

Introduction to Quantum Computing

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Quantum bits or Qubits

- The set $\{0, 1\}$ is a classical bit. If $x \in \{0, 1\}$, we say that x is the state of a classical **bit**.
- The set $\left\{ \begin{pmatrix} a \\ b \end{pmatrix} : a, b \in \mathbb{C}, |a|^2 + |b|^2 = 1 \right\}$ is a quantum bit, or a **qubit**.
- Example 1: the space of all possible polarization states of a photon is a qubit. *spin*
- Example 2: the space of all possible spins of an electron is said to be a qubit.

- ✓ Superposition - quantum state
 - ✓ Entanglement - of qubit
- A single

bit is a system that can exist in two states.

- 0 'zero' state
- 1 'one' state

~~1 1 1 1 1 1 1~~

64

qubits

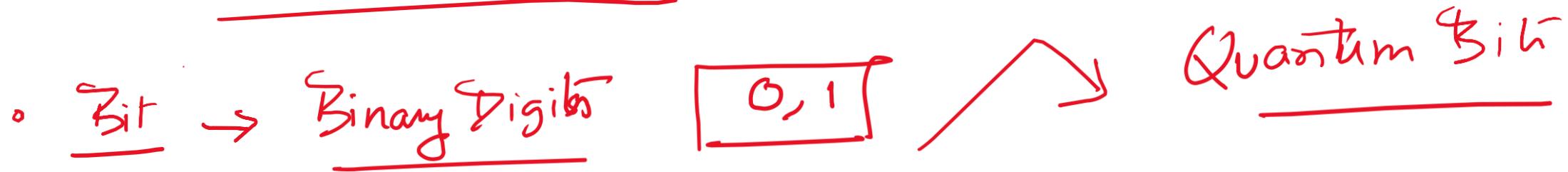
~~1 1 1~~ bits

What can we do with
bits?

x Read
x Write

x Store → without
and over
in a very tiny

- On a quantum computer there are issues like
 - × readers . × stay



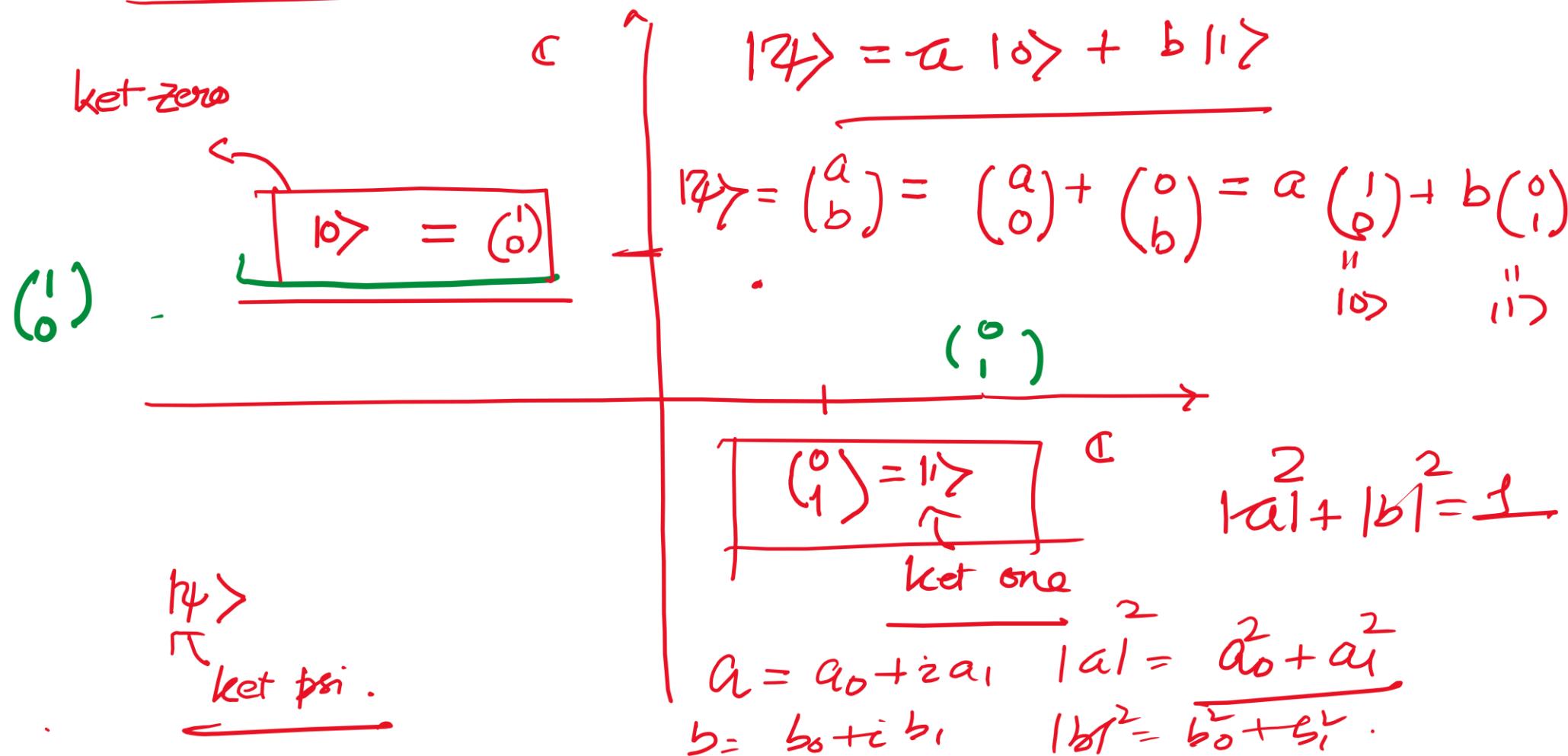
A quantum bit is a quantum mechanical system that exist in infinite number of states. A single qubit state can be specified by a pair of complex number a, b . $\begin{pmatrix} a \\ b \end{pmatrix}$ when $|a|^2 + |b|^2 = 1$

$$\begin{pmatrix} a \\ b \end{pmatrix}$$

$$|a|^2 + |b|^2 = 1$$

Unit of information

The single qubit state space



Vector spaces over \mathbb{C} and qubits

- A single-qubit state space is a two-dimensional vector space over the field of complex numbers \mathbb{C} .
- We represent it as

$$\mathbb{C}^2 = \left\{ \begin{pmatrix} a \\ b \end{pmatrix} : a, b \in \mathbb{C}, |a|^2 + |b|^2 = 1 \right\}$$

- The computational basis is

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- A qubit state is written as

$$|\psi\rangle = a|0\rangle + b|1\rangle$$

We can write
and single-qubit
state as a

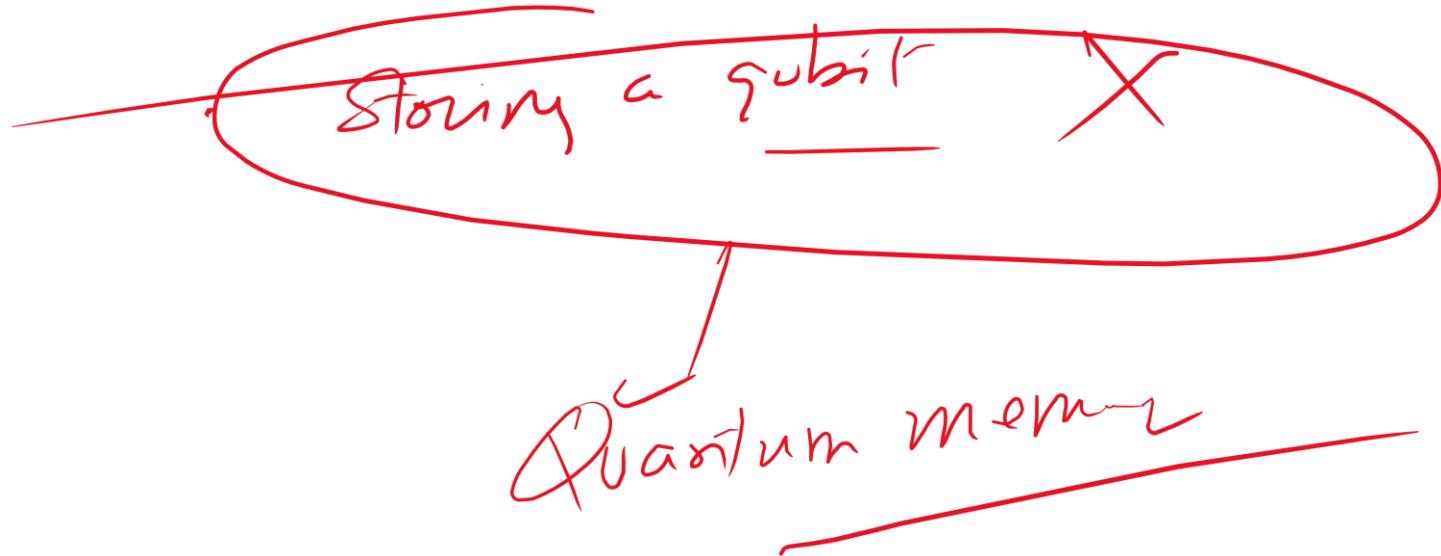
$$\left\{ \begin{pmatrix} a \\ b \end{pmatrix} : a, b \in \mathbb{C}, |a|^2 + |b|^2 = 1 \right\}$$

linear combination of
 $|0\rangle$ (ket 0) and $|1\rangle$ (ket 1)

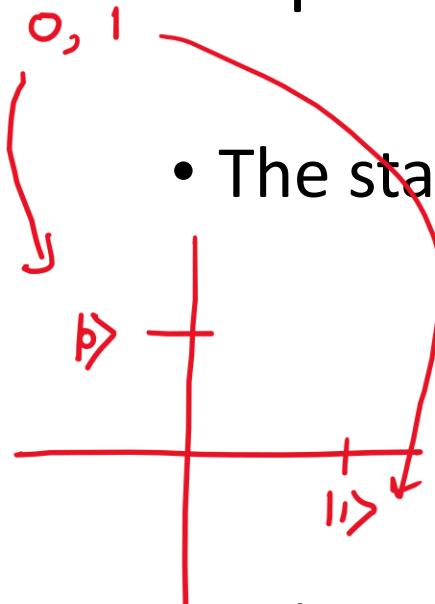
- It is possible to write a single qubit to a particular single qubit state.

Write ✓

- Reading a qubit X - Measurement



Superposition of states



- The state of a single-qubit is of the form

$$|\psi\rangle = \begin{pmatrix} a \\ b \end{pmatrix} = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix} = a|0\rangle + b|1\rangle$$

where $|a|^2 + |b|^2 = 1$.

$|0\rangle$, $|1\rangle$, and $|\psi\rangle = a|0\rangle + b|1\rangle$
when $a \neq 0, b \neq 0$. Then

$|\psi\rangle$ is called a SUPERPOSITION of the states $|0\rangle$ and $|1\rangle$

- If $a \neq 0$ and $b \neq 0$ the qubit is said to be in the superposition of two states $|0\rangle$ and $|1\rangle$.

What is a superposition of states?

If I am construction a single qubit state by taking a (linear) combintion of

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad |2\rangle = a|0\rangle + b|1\rangle$$

The pair $\{|0\rangle, |1\rangle\}$ is called a basis of the single-qubit state space.

This is a very important basis, so much so that it has special name. It is called **COMPUTATIONAL BASIS**.

Hadamard Basis

$$|H\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}$$

$$|+\rangle = \frac{|0\rangle - |1\rangle}{\sqrt{2}}$$

$$\frac{c+d}{\sqrt{2}} = a, \frac{c-d}{\sqrt{2}} = b$$

$$c+d = a\sqrt{2} \quad c-d = b\sqrt{2}$$

$$2c = (a+b)\sqrt{2}$$

$$c = \frac{a+b}{\sqrt{2}}$$

$$d = \frac{a-b}{\sqrt{2}}$$

$$|\psi\rangle = a|0\rangle + b|1\rangle$$

$$|\psi\rangle = c|+\rangle + d|-\rangle$$

$$= \frac{c}{\sqrt{2}}(|0\rangle + |1\rangle) + \frac{d}{\sqrt{2}}(|0\rangle - |1\rangle)$$

$$= \left(\frac{c+d}{\sqrt{2}}\right)|0\rangle + \left(\frac{c-d}{\sqrt{2}}\right)|1\rangle = (a)|0\rangle + b|1\rangle$$

$$|\psi\rangle = \frac{a+b}{\sqrt{2}}|+\rangle + \frac{a-b}{\sqrt{2}}|-\rangle$$

We have a qubit state $|\psi\rangle$ written in the computational basis.

Suppose there is a quantum state $|24\rangle$ which is in superposition with respect to the computational basis. Is it in superpositn with all other basis?

$$\underline{|24\rangle} = \frac{1}{\sqrt{2}}|10\rangle + \frac{1}{\sqrt{2}}|11\rangle = \frac{|10\rangle + |11\rangle}{\sqrt{2}}$$

With respect to the Hadamard basis $|24\rangle$ is just $|+\rangle$. So it is not in superpositn.

Once a superposition, always a superposition?

NO

- $|\psi\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}$ is a superposition of two states $|0\rangle$, and $|1\rangle$.
- We say that $|\psi\rangle$ is in superposition with respect to the basis $\{|0\rangle, |1\rangle\}$.
- However, the representation of $|\psi\rangle$ with respect to the basis $\mathcal{H} = \{|+\rangle, |-\rangle\}$ is $|\psi\rangle = |+\rangle$.
- Therefore, $|\psi\rangle$ is not in superposition with respect to the basis \mathcal{H} .

Changing a Qubit representation from computational to Hadamard basis

- $|\psi\rangle = a|0\rangle + b|1\rangle$ is a single-qubit state written in computational basis.
- The Hadamard basis vectors in terms of computational basis vectors are:

$$|+\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}, \quad |-\rangle = \frac{|0\rangle - |1\rangle}{\sqrt{2}}.$$

- Solving for $|0\rangle$ and $|1\rangle$ yields:

$$|0\rangle = \frac{|+\rangle + |-\rangle}{\sqrt{2}}, \quad |1\rangle = \frac{|+\rangle - |-\rangle}{\sqrt{2}}.$$

- $|\psi\rangle = a\left(\frac{|+\rangle + |-\rangle}{\sqrt{2}}\right) + b\left(\frac{|+\rangle - |-\rangle}{\sqrt{2}}\right) = \frac{a+b}{\sqrt{2}} |+\rangle + \frac{a-b}{\sqrt{2}} |-\rangle.$

Global phase versus relative phase

- Two single-qubit states $|\psi\rangle = a|0\rangle + b|1\rangle$ and $|\phi\rangle = c|0\rangle + d|1\rangle$ are said to differ by the global phase θ if

$$|\psi\rangle = a|0\rangle + b|1\rangle = e^{i\theta}(c|0\rangle + d|1\rangle) = e^{i\theta} |\phi\rangle.$$

- If two quantum states differ by a global phase, they are considered to be same. We write $|\psi\rangle \sim |\phi\rangle$.
- The relative phase of a single-qubit state $|\psi\rangle = a|0\rangle + b|1\rangle$ is a number φ which satisfies the equation

$$\frac{a}{b} = e^{i\varphi} \frac{|a|}{|b|}.$$

- Two quantum states with different relative phases are not the same quantum state.

Examples of qubits differing by a global phase

- Consider: $\frac{1}{\sqrt{2}}(|0\rangle + e^{\frac{i\pi}{4}}|1\rangle)$ and $\frac{1}{\sqrt{2}}(e^{-\frac{i\pi}{4}}|0\rangle + |1\rangle)$
- The qubit state $\frac{1}{\sqrt{2}}(e^{-\frac{i\pi}{4}}|0\rangle + |1\rangle) = \frac{e^{-\frac{i\pi}{4}}}{\sqrt{2}}(|0\rangle + e^{\frac{i\pi}{4}}|1\rangle)$
- Therefore, these two quantum states are the same.

Examples of qubits differing by relative phases

- Consider: $\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ and $\frac{1}{\sqrt{2}}(-|0\rangle + \mathbf{i}|1\rangle)$

- Let $a|0\rangle + b|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$
and $a'|0\rangle + b'|1\rangle = \frac{1}{\sqrt{2}}(-|0\rangle + \mathbf{i}|1\rangle).$
$$\frac{a}{b} = \frac{1}{\sqrt{2}} \frac{\sqrt{2}}{1} = e^{0\mathbf{i}} \frac{|a|}{|b|}, \quad \text{and} \quad \frac{a'}{b'} = -\frac{1}{\sqrt{2}} \frac{\sqrt{2}}{\mathbf{i}} = -\frac{1}{\mathbf{i}} = \mathbf{i} = e^{\frac{\pi\mathbf{i}}{2}} \frac{|a'|}{|b'|}.$$

By definition the relative phase of the first qubit is 0 and the relative phase of the second qubit is $\frac{\pi}{2}$. Since they have different relative phases they are different quantum states.

Complex Inner Product

- Let V be an n -dimensional \mathbb{C} -vector space.

- Let $|a\rangle = \begin{pmatrix} a_1 \\ a_2 \\ \vdots \\ a_n \end{pmatrix}$ and $|b\rangle = \begin{pmatrix} b_1 \\ b_2 \\ \vdots \\ b_n \end{pmatrix}$

- $\langle a|b \rangle = (\bar{a}_1, \bar{a}_2, \dots, \bar{a}_n) \begin{pmatrix} b_1 \\ b_2 \\ \vdots \\ b_n \end{pmatrix} = \sum_{i \in [n]} \bar{a}_i b_i$

$$\underline{ac + bd = 0}$$

$$\|u\|^2 = u \cdot u = a^2 + b^2$$

$$\|u\| = \sqrt{a^2 + b^2}$$

$u \cdot v = ac + bd$

$$(\begin{matrix} a \\ b \end{matrix})^T (\begin{matrix} c \\ d \end{matrix}) = (a \ b) (\begin{matrix} c \\ d \end{matrix})$$

$$= ac + bd$$

$$\mathbb{R}^n = \left\{ \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix} : a_i \in \mathbb{R} \right\}$$

$$\bar{a}, \bar{b} \in \mathbb{R}^n$$

$$\bar{a} \cdot \bar{b} = \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix}^T \begin{pmatrix} b_1 \\ \vdots \\ b_n \end{pmatrix}$$

$$= (a_1 \dots a_n) \begin{pmatrix} b_1 \\ \vdots \\ b_n \end{pmatrix}$$

$$= a_1 b_1 + \dots + a_n b_n$$

Dot product in Complex
vector spaces \rightarrow Inner product.

$$\bar{A}^n = \left\{ \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix} : a_i \in \mathbb{C} \right\}$$

$$\bar{a} = \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix} \quad \begin{matrix} 1+i \\ -1-i \end{matrix}$$

$$\bar{b} = \begin{pmatrix} b_1 \\ \vdots \\ b_n \end{pmatrix}$$

$$\bar{a} \cdot \bar{b} = (a_1 \dots a_n) \begin{pmatrix} b_1 \\ \vdots \\ b_n \end{pmatrix}$$

$$= a_1 b_1 + \dots + a_n b_n$$

$$\bar{a} \cdot \bar{a} = \sum_{i=1}^n \underbrace{(a_i)^2}_{a_i \in \mathbb{C}} \quad a_i^2 \in \mathbb{C}$$

$$\bar{a} \cdot \bar{a}$$

$$x = \begin{pmatrix} x_1 \\ \vdots \\ x_n \end{pmatrix}$$

$$x^T = (\bar{x}_1, \bar{x}_2, \dots, \bar{x}_n)$$

$$\langle x, y \rangle = x^T y = (\bar{x}_1, \dots, \bar{x}_n) \begin{pmatrix} y_1 \\ \vdots \\ y_n \end{pmatrix}$$

$$= \bar{x}_1 y_1 + \dots + \bar{x}_n y_n$$

$$\boxed{\begin{aligned} \bar{a} &= a_0 + i a_1 \\ \overline{\bar{a}} &= a_0 - i a_1 \end{aligned}}$$

$$\begin{aligned} \langle x, x \rangle &= (\bar{x}_1, \dots, \bar{x}_n) \begin{pmatrix} x_1 \\ \vdots \\ x_n \end{pmatrix} = \bar{x}_1 x_1 + \dots + \bar{x}_n x_n = \\ &= |x_1|^2 + \dots + |x_n|^2. \end{aligned}$$

ket a

$$|a\rangle = \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix}$$

$\langle a | b \rangle$

$$\begin{aligned} &= (\bar{a}_1, \bar{a}_2, \dots, \bar{a}_n) \begin{pmatrix} b_1 \\ b_2 \\ \vdots \\ b_n \end{pmatrix} \\ &= \bar{a}_1 \cdot b_1 + \dots + \bar{a}_n b_n = \sum_{i=1}^n \bar{a}_i b_i \end{aligned}$$

bare a

$\langle a | = (\bar{a}_1 \dots \bar{a}_n)$

$$= \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix}^\dagger$$

bra' a' ket' b'

↙ ↘

bra ket a, b

(c)

$$\langle \bar{a} | a \rangle = (\bar{a}_1, \dots, \bar{a}_n) \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix} = \sum_{i=1}^n \bar{a}_i a_i$$

$$= \sum_{i=1}^n |a_i|^2$$

$$a_i = x + iy$$

$$\tilde{z} = -i$$

$$\bar{a}_i = x - iy$$

$$\bar{a}_i \cdot a_i = (x + iy)(x - iy) = x^2 - ixy + ixy - i^2 y^2$$

$$= \cancel{x^2 + y^2} = |a_i|^2.$$

Measuring a qubit

Measurement of a Single-Qubit System

How do we "read" a qubit state??

- Any measurement of a quantum system is associated to an orthonormal basis of its state space.
- Two orthonormal bases of \mathbb{C}^2 are

Orthonormal Basis

$$\mathcal{B} = \{|0\rangle, |1\rangle\}$$

$$\mathcal{H} = \{|+\rangle, |-\rangle\} = \left\{ \frac{|0\rangle + |1\rangle}{\sqrt{2}}, \frac{|0\rangle - |1\rangle}{\sqrt{2}} \right\}$$

- \mathcal{B}_1 is said to be the computational basis, \mathcal{B}_2 is said to be the Hadamard basis of \mathbb{C}^2 .

o Any measurement corresponds to an orthonormal basis.

↓
of a quantum state

↓
single qubit state

$$|4\rangle = a|0\rangle + b|1\rangle$$

If I measure $|4\rangle$ by the measurment having the basis $\{|0\rangle, |1\rangle\}$

the measurement outcome is

$|0\rangle$ with probability

$|1\rangle$ with probability

$$|\langle 0|4\rangle|^2 \text{ and}$$

$$|\langle 1|4\rangle|^2$$

↓
of the state space

↓
single qubit state

$$\mathbb{C}^2 \quad \left| \begin{array}{l} \{|0\rangle, |1\rangle\} \xrightarrow{\rho_0} \\ \xrightarrow{\hspace{1cm}} \end{array} \right.$$

$$\left| \begin{array}{l} \{|0\rangle, |1\rangle\} \\ \xrightarrow{\hspace{1cm}} \end{array} \right. = \mathcal{H}$$

Single qubit measurement

- A single-qubit measurement, M is associated to an orthonormal basis $\{|\Phi_1\rangle, |\Phi_2\rangle\}$
- Measuring $|\Psi\rangle = a|0\rangle + b|1\rangle$ by M outputs either $|\Phi_1\rangle$ or $|\Phi_2\rangle$.
- The probability of outcome $|\Phi_1\rangle$ is $|\langle\Phi_1|\Psi\rangle|^2$
- The probability of outcome $|\Phi_2\rangle$ is $|\langle\Phi_2|\Psi\rangle|^2$

Example 1

- Consider the single-qubit state $|\Psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + \mathbf{i}|1\rangle)$ and the measurement basis $\{|0\rangle, |1\rangle\}$.
- The measurement outcome is $|0\rangle$ with probability

$$|\langle 0|\Psi\rangle|^2 = \left| \frac{1}{\sqrt{2}} \right|^2 = \frac{1}{2}$$

- The measurement outcome is $|1\rangle$ with probability

$$|\langle 1|\Psi\rangle|^2 = \left| \mathbf{i} \frac{1}{\sqrt{2}} \right|^2 = \frac{1}{2}$$

Calculations

$$\bullet \langle 0 | \Psi \rangle = \langle 0 | \left(\frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} \mathbf{i} |1\rangle \right) = \frac{1}{\sqrt{2}} \langle 0 | 0 \rangle + \frac{1}{\sqrt{2}} \mathbf{i} \langle 0 | 1 \rangle = \frac{1}{\sqrt{2}}.$$

$$\bullet \langle 0 | \Psi \rangle = \langle 1 | \left(\frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} \mathbf{i} |1\rangle \right) = \frac{1}{\sqrt{2}} \langle 1 | 0 \rangle + \frac{1}{\sqrt{2}} \mathbf{i} \langle 1 | 1 \rangle = \frac{1}{\sqrt{2}} \mathbf{i}.$$

Example 2

- Consider the single-qubit state $|\Psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + \mathbf{i}|1\rangle)$ and the measurement basis $\{|+\rangle, |-\rangle\}$.

- The measurement outcome is $|+\rangle$ with probability

$$|\langle +|\Psi\rangle|^2 = \left| \frac{1}{2} (1 + \mathbf{i}) \right|^2 = \frac{1}{2}.$$

- The measurement outcome is $|-\rangle$ with probability

$$|\langle -|\Psi\rangle|^2 = \left| \frac{1}{2} (1 - \mathbf{i}) \right|^2 = \frac{1}{2}.$$

Calculations

- $\langle +|\Psi\rangle = \left(\frac{1}{\sqrt{2}}(\langle 0| + \langle 1|)\right) \left(\frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle)\right) = \frac{1}{2}(1 + i).$
- $\langle -|\Psi\rangle = \left(\frac{1}{\sqrt{2}}(\langle 0| - \langle 1|)\right) \left(\frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle)\right) = \frac{1}{2}(1 - i).$
- $|\langle +|\Psi\rangle|^2 = \left|\frac{1}{2}(1 + i)\right|^2 = \frac{1}{2}.$
- $|\langle -|\Psi\rangle|^2 = \left|\frac{1}{2}(1 - i)\right|^2 = \frac{1}{2}.$

Outer product

- Let $|\psi\rangle$ and $|\Phi\rangle$ be two vectors.
- $|\psi\rangle = a|0\rangle + b|1\rangle$ and $|\Phi\rangle = c|0\rangle + d|1\rangle$.
- The outer product of $|\psi\rangle$ and $|\Phi\rangle$ is
$$|\Psi\rangle\langle\Phi| = \begin{pmatrix} a \\ b \end{pmatrix} \begin{pmatrix} c \\ d \end{pmatrix}^\dagger = \begin{pmatrix} a \\ b \end{pmatrix} (\bar{c} \quad \bar{d})$$
$$= \begin{pmatrix} a\bar{c} & a\bar{d} \\ b\bar{c} & b\bar{d} \end{pmatrix}$$

Quantum state transformations

- Quantum computers have the capability of transforming one quantum state to another by applying unitary transformations on the former.
- A linear transformation T is said to be unitary if

$$T T^\dagger = I$$

where I is the identity operator.

The Pauli Transformations

- $I : |0\rangle\langle 0| + |1\rangle\langle 1| = \begin{pmatrix} 1 \\ 0 \end{pmatrix} (1 \quad 0) + \begin{pmatrix} 0 \\ 1 \end{pmatrix} (0 \quad 1)$
 $= \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$
- $X : |1\rangle\langle 0| + |0\rangle\langle 1| = \begin{pmatrix} 0 \\ 1 \end{pmatrix} (1 \quad 0) + \begin{pmatrix} 1 \\ 0 \end{pmatrix} (0 \quad 1)$
 $= \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

The Pauli Transformations

- $Y : -|1\rangle\langle 0| + |0\rangle\langle 1| = -\begin{pmatrix} 0 \\ 1 \end{pmatrix}(1 \quad 0) + \begin{pmatrix} 1 \\ 0 \end{pmatrix}(0 \quad 1)$
 $= -\begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$
- $Z : |0\rangle\langle 0| - |1\rangle\langle 1| = \begin{pmatrix} 1 \\ 0 \end{pmatrix}(1 \quad 0) - \begin{pmatrix} 0 \\ 1 \end{pmatrix}(0 \quad 1)$
 $= \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} - \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

Action of the Pauli Transformations

- I = identity transformation
- X = negation, it is similar to the classical not operation
- Z = changing the relative phase of a superposition in the standard basis.
- $Y = ZX$.

The Hadamard Transformation

- $H = \frac{1}{\sqrt{2}}(|0\rangle\langle 0| + |1\rangle\langle 0| + |0\rangle\langle 1| - |1\rangle\langle 1|) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}.$

Two qubit states

- Consider two qubits

$$|\Phi_1\rangle = a |0\rangle + b |1\rangle$$

and

$$|\Phi_2\rangle = c |0\rangle + d |1\rangle$$

If these two qubits exist side by side, then we have a two-qubit state

$$(|\Phi_1\rangle, |\Phi_2\rangle) = (a |0\rangle + b |1\rangle, \quad c |0\rangle + d |1\rangle)$$

Two qubit states: All measurements are with respect to $\{|0\rangle, |1\rangle\}$

$$(|\Phi_1\rangle, |\Phi_2\rangle) = (a|0\rangle + b|1\rangle, c|0\rangle + d|1\rangle)$$

- If we measure $|\Phi_1\rangle$ and $|\Phi_2\rangle$ the outcomes are

$$|0\rangle|0\rangle, |0\rangle|1\rangle, |1\rangle|0\rangle, |1\rangle|1\rangle$$

or

$$|00\rangle, |01\rangle, |10\rangle, |11\rangle$$

Two qubit states: All measurements are with respect to $\{|0\rangle, |1\rangle\}$

$$(|\Phi_1\rangle, |\Phi_1\rangle) = (a|0\rangle + b|1\rangle, c|0\rangle + d|1\rangle)$$

- Probability of observing $|0\rangle |0\rangle$ is $= |ac|^2$
- Probability of observing $|0\rangle |1\rangle$ is $= |ad|^2$
- Probability of observing $|1\rangle |0\rangle$ is $= |bc|^2$
- Probability of observing $|1\rangle |1\rangle$ is $= |bd|^2$

Two qubit states:

All measurements are with respect to $\{|0\rangle, |1\rangle\}$

$$(|\Phi_1\rangle, |\Phi_1\rangle) = (a|0\rangle + b|1\rangle, c|0\rangle + d|1\rangle)$$

- Probability of observing $|0\rangle |0\rangle$ is $= |ac|^2$ $\cdot \begin{pmatrix} a \\ b \end{pmatrix} \otimes \begin{pmatrix} c \\ d \end{pmatrix} = \begin{pmatrix} a \begin{pmatrix} c \\ d \end{pmatrix} \\ b \begin{pmatrix} c \\ d \end{pmatrix} \end{pmatrix} = \begin{pmatrix} ac \\ ad \\ bc \\ bd \end{pmatrix}$
- Probability of observing $|0\rangle |1\rangle$ is $= |ad|^2$
- Probability of observing $|1\rangle |0\rangle$ is $= |bc|^2$
- Probability of observing $|1\rangle |1\rangle$ is $= |bd|^2$ $\cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ 0 \begin{pmatrix} 1 \\ 0 \end{pmatrix} \end{pmatrix} = \begin{pmatrix} 1 \times 1 \\ 1 \times 0 \\ 0 \times 1 \\ 0 \times 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}$
- $\begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = |0\rangle \otimes |0\rangle = |0\rangle|0\rangle = |00\rangle$

Two qubit states:

All measurements are with respect to $\{|0\rangle, |1\rangle\}$

$$\cdot \begin{pmatrix} a \\ b \end{pmatrix} \otimes \begin{pmatrix} c \\ d \end{pmatrix} = \begin{pmatrix} a \begin{pmatrix} c \\ d \end{pmatrix} \\ b \begin{pmatrix} c \\ d \end{pmatrix} \end{pmatrix} = \begin{pmatrix} ac \\ ad \\ bc \\ bd \end{pmatrix} = |\Phi\rangle \otimes |\Psi\rangle$$

$$\cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\ 0 \begin{pmatrix} 0 \\ 1 \end{pmatrix} \end{pmatrix} = \begin{pmatrix} 1 \times 0 \\ 1 \times 1 \\ 0 \times 0 \\ 0 \times 1 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}$$

$$\cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix} = |0\rangle \otimes |1\rangle = |0\rangle|1\rangle = |01\rangle$$

$$\cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ 1 \begin{pmatrix} 1 \\ 0 \end{pmatrix} \end{pmatrix} = \begin{pmatrix} 0 \times 1 \\ 0 \times 0 \\ 1 \times 1 \\ 1 \times 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

$$\cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = |1\rangle \otimes |0\rangle = |1\rangle|0\rangle = |10\rangle$$

$$\cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\ 1 \begin{pmatrix} 0 \\ 1 \end{pmatrix} \end{pmatrix} = \begin{pmatrix} 0 \times 0 \\ 0 \times 1 \\ 1 \times 0 \\ 1 \times 1 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix} = |1\rangle \otimes |1\rangle = |1\rangle|1\rangle = |11\rangle$$

Two-qubit states

- $|\Phi\rangle|\Psi\rangle = ac|0\rangle \otimes |0\rangle + ad|0\rangle \otimes |1\rangle + bc|1\rangle \otimes |0\rangle + bd|1\rangle \otimes |1\rangle$
 $= ac|00\rangle + ad|01\rangle + bc|01\rangle + bd|11\rangle$
- $|ac|^2 + |ad|^2 + |bc|^2 + |bd|^2$
 $= |a|^2|c|^2 + |a|^2|d|^2 + |b|^2|c|^2 + |b|^2|d|^2$
 $= |a|^2(|c|^2 + |d|^2) + |b|^2(|c|^2 + |d|^2) = (|a|^2 + |b|^2)(|c|^2 + |d|^2)$
 $= 1 \times 1 = 1$

Two-qubit states

- $|\Psi\rangle = a_{00}|0\rangle \otimes |0\rangle + a_{01}|0\rangle \otimes |1\rangle + a_{10}|1\rangle \otimes |0\rangle + a_{11}|1\rangle \otimes |1\rangle$
 $= a_{00} |00\rangle + a_{01} |01\rangle + a_{10} |01\rangle + a_{11} |11\rangle$

where $|a_{00}|^2 + |a_{01}|^2 + |a_{10}|^2 + |a_{11}|^2 = 1$

- Any vector of the above type is a two-qubit state.
- All such vectors are not (tensor) products of single-qubit states.

Entangled states

- Consider the state

$$|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

$$(a|0\rangle + b|1\rangle)(c|0\rangle + d|1\rangle)$$

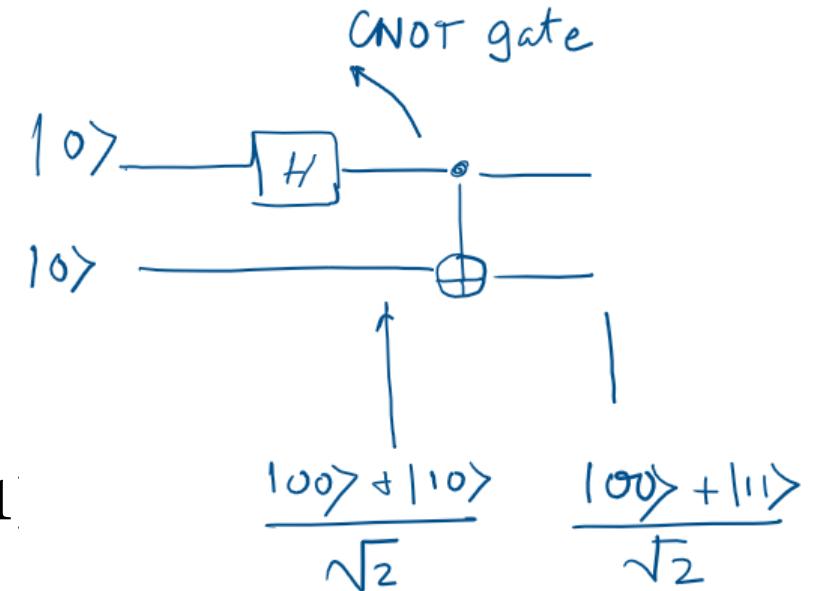
$$= ac|0\rangle \otimes |0\rangle + ad|0\rangle \otimes |1\rangle + bc|1\rangle \otimes |0\rangle + bd|1\rangle \otimes |1\rangle$$

$$= ac|00\rangle + ad|01\rangle + bc|01\rangle + bd|11\rangle$$

- $ac = \frac{1}{\sqrt{2}}, ad = 0, bc = 0, bd = \frac{1}{\sqrt{2}}$

- $ad = 0 \Rightarrow a = 0$ or $d = 0$. Both options lead to a contradiction.

- Therefore, the quantum state $|\Phi^+\rangle$ cannot be written as a tensor product of two single-qubit states.



Multiple qubit states

- An n -qubit state is

$$|\Psi\rangle = a_0|0\rangle + a_1|1\rangle + a_2|2\rangle + \cdots + a_{2^n-1}|2^n - 1\rangle$$

where $|a_0|^2 + |a_1|^2 + \cdots + |a_{2^n-1}|^2 = 1$.

- For any number, m , between $0 \leq m \leq 2^n - 1$, its binary representation is denoted by **m**.

Multiple qubit states

- An n -qubit state is

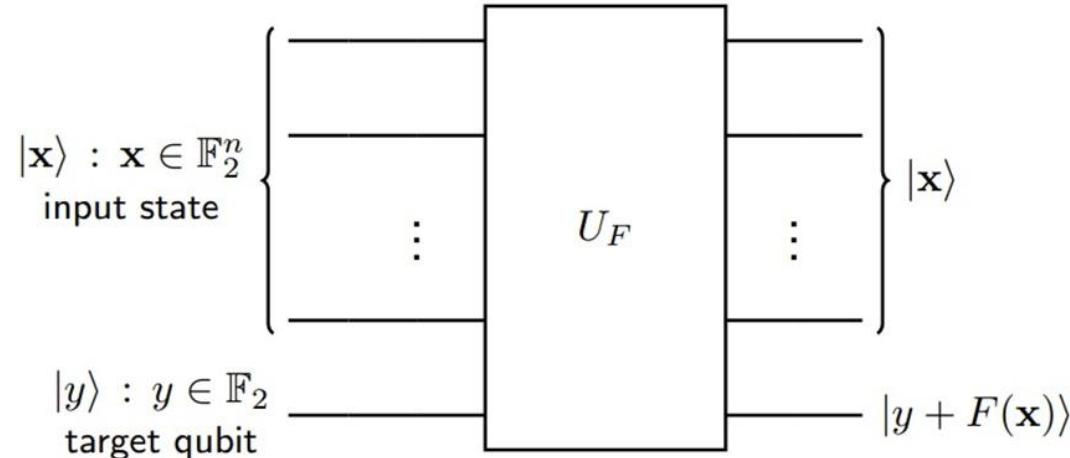
$$\begin{aligned} |\Psi\rangle = & a_0|000\rangle + a_1|001\rangle + a_2|010\rangle + a_3|011\rangle \\ & + a_4|100\rangle + a_5|101\rangle + a_6|110\rangle + a_7|111\rangle \end{aligned}$$

where

$$|a_0|^2 + |a_1|^2 + |a_2|^2 + |a_3|^2 + |a_4|^2 + |a_5|^2 + |a_6|^2 + |a_7|^2 = 1.$$

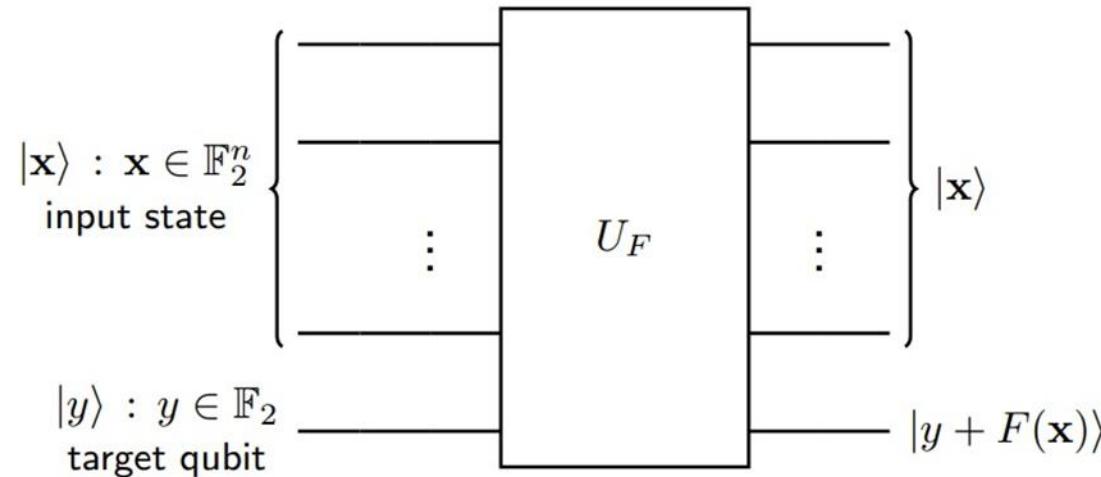
Quantum implementation of Boolean functions

- A Boolean function in n variables is a mapping from $\{0,1\}^n$ to $\{0, 1\}$.
- Suppose f is an n -variable Boolean function.
- On a quantum computer f is implemented as a transformation U_f as follows: ($x_i, y \in \{0, 1\}$ for all $i \in [n]$)

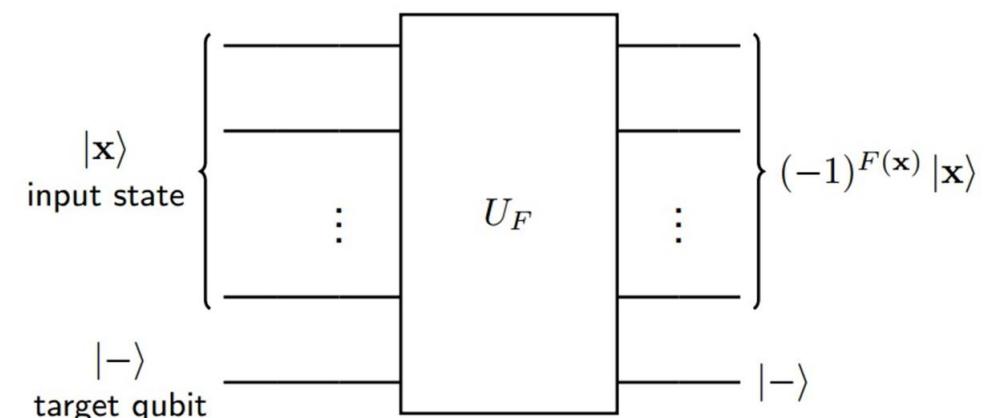
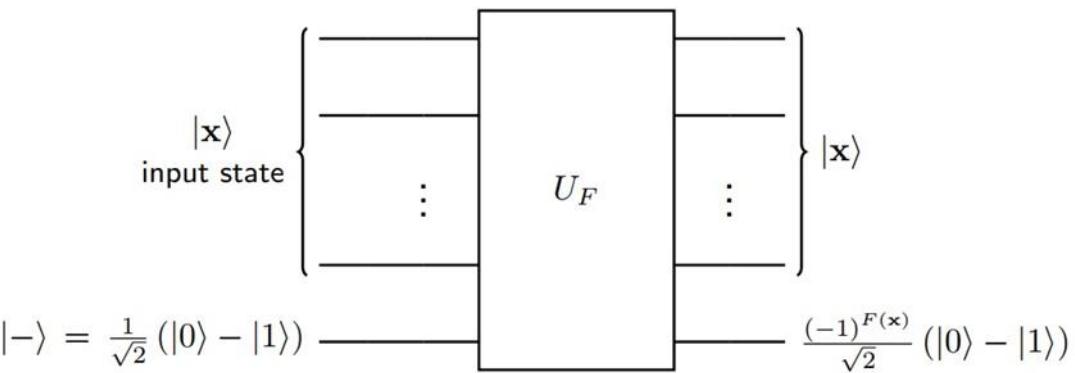
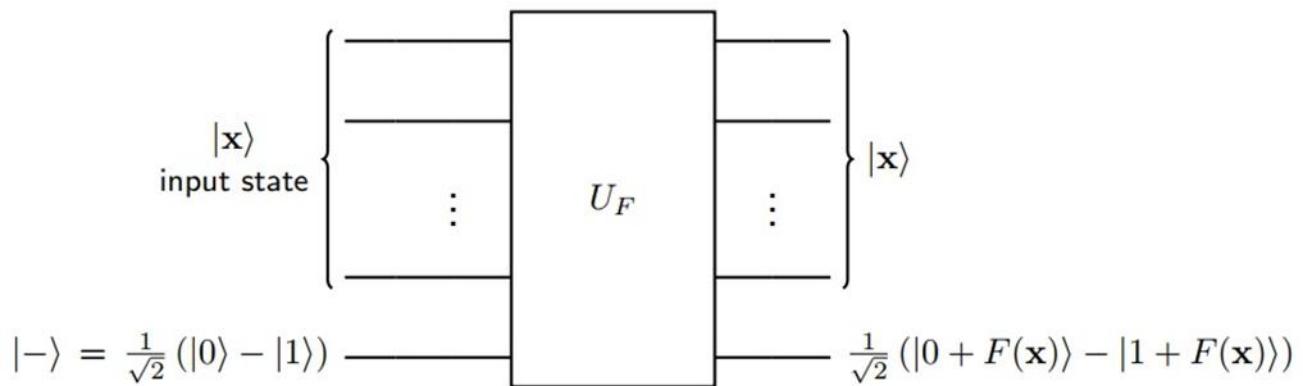


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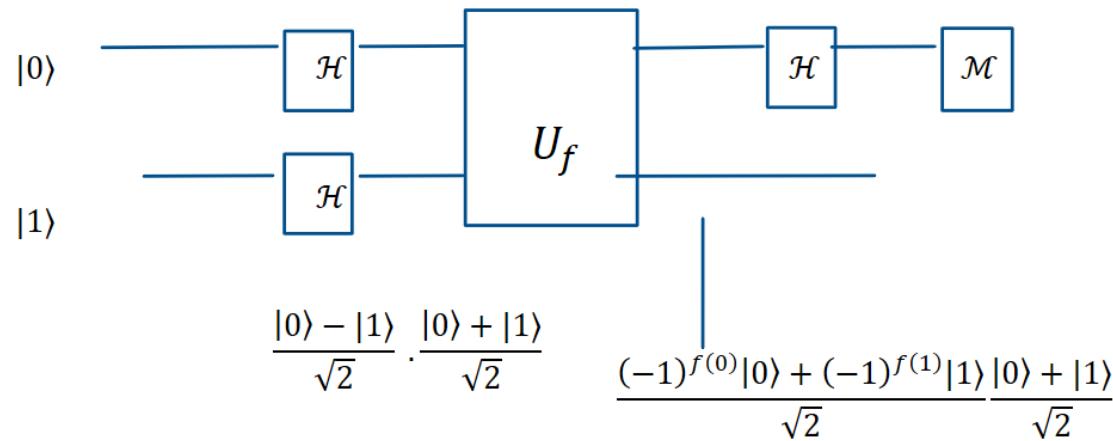
Bit Oracle to Phase Oracle



Deutsch Algorithm

- Consider 1-variable Boolean functions

- $f_0(0) = 0, f_0(1) = 0$
- $f_1(0) = 0, f_1(1) = 1$
- $f_2(0) = 1, f_2(1) = 0$
- $f_3(0) = 1, f_3(1) = 1$

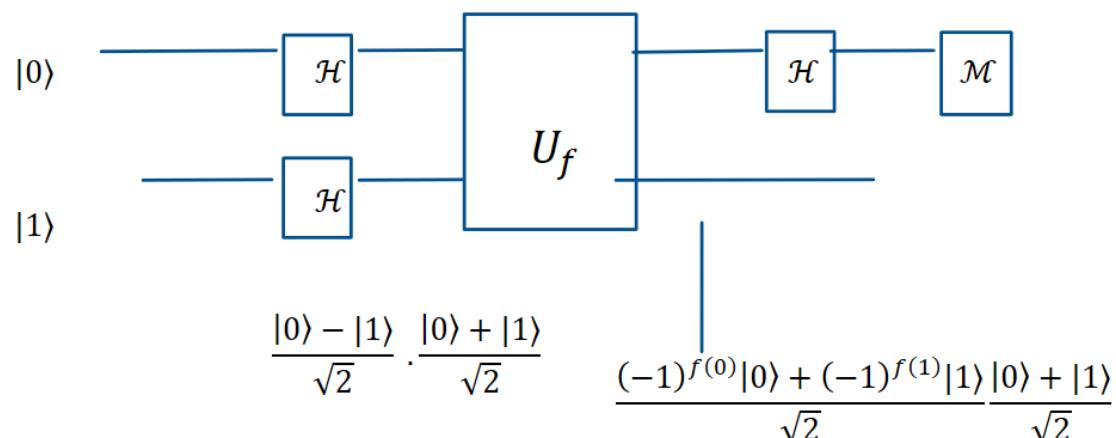


Deutsch Algorithm

- After the final Hadamard transformation we have

$$(-1)^{f(0)}|0\rangle + (-1)^{f(1)}|1\rangle \rightarrow (-1)^{f(0)} \frac{|0\rangle + |1\rangle}{\sqrt{2}} + (-1)^{f(1)} \frac{|0\rangle - |1\rangle}{\sqrt{2}}$$

$$= \frac{(-1)^{f(0)} + (-1)^{f(1)}}{\sqrt{2}} |0\rangle + \frac{(-1)^{f(0)} - (-1)^{f(1)}}{\sqrt{2}} |1\rangle$$



Deutsch-Jozsa Algorithm

- Let $\mathbf{x} = x_1 \cdots x_n \in \{0, 1\}^n$
- $|x_i\rangle \xrightarrow[\sqrt{2}]{H} \frac{|0\rangle + (-1)^{x_i}|1\rangle}{\sqrt{2}}$
- $|\mathbf{x}\rangle \xrightarrow{H^{\otimes n}} 2^{-n/2} \sum_{\mathbf{y} \in \{0,1\}^n} (-1)^{\mathbf{x} \cdot \mathbf{y}} |\mathbf{y}\rangle$
- $|\mathbf{0}_n\rangle \xrightarrow{H^{\otimes n}} 2^{-n/2} \sum_{\mathbf{x} \in \{0,1\}^n} |\mathbf{x}\rangle \xrightarrow{U_f} 2^{-n/2} \sum_{\mathbf{x} \in \{0,1\}^n} (-1)^{f(\mathbf{x})} |\mathbf{x}\rangle$
- $2^{-n/2} \sum_{\mathbf{x} \in \{0,1\}^n} (-1)^{f(\mathbf{x})} |\mathbf{x}\rangle \xrightarrow{H^{\otimes n}} 2^{-n} \sum_{\mathbf{x} \in \{0,1\}^n} \sum_{\mathbf{y} \in \{0,1\}^n} (-1)^{f(\mathbf{x}) + \mathbf{x} \cdot \mathbf{y}} |\mathbf{y}\rangle$

Deutsch-Jozsa Algorithm

- $|\psi\rangle = 2^{-n} \sum_{\mathbf{x} \in \{0,1\}^n} \sum_{\mathbf{y} \in \{0,1\}^n} (-1)^{f(\mathbf{x}) + \mathbf{x} \cdot \mathbf{y}} |\mathbf{y}\rangle$
 $= \sum_{\mathbf{y} \in \{0,1\}^n} (2^{-n} \sum_{\mathbf{x} \in \{0,1\}^n} (-1)^{f(\mathbf{x}) + \mathbf{x} \cdot \mathbf{y}}) |\mathbf{y}\rangle$
- Suppose we measure $|\psi\rangle$ using the computational basis.
- The state $|0_n\rangle$ appears with probability $2^{-n} \left| \sum_{\mathbf{x} \in \{0,1\}^n} (-1)^{f(\mathbf{x})} \right|^2$.
 - If f is balanced $2^{-n} \left| \sum_{\mathbf{x} \in \{0,1\}^n} (-1)^{f(\mathbf{x})} \right|^2 = 0$. So $|0_n\rangle$ will never appear.
 - If f is a constant $2^{-n} \left| \sum_{\mathbf{x} \in \{0,1\}^n} (-1)^{f(\mathbf{x})} \right|^2 = 1$. So $|0_n\rangle$ will always be the result of the measurement.

Thank You

Questions Please!?