studip)**

## 1 Phase transition and large deviations in the infinite range Ising model (20 pts. + 6 bonus pts.)

The Hamiltonian for the infinite range Ising model (in zero field) is

$$H = -\frac{1}{2N} \sum_{i,j=1}^N J \sigma_i \sigma_j$$

with Ising spins  $\sigma_i = \pm 1$ . The aim of this question is to analyse the (equilibrium) fluctuations of the magnetization  $M = \sum_{i=1}^N \sigma_i$  in this system, as a function of temperature  $T = 1/\beta$ .

- a) Show that the probability of observing magnetization  $M$  is

$$P(M) = \frac{1}{Z} \frac{N!}{((N+M)/2)!((N-M)/2)!} e^{\beta J M^2 / (2N)}$$

(Hint: Count how many spin configurations have a given  $M$ .) (3 pts.)

- b) Use the Stirling approximation to show from (a) that the rate function for the magnetization is, with  $m = M/N$ ,

$$I(m) = -\beta J m^2 / 2 + \frac{1+m}{2} \ln \frac{1+m}{2} + \frac{1-m}{2} \ln \frac{1-m}{2} + \text{const}$$

and sketch  $I(m)$  for  $\beta J < 1$  and  $\beta J > 1$ . Why does the form in the latter case indicate a phase transition? (6 pts.)

- c) Derive the “Hubbard-Stratonovich” transformation

$$e^{\beta J M^2 / (2N)} = \int \frac{dx}{\sqrt{2\pi\beta J/N}} e^{-Nx^2/(2\beta J) + Mx}$$

which is used to factorize the sum defining the partition function  $Z = \text{Tr } e^{-\beta(H-hM)}$  in a field  $h$  into independent sums over individual spins. Carry out these sums to express  $Z$  as an integral over  $x$ , and deduce that the free energy per spin becomes for  $N \rightarrow \infty$

$$f(h) = \min_x \left[ \frac{x^2}{2\beta^2 J} - \frac{1}{\beta} \ln(2 \cosh(\beta h + x)) \right]$$

Find the condition for the position  $x^*(h)$  of the minimum. Sketch function to be minimized for  $\beta J < 1$  and  $> 1$  and a few  $h$ ? (5 pts.)

- d) Find the average magnetization  $m = -f'(h)$  and show that it obeys  $m = \tanh(\beta(Jm + h))$ . (Hint:  $f(h)$  depends on  $h$  both explicitly and via  $x^*(h)$ ; show using the chain rule that only the explicit dependence is needed to calculate  $f'(h)$ .) (3 pts.)
- e) Discuss qualitatively, and sketch, how the rate function  $I^*(m)$  that is obtained by applying the Gärtner-Ellis-Theorem to  $f(h)$  differs from the result of (b). (3 pts.)
- f) (bonus) Verify explicitly that  $I^*(m)$  agrees with  $I(m)$  in the appropriate range of  $m$ . (2 pts.)
- g) (bonus) Explain how the rate function  $I(m)$  for  $\beta = 0$  can be used to write down the rate function for the fluctuations of the end-to-end distance of a one-dimensional polymer, consisting of  $N$  chain segments of fixed length  $l$  that can randomly point left or right. (4 pts.)

## 2 Properties of equilibrium correlation functions (10 pts.)

Correlation functions are important for describing time-dependent fluctuations of a system and can be used e.g. to deduce transport properties. In this exercise, we want to prove some general properties of correlation functions, defined as in the lectures as  $C_{AB}(t, t') = \text{Tr } A(t)B(t')\rho(0)$  where  $A(t)$ ,  $B(t')$  are time-dependent phase space observables (classically) or Hermitian operators.

- a) In the classical case one has  $C_{BA}(t', t) = C_{AB}(t, t')$ . Prove that in the quantum case,  $C_{BA}^*(t', t) = C_{AB}(t, t')$ . (Hint: Use that  $(\text{Tr } M)^* = \text{Tr}(M^\dagger)$  where  $\dots^\dagger$  indicates Hermitian conjugation.) (2 pts.)
- b) From now on, consider equilibrium correlators, which can be expressed as  $C_{AB}(t, t') = (A(t), B(t'))$  in terms of the product  $(A, B) = \text{Tr } AB\rho^{\text{eq}}$ .  
Prove in the classical case that  $(A, \mathcal{L}B) = -(\mathcal{L}A, B)$ . Deduce that  $(A, e^{\mathcal{L}t'} B) = (e^{-\mathcal{L}t'} A, B)$  and therefore that equilibrium correlation functions are time-translation invariant. (5 pts.)
- c) Show that in the classical case,  $C_{AA}(\Delta t) \leq C_{AA}(0)$ , i.e. the autocorrelation function is bounded from above by its equal-time value. (Hint: Start from the average of  $[A(\Delta t) - A(0)]^2$ .) (3 pts.)

### 3 Liouville operator for the harmonic oscillator (20 pts. + 5 bonus pts.)

Consider a simple harmonic oscillator in one-dimension. The Hamiltonian of the system is given by

$$H(x, p) = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2 \quad (1)$$

- a) Write down the Liouville operator  $\mathcal{L}$  as defined in the lectures. (2 pts.)
- b) Consider the coordinate transformation

$$x = \sqrt{\frac{2J}{m\omega}} \cos \phi, \quad \text{and} \quad p = \sqrt{2Jm\omega} \sin \phi$$

Show that in these coordinates, the Liouvillian becomes simply

$$\mathcal{L} = -\omega \frac{\partial}{\partial \phi}$$

Hint: Use the chain rule to transform from  $(\partial/\partial x, \partial/\partial p)$  to  $(\partial/\partial \phi, \partial/\partial J)$ ; it is convenient to express the Jacobian via the inverse function relation

$$\begin{pmatrix} \frac{\partial \phi}{\partial x} & \frac{\partial J}{\partial x} \\ \frac{\partial \phi}{\partial p} & \frac{\partial J}{\partial p} \end{pmatrix} = \begin{pmatrix} \frac{\partial x}{\partial \phi} & \frac{\partial p}{\partial \phi} \\ \frac{\partial x}{\partial J} & \frac{\partial p}{\partial J} \end{pmatrix}^{-1}$$

(10 pts.)

- c) Consider the eigenvalue problem  $\mathcal{L}f = \lambda f$ . Show that the eigenvalues are of the form  $\lambda = -in\omega$  with  $n$  an integer, and find the corresponding eigenfunctions  $f$ . (8 pts.)
- d) (bonus) Show that the eigenvalues of  $\mathcal{L}$  are always imaginary in deterministic dynamics, given the anti-self-adjointness of  $\mathcal{L}$ . (Hint: Anti-self-adjointness for a matrix  $\mathbf{L}$  and the conventional scalar product would be the property that  $\mathbf{a}^T \mathbf{L} \mathbf{b} = -(\mathbf{L} \mathbf{a})^T \mathbf{b}$  for arbitrary vectors  $\mathbf{a}, \mathbf{b}$ . Deduce that  $\mathbf{L}$  is antisymmetric (also called skew-symmetric). What does this imply for the eigenvalues?) (5 pts.)

### 4 Rate function in nonzero field (bonus) (15 pts.)

Consider an extensive observable  $A$  with conjugate field  $h$ . We discussed in the lectures the Gärtner-Ellis theorem that gives the rate function for the distribution of  $A$  at zero field.

- a) Generalize the derivation from the lectures to show that the rate function for the distribution of  $A$  in a nonzero field  $h_0$  is given by

$$I_{h_0}(a) = \sup_h \beta [ha + f(h_0 + h) - f(h_0)]$$

(5 pts.)

- b) By shifting  $h$  appropriately, show that this result can be rewritten as

$$I_{h_0}(a) = I_0(a) - \beta h_0 a + \beta[f(0) - f(h_0)]$$

where  $I_0(a)$  is the zero field rate function derived in the lectures. Interpret this result graphically. (5 pts.)

- c) Show that for arbitrary observables,

$$P_{h_0}(A) \propto P_0(A)e^{\beta h_0 A}$$

Find the proportionality constant. Explain how the relation between  $P_{h_0}(A)$  and  $P_0(A)$  can be used to give an alternative derivation of the result of (b). (5 pts.)