

Ih. Astro - Radiation transfer

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Definitions:

- Spec. radiation intensity:

$$I_\nu \rightarrow I_\nu = \frac{dE}{dA d\Omega d\nu dt}$$

- Spec. radiation energy density:

$$u_\nu = \frac{dE}{dA dt \cdot c} \Rightarrow u_\nu = \frac{I_\nu}{c}$$

- total radiation energy density:

$$E_\nu = \int_{4\pi} u_\nu d\Omega = \int \frac{I_\nu}{c} d\Omega$$

- mean Spec. radiation intensity:

$$\langle I_\nu \rangle = \frac{1}{4\pi} \int_0^{4\pi} I_\nu d\Omega$$

$$J_v = \frac{1}{4\pi} \int_{4\pi} I_v d\Omega$$

- Energy conservation in vacuum:

$$\frac{dI_v}{ds} = 0 \xrightarrow{\text{Lagrange derivative}} \frac{\partial I_v}{\partial t} + \hat{c} \cdot \vec{\nabla} I_v = 0$$

- Different derivation:

$$dN = f(\vec{x}, \vec{p}) d^3x d^3p$$

$$\hookrightarrow \begin{cases} d^3x = ds c dt \\ d^3p = p^2 dp d\Omega = \left(\frac{h\nu}{c}\right)^2 \left(\frac{h d\nu}{c}\right) d\Omega \end{cases}$$

$$\hookrightarrow dE = 2 \cdot h\nu \cdot dN$$

d

$$\Rightarrow dE = f \, ds \, c \, dt \left(\frac{hv}{c} \right)^2 \left(\frac{u_{vv}}{c} \right) dr \cdot hv$$

$$= I_v \, ds \, dt \, dr \, dv$$

$$\Rightarrow I_v = \frac{h^4 v^3}{c^2} f$$

- Boltzmann equation:

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \vec{\nabla} f = 0 \quad (\Leftrightarrow)$$

$$\frac{\partial I_v}{\partial t} + c \cdot \hat{n} \cdot \vec{\nabla} I_v = 0$$

- Interaction with matter:

$$\frac{1}{c} \frac{\partial I_v}{\partial t} + \hat{n} \cdot \vec{\nabla} I_v = j_v - \alpha_v I_v$$

$$j_v = [\text{cm}^{-2} \text{s}^{-1} \text{sr}^{-1}]$$

Radiative transfer equation

$$\frac{dI_v}{ds} = j_v - \alpha_v I_v$$

induced emission
Same direction as
incoming r

$ds = j^v \sim \sim$ Same direction as incoming radiation
→ quantum effect!

optical depth:

$$d\tau_v \equiv \alpha_v ds \Rightarrow \frac{dI_v}{d\tau_v} = \underbrace{\frac{j_v}{\alpha_v}}_{= S_v} - I_v$$

$$\Rightarrow \frac{dI_v}{d\tau_v} = S_v - I_v$$

↳ formal solution:

$$I_v = \int_0^{\tilde{\tau}_v} S_v e^{(\tilde{\tau} - \tilde{\tau}_v)} d\tilde{\tau} + I_v(0) e^{-\tilde{\tau}_v}$$

- optically thin $\tilde{\tau}_v \ll 1$: $I_v \approx I_v(0) + [S_v - I_v(0)] \tilde{\tau}_v$
- optically thick $\tilde{\tau}_v \gg 1$: $I_v \approx S_v$
- Moments of the radiative transfer equation:

∫

$$1) \int d\Omega : \int \frac{1}{c} \frac{\partial I_v}{\partial t} d\Omega + \int \hat{u} \cdot \nabla I_v d\Omega = \int j_v d\Omega - \int \alpha_v I_v d\Omega$$

$$\Leftrightarrow \underbrace{\frac{\partial}{\partial t} \left(\int \frac{I_v}{c} d\Omega \right)}_{\equiv E_v} + \underbrace{\nabla \left(\int \hat{u} \cdot I_v d\Omega \right)}_{\equiv \vec{F}_v} = 4\pi j_v - \underbrace{\alpha_v c \int \frac{I_v}{c} d\Omega}_{\equiv E_v}$$

$$\Rightarrow \boxed{\frac{\partial E_v}{\partial t} + \nabla \cdot \vec{F}_v = 4\pi j_v - \alpha_v c E_v} \quad (\text{Sometimes } S_v = \int_{4\pi} j_v d\Omega)$$

radiation energy conservation equation

$$2) \int \hat{u} d\Omega : \int \frac{1}{c} \hat{u} \frac{\partial I_v}{\partial t} d\Omega + \int \nabla (\hat{u} \otimes \hat{u}) I_v d\Omega = \underbrace{\int \hat{u} j_v d\Omega}_{=0 \text{ (isotropic)}} - \int \alpha_v \hat{u} I_v d\Omega \equiv \vec{P}_v$$

$$\Rightarrow \boxed{\frac{1}{c} \frac{\partial \vec{P}_v}{\partial t} + c \nabla \cdot \vec{P}_v = -\alpha_v \vec{F}_v} \quad \text{or} \quad \boxed{\frac{1}{c^2} \frac{\partial \vec{F}_v}{\partial t} + \nabla \cdot \vec{P}_v = -\frac{\alpha_v \vec{F}_v}{c}}$$

radiation momentum conservation equation

Radiation momentum conservation

No need for closure relation!

\tilde{F}_v : diffusion limit

\underline{P}_v : M1-closure (\rightarrow analogously to EOS)

Coupling to hydrodynamics:

exchange of momentum and energy via radiation-fluid interaction

\hookrightarrow sink-source terms for Euler equation:

$$1) \int d\Omega : \frac{\partial E_{rad}}{\partial t} + \vec{\nabla} \cdot \vec{F}_{rad} = \Lambda_{rad} - P_{rad}$$

$$\text{with } \Lambda_{rad} = \int d\Omega c E_v d\Omega \quad (\text{cooling})$$

$$P_{rad} = \int S_v d\Omega = \iint j_v dS d\Omega \quad (\text{heating})$$

$$\vec{E}_{rad} = \int \vec{E}_v d\Omega$$

$$\tilde{F}_{\text{rad}} = \int \tilde{F}_v dv$$

\Rightarrow Euler equation:
(energy conservation)

$$\frac{\partial E}{\partial t} + \bar{\nabla}(E + \underline{P}) \bar{v} = \underline{F}_{\text{rad}} - \underline{\Lambda}_{\text{rad}} + \rho \bar{g} \bar{v}$$

$$2) \int dv : \frac{1}{c^2} \frac{\partial \tilde{F}_{\text{rad}}}{\partial t} + \bar{\nabla} \cdot \underline{\underline{P}_{\text{rad}}} = - \tilde{F}_{\text{rad}}$$

with $\underline{\underline{P}_{\text{rad}}} = \int \underline{\underline{P}_v} dv$

$$\tilde{F}_{\text{rad}} = \int \frac{\alpha_v \tilde{F}_v}{c} dv$$

$$\Rightarrow \text{Euler equation: } \frac{\partial(\rho \bar{v})}{\partial t} + \bar{\nabla}(\rho(\bar{v} \otimes \bar{v}) + \underline{P}) = \rho \bar{g} + \tilde{F}_{\text{rad}}$$

• Equilibria / LTE:

1) thermal radiation: gas at equilibrium

$$j_v = B_v \alpha_v = I_v \alpha_v$$

2) Blackbody radiation: radiation at equilibrium

$$I_v = B_v(T_{\text{rad}}) \rightarrow \text{isotropic}$$

$$\hookrightarrow \bar{F}_v = \int \hat{n} \cdot \vec{I}_v d\Omega = B_v \int \hat{n} d\Omega = 0 \Rightarrow \bar{F}_{\text{rad}} = 0$$

$$\hookrightarrow \underline{P}_v = \int \frac{\vec{I}_v}{c} \hat{n} \otimes \hat{n} d\Omega = \frac{B_v}{c} \int \hat{n} \otimes \hat{n} d\Omega = \frac{B_v}{c} \cdot \frac{4\pi}{3} \underline{1}$$

$$\hookrightarrow E_v = \int \frac{\vec{I}_v}{c} d\Omega = \frac{B_v}{c} \int d\Omega = \frac{4\pi}{c} B_v$$

$$\Rightarrow \underline{P}_v = \frac{E_v}{3} \underline{1} \Rightarrow P_{\text{rad}} = \frac{E_{\text{rad}}}{3}$$

Eddington's approximation

$$\left. \begin{aligned} B_v|_{RJ} &\approx \frac{2v^2}{c^2} k_B T \\ B_v|_{Wien} &\approx \frac{2v^2}{c^2} h v e^{-\frac{hv}{k_B T}} \end{aligned} \right\} \text{Blackbody: } I_v = \frac{h^4 v^3}{c^2} f$$

$$f = \frac{N}{h^3} \quad N = \frac{2}{e^{\frac{hv}{kT}} - 1} \quad \left\{ \begin{array}{l} \downarrow \\ B_v = \frac{2v^2}{c^2} \frac{hv}{e^{\frac{hv}{kT}} - 1} \\ \hookrightarrow E_{rad} = \int \int \frac{B_v}{c} dS dv = a T^4 \end{array} \right.$$

Diffusion limit:

$$\frac{1}{c} \frac{\partial F_v}{\partial t} + c \nabla \underline{\underline{P}_v} = -\alpha_v F_v$$

$L \leq cT$ and $\alpha_v = \frac{1}{\lambda_v}$ with $\lambda_v \ll L$

\hookrightarrow discard $\frac{\partial F_v}{\partial t}$ and use Eddington tensor $\underline{\underline{P}_v} = \underline{\underline{E}_v} \underline{\underline{D}_v} \stackrel{\text{isotropic}}{\simeq} \frac{1}{3} \underline{\underline{E}_v}$

$$\Rightarrow \boxed{\nabla \frac{c}{3} \underline{\underline{E}_v} = -\alpha_v \underline{\underline{F}_v}} \quad \text{with } \alpha_v = \frac{1}{\lambda_v} \quad \text{diffusion limit}$$

$$\Rightarrow \boxed{\frac{\partial E_v}{\partial t} - \nabla \left(\frac{c}{3\alpha_v} \nabla E_v \right) = 4\pi j_v - \alpha_v c E_v} \quad \text{diffusion approximation}$$

Diffusion equation

Rosseland mean: $\bar{T}_v = -\frac{c}{3\alpha_v} \bar{\nabla} E_v$

$$\Rightarrow \bar{F}_{rad} = \frac{c}{3\alpha_R} \bar{\nabla} (a T^4) = K_{rad} \bar{\nabla} T$$

with $K_{rad} = \frac{c}{3\alpha_R} 4 a T^3$

thermal radiation:

$$j_v = \alpha_v B_v$$

$$\hookrightarrow \frac{\partial E_v}{\partial t} - \bar{\nabla} \left(\frac{c}{3\alpha_v} \bar{\nabla} E_v \right) = \alpha_v c \left(\frac{4\pi}{c} B_v - E_v \right)$$

Planck mean:

$$\frac{\partial E_v}{\partial t} - \bar{\nabla} \left(\frac{c}{3\alpha_R} \bar{\nabla} E_v \right) \approx \alpha_p c a \underbrace{\left(T_{gas}^4 - T_{rad}^4 \right)}_{\text{cooling / heating equilibrium}}$$

cooling / heating equilibrium

- Planar geometry: