Report 4: Partial Differnetial Equations (PDEs)

Simon Blaue

December 15, 2022

Universität Göttingen
Faculty of Physics
Instructor: Prof. Dr. S. Schumann
Tutors: Dr. E. Bothmann, M. Knobbe

1   Laplace Equation

1.1   Iterator methods

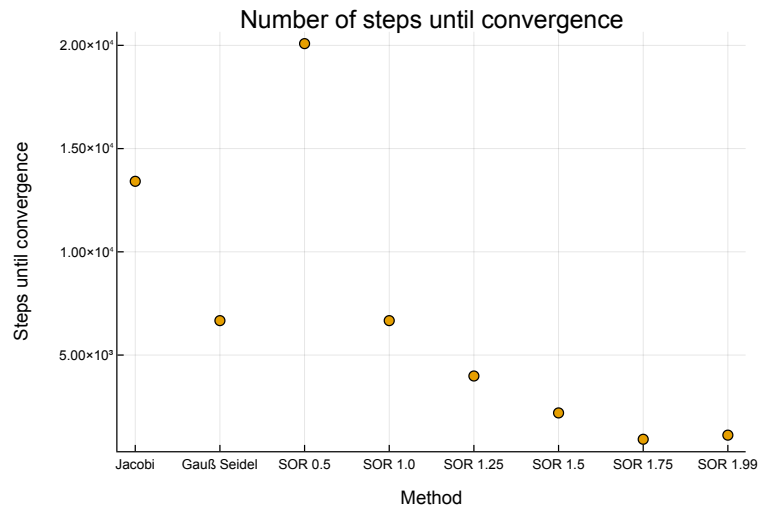First lets check which method converges the fastest:

Figure 1: Number of steps until convergence concerning the Laplace error $\max \epsilon < 1 \times 10^{-3}$.

I observe that Gausß-Seidel and SOR wit $\alpha = 1.0$ need the same timestep as expected. The fastest method is SOR with $\alpha = 1.75$, because with $\alpha = 1.99$ we get to close to unstable regimes. Obivoulsly SOR with $\alpha = 0.5$ takes the longest, as the updating step is damped with the factor α. Now I will observe the evolution of the maximal error and the average error of the different methods.
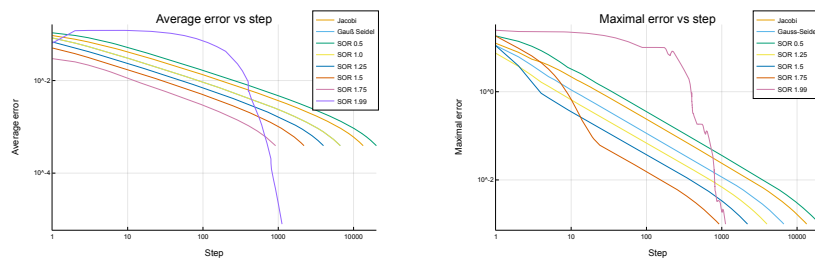

Figure 2: Maximal and average error for different iteration methods.

As discussed before, SOR with $\alpha = 1.0$ and Gauß-Seidel method are the same, hence I plotted only the latter. The error development for all methods besides

SOR with$\alpha \geq 1.5$ seem qualitatively the same. However for higher α I observe indents in the curve, which seem to boost the algorithm. This is due to the algorithm taking bigger steps in the right direction once it found this direction. For $\alpha = 1.99$ the curve is very rigid, and I was not sure, if it realy converged to the right solution, hence I plotted the result below. As it turns out the method also converges to the expected result.
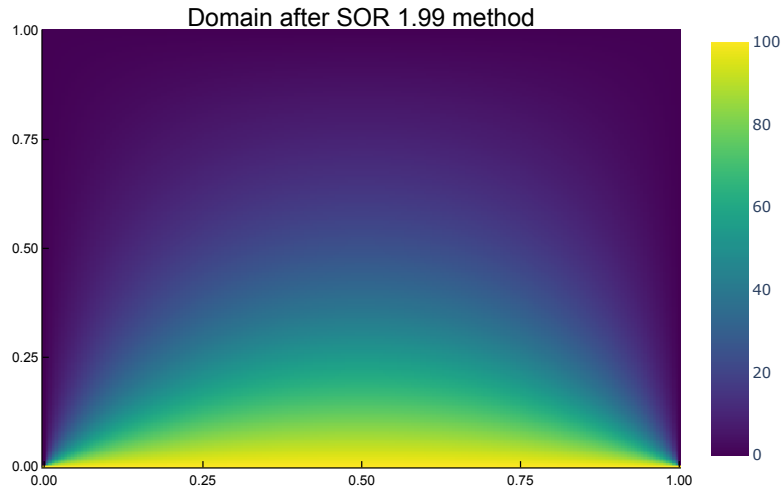

Figure 3: Domain after iterating with the SOR 1.99 method to varify right convergence.

1.2   SOR for $\alpha = 2.0$

The natural question to ask is what happens for an even further boosted SOR method? I found that for $\alpha = 2.0$ the algorithm does not converge in 50000 steps. It seems that it would not have converged to the right domain, as the result shows stripes and the maximal error starts to fluctuate a lot after 100 timesteps. The system can not recover from that.
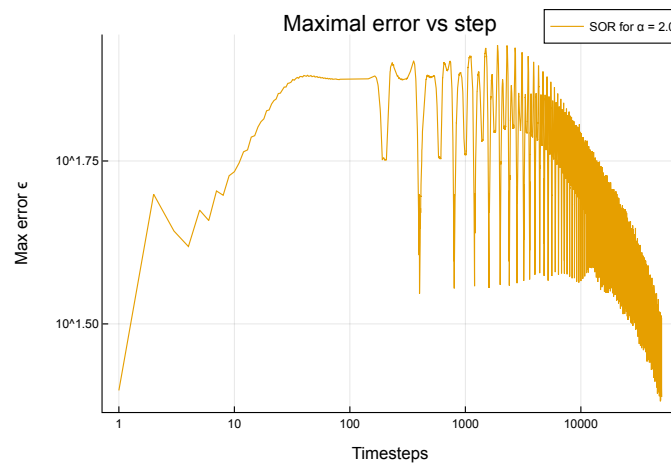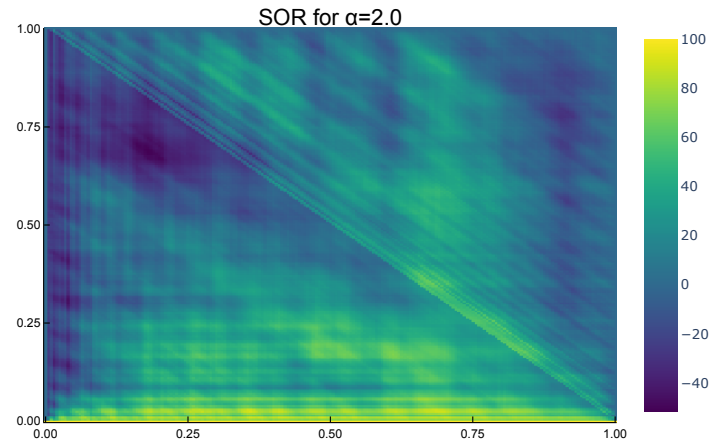
1.3   Infinite sum solution

Now I will cross verify the iterative solution with an aproximated analytical solution. First have a look at the analytical "infinite" sum solution for different numbers of summands.
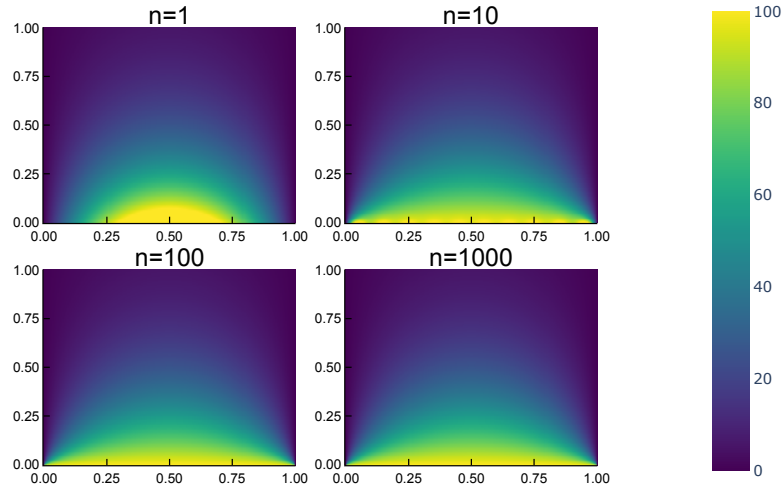
Figure 4: Results for the "infinite" sum solution to the Laplace equation for different numbers of terms n.

As I increse the number of terms more and more sine functions are added left and right to the first one, and towards the top an exponential decay fills in. Comparing the "infinite" sum solution for 1000 terms with the Gauß-Seidel method reveals that the exponential decay is to weak, or the Gauß-Seidel method did not propergate long enough. Substracting both results displays that at the edges the solutions are identical (yellow) in the middle-top they dont fit together as discussed before.
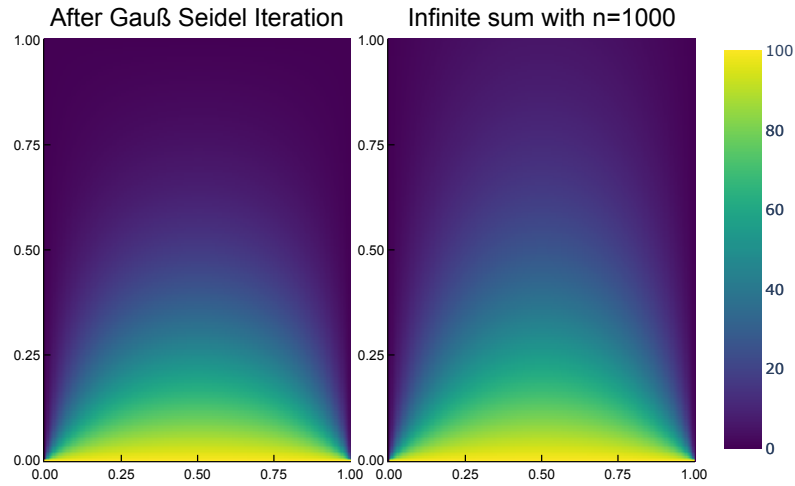
(a) Comparison of the infinite sum solution with 1000 terms and the Gauß Seidel iterator solution.
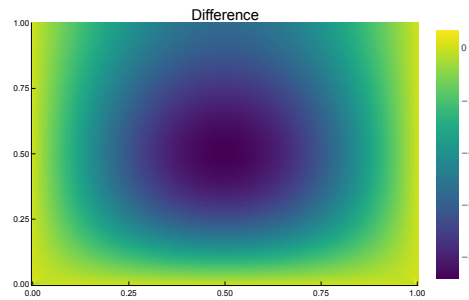
(b) Difference between the infinite sum solution with 1000 terms and the Gauß Seidel iterator solution.

2   Diffusion

In this task I implemented four different method for solving the diffusion equation. To varyfiy they all converged to the same solution I plotted a time evolution of all of them. The time domain is in the y-axis, the rod domain in the x-axis, displayed in Figure 6. As expected, th rod cools down, but holds the qualitative temperature profile.
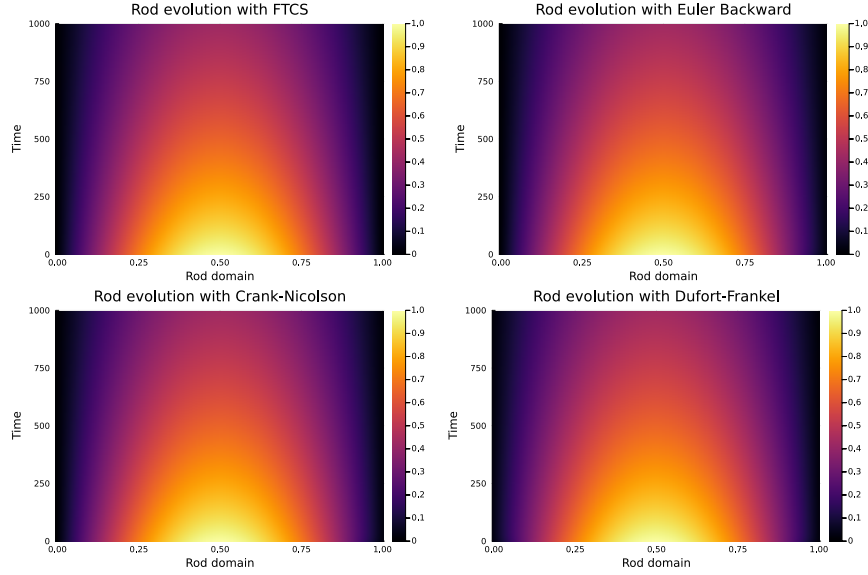
2.1   Integration methods

Figure 6: Temporal evolution of the rod's temperature along the y-axis. The evolution seems to be the same for all four integration methods.

2.2   Error comparison

To investigate the stabiliy of the different method, I plotted the error against the analytical soluution for different timestep widths. I expect unconditionall stability for the implicit methods Euler Backwards and Crank-Nicelson and the Dufort-Frankel method. The FCTS method is not unconditionally stable and should result in a huge error at some large enough Δt, depending on Δx ($\Delta x = 0.01$). This can be observed in Figure 7.
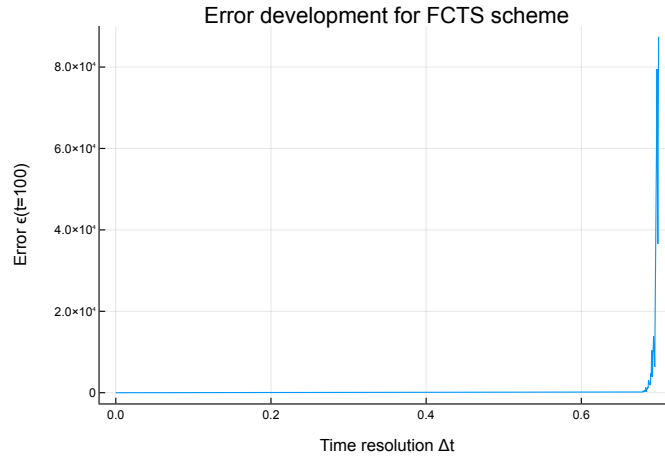
Figure 7: Error development for the FCTS method

For comparing the other erros I will cut off the instbaility of the FCTS scheme at an error of $\epsilon = 0.007$. I observe that the other schemes seem stable in the tested regime, however all of them have a relatively high error. An intersting result is that the error for the implicit methods is the same, while the Dufort-Frankel error is approximatly twice as large.
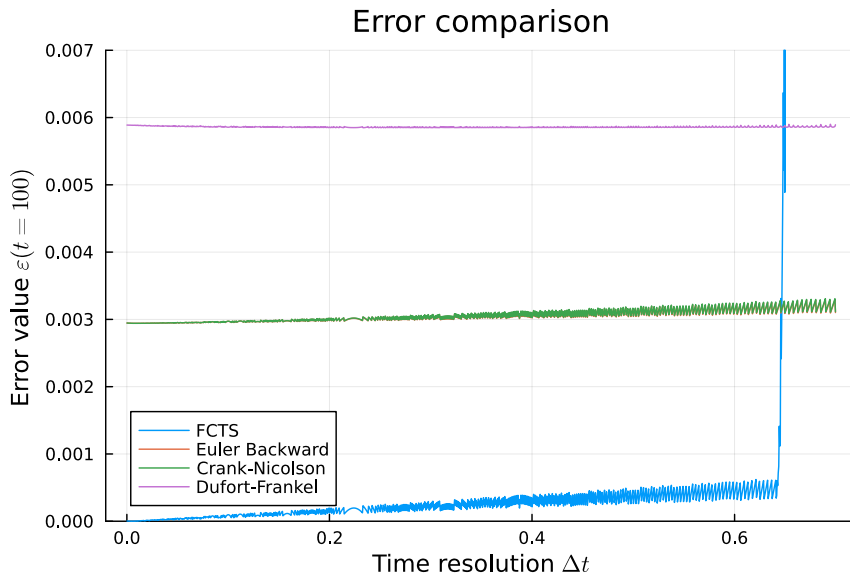

Figure 8: Error comparison for the different methods

3 Solitons