

# Inverse Design and Automatic Differentiation for Optical Devices



**Ian Williamson and Momchil Minkov**

Stanford University Optical Society Workshop

Nov 21, 2019

<https://github.com/fancompute/workshop-invdesign>

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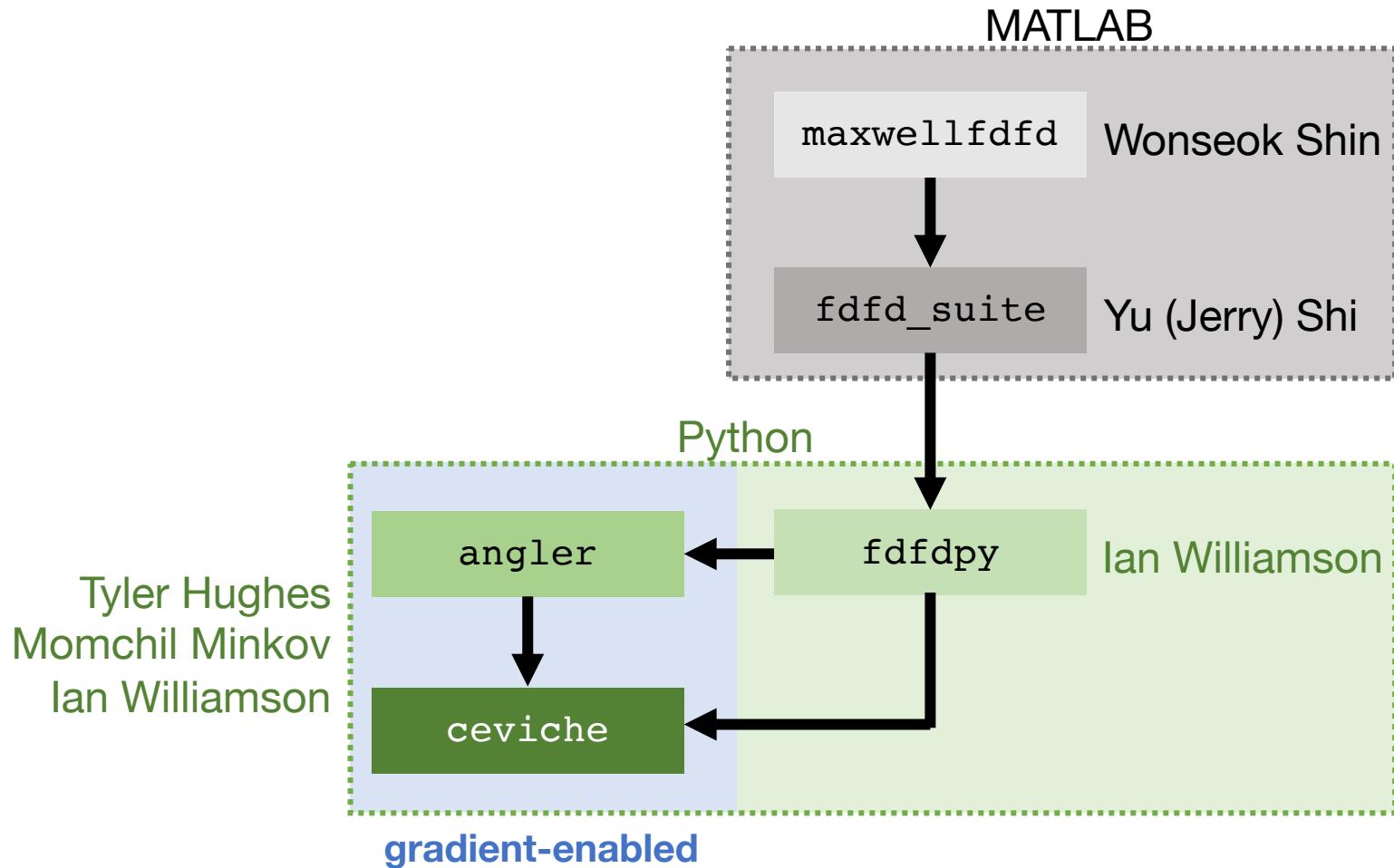
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# Workshop outline

- Brief crash course on optical simulation
  - Theory of numerically solving Maxwell's equations
  - Practical demonstration with our open source code
- Inverse design theory / practical introduction
  - Adjoint gradient computation
  - Automatic differentiation
  - Connection to machine learning
- Device parameterization techniques
  - Binarizing features
  - Optimization penalties

# A history of optical simulation packages *(from our group)*



# A quick crash course in optical simulation

# How do we simulate optical devices?

... we solve Maxwell's equations!

The diagram shows two equations for Maxwell's equations within a rectangular region representing a device. The top equation is  $\nabla \times E = -j\omega\mu_0 H$ , with a blue arrow pointing down from the text "electric field". The bottom equation is  $\nabla \times H = j\omega\epsilon_0\epsilon_r E + J$ , with a blue arrow pointing up from the text "current source". A vertical blue arrow points upwards from the bottom equation through the center of the rectangle to the text "relative permittivity". Below the rectangle, the text "this represents our device!" is displayed.

$$\nabla \times E = -j\omega\mu_0 H$$
$$\nabla \times H = j\omega\epsilon_0\epsilon_r E + J$$

electric field

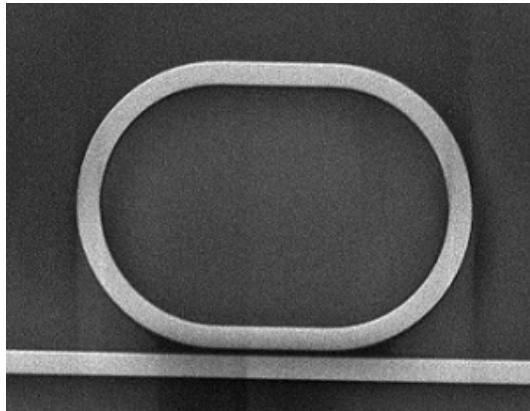
magnetic field

current source

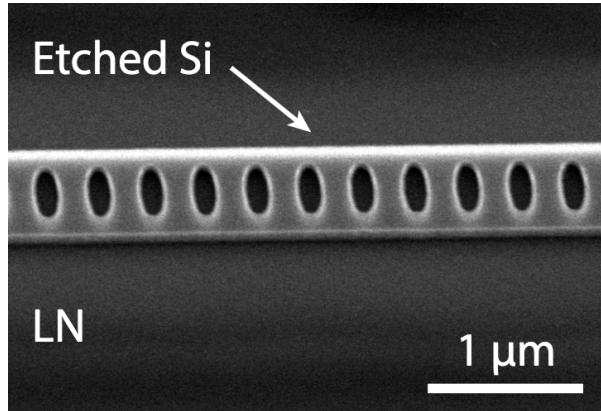
relative permittivity

**this represents our device!**

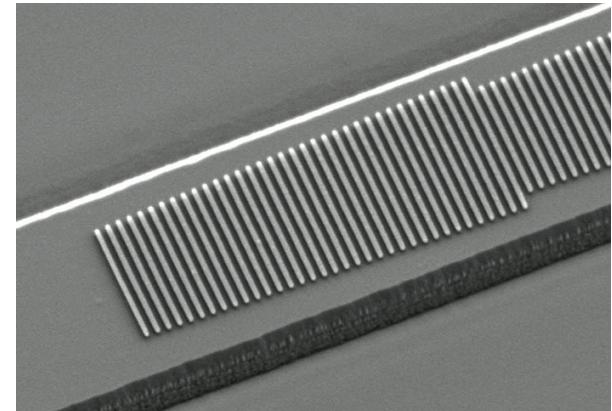
# Welcome to the optical device zoo



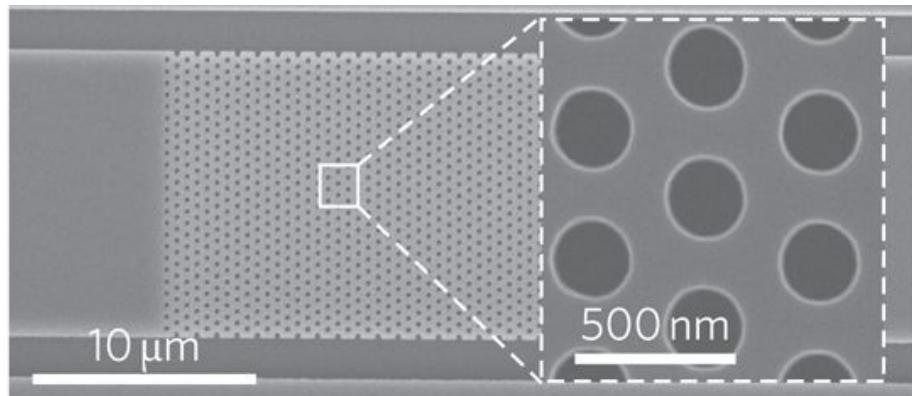
Stern Group (Northwestern Univ)



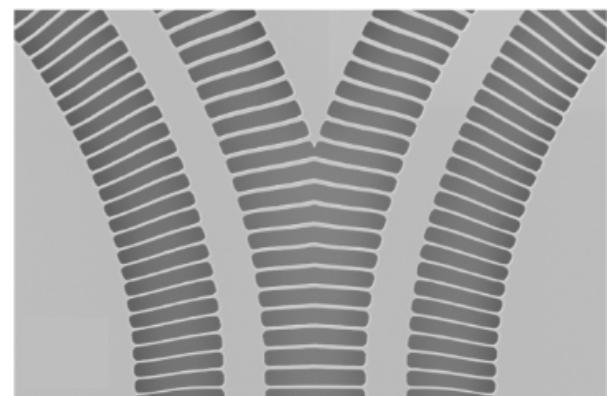
Witmer et al. Opt. Express 24, 5876-5885 (2016)



Wang et al. Nat. Comm. 8, 2098 (2017)



Kocaman et al. Nat. Pho. 5, 499–505 (2011)



Soler Penades et al. Opt. Express 24, 22908-22916 (2016)

# Constructing a set of equations to solve

$$\nabla \times E = -j\omega\mu_0 H$$
$$\nabla \times H = j\omega\epsilon_0\epsilon_r E + J$$

## Some assumptions we make:

- A two-dimensional domain (invariant along z-direction)
- All fields are also invariant along z-direction
- Electric-field polarized only along z-direction

$$\partial_y E_z = -j\omega\mu_0 H_x$$

$$\partial_x E_z = j\omega\mu_0 H_y$$

$$\partial_x H_y - \partial_y H_x = j\omega\epsilon E_z + J_z$$

# Constructing a set of equations to solve

$$\partial_y E_z = -j\omega\mu_0 H_x$$

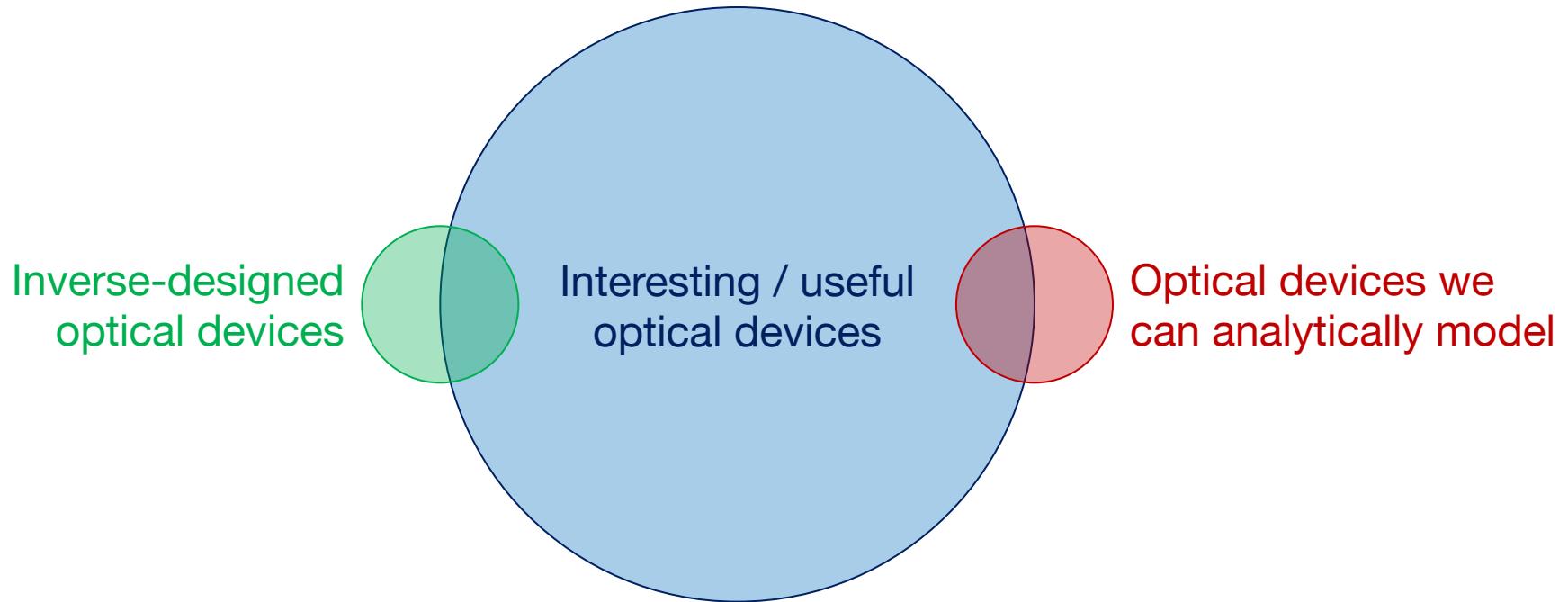
$$\partial_x E_z = j\omega\mu_0 H_y$$

$$\partial_x H_y - \partial_y H_x = j\omega\epsilon E_z + J_z$$



$$\partial_x \partial_x E_z + \partial_y \partial_y E_z + \omega^2 \mu_0 \epsilon E_z = j\omega\mu_0 J_z$$

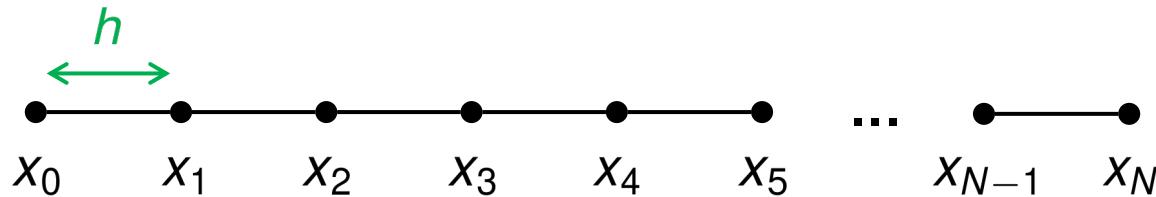
# Where can these equations take us?



$$\partial_x \partial_x E_z + \partial_y \partial_y E_z + \omega^2 \mu_0 \epsilon E_z = j\omega \mu_0 J_z$$

# The key idea behind finite differences

$$\frac{\partial y}{\partial x} = F(x) \rightarrow \frac{\partial y}{\partial x} \approx \frac{y(x + h) - y(x)}{h}$$

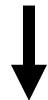


- Very simple concept that can be scaled up to solve complex PDEs
- Convergence to the exact solution at a rate proportional to  $h$  (depends on “stencil”)

# Discretizing Maxwell's equations

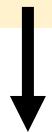
$$\partial_x \mu_0^{-1} \partial_x E_z + \partial_y \mu_0^{-1} \partial_y E_z + \omega^2 \epsilon E_z = j\omega J_z$$

continuous PDE



$$\left[ D_x^b \mu_0^{-1} D_x^f + D_y^b \mu_0^{-1} D_y^f + \omega^2 \text{diag}(\epsilon) \right] e_z = (j\omega) j_z$$

discretized PDE

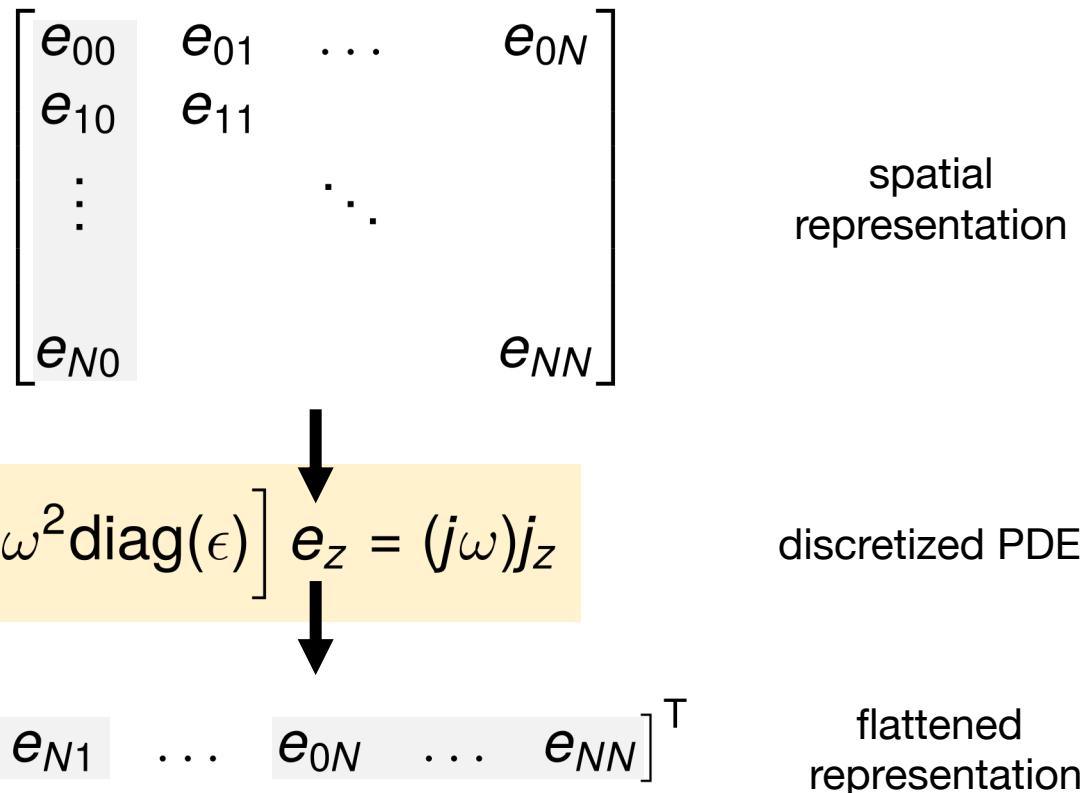


Along just one dimension,  
this looks like...

$$D_x^f = \frac{1}{\Delta x} \begin{bmatrix} -1 & 1 & & & \\ & -1 & 1 & & \\ & & \ddots & \ddots & \\ & & & -1 & 1 \\ 1 & & & & -1 \end{bmatrix}$$

Along two dimensions, the  
representation is a bit more  
complex, *but the same idea  
applies!*

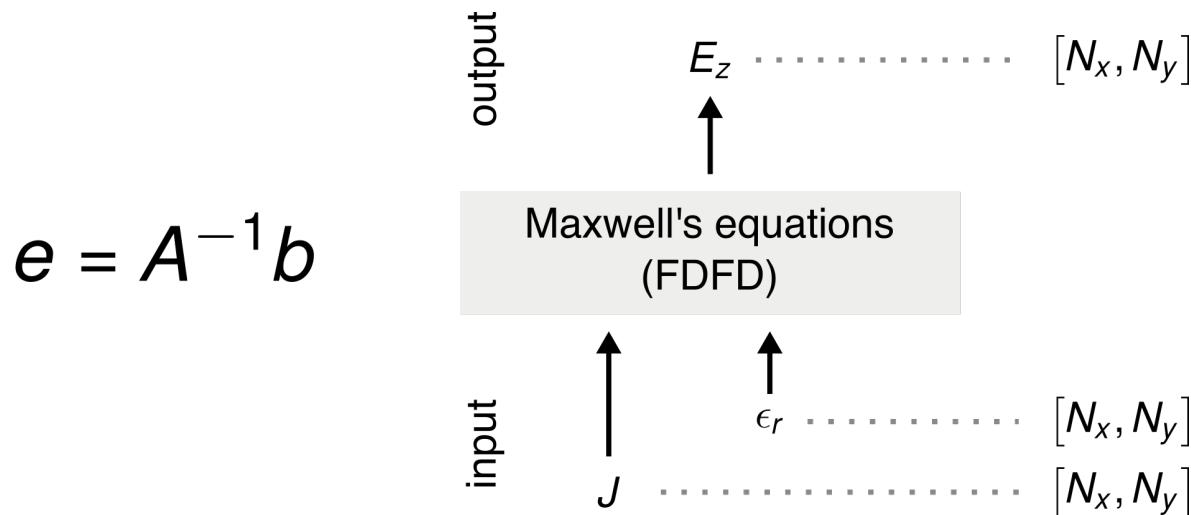
# Discretizing Maxwell's equations



sparse linear problem:  $A \cdot e = b \rightarrow A = f(\epsilon) \rightarrow e = A^{-1}b$

# The `ceviche` package (at a high level)

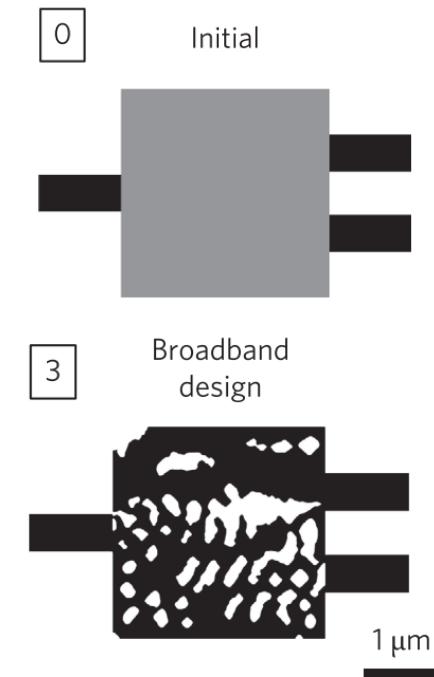
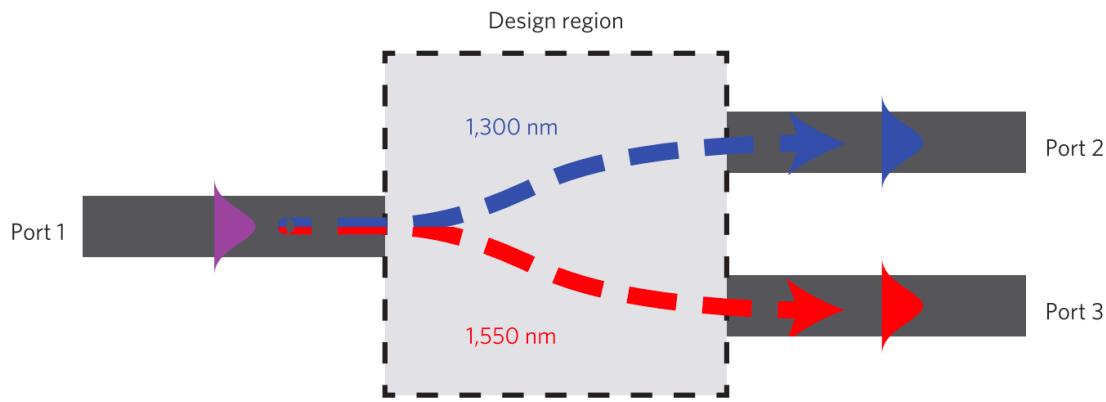
$$\left[ D_x^b \mu_0^{-1} D_x^f + D_y^b \mu_0^{-1} D_y^f + \omega^2 \text{diag}(\epsilon) \right] e_z = (j\omega) j_z$$



# Notebook #01

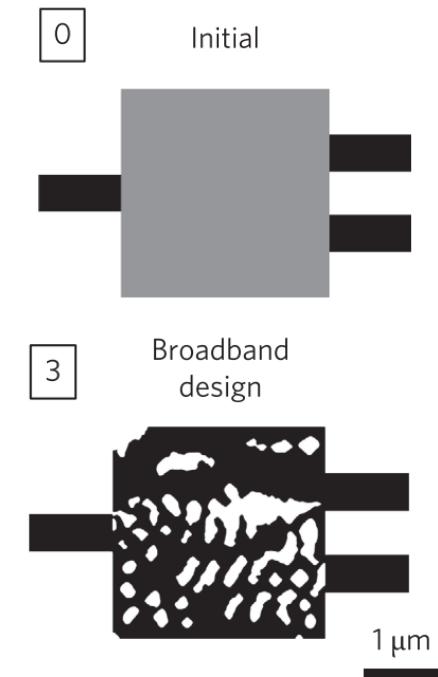
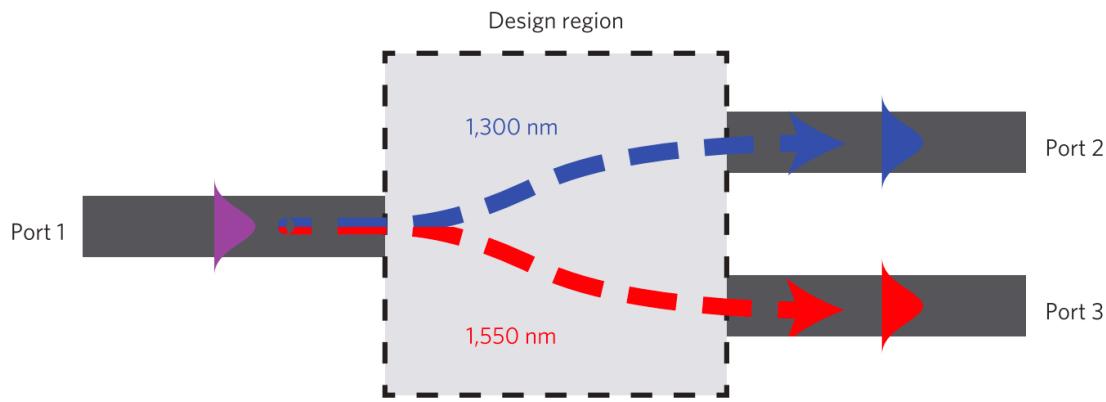
# Inverse design

- Assume that we can “mold” the permittivity
- Try to get the **best** possible structure for a particular desired functionality.



# Inverse design

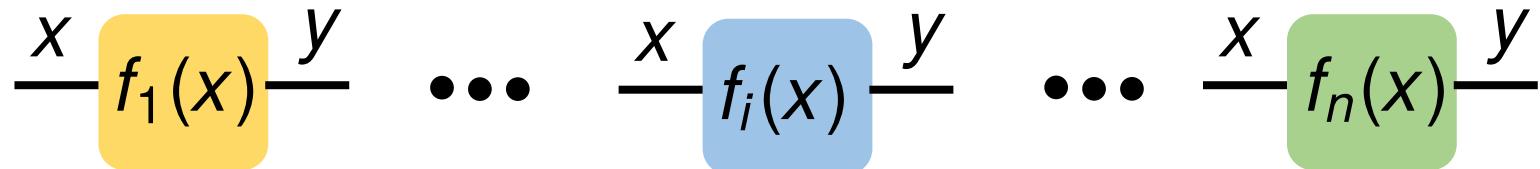
- We want the gradient of some objective function w.r.t. the permittivity at every point
- Efficiently!



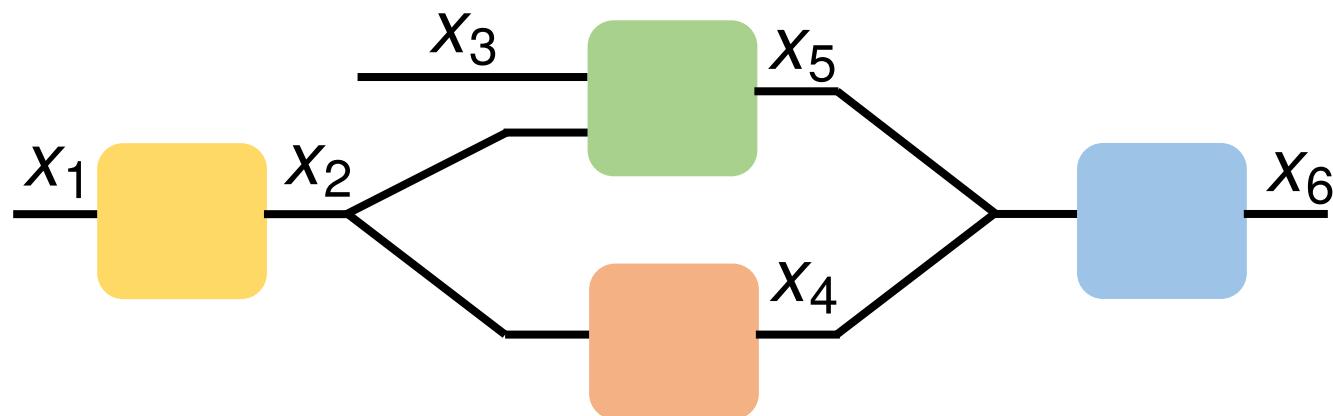
# Automatic differentiation

or: differentiable programming

Usual programming: library of functions

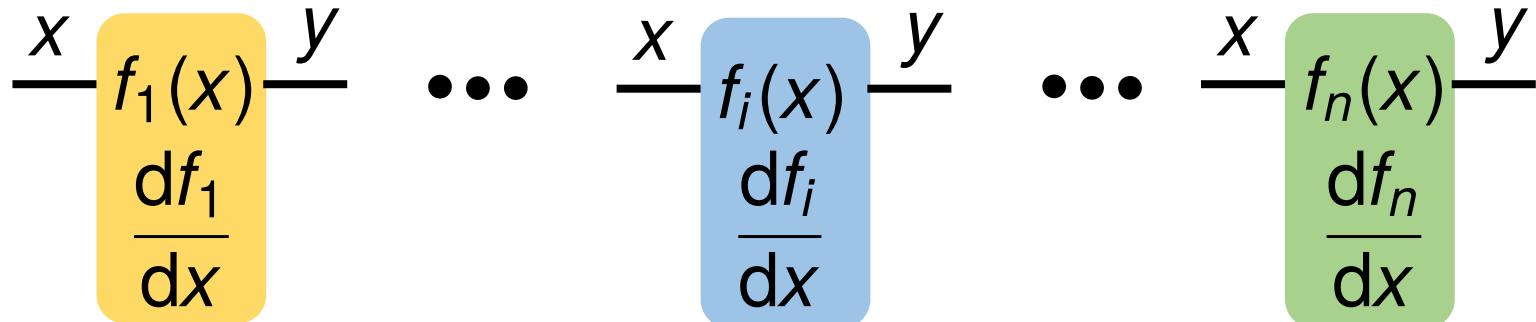


Program = “computational graph”

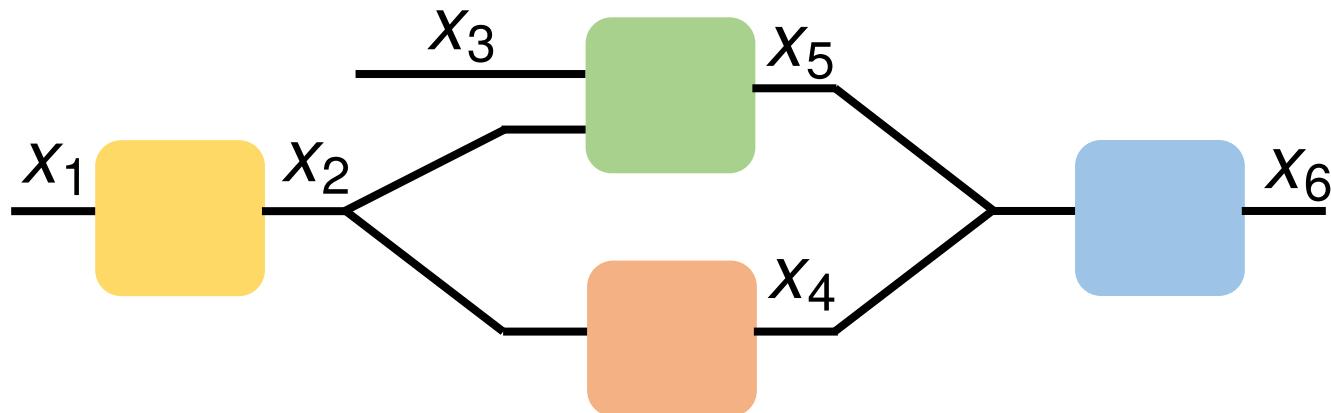


# Automatic differentiation

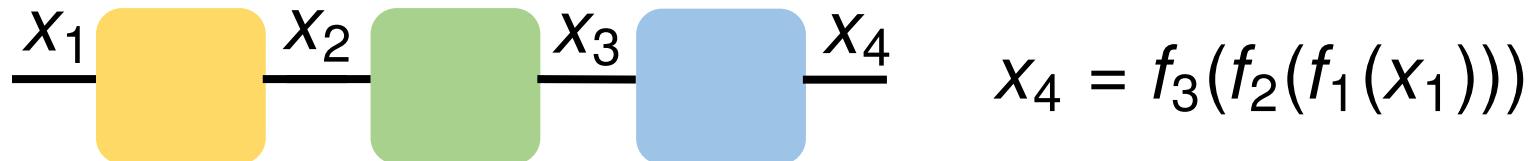
Differentiable programming library:



Track a way to compute any  $dx_i/dx_j$



# Forward/reverse autodiff



$$\frac{dx_4}{dx_1} = \frac{dx_4}{dx_3} \frac{dx_3}{dx_2} \frac{dx_2}{dx_1}$$

The diagram shows the structure of the Jacobian matrix resulting from the forward pass. The top row consists of three diagonal lines connecting the inputs to the outputs of each operation. Below these lines are three labels indicating the dimensions of the resulting matrix columns:  $[N_4 \times N_3]$ ,  $[N_3 \times N_2]$ , and  $[N_2 \times N_1]$ . A vertical line connects the center of the first two columns to the center of the third column.

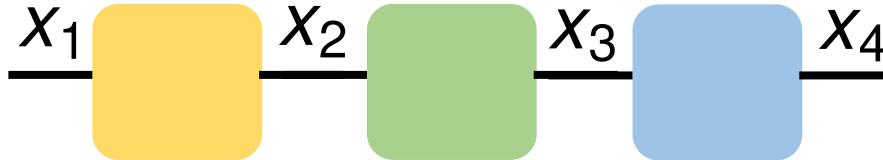
# Forward/reverse autodiff

$$[N_4 \times N_3] \quad [N_3 \times N_2] \quad [N_2 \times N_1]$$

$$\frac{dx_4}{dx_1} = \frac{dx_4}{dx_3} \frac{dx_3}{dx_2} \frac{dx_2}{dx_1}$$

$$\frac{dx_4}{dx_1} = \frac{dx_4}{dx_3} \left( \frac{dx_3}{dx_2} \frac{dx_2}{dx_1} \right) \quad \mathcal{O}(N_1 N_2 N_3 + N_1 N_3 N_4)$$

$$\frac{dx_4}{dx_1} = \left( \frac{dx_4}{dx_3} \frac{dx_3}{dx_2} \right) \frac{dx_2}{dx_1} \quad \mathcal{O}(N_2 N_3 N_4 + N_1 N_2 N_4)$$



$$\frac{dx_4}{dx_1} = \frac{dx_4}{dx_3} \left( \frac{dx_3}{dx_2} \frac{dx_2}{dx_1} \right) \quad \text{— Forward-mode} \quad \mathcal{O}(N_1 N_2 N_3 + N_1 N_3 N_4)$$

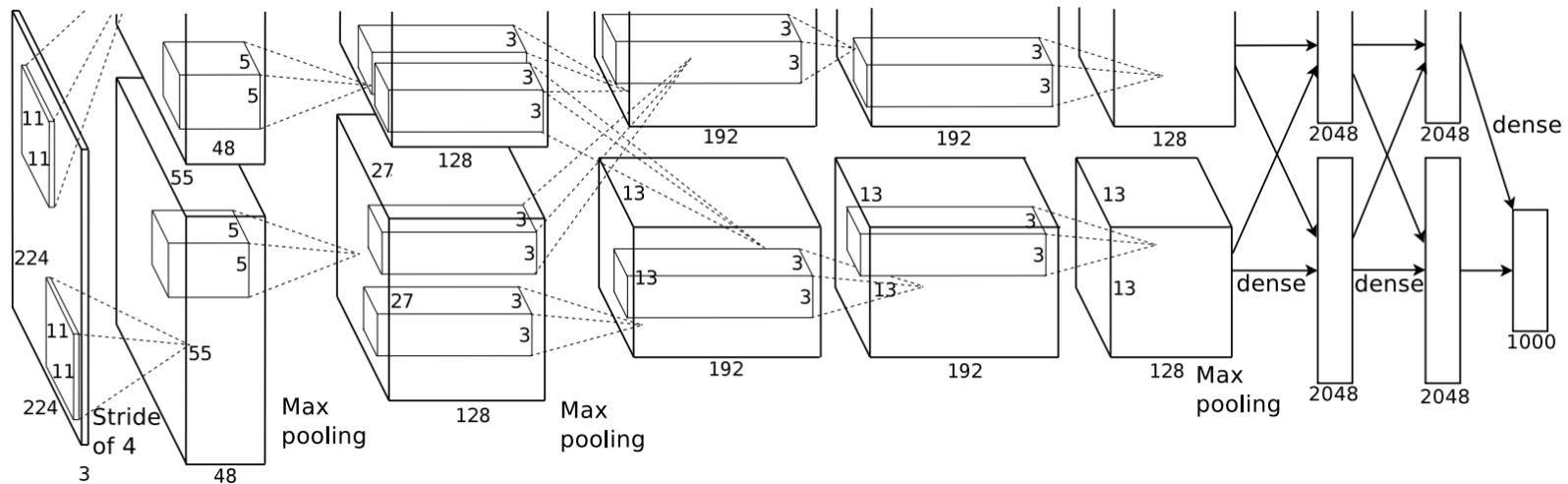
$$\frac{dx_4}{dx_1} = \left( \frac{dx_4}{dx_3} \frac{dx_3}{dx_2} \right) \frac{dx_2}{dx_1} \quad \text{— Reverse-mode} \quad \mathcal{O}(N_2 N_3 N_4 + N_1 N_2 N_4)$$

Reverse-mode autodiff always better when  
 $N_4 \ll N_1$  (few outputs, many inputs)

In optimization problems, this is always the case! Final objective function is just one number.

# Reverse-mode autodiff

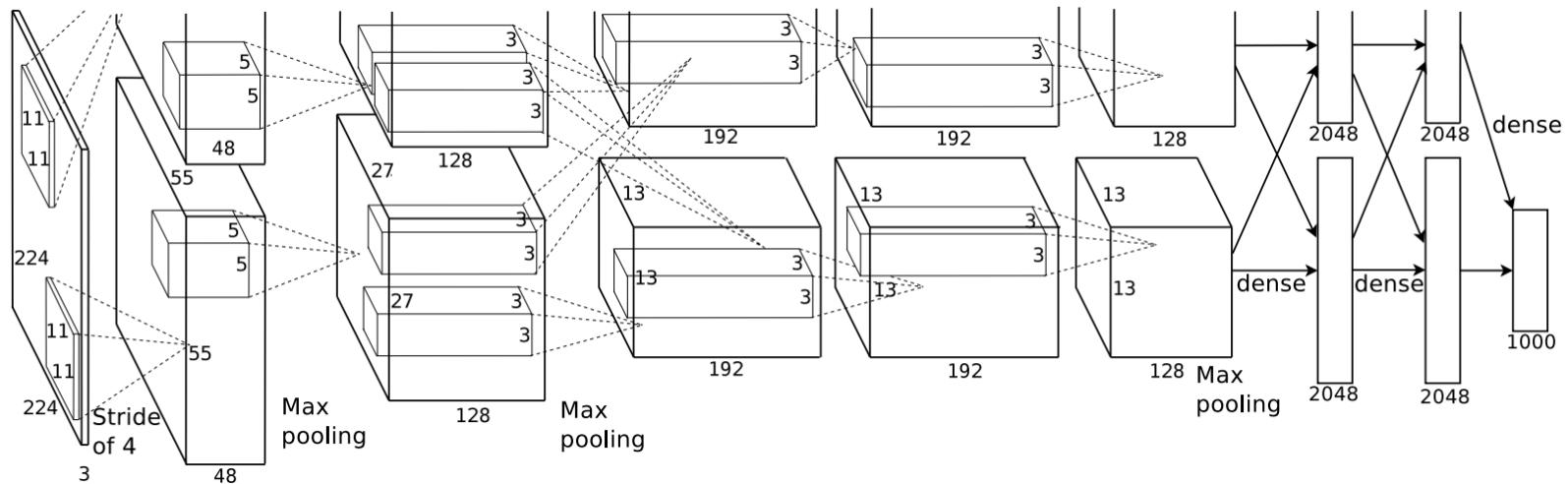
It's backpropagation in machine learning lingo!



**62 million parameters... one final output!**

# Reverse-mode autodiff

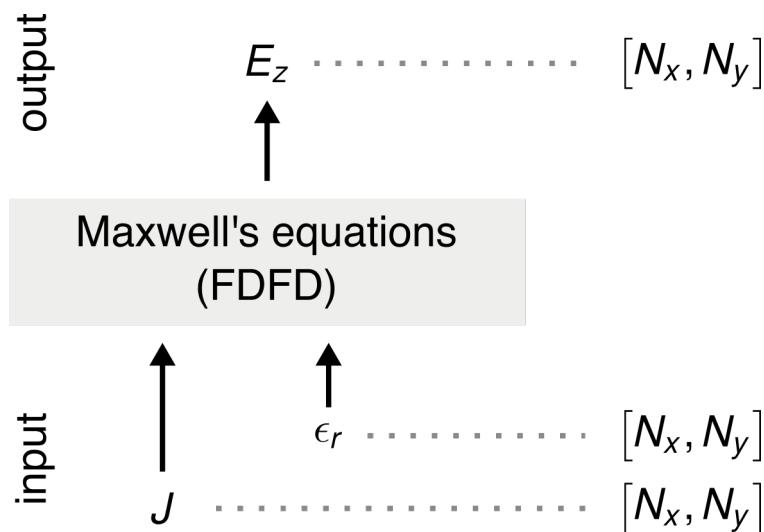
It's backpropagation in machine learning lingo!



TensorFlow, PyTorch, Autograd et al. are all, in their heart, automatic differentiation libraries

# Adjoint variable method

In photonics inverse design, it refers to how we “backprop” through the field computation



$$Ae = b \quad e = A^{-1}b$$

$$\frac{d\mathcal{L}}{d\epsilon_r} = \frac{d\mathcal{L}}{de} \frac{de}{d\epsilon_r}$$

# Adjoint variable method

In photonics inverse design, it refers to how we “backprop” through the field computation

“dude, trust me”

$$\begin{aligned} \frac{d\mathcal{L}}{d\epsilon_r} &= \frac{d\mathcal{L}}{de} \frac{dA^{-1}}{d\epsilon_r} b = - \underbrace{\frac{d\mathcal{L}}{de}}_{e^T} A^{-1} \underbrace{\frac{dA}{d\epsilon_r}}_e A^{-1} b = \\ &= e_{aj}^T \frac{dA}{d\epsilon_r} e \end{aligned}$$

diagonal tensor

$$e_{aj} = (A^T)^{-1} \left[ - \frac{d\mathcal{L}}{de} \right]^T$$

# Adjoint variable method

$$\frac{d\mathcal{L}}{d\epsilon_r} = e_{aj}^T \frac{dA}{d\epsilon_r} e \quad e_{aj} = (A^T)^{-1} \left[ -\frac{d\mathcal{L}}{de} \right]^T$$

For a “reciprocal” system (very often),  $A^T = A$

$$e = A^{-1}b \quad e_{aj} = A^{-1}b_{aj}$$

The adjoint field is a solution to **the same Maxwell** problem, but with a **different source** that depends on the objective function and the original field.

# Autograd

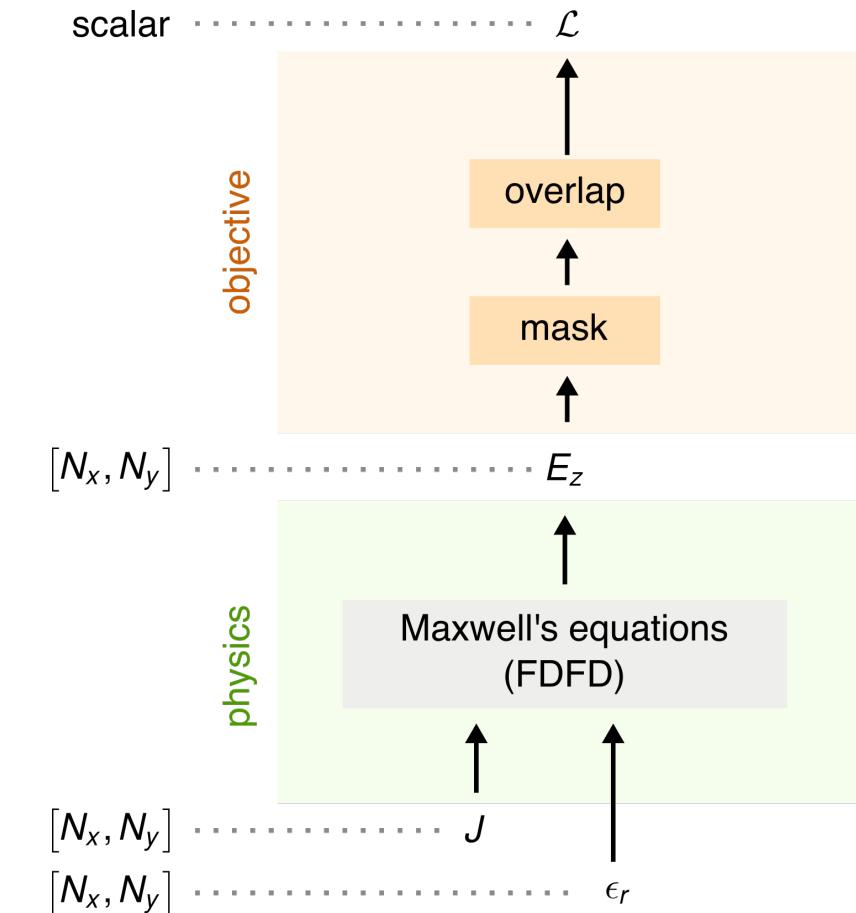
[github.com/HIPS/autograd](https://github.com/HIPS/autograd)

- The *autograd* package from Harvard (HIPS) is an autodiff library for NumPy and SciPy
- In many cases, you can just use it blindly, and then get all the gradients you need!
- You can also extend it by telling *autograd* how to differentiate through custom functions

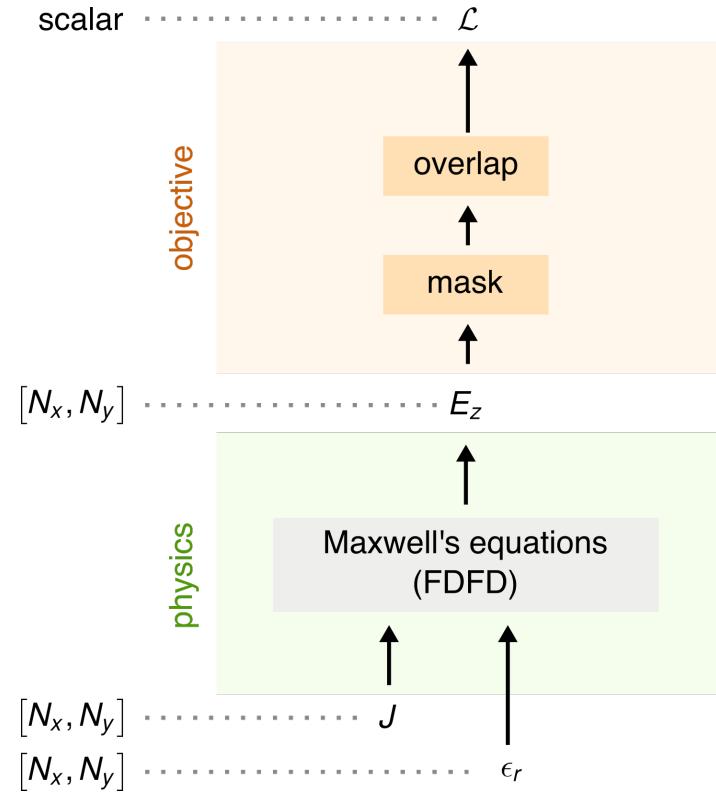
# Ceviche

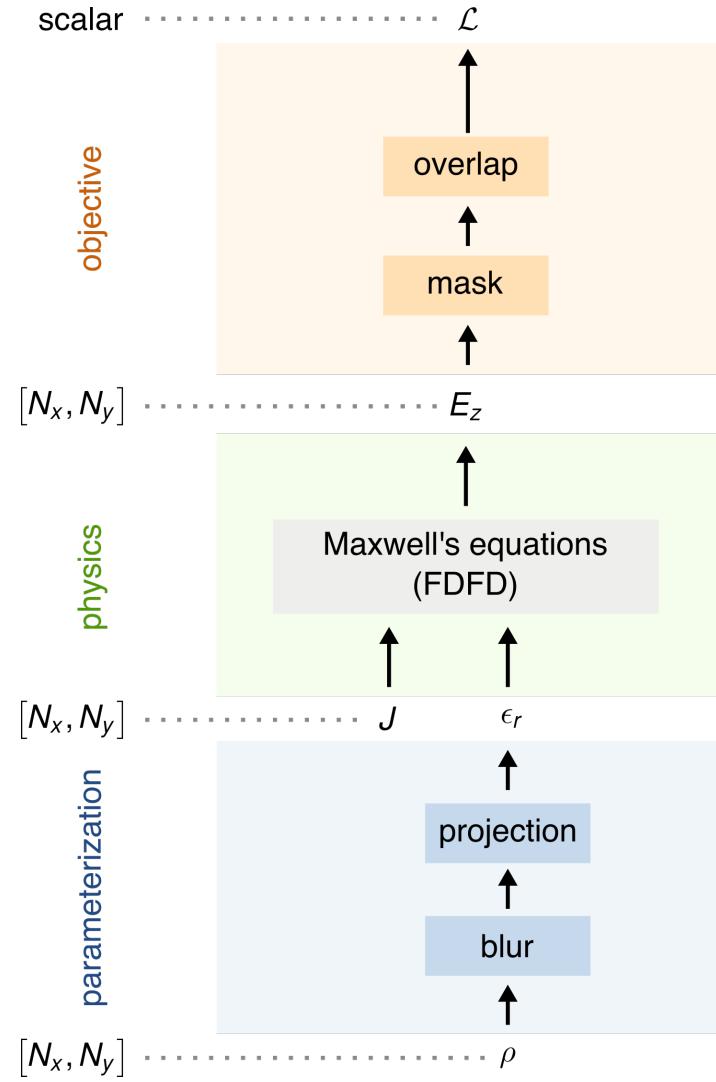
In *ceviche* we've interfaced the AVM with *autograd*

This allows great flexibility in building optimization graphs!

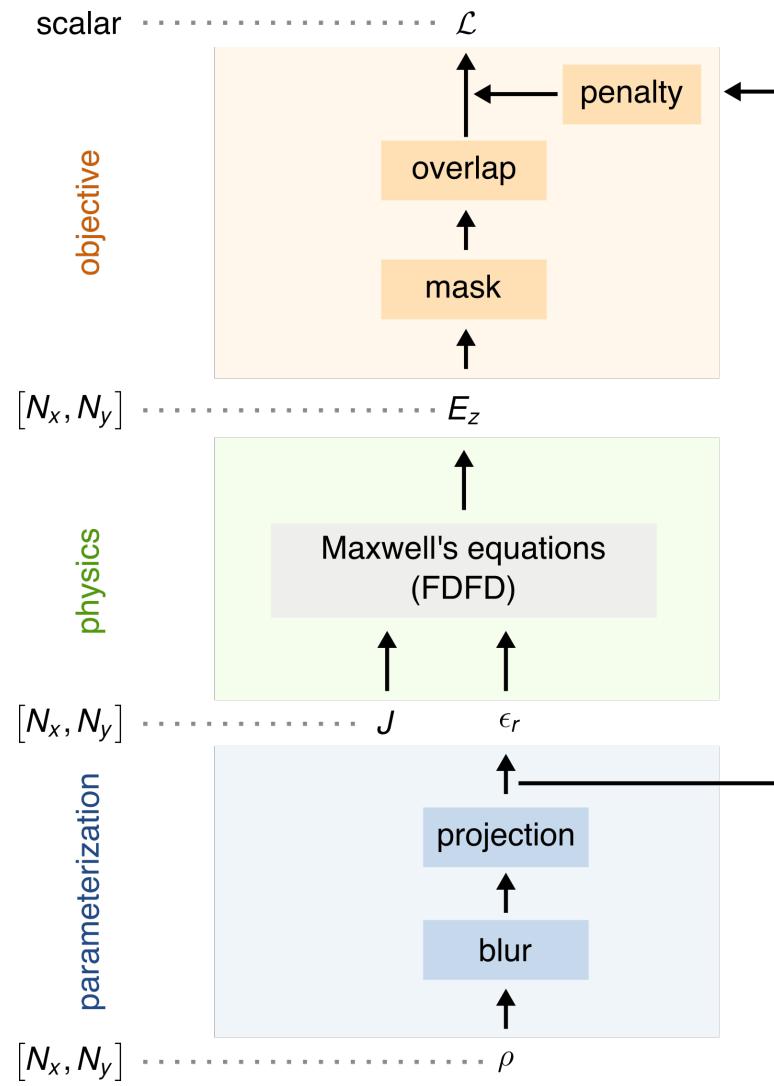


# Notebook #02



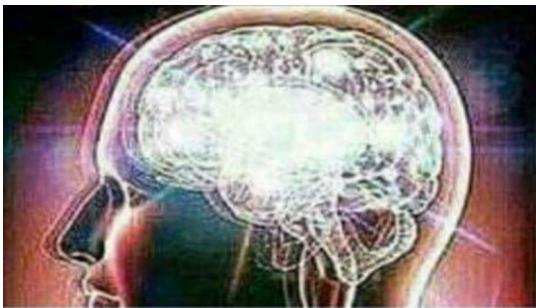


# Notebook #03





The **adjoint variable method** lets us efficiently compute the gradients in *PDE-constrained optimization problems* (inverse design)



The **backpropagation** algorithm is the equivalent of the *adjoint method* for training neural networks



**Automatic differentiation** is the generalization of both *backpropagation* and the *adjoint method* to arbitrary *computational graphs (programs)*

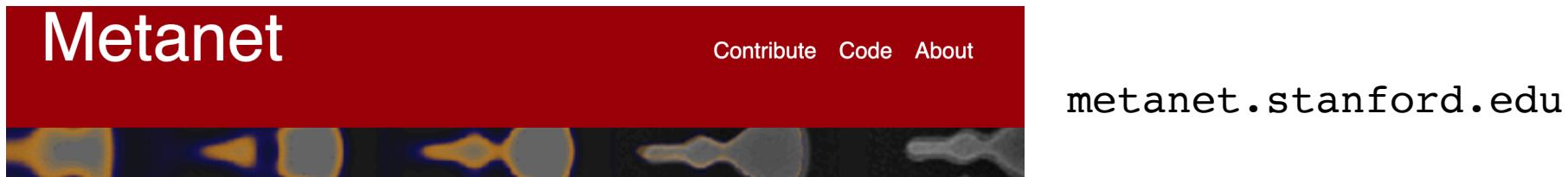


We can utilize **automatic differentiation** to efficiently optimize complex models that freely combine physics, machine learning, and other computations

# Final thoughts (homework?)

- Viewing parameterization as a sequence of ***composable transformations*** is powerful!
  - ML libraries can be useful for things other than ML!
- Hyper parameters are important!
  - *penalty term, binarization*, etc.
  - *Scheduling* (adjusting hyper params during progress of optimization)
- Must go to simulations which capture 3D features to make real devices
  - but doesn't necessarily mean that you can't prototype in 2D!

# Other resources



A collection of resources being setup by Jiaqi Jiang / Jonathan Fan

## Stanford University

- Shanhui Fan
- Jonathan Fan
- Jelena Vuckovic

## UC Berkeley

- Eli Yablonovitch

## Princeton University

- Alejandro Rodriguez

## MIT

- Steven Johnson

## Technical University of Denmark

- Ole Sigmund

... and **many** others!

<https://github.com/fancompute/workshop-invdesign>

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