Chapter 7: Ordinary Differential Equations

7.5 Exercises

7.5.1

7.5.7 A quantum mechanical analysis of the Stark effect (parabolic coordinates) leads to the differential equation

$$\frac{d}{d\xi} \left(\xi \frac{du}{d\xi} \right) + \left(\frac{1}{2} E \xi + \alpha - \frac{m^2}{4\xi} - \frac{1}{4} F \xi^2 \right) u = 0 \tag{1}$$

Here α is a constant, E is the total energy, and F is a constant such that Fz is the potential energy added to the system by the introduction of an electric field.

Using the larger root of the indicial equation, develop a power-series solution about $\xi = 0$. Evaluate the first three coefficients in terms of a_0 .

Solution:

$$0 = \frac{d}{d\xi} \left(\xi \frac{du}{d\xi} \right) + \left(\frac{1}{2} E \xi + \alpha - \frac{m^2}{4\xi} - \frac{1}{4} F \xi^2 \right) u \tag{2}$$

$$0 = \xi \frac{d^2 u}{d\xi^2} + \frac{du}{d\xi} + \left(\frac{1}{2}E\xi + \alpha - \frac{m^2}{4\xi} - \frac{1}{4}F\xi^2\right)u \tag{3}$$

applying our series solution of the form: $\sum_{j=0}^{\infty} a_j \xi^{s+j}$. We have,

$$0 = \xi \frac{d^2 u}{d\xi^2} + \frac{du}{d\xi} + \left(\frac{1}{2}E\xi + \alpha - \frac{m^2}{4\xi} - \frac{1}{4}F\xi^2\right)u \tag{4}$$

$$0 = \xi \sum_{j=0}^{\infty} a_j (s+j)(s+j-1) \xi^{s+j-2} + \sum_{j=0}^{\infty} a_j (s+j) \xi^{s+j-1} + \left(\frac{1}{2} E \xi + \alpha - \frac{m^2}{4\xi} - \frac{1}{4} F \xi^2\right) \sum_{j=0}^{\infty} a_j \xi^{s+j}$$
 (5)

we can simplify this as follows,

$$0 = \xi \sum_{j=0}^{\infty} a_j(s+j)(s+j-1)\xi^{s+j-2} + \sum_{j=0}^{\infty} a_j(s+j)\xi^{s+j-1} + \left(\frac{1}{2}E\xi + \alpha - \frac{m^2}{4\xi} - \frac{1}{4}F\xi^2\right)\sum_{j=0}^{\infty} a_j\xi^{s+j}$$
 (6)

$$0 = \sum_{j=0}^{\infty} a_j(s+j)(s+j-1)\xi^{s+j-1} + \sum_{j=0}^{\infty} a_j(s+j)\xi^{s+j-1} + \frac{1}{2}E\sum_{j=0}^{\infty} a_j\xi^{s+j+1} + \alpha\sum_{j=0}^{\infty} a_j\xi^{s+j}$$
 (7)

$$-\frac{m^2}{4}\sum_{j=0}^{\infty}a_j\xi^{s+j-1} - \frac{F}{4}\sum_{j=0}^{\infty}a_j\xi^{s+j+2}$$
(8)

$$0 = \sum_{i=0}^{\infty} a_j \left[(s+j)(s+j-1) + (s+j) - \frac{m^2}{4} \right] \xi^{s+j-1} + \frac{1}{2} E \sum_{i=0}^{\infty} a_i \xi^{s+j+1}$$
(9)

$$+ \alpha \sum_{i=0}^{\infty} a_j \xi^{s+j} - \frac{F}{4} \sum_{i=0}^{\infty} a_j \xi^{s+j+2}$$
 (10)

$$0 = a_0 \left\{ s(s-1) + s - \frac{m^2}{4} \right\} \xi^{s-1} + \sum_{j=1}^{\infty} a_j \left[(s+j)(s+j-1) + (s+j) - \frac{m^2}{4} \right] \xi^{s+j-1}$$
(11)

$$+\frac{1}{2}E\sum_{j=0}^{\infty}a_{j}\xi^{s+j+1} + \alpha\sum_{j=0}^{\infty}a_{j}\xi^{s+j} - \frac{F}{4}\sum_{j=0}^{\infty}a_{j}\xi^{s+j+2}$$
(12)

$$0 = a_0 \{ s(s-1) + s - \frac{m^2}{4} \} \xi^{s-1} + \sum_{i=0}^{\infty} a_{j+1} \left[(s+j+1)(s+j) + (s+j+1) - \frac{m^2}{4} \right] \xi^{s+j}$$
 (13)

$$+\frac{1}{2}E\sum_{j=1}^{\infty}a_{j-1}\xi^{s+j} + \alpha\sum_{j=0}^{\infty}a_{j}\xi^{s+j} - \frac{F}{4}\sum_{j=2}^{\infty}a_{j+2}\xi^{s+j}$$
(14)

Our indicial equation is given as,

$$s^2 - \frac{m^2}{4} = 0 ag{15}$$

the solutions are $s=\pm rac{m}{2}$

The remaining terms can be simplified as follows,

$$0 = a_0 \{ s(s-1) + s - \frac{m^2}{4} \} + \sum_{j=0}^{\infty} a_{j+1} \left[(s+j+1)(s+j) + (s+j+1) - \frac{m^2}{4} \right] \xi^{s+j} + \frac{1}{2} E \sum_{j=1}^{\infty} a_{j-1} \xi^{s+j}$$
 (16)

$$+ \alpha \sum_{i=0}^{\infty} a_j \xi^{s+j} - \frac{F}{4} \sum_{i=2}^{\infty} a_{j+2} \xi^{s+j}$$
 (17)

$$0 = a_0 \{ s(s-1) + s - \frac{m^2}{4} \} \xi^{s-1} + a_1 [(s+1)s + (s+1) - \frac{m^2}{4}] \xi^s + a_2 [(s+2)(s+1)]$$
(18)

$$+(s+2) - \frac{m^2}{4}]s^{s+1} + a_0 \frac{E}{2} \xi^{s+1} + \alpha a_0 \xi^s + \alpha a_1 \xi^{s+1} + \cdots$$
(19)

And combining like power coefficients, first for ξ^s we have,

$$0 = a_1[(s+1)s + (s+1) - \frac{m^2}{4}] + \alpha a_0 \qquad @s = m/2$$
 (20)

$$0 = a_1[(s^2 + 2s + 1 - \frac{m^2}{4})] + \alpha a_0$$
 (21)

$$0 = a_1[(m/2)^2 + 2(m/2) + 1 - \frac{m^2}{4}] + \alpha a_0$$
 (22)

$$0 = a_1[m+1] + \alpha a_0 \tag{23}$$

So our first coefficient $a_1=-a_0rac{lpha}{m+1}.$ For the coefficient of ξ^{s+1} we have,

$$0 = a_2[s+2)(s+1) + (s+2) - \frac{m^2}{4}] + a_0 \frac{E}{2} + \alpha a_1$$
 (24)

$$0 = a_2 2[m+2] + a_0 \frac{E}{2} - a_0 \alpha \frac{\alpha}{m+1}$$
 (26)

Hence for our third coefficient we have,

$$a_2 = a_0 \left[\frac{\alpha}{2(m+2)(m+1)} - \frac{E}{4(m+2)} \right]$$
 (27)

Our solution with the first three coefficients in terms of a_0 is given:

$$y(x) = x^{m/2} \left(a_0 + \left[-a_0 \frac{\alpha}{m+1}\right] \xi + a_0 \left[\frac{\alpha}{2(m+2)(m+1)} - \frac{E}{4(m+2)}\right] \xi^2 + \cdots\right)$$
(29)

$$y(x) = a_0 x^{m/2} \left\{ 1 - \frac{\alpha}{m+1} \xi + \left[\frac{\alpha}{2(m+2)(m+1)} - \frac{E}{4(m+2)} \right] \xi^2 + \cdots \right\}$$
 (30)

Note that the perturbation F does not appear until a_3 is included.

7.5.8 For the special case of no azimuthal dependence, the quantum mechanical analysis of the hydrogen molecular ion leads to the equation

$$\frac{d}{d\eta} \left[(1 - \eta^2) \frac{du}{d\eta} \right] + \alpha u + \beta \eta^2 u = 0 \tag{31}$$

Develop a power-series solution for $u(\eta)$. Evaluate the first three nonvanishing coefficients in terms of a_0 .

ANS: indicial equation s(s-1)=0

Solution:

$$u_{k=1} = a_0 \eta \left\{ 1 + \frac{2 - \alpha}{6} \eta^2 + \left[\frac{(2 - \alpha)(12 - \alpha)}{120} - \frac{\beta}{20} \right] \eta^4 + \cdots \right\}$$
 (32)