1 Abstract

- Derive stability condition on FTCS solution of heat equation;
- Give Von-Neumann condition. (no proof)

2 Problem

We already observed the danger of FTCS scheme for the heat equation. Find stability condition of FTCS solution for the heat equation

$$u_t = u_{xx}, \quad t > 0, x \in \mathbb{R}$$

with initial data

$$u(x,0) = \phi(x), \quad x \in \mathbb{R}.$$

3 Analysis

3.1 Setup

First, we recall FTCS scheme. FTCS means that we use finite difference form of

$$u_t(x,t) \simeq \frac{u(x,t+\theta) - u(x,t)}{\theta} := \delta_{\theta}^t u(x,t)$$

and

$$u_{xx}(x,t) \simeq \frac{u(x+h,t) - 2u(x,t) + u(x-h,t)}{h^2} := \delta_h^{xx} u(x,t).$$

where h and θ are some positive mesh size in space h and in time, respectively.

Discrete domain is accordingly a grid of

$$\{(jh,n\theta): j+1\in\mathbb{N}, j\in\mathbb{Z}\}.$$

We denote by u_j^n is the FTCS solution at a grid point $(jh, n\theta)$, then we shall have

$$u_t(jh, n\theta) \simeq \frac{u_j^{n+1} - u_j^n}{\theta}, \quad u_{xx}(jh, n\theta) \simeq \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{h^2}.$$

Plug it into the heat equation, we obtain discrete heat equation

$$u_j^{n+1} = su_{j+1}^n + (1-2s)u_j^n + su_{j-1}^n, \quad \forall j \in \mathbb{Z}, n+1 \in \mathbb{N}.$$
 (1)

with the initial condition

$$u_j^0 = \phi(jh), \quad \forall j \in \mathbb{Z}.$$
 (2)

where

$$s = \frac{\theta}{h^2}.$$

By the FTCS solution of heat equation, we mean

$$\{u_i^n: \forall j \in \mathbb{Z}, n \in \mathbb{N}\}$$

satisfying equations (2) - (1).

3.2 Solution

In this below, we solve for FTCS solution in two steps. First, we use the technic of the separation in variable to find all possible solutions satisfying (1). Second step is to choose specific solution by fitting the initial condition (2).

We first search for the solution of (1) given by the product of j-function and n-function:

$$u_j^n = X_j T_n$$
.

Of course, any linear combination of such solutions shall give another solution of (1).

Plug above form into (1), it writes

$$\frac{T_{n+1}}{T_n} = 1 - 2s + s \frac{X_{j+1} - X_j}{X_j}.$$

Note that, left hand side is a function of n while right hand side is a function of j for all (n, j). Thus, to be equal, they must be equal to a constant, say ξ , i.e.

$$\frac{T_{n+1}}{T_n} = 1 - 2s + s \frac{X_{j+1} - X_j}{X_j} := \xi.$$

Therefore, we have

$$T_n = \xi^n T_0 = \xi^n, \forall n \in \mathbb{N}$$

and

$$1 - 2s + s \frac{X_{j+1} - X_j}{X_j} := \xi, \forall j \in \mathbb{Z}.$$

In the above, $T_0 = 1$ is assumed w.l.o.g. (why?) Another trick is to postulate X in the form of

$$X_i = (e^{ikh})^j$$

for some $k \in \mathbb{Z}$. The reason is that, we expect the solution X is in L^2 , and our solution could be linear combination of the above Fourier basis functions.

Then, we can solve for ξ by plugging into X-equation above, that

$$\xi(k) := \xi = 1 - 2s + 2s \cos kh j$$
.

At last, we have general representation of the numerical solution:

$$u_j^n = \sum_{k \in \mathbb{Z}} b_k e^{ikhj} \xi^n(k).$$

The rest is to determine coefficients $(b_k : k \in \mathbb{Z})$ from the initial condition using orthogonal basis functions.

3.3 Stability condition

From the above solution representation, we shall have

if $|\xi(k)| \leq 1$ for all $k \in \mathbb{Z}$, then FTCS is stable.

Since s > 0, the inequality $\xi(k) \le 1$ holds automatically. The lower bound of $\xi(k)$ is 1 - 4s. So we shall require $-1 \le 1 - 4s$ for its stability, which at last gives the **sufficient condition of the stability** by

$$s \le 1/2$$
.

The example discussed above indicates that the general procedure to determine stability in a diffusion or wave problem is to separate variables in the difference equation. For the time factor we obtain a simple equation with an amplication factor $\xi(k)$. In the analysis above, we use $|\xi(k)| \leq 1$ for the stability. More precisely, it can be shown that the correct condition necessary for stability is

$$|\xi(k)| \le 1 + O(\theta), \forall k \in \mathbb{Z}.$$

This is the von Neumann stability condition.