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Energy in units of
 Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$\begin{aligned}
 V(x) &= \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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Energy in units of V_0

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

$E < V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
 $-a \leq x \leq a$

Figure 4

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

$$\text{inside}$$

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n = \frac{\hbar^2 k_n^2}{2m}$$

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$

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, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega$

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$2mE/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar \text{ outside.}$$

$$2m(E - V_0)/\hbar$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

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$$V_0 \quad (|x| > a).$$

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$$E = 0$$

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$$-a \leq x \leq a$$

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$$k = \sqrt{2m(E - V_0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto n^2$$

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E = 0$

$V = V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

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$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V(x))/\hbar^2} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar^2} \text{ outside.}$$

$$2m(E - V_0)/\hbar^2$$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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by

$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

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$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 > 0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$$-V_0 < E < 0$$

$$E > 0$$

Figure 4

• $E > 0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

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$$\text{to } \psi(0) = 1$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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Energy in units of

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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Energy in units of V_0

Figure 3

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Energy in units of

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$$E_n \propto n^2$$

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $\sqrt{2m(E - V(x))} / \hbar$. This equals $\sqrt{2m(E - V_0)} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

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Energy in units of V_0

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Figure 5

Energy in units of V_0

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

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-a a

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

0 ($|x| \leq a$)

V_0 ($|x| > a$).

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Energy in units of E_1

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

by

$$V(x) = \begin{cases} \frac{1}{2} \\ 0 \end{cases}$$

$$0 \quad (|x| \leq a)$$

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

$$\text{inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter}$$

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$E < V_0$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar.$$

This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

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to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

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$-a \ a$

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Energy in units of

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Energy in units of

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are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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0 ($|x| \leq a$)

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E

Energy in units of

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inside
and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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E
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Figure 3
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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$
 $E < V_0$

-a a
 Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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$$p = \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } \sqrt{2mE}/\hbar \text{ outside.}$$

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$$k \text{ is larger inside the well, so the wavelength is shorter there.}$$

A possible wavefunction might look something like the one in Fig. 5. It is customary to

to $-a \quad a$

E

Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string

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$$n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}.$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$$V=V_0$$

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

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Energy in units of V_0

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$$V_0 \quad (|x| > a).$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

$E < V_0$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n

ω

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$E < V_0$

$E = V_0$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of

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E

Energy in units of E_1

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by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2mE}/\hbar \text{ outside.}$$

$$2mE/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar \text{ outside.}$$

$$2m(E - V_0)/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

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$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

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$$E$$

Energy in units of

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$$V_0 \quad (|x| \leq a).$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of $p = \hbar k$. This equals $\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter

there. A

possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \leq x \leq a$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

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$$V_0 (|x| \leq a)$$

$$0 (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar$$

$$2m(E - V_0)/\hbar^2$$

$$2m(E - 0)/\hbar^2$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

n^2

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$$E > -V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > -V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E + V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

$$\text{inside}$$

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto n^2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

$\psi(x)$ is given

by

$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} \frac{1}{2} \\ 0 \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE/\hbar} \text{ inside the well and } p$$

$$2m(E - V_0)/\hbar$$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$
 $E < V_0$

-a a
 Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

$2m(E - V_0)/\hbar$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))$. This equals $\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter

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$\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

$E < V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of

Figure 3

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Figure 3

E

Energy in units of $\hbar^2/(2ma^2)$

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
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Energy in units of

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Energy in units of V_0

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Energy in units of E

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of V_0

Figure 3

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Energy in units of

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$$p = \sqrt{2m(E - V_0)}/\hbar$$

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Energy in units of V_0

Figure 3

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E

Energy in units of

Figure 3

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Energy in units of

Figure 3

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Energy in units of

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$\begin{aligned}
 V(x) &= \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \\
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 \end{aligned}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$E = V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of
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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

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by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

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-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

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apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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Energy in units of

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E

Energy in units of E_1

Figure 3

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$E_n \propto n^2$$

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, because although $\omega = ck \propto n$, the energy is proportional to $\omega^2(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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E

Energy in units of

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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-a a

Figure 4

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Energy in units of

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$\begin{aligned}
 V(x) &= \frac{1}{2} \\
 0 & \quad (|x| \leq a) \\
 V_0 & \quad (|x| > a).
 \end{aligned}
 \tag{19}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

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Energy in units of

Figure 3

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Energy in units of

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Energy in units of

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$2m(E - V_0)/\hbar$

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & 0 \leq |x| \leq a \\ 0 & |x| > a \end{cases}$$

$V_0 (|x| > a)$.

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \quad a$$

$$E$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Energy in units of
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E

Energy in units of V_0

Figure 3

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Energy in units of

Figure 3

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Energy in units of

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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10.3.3 Finite square well

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$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 > 0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$$V_0 = 1$$

Energy in units of E

Figure 3

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$$E_n \propto n^2$$

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Energy in units of

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \quad a$$

$$E$$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

$$E,$$

2

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

2

That wouldn't be much of a state, because the probability would be zero everywhere.

The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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Energy in units of

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$$V_0 \quad (|x| > a).$$

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$$E$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

$$n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well.

k is larger inside the well, so the wavelength is shorter there.

A possible wavefunction might look something like the one in Fig. 5. It is customary to

normalize $\psi(x)$ to 1.

Figure 3

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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E

Energy in units of

Figure 3

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Energy in units of

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(unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))} / \hbar$. This equals $\sqrt{2m(E - V_0)} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

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E
Energy in units of

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E

Energy in units of V_0

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string
are also proportional to n^2
, because although $\omega = ck \propto n$, the energy is proportional to ω^2
(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both
apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

E,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E:

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

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$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{inside the well and } p = \sqrt{2mE}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

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Energy in units of

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Energy in units of

Figure 3

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are also proportional to n

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

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E

Energy in units of

Figure 3

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$2m(E - V_0)/\hbar$

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
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 &2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar \\
 &\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a
 \end{aligned}$$

Energy in units of
 Figure 3
 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| \leq a).$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

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$$E$$

Energy in units of V_0

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Energy in units of V_0

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are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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Energy in units of

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Energy in units of

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E

Energy in units of

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A possible wavefunction might look something like the one in Fig. 5. It is customary to

plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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10.3.3 Finite square well

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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 $E < V_0$

-a a
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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of V_0

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Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside}$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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Energy in units of

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Energy in units of V_0

Figure 3

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Energy in units of

Figure 3

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$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$
 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
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$$\begin{aligned}
 V(x) &= \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of

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$$E$$

Energy in units of

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to $-a \leq x \leq a$

Energy in units of V_0

Figure 3

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Energy in units of V_0

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Energy in units of

Figure 3

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Energy in units of V_0

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to -a a

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E

Energy in units of

Figure 3

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$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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That wouldn't be much of a state, because the probability would be zero everywhere.

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}.$$

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Figure 5

Energy in units of V_0

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Energy in units of

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$\sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } p$$

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Energy in units of
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$2m(E - V_0)/\hbar$

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E

Energy in units of E_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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E

Energy in units of

Figure 3

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are also proportional to n

2

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Energy in units of
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 $V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E = 0$
 $E = V_0$
-a a

Figure 4

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Energy in units of
 Figure 3
 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$\begin{aligned}
 V(x) &= \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

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 $V=V_0$
 -a a
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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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Energy in units of

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$$E$$

$$E$$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n$$

$$E_n$$

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}.$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2 . This equals $\sqrt{2mE}/\hbar$ inside the well and p

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10.3.3 Finite square well

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}.$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

ω^2

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Energy in units of

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$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of

Figure 3

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2

2

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E

Energy in units of

Figure 3

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A possible wavefunction might look something like the one in Fig. 5. It is customary to

plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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E
Energy in units of

Figure 3
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Energy in units of
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- $V = V_0$

-a a

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω

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$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside}$$

$$\text{and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A}$$

$$\text{possible wavefunction might look something like the one in Fig. 5. It is customary to -a a}$$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{}$

$$2mE/\hbar \text{ inside the well and } p$$

$$2m(E - V_0)/\hbar$$

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Energy in units of

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω

$$2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

E,

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$$V_0 \quad (|x| > a).$$

(19)

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$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

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$$k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well.}$$

$$k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

to $-a \leq x \leq a$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar \text{ outside.}$$

$$k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well.}$$

$$k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

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$0 \quad (|x| \leq a)$
 $V_0 \quad (|x| > a).$
 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

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$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

$2m(E - V_0)/\hbar$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p
 $\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a
 E
 Energy in units of
 Figure 3
 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & 0 \leq |x| \leq a \\ 0 & |x| > a \end{cases}$$

V_0 ($|x| > a$).

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

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$V(x) = \frac{1}{2}$

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of
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Energy in units of

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Energy in units of E
 Figure 3

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$E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
 Figure 3
 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$\begin{aligned}
 V(x) &= \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
 $-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$\begin{aligned}
 &2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar \\
 &\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a
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$$V_0 \quad (|x| \leq a).$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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$E < V_0$

$E = V_0$

$E < 0$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$p = \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar \text{ outside.}$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n^2

$E_n \propto n^2$

E_n

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of V_0

Figure 3

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll
let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given
by

$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$
(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p
 $\sqrt{2m(E - V(x))/\hbar}$. This equals $\sqrt{2m(E - V_0)/\hbar}$ inside the well and $\sqrt{2m(E - V(x))/\hbar}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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 $E < V_0$

-a a
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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$

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Energy in units of
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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

$$p = \sqrt{2m(E - V_0)}/\hbar$$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega = \hbar ck = \hbar c k$. This equals $\sqrt{2mE}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$0 \quad (|x| \leq a)$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

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$$-a \leq x \leq a$$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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E

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n = \frac{1}{2} \hbar \omega_n$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$-a \quad a$$

Figure 4

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$$E$$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy

E:

$V=0$

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well.

k is larger inside the well, so the wavelength is shorter there.

A possible wavefunction might look something like the one in Fig. 5. It is customary to

normalize $\psi(x)$ to 1.

Figure 3

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

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$$2m(E - V_0)/\hbar$$

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E

Energy in units of

Figure 3

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega \hbar (E - V(x)) / \hbar$. This equals $\sqrt{2m(E - V(x))} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $\sqrt{2m(E - V(x))} / \hbar$. This equals $\sqrt{2m(E - V_0)} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

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-a a

E

Energy in units of

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$E > V_0$

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Figure 5

Energy in units of V_0

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$E < V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and $\sqrt{2m(E - V_0)/\hbar}$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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Figure 5

Energy in units of V_0

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Energy in units of E_1

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Energy in units of
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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$
 $V=V_0$
 -a a
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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a). \end{cases}$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$V_0 \quad (|x| \leq a).$$

$$(19)$$

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Energy in units of V_0

Figure 3

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Energy in units of

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(19)

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Figure 4

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$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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Energy in units of

Figure 3

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Energy in units of V_0

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}.$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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normalize ψ to 1.

Energy in units of E_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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Energy in units of E_0

Figure 3

$$2m(E - V_0)/\hbar$$

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Energy in units of

Figure 3

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Energy in units of

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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E
Energy in units of

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$$V(x) = \frac{1}{2}$$

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by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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$$E$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$E$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

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$2m(E - V_0)/\hbar$

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E ,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
 Figure 3
 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$\begin{aligned}
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 (19)
 \end{aligned}$$

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Energy in units of

Figure 3

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Energy in units of

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Energy in units of

Figure 3

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

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Energy in units of V_0

Figure 3

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$E = 0$

$E = V_0$

$-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere.

The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

$\psi(x)$ is given

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$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$$E > -V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > -V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of E_0

Figure 3

gap

between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

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Energy in units of E_0

Figure 3

gap

between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of

Figure 3

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$$\text{inside the well and } p = \sqrt{\frac{2m(E - V_0)}{\hbar^2}}$$

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inside

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

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2

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Energy in units of

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E
Energy in units of
Figure 3

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$V=0$
 $V=V_0$

-a a
 Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$

inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))$. This equals $\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of V_0

Figure 3

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Figure 4

there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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$V = V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2

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$$E > V_0$$

$$E < V_0$$

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

$$p = \sqrt{2m(E - V_0)}/\hbar$$

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Energy in units of

Figure 3

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$$E < V_0$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

$$E_n = \frac{1}{2} \hbar \omega_n$$

Energy in units of V_0

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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 $V = V_0$
 $-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

Energy in units of V_0

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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E

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega = \hbar ck = \hbar c \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

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Energy in units of V_0

Figure 3

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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$E < V_0$

$E = V_0$

$E < 0$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

ω^2

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Energy in units of

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Energy in units of

Figure 3

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Energy in units of E_0

Figure 3

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$$E$$

Energy in units of V_0

Figure 3

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Energy in units of V_0

Figure 3

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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$$\begin{aligned}
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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2mE}/\hbar \text{ outside.}$$

$$p = \sqrt{2m(E - V_0)}/\hbar$$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 . This equals $\sqrt{2m(E - V(x))}/\hbar$. This

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a$ a

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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E

Energy in units of

Figure 3

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2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

$$\text{inside}$$

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n = \frac{1}{2} \mu \omega_n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$V=V_0$

$-a \quad a$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Figure 4

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$$V_0 \quad (|x| > a).$$

(19)

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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E ,

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of

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Energy in units of E_0

Figure 3

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$$k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well.}$$

k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to

-a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

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$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Figure 5

Energy in units of V_0

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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10.3.3 Finite square well

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
 -a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2

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$$2mE/\hbar^2 \text{ inside the well and } 2m(E - V_0)/\hbar^2 \text{ outside.}$$

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Energy in units of

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Energy in units of V_0

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Energy in units of

Figure 3

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$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

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Energy in units of

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Energy in units of V_0

Figure 3

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Energy in units of V_0

Figure 3

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary

to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future

convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy

E:

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E
Energy in units of
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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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E
Energy in units of
Figure 3
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, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V(x))/\hbar}$ inside the well and $\sqrt{2m(E - V_0)/\hbar}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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 $V = V_0$
-a a

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$V_0 \quad (|x| > a).$$

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$E > V_0$

$E < V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside.

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

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to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$0 \quad (|x| \leq a)$$

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

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Figure 4

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$

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Energy in units of

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E

Energy in units of

Figure 3

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Figure 3

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ V_0 & (|x| > a). \end{cases} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy

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$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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Energy in units of

Figure 3
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Energy in units of V_0

Figure 3

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Energy in units of

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of V_0

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Energy in units of V_0

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p$$

$$2m(E - V_0)/\hbar$$

$$2m(E - V_0)/\hbar$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

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E ,

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$2m(E - V_0)/\hbar$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
 Figure 3
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Figure 4

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 $\frac{2m(E - V(x))}{\hbar}$. This equals $\sqrt{\frac{2mE}{\hbar}}$ inside the well and $\sqrt{\frac{2m(E - V_0)}{\hbar}}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$-a \leq x \leq a$$

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$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

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E

Energy in units of

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Figure 4

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$$2m(E - V_0)/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of
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$V = 0$
 $V = V_0$
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Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$\begin{aligned}
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Energy in units of
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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of

Figure 3

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are also proportional to n^2

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$E = 0$

$E = V_0$

$-a \leq x \leq a$

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Energy in units of

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That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2mE}/\hbar \text{ outside.}$$

$$k = \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar \text{ outside.}$$

$$k = \sqrt{2m(E - V_0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto n^2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E ,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E = 0$

$V = V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of p

$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V(x))/\hbar^2} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar^2} \text{ outside.}$$

$$2m(E - V_0)/\hbar^2$$

$$2m(E - V_0)/\hbar^2$$

$$2m(E - V_0)/\hbar^2$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

\sqrt{E} ,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section

10.4 below),

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by

$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E = 0$

$V = V_0$

$-a \leq x \leq a$

Figure 4

$V=0$

$V=V_0$

$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E ,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \ (|x| \leq a)$

$V_0 \ (|x| > a)$.

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside.

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future

convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

(x) is given

by

$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$

V_0 ($|x| \leq a$).

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy

E:

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

inside
and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case

of a string
are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

$-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by
 $V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$
 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$
 $E < V_0$

-a a
 Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

$$p = \sqrt{2m(E - V_0)}/\hbar$$

$$p = \sqrt{2m(E - 0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V = 0$$

$$V = V_0$$

$$-a \leq x \leq a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V(x))/\hbar^2} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar^2} \text{ outside.}$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$n^2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E ,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } p = \sqrt{2mE}/\hbar \text{ outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

$$p = \sqrt{2mE}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar \text{ outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

$$p = \sqrt{2m(E - V_0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2 . This equals $\sqrt{2m(E - V(x))}/\hbar$. This

$$p = \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p = \sqrt{2mE}/\hbar \text{ outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

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Energy in units of

Figure 3

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a). \end{cases} \quad (19)$$

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Figure 4

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Energy in units of

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Energy in units of

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Energy in units of E_0

Figure 3

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Energy in units of E_0

Figure 3

$$2m(E - V_0)/\hbar$$

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E

Energy in units of

Figure 3

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E

Energy in units of V_0

Figure 3

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E

Energy in units of

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of V_0

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gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE/\hbar} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar} \text{ outside.}$$

$$2m(E - V_0)/\hbar$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

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$E = V_0$

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$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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$$\text{to } -a \leq x \leq a$$

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Energy in units of

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$$V_0 \quad (|x| \leq a).$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of $p = \hbar k$. This equals $\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside.

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E = 0$

$E = V_0$

$-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$V_0 > 0$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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Energy in units of

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by

$$V(x) = \begin{cases} \frac{1}{2} \\ 0 \end{cases}$$

$$0 \quad (|x| \leq a)$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of

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Energy in units of

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$2m(E - V_0)/\hbar$

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Energy in units of

Figure 3

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of

Figure 4

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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E

Energy in units of V_0

Figure 3

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plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

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$0 \ (|x| \leq a)$

$V_0 \ (|x| > a)$.

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of

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Energy in units of V_0

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$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

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(19)

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E

Energy in units of V_0

Figure 3

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Energy in units of

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Energy in units of V_0

Figure 3

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Energy in units of

Figure 3

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < 0$$

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$V=V_0$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

Figure 3

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$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

2

2

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$V(x) = \frac{1}{2}$

0 ($|x| \leq a$)

V_0 ($|x| > a$).

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E:

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega^2(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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E:

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A possible wavefunction might look something like the one in Fig. 5. It is customary to

plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of

Figure 3
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Figure 4

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Energy in units of

Figure 3

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Figure 4

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
 $-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \hbar k$
 $k = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2mE}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string
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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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That wouldn't be much of a state, because the probability would be zero everywhere. The

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Energy in units of V_0

Figure 3

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$0 \quad (|x| \leq a)$
 $V_0 \quad (|x| > a).$
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Energy in units of

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Energy in units of V_0

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$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto n^2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E ,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

$E < V_0$

$E = V_0$

$E < V_0$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar.$$

This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

$$k = \sqrt{2m(E - V_0)}/\hbar$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well.

k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$V=V_0$

$-a$ a

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2

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$V = 0$
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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and p

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$$E > V_0$$

Energy in units of V_0

Figure 3

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Energy in units of

Figure 3

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Energy in units of

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n = \frac{1}{2} \hbar \omega_n$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Figure 3

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of

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E

Energy in units of E_1

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

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to $-a \leq x \leq a$

Energy in units of E_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$2m(E - V_0)/\hbar$$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω

$$2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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E Energy in units of

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$-a \ a$

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A possible wavefunction might look something like the one in Fig. 5. It is customary to

normalize $\psi(x)$ to 1.

E

Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string

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E

Energy in units of

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Energy in units of
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$E < V_0$

$E > V_0$

$E = V_0$

Figure 4

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to $-a$ a

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

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-a a
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E
Energy in units of
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k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$V(x) = \frac{1}{2}$$

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

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Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E:

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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Energy in units of

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$E < V_0$

$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of $p = \hbar k$. This equals $\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside.

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$V=V_0$

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Energy in units of

Figure 3

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by

$V(x) = \frac{1}{2}$

$0 \ (|x| \leq a)$

$V_0 \ (|x| > a)$.

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of

Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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- $V = V_0$

-a a

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p
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Energy in units of V_0

Figure 3

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E

Energy in units of

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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E

Energy in units of

Figure 3

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Energy in units of

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$V(x) = \frac{1}{2}$

$0 \ (|x| \leq a)$

$V_0 \ (|x| > a)$.

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

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Energy in units of

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega \hbar (E - V(x)) / \hbar$. This equals $\sqrt{2m(E - V_0) / \hbar}$ inside the well and p

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$\sqrt{2mE / \hbar}$$

$$2m(E - V_0)/\hbar$$

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Energy in units of

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Energy in units of

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$$-a \leq x \leq a$$

Figure 4

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$$\text{inside the well and } p = \sqrt{2mE}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$0 \text{ for } |x| > a$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar$$

$$p = \sqrt{2mE}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

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Energy in units of

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$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of

Figure 3

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Energy in units of

Figure 3

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$$-a \quad a$$

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$$k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a}$$

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Energy in units of

Figure 3

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| > a) \\ 0 & (|x| \leq a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of E_1

Figure 3

gap

between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

Figure 3

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-a a

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E

Energy in units of V_0

Figure 3

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 Energy in units of

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Energy in units of

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E_n

Energy in units of E_1

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p
 $\frac{2m(E - V(x))}{\hbar}$. This equals $\sqrt{\frac{2mE}{\hbar^2}}$ inside the well and $\sqrt{\frac{2m(E - V_0)}{\hbar^2}}$ outside.

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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Energy in units of V_0

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are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This

$\sqrt{2m(E - V_0)}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of E
 Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$\begin{aligned}
 V(x) &= \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \\
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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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Energy in units of V_0

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Energy in units of E_0

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V(x))}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

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Energy in units of

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

$$\text{inside}$$

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

E_n

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

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Energy in units of

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E_n

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$$V=0$$

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Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE/\hbar} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar} \text{ outside.}$$

$$2m(E - V_0)/\hbar$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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10.3.3 Finite square well

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \quad (19)$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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Energy in units of V_0

Figure 3

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Energy in units of

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E

Energy in units of V_0

Figure 3

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E

Energy in units of E_1

Figure 3

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Energy in units of

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E:

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

$$2m(E - V_0)/\hbar$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega \hbar (E - V(x)) / \hbar$. This equals $\sqrt{2m(E - V(x))} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

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-a a

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-a a

E

Energy in units of

Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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-a a

E

Energy in units of

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$$-a \leq x \leq a$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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√

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

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-a a

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E

Energy in units of

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$n^2$$

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Energy in units of

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

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Energy in units of

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Figure 5

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Energy in units of

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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2

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \frac{1}{2}$$

$0 \quad (|x| \leq a)$
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 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$V=V_0$

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a). \end{cases}$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ outside the well and } \sqrt{2m(E - 0)}/\hbar \text{ inside the well.}$$

$$k = \sqrt{2m(E - V_0)}/\hbar \text{ outside the well and } \sqrt{2m(E - 0)}/\hbar \text{ inside the well.}$$

$$k = \sqrt{2m(E - V_0)}/\hbar \text{ outside the well and } \sqrt{2m(E - 0)}/\hbar \text{ inside the well.}$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

$\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

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2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E ,

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$$V_0 \quad (|x| > a).$$

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of

$\sqrt{2m(E - V(x))/\hbar}$. This equals $\sqrt{2mE/\hbar}$ inside the well and

$\sqrt{2m(E - V_0)/\hbar}$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there.

A possible wavefunction might look something like the one in Fig. 5. It is customary to

plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

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where as the quantum mechanical system has $\omega \propto E$.

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$0 \ (|x| \leq a)$

$V_0 \ (|x| > a)$.

(19)

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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possible wavefunction might look something like the one in Fig. 5. It is customary

to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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convenience, we'll

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by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside.

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inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case

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$$\begin{aligned} V(x) &= \frac{1}{2} \\ 0 &(|x| \leq a) \\ V_0 &(|x| > a). \end{aligned} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

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Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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 $V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$
 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$
 $E < V_0$

-a a
 Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of
Figure 3

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$$E > V_0$$

$$-a \leq x \leq a$$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$p = \sqrt{2m(E - V_0)}/\hbar$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$\sqrt{2m(E - V_0)}/\hbar \text{ outside the well and } \sqrt{2m(E - 0)}/\hbar \text{ inside the well.}$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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n^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 . This equals $\sqrt{2m(E - V(x))}/\hbar$. This

$$p = \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } \sqrt{2mE}/\hbar \text{ outside.}$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a). \end{cases} \quad (19)$$

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Energy in units of E_0

Figure 3

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Energy in units of

Figure 3

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Energy in units of E_1

Figure 3

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of V_0

Figure 3

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-a a

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Energy in units of

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Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

$$E_n = \frac{\hbar^2 k_n^2}{2m}$$

Energy in units of V_0

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Energy in units of V_0

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Energy in units of V_0

$$E/V_0$$

Figure 3

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E

Energy in units of

Figure 3

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Energy in units of

Figure 3

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normalize ψ to 1. The energy E is in units of V_0 .

Figure 3 shows the energy levels for a string fixed at both ends. The gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 .

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 .

$E_n = \frac{1}{2} \hbar \omega_n$, because although $\omega = ck \propto n$, the energy is proportional to $\hbar \omega$.

E_n

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

consider a finite square well of width L and height V_0 . The potential is zero outside the well and V_0 inside.

The energy levels for a finite square well are shown in Fig. 4. The gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 .

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Energy in units of V_0

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \quad (19)$$

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $\sqrt{2m(E - V(x))/\hbar}$. This equals $\sqrt{2m(E - V(x))/\hbar}$ inside the well and $\sqrt{2m(E - V_0)/\hbar}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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Energy in units of V_0

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)/\hbar}$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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Energy in units of V_0

Figure 3

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Energy in units of V_0

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Energy in units of

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 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$\begin{aligned}
 V(x) &= \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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 $V = V_0$
 -a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p
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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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10.3.3 Finite square well

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

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Energy in units of

Figure 3

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$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

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$$V_0 > 0 \quad (|x| > a).$$

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$$-V_0 < E < 0$$

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$$\text{to } \psi(0) = 1$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 > 0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$$-V_0 < E < 0$$

$$E > 0$$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

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-a a

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$E > V_0$

-a a

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

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Energy in units of

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to $-a$ a

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Figure 4

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Energy in units of

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Energy in units of

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Energy in units of

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar \text{ outside.}$$

$$2m(E - V_0)/\hbar$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$\begin{aligned}
 V(x) &= \frac{1}{2} \\
 0 & \quad (|x| \leq a) \\
 V_0 & \quad (|x| > a).
 \end{aligned}
 \tag{19}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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by

$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Figure 4

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Energy in units of

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by

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$$0 \quad (|x| \leq a)$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \quad a$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V = 0$$

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$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar$$

$$\sqrt{2m(E - V_0)}/\hbar \text{ outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \quad a$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω

$$2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

E,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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Energy in units of E_0

Figure 3

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Energy in units of E_0

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gap between the energies grows as n increases. Note that the energies in the case of a string

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n^2

n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

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$$-a \quad a$$

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Energy in units of

Figure 3

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$$E$$

Energy in units of V_0

Figure 3

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Energy in units of V_0

Figure 3

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Energy in units of

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Energy in units of
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Energy in units of
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Energy in units of

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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E

Energy in units of

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E_n

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Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE/\hbar} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar} \text{ outside.}$$

$$2m(E - V_0)/\hbar$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$k = \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

$$k = \sqrt{2m(E - V_0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto n^2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

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$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E = 0$

$V = V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of p

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$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

$$\text{inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter}$$

$$\text{there. A possible wavefunction might look something like the one in Fig. 5. It is customary}$$

to $-a \leq x \leq a$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section

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10.3.3 Finite square well

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by

$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy

E :

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

$\psi(x)$ is given

by

$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2mE}/\hbar \text{ outside.}$$

$$k = \sqrt{2m(E - V_0)}/\hbar$$

$$k = \sqrt{2mE}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

$E_n \propto n^2$

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto n^2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

$$\omega \propto \sqrt{E}$$

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future

convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

inside
and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case

of a string
are also proportional to n^2
, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V(x))/\hbar}$ inside the well and $\sqrt{2m(E - V_0)/\hbar}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string
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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.
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$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$
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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p
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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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 $V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$
 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$
 $E < V_0$

-a a
 Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside.

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E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

$$p = \sqrt{2m(E - V_0)}/\hbar$$

$$p = \sqrt{2m(E - 0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega = \hbar ck = \hbar c \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of V_0

Figure 3

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$E_n = \frac{1}{2} \hbar \omega_n$$

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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Energy in units of
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Energy in units of V_0

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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10.3.3 Finite square well

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Figure 5

Energy in units of V_0

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E

Energy in units of

Figure 3

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Energy in units of E_1

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E:

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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$$E = \frac{1}{2} \hbar \omega^2 = \frac{1}{2} \hbar c^2 k^2 = \frac{1}{2} \hbar c^2 \left(\frac{n\pi}{L} \right)^2$$

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Energy in units of V_0

Figure 3

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$$2mE/\hbar^2 \text{ inside the well and } 2m(E - V_0)/\hbar^2$$

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E

Energy in units of

Figure 3

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of

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plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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Energy in units of

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Energy in units of

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy

E:

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well.

k is larger inside the well, so the wavelength is shorter there.

A possible wavefunction might look something like the one in Fig. 5. It is customary to

normalize $\psi(x)$ to 1.

Figure 3 shows the energy levels for a string fixed at both ends. The gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 .

Figure 3

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 .

are also proportional to n^2 .

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 .

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere.

The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$V(x) = \begin{cases} V_0 & \text{for } |x| \leq a \\ 0 & \text{for } |x| > a \end{cases}$

by

$V(x) = \begin{cases} V_0 & \text{for } |x| \leq a \\ 0 & \text{for } |x| > a \end{cases}$

0 ($|x| \leq a$)

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well.

$$2m(E - V_0)/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside}$$

$$\text{and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A}$$

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

equals $\sqrt{2m(E - V(x))}/\hbar$. This

$$\text{inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

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-a a

E

Energy in units of

Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

E,

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$(19)$$

$$V_0$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

$E < V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

$$2m(E - V_0)/\hbar^2 \text{ outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

√

E,

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(x) is given

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$$V_0 \quad (|x| > a).$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

$E < V_0$

$-a \leq x \leq a$

$x < -a$

$x > a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

inside

and outside the well. k is larger inside the well, so the wavelength is shorter

there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n

ω

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

ω

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

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E,

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$-a \quad a$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

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Energy in units of

Figure 3

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Energy in units of

Figure 3

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Energy in units of

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 E
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$$E$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of

Figure 3

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E

Energy in units of E_1

Figure 3

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Energy in units of V_0

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$$\text{to -a a}$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere.

The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 > 0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$$-V_0 < E < 0$$

$$E > 0$$

Figure 4

• $E > 0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$V=V_0$

-a a

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$\propto (E - V(x))/\hbar^2$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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$$-a \quad a$$

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Energy in units of

Figure 3

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Energy in units of
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 $E < V_0$

-a a
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E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of V_0

Figure 3

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E

Energy in units of

Figure 3

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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$$\sqrt{2m(E - V_0)}/\hbar$$

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E

Energy in units of V_0

Figure 3

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$$E$$

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

$$E_n = \frac{1}{2} \hbar \omega_n$$

Energy in units of E_1

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$$E$$

$$E$$

Energy in units of E_1

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n^2

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto \omega_n^2$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω

$$2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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10.3.3 Finite square well

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$$0 \quad (|x| \leq a)$$

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(19)

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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n^2

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Energy in units of E_1

Figure 3

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Energy in units of V_0

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

√

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere.

The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section

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$-a \leq x \leq a$

$x < -a$

$x > a$

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Energy in units of

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of V_0

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } p = \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

$$2m(E - V_0)/\hbar^2 \text{ inside the well and } p = \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

$$2m(E - 0)/\hbar^2 \text{ outside.}$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

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E Energy in units of

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$$E > V_0$$

$$-a \leq x \leq a$$

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Energy in units of

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$$\text{to } \psi(0) = 1$$

Energy in units of V_0

Figure 3

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to -a a

E

Energy in units of

Figure 3

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$V(x) = \frac{1}{2}$

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Energy in units of
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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p $\sqrt{2m(E - V_0)/\hbar}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$\text{equals } \sqrt{2m(E - V_0)}/\hbar$$

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-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$\sqrt{2m(E - V_0)}/\hbar$$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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to $-a \leq x \leq a$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n = \frac{1}{2} \mu \omega_n^2$$

$$= \frac{1}{2} \mu c^2 k^2 \propto n^2, \text{ because although } \omega = ck \propto n, \text{ the energy is proportional to } \omega^2$$

$$(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both$$

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 $V=V_0$
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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 . This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

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E

Energy in units of

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$V=V_0$

$-a \ a$

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Energy in units of

Figure 3

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A possible wavefunction might look something like the one in Fig. 5. It is customary to

normalize $\psi(x)$ to 1.

Figure 3

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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plot $\psi(x)$ versus E .

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{inside the well and } p = \sqrt{2m(E - 0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to

-a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

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gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

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$$k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

Figure 5

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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E

Energy in units of

Figure 3

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Energy in units of

Figure 3

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$E < V_0$

$-a \leq x \leq a$

$x < -a$

$x > a$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

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$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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 E
 Energy in units of
 Figure 3
 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of
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10.3.3 Finite square well

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E:

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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2

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Energy in units of V_0

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter

there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

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Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E : $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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Figure 4

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Energy in units of

Figure 3

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Energy in units of

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are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E ,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

$\psi(x)$ is given

by

$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$$E > -V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > -V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } p = \sqrt{2mE}/\hbar \text{ outside.}$$

$$k = \sqrt{2m(E - V_0)}/\hbar$$

$$k = \sqrt{2m(E - V_0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto n^2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

E ,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

$\psi(x)$ is given

by

$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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zero.

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because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future

convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

inside
and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case

of a string
are also proportional to n^2
, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V(x))/\hbar}$ inside the well and $\sqrt{2m(E - V_0)/\hbar}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
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are also proportional to n^2
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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given
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$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$
(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p
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E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

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by
 $V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$
 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$
 $E < V_0$

-a a
Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside.

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

$$p = \sqrt{2m(E - V_0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ outside the well and } \sqrt{2m(E - 0)}/\hbar \text{ inside the well.}$$

$$\sqrt{2m(E - V_0)}/\hbar$$

$$\sqrt{2m(E - 0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

$$E$$

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}.$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \quad a$$

$$E$$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$E,$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E ,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section

10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future

convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ V_0 & (|x| > a). \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

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k is larger inside the well, so the wavelength is shorter there.

A possible wavefunction might look something like the one in Fig. 5. It is customary to

normalize $\psi(x)$ to 1.

Figure 3

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$2m(E - V_0)/\hbar$$

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Energy in units of

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$$0 \quad (|x| \leq a)$$

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-a a

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega \hbar (E - V(x)) / \hbar$. This equals $\sqrt{2m(E - V(x))} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of

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, because although $\omega = ck \propto n$, the energy is proportional to ω (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $\sqrt{2m(E - V(x))} / \hbar$. This equals $\sqrt{2m(E - V_0)} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $\sqrt{2m(E - V(x))} / \hbar$. This equals $\sqrt{2m(E - V_0)} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

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(19)

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$E > V_0$

$E < V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $\sqrt{2m(E - V(x))} / \hbar$. This equals $\sqrt{2m(E - V_0)} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

$\sqrt{2m(E - V(x))} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

$\sqrt{2m(E - V_0)} / \hbar$ inside the well and $\sqrt{2m(E - V_0)} / \hbar$ outside.

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of

Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2
, because although $\omega = ck \propto n$, the energy is proportional to ω

2
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apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

2
 2
outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well.}$$

$$k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

to $-a \leq x \leq a$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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√

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to $\omega^2 m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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Energy in units of

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- $V = 0$
- $V = V_0$

-a a

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E

Energy in units of V_0

Figure 3

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are also proportional to n^2
, because although $\omega = ck \propto n$, the energy is proportional to ω^2
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Energy in units of

Figure 3

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$$2m(E - V_0)/\hbar$$

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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10.3.3 Finite square well

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by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$V_0 (|x| > a)$.

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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Energy in units of

Figure 3

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$$\begin{aligned} V(x) &= \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \end{aligned} \quad (19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

$E < V_0$

-a a

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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 Energy in units of
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$$\begin{aligned}
 &\text{by} \\
 &V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \\
 &(19)
 \end{aligned}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$2mE/\hbar$ inside the well and p

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and p

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

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Energy in units of

Figure 3

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \ a$

E

Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V(x))/\hbar^2} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar^2} \text{ outside.}$$

$$k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well.}$$

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Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

$$E = V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V(x))/\hbar^2} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar^2} \text{ outside.}$$

$$k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well.}$$

k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

E

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

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, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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E

Energy in units of

Figure 3

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E

Energy in units of E_1

Figure 3

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$E = 0$

$E = V_0$

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Figure 4

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

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Energy in units of V_0

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by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \quad a$$

$$E$$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

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$\sqrt{2m(E - V_0)/\hbar}$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \quad a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$\sqrt{2m(E - V_0)/\hbar}$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega$

n^2

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$$V_0 (|x| \leq a)$$

$$0 (|x| > a).$$

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E

Energy in units of

Figure 3

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are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

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$$2m(E - V_0)/\hbar^2$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$-a \ a$

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Energy in units of

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Energy in units of

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0 ($|x| \leq a$)

V_0 ($|x| > a$).

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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-a a

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E
Energy in units of
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-a a
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E

Energy in units of

Figure 3

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Energy in units of E_1

Figure 3

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Energy in units of V_0

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by
 $V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$
(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E
Energy in units of
Figure 3

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E

Energy in units of V_0

Figure 3

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Energy in units of

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Energy in units of

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$\begin{aligned}
 V(x) &= \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
 $-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

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Energy in units of V_0

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Energy in units of V_0

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Energy in units of

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Energy in units of E_1

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 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$\begin{aligned}
 V(x) &= \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$
 $V=V_0$
 -a a
 Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

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$$E = 0$$

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$$-a \leq x \leq a$$

Figure 4

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$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))/\hbar} \text{ inside the well and } p$$

$$2m(E - V_0)/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

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$$V_0 \quad (|x| \leq a).$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of $p = \hbar k$. This equals $\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter

there. A possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \leq x \leq a$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E = 0$$

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$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to

to $-a \leq x \leq a$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

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Figure 4

• $E > 0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$\text{to } \psi(0) = 1$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

$$E_n \propto n^2$$

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$$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E = 0$

$V = V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

$$\text{inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter}$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere.

The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

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Energy in units of

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Energy in units of

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $\sqrt{2m(E - V(x))/\hbar}$. This equals $\sqrt{2m(E - V_0)/\hbar}$ inside the well and $\sqrt{2m(E - V_0)/\hbar}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
 $-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \hbar k$
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Energy in units of V_0

Figure 3

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E

Energy in units of

Figure 3

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E

Energy in units of V_0

Figure 3

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Energy in units of

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Energy in units of

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$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \quad a$$

$$E$$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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Energy in units of

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Energy in units of

Figure 3

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(19)

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Figure 4

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Figure 5

Energy in units of V_0

Figure 3

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$$V(x) = \begin{cases} V_0 & |x| \leq a \\ 0 & |x| > a \end{cases}$$

$$V_0 \text{ for } |x| \leq a$$

$$0 \text{ for } |x| > a$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar$$

$$p = \sqrt{2mE}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of V_0

Figure 3

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, because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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E

Energy in units of

Figure 3

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of E

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Energy in units of

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Energy in units of V_0

Figure 3

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Energy in units of

Figure 3

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Energy in units of

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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That wouldn't be much of a state, because the probability would be zero everywhere.

The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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$E = V_0$

$-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of V_0

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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Energy in units of V_0

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are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

ω^2

$\omega^2(E - V(x))/\hbar^2$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

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apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

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Energy in units of

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A possible wavefunction might look something like the one in Fig. 5. It is customary to

normalize ψ to 1.

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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E
Energy in units of

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-a a

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Energy in units of V_0

Figure 3

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$V = 0$
 $V = V_0$
 $-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \hbar k$
 $p = \sqrt{2m(E - V(x))}$. This equals $\sqrt{2m(E - V_0)}$ inside the well and $p = \sqrt{2m(E - 0)}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

Energy in units of V_0

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$$V=0$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

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Energy in units of E

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Energy in units of V_0

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Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of E_1

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$$V_0 \quad (|x| > a).$$

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$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar$$

$$\sqrt{2m(E - V_0)}/\hbar \text{ outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically

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That wouldn't be much of a state, because the probability would be zero everywhere. The

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

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Figure 4

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Energy in units of

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{inside the well and } p = \sqrt{2mE}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string

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$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega = \hbar ck = \hbar c k$. This equals $\sqrt{2mE}$ inside the well and $\sqrt{2m(E - V_0)}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$\sqrt{2m(E - V_0)}/\hbar$$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Figure 4

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$E$$

Energy in units of

Figure 3

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Figure 3

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

2

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Energy in units of

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Energy in units of

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Figure 4

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of

Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2
, because although $\omega = ck \propto n$, the energy is proportional to ω

2
(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.
There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$
(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside.

$2m(E - V_0)/\hbar$

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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E

Energy in units of

Figure 3

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$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to $\omega^2 m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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Energy in units of V_0

Figure 3

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both
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E

Energy in units of

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Figure 5

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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E

Energy in units of

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0 ($|x| \leq a$)

V_0 ($|x| > a$).

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E:

$V=0$

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$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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That wouldn't be much of a state, because the probability would be zero everywhere.

The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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Energy in units of V_0

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E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of
 Figure 3
 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$\begin{aligned}
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 (19)
 \end{aligned}$$

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a). \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

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 $\frac{2m(E - V(x))}{\hbar}$. This equals $\sqrt{\frac{2mE}{\hbar}}$ inside the well and $\sqrt{\frac{2m(E - V_0)}{\hbar}}$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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Energy in units of

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E

Energy in units of

Figure 3

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Energy in units of V_0

Figure 3

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to $\omega m(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

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Energy in units of

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 $V = V_0$
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Energy in units of V_0

Figure 3

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$$V_0 \quad (|x| > a).$$

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$$E > V_0$$

$$-a \leq x \leq a$$

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$$E$$

Energy in units of

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Energy in units of

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E

Energy in units of

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E

Energy in units of

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A possible wavefunction might look something like the one in Fig. 5. It is customary to

plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

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$V_0 \ (|x| > a)$.

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of

$\sqrt{2m(E - V(x))/\hbar}$. This equals $\sqrt{2mE/\hbar}$ inside the well and

$\sqrt{2m(E - V_0)/\hbar}$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

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Energy in units of
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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$\begin{aligned} V(x) &= V_0 & (|x| \leq a) \\ V(x) &= 0 & (|x| > a). \end{aligned} \quad (19)$$

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$E < V_0$

$E > V_0$

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 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$
 $E < V_0$

-a a
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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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E

Energy in units of V_0

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Energy in units of

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$$E_n = \frac{1}{2} \rho \omega^2 L^2$$

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \ (|x| \leq a)$

$V_0 \ (|x| > a)$.

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

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Energy in units of V_0

Figure 3

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{}$$

$$2mE/\hbar \text{ inside the well and } p$$

$$2m(E - V_0)/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{}$

E,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

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$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \quad a$$

$$E$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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$$E$$

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and p

$\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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$0 \quad (|x| \leq a)$
 $V_0 \quad (|x| > a).$
 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

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$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
 Figure 3
 gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

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$$\begin{aligned}
 V(x) &= \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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2

, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of

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Energy in units of

Figure 3

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Energy in units of

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Energy in units of

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Figure 3

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are also proportional to n^2

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, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then V

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by

$$V(x) = \begin{cases} -V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < -V_0$$

$$E > -V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > -V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

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Energy in units of
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E
Energy in units of
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Energy in units of V_0

Figure 3

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string

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2

2

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Figure 3

E

Energy in units of $\hbar^2/(2ma^2)$

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Energy in units of

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A possible wavefunction might look something like the one in Fig. 5. It is customary to

plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω

$$2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V(x))}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

$$2m(E - V_0)/\hbar^2 \text{ outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both inside and outside the well. } k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \leq x \leq a$$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

√

E,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere.

The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section

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$x < -a$

$x > a$

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$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

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Energy in units of

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$$V(x) = \begin{cases} \frac{1}{2}V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a$ a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere.

The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

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$E < V_0$

$E = V_0$

$E < 0$

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$-a \leq x \leq a$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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$$E$$

Energy in units of V_0

Figure 3

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E:

$V=0$

$V=V_0$

$-a$ a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and p

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outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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Energy in units of V_0

Figure 3

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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10.3.3 Finite square well

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$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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Energy in units of

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Energy in units of

Figure 3

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Energy in units of V_0

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Energy in units of

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and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$. There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero. That wouldn't be much of a state, because the probability would be zero everywhere. The lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below), because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

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Energy in units of V_0

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$

$E < V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar^2. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar$$

$$2mE/\hbar^2 \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$2m(E - V_0)/\hbar^2$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Figure 5

Energy in units of V_0

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gap between the energies grows as n increases. Note that the energies in the case of a string

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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E

Energy in units of

Figure 3

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Energy in units of V_0

Figure 3

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E

Energy in units of E_1

Figure 3

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 E
 Energy in units of
 Figure 3
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E

Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

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 V(x) &= \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases} \\
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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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10.3.3 Finite square well

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let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

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Energy in units of V_0

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case

of a string
are also proportional to n^2
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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both
apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,
where as the quantum mechanical system has $\omega \propto E$.
There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.
That wouldn't be much of a state, because the probability would be zero everywhere. The
lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),
because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll
let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given
by

$$V(x) = \begin{cases} \frac{1}{2} & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases}$$
(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$
 $V = V_0$
-a a

Figure 4

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

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$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and p

$2m(E - V_0)/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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Energy in units of

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$p = \sqrt{2m(E - V_0)}/\hbar$$

$$p = \sqrt{2m(E - 0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

n^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega = \hbar ck = \hbar c k$. This equals $\sqrt{2mE}/\hbar$ inside the well and $p = \sqrt{2m(E - V_0)}/\hbar$

$$\sqrt{2mE}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$\sqrt{2m(E - V_0)}/\hbar$$

$$\sqrt{2m(E - 0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

E,

where as the quantum mechanical system has $\omega \propto E$.

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$$V = 0$$

$$V = V_0$$

$$-a \quad a$$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$2m(E - V(x))/\hbar. \text{ This equals } \sqrt{2mE/\hbar} \text{ inside the well and } \sqrt{2m(E - V_0)/\hbar} \text{ outside.}$$

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$$k \text{ is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to } -a \quad a$$

E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$E$$

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

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2

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k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \ a$

E

Energy in units of E_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar\omega(E - V(x))/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

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$$2m(E - V_0)/\hbar$$

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Energy in units of

Figure 3

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n^2

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Energy in units of

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2

2

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Energy in units of

Figure 3

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(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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Energy in units of V_0

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10.3.3 Finite square well

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$V(x) = \frac{1}{2}$

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V_0 ($|x| > a$).

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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Figure 5

Energy in units of V_0

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Energy in units of E_1

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
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Energy in units of V_0

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

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$V=0$

$V=V_0$

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

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, because although $\omega = ck \propto n$, the energy is proportional to ω

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Energy in units of V_0

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Figure 3

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$V(x) = \begin{cases} \frac{1}{2} & 0 \leq |x| \leq a \\ V_0 & |x| > a \end{cases}.$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V = 0$

$V = V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V(x))/\hbar}$ inside the well and $\sqrt{2m(E - V_0)/\hbar}$ outside.

$2m(E - V_0)/\hbar$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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Energy in units of

Figure 3

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k is larger inside the well, so the wavelength is shorter there.

A possible wavefunction might look something like the one in Fig. 5. It is customary to

normalize $\psi(x)$ to 1.

Energy in units of E_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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Energy in units of

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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Energy in units of V_0

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Energy in units of

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 . This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

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(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \frac{1}{2}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \ a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2mE/\hbar}$ inside the well and $\sqrt{2m(E - V_0)/\hbar}$

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Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

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$$E$$

Energy in units of V_0

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$
 $V_0 \quad (|x| > a).$
 (19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$2m(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside.

$2m(E - V_0)/\hbar$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \quad a$

E

Energy in units of E_1

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

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$V(x) = \frac{1}{2}$

$0 \quad (|x| \leq a)$

$V_0 \quad (|x| > a).$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$

$V=V_0$

$-a \quad a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

Energy in units of
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$$\begin{aligned}
 V(x) &= \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \\
 (19)
 \end{aligned}$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$V=0$
 $V=V_0$
 -a a
 Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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Energy in units of

Figure 3

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E < V_0$$

$$E > V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } p = \sqrt{2mE}/\hbar \text{ outside.}$$

$$p = \sqrt{2mE}/\hbar \text{ inside the well and } p = \sqrt{2m(E - V_0)}/\hbar \text{ outside.}$$

$$p = \sqrt{2m(E - V_0)}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

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E
Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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-a a

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to $-a \quad a$

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(19)

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$$E = 0$$

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$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n = \frac{1}{2} \hbar \omega_n$$

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar \omega(E - V(x))/\hbar$. This equals $\sqrt{2m(E - V_0)}/\hbar$ inside the well and $\sqrt{2m(E - 0)}/\hbar$ outside.

$$k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

$$\text{inside}$$

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n = \frac{1}{2} \hbar \omega_n$$

, because although $\omega = ck \propto n$, the energy is proportional to $\hbar \omega$

$$E_n = \frac{1}{2} \hbar \omega_n$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$V_0 \quad (|x| \leq a).$$

$$(19)$$

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E < V_0$

$E > V_0$

$-a \leq x \leq a$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form

of $p = \hbar k$. This equals $\sqrt{2m(E - V(x))}$ inside the well and $\sqrt{2m(E - V_0)}$ outside.

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

inside and outside the well. k is larger inside the well, so the wavelength is shorter

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possible wavefunction might look something like the one in Fig. 5. It is customary

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Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

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$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$E = 0$

$E = V_0$

$-a \leq x \leq a$

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$k = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2mE}/\hbar \text{ inside the well and } \sqrt{2m(E - V_0)}/\hbar$$

$$\text{outside. } k \text{ is therefore real everywhere, so } \psi(x) \text{ is an oscillatory function both}$$

inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$,

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$\psi(x)$ is given

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$$V_0 \quad (|x| > a).$$

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$$E < -V_0$$

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$$-a \leq x \leq a$$

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• $E > -V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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$$k = \sqrt{2m(E - V_0)}/\hbar$$

$$k = \sqrt{2mE}/\hbar$$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

$$E_n \propto n^2$$

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

$$E_n \propto n^2$$

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$

E ,

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$$0 \quad (|x| \leq a)$$

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-a a

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• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E
Energy in units of
Figure 3
gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

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gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2

are also proportional to n^2

by

$$V(x) = \begin{cases} \frac{1}{2} \\ 0 \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$V=0$$

$$V=V_0$$

$$-a \quad a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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Energy in units of E

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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That wouldn't be much of a state, because the probability would be zero everywhere. The

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possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

n^2

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$0 \quad (|x| \leq a)$

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(19)

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-a a

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E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2 , because although $\omega = ck \propto n$, the energy is proportional to ω

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E
Energy in units of
Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of V_0

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are also proportional to n^2

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there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n^2 , because although $\omega = ck \propto n$, the energy is proportional to ω^2 (because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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$V = 0$

$V = V_0$

-a a

Figure 4

- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of $p = \sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$ outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside and outside the well. k is larger inside the well, so the wavelength is shorter there. A possible wavefunction might look something like the one in Fig. 5. It is customary to -a a

E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string are also proportional to n

2

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2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

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possible wavefunction might look something like the one in Fig. 5. It is customary

to $-a \leq x \leq a$

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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n

2

, because although $\omega = ck \propto n$, the energy is proportional to ω

2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2 and 3 both

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10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given by

$$V(x) = \begin{cases} 0 & (|x| \leq a) \\ V_0 & (|x| > a) \end{cases} \quad (19)$$

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$$E$$

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Figure 3

E

Energy in units of V_0

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plot $\psi(x)$ versus x .

E

Energy in units of V_0

Figure 3

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This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

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- $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p
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Energy in units of E

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string
are also proportional to n^2
, because although $\omega = ck \propto n$, the energy is proportional to ω^2
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E

Energy in units of

Figure 3

gap between the energies grows as n increases. Note that the energies in the case

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$$-a \leq x \leq a$$

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

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Energy in units of

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Energy in units of

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where as the quantum mechanical system has $\omega \propto E$.

There is no $n = 0$ state, because from Eq. (18) this would make ψ be identically zero.

That wouldn't be much of a state, because the probability would be zero everywhere. The

lack of a $n = 0$ state is consistent with the uncertainty principle (see Section 10.4 below),

because such a state would have $\Delta x \Delta p = 0$ (since $\Delta x < L$, and $\Delta p = 0$ because $n = 0 \Rightarrow k = 0 \Rightarrow p = \hbar k = 0$), which would violate the principle.

10.3.3 Finite square well

Things get more complicated if we have a finite potential well. For future convenience, we'll

let $x = 0$ be located at the center of the well. If we label the ends as $\pm a$, then $V(x)$ is given

by

$$V(x) = \begin{cases} V_0 & (|x| \leq a) \\ 0 & (|x| > a) \end{cases}$$

$$0 \quad (|x| \leq a)$$

$$V_0 \quad (|x| > a).$$

(19)

This is shown in Fig. 4. Given V_0 , there are two basic possibilities for the energy E :

$$E > V_0$$

$$E < V_0$$

$$-a \leq x \leq a$$

Figure 4

• $E > V_0$ (unbound state): From Eq. (11), the wavenumber k takes the general form of p

$$p = \sqrt{2m(E - V(x))}/\hbar. \text{ This equals } \sqrt{2m(E - V_0)}/\hbar \text{ inside the well and } \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

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$$p = \sqrt{2m(E - 0)}/\hbar \text{ outside.}$$

k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both inside

and outside the well. k is larger inside the well, so the wavelength is shorter there. A

possible wavefunction might look something like the one in Fig. 5. It is customary to $-a \leq x \leq a$

E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

(because the time derivative in Eq. (4.50) brings down a factor of ω). So Figs. 2

and 3 both

apply to both systems. The difference between the systems is that a string has $\omega \propto \sqrt{E}$, where as the quantum mechanical system has $\omega \propto E$.

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$\sqrt{2m(E - V(x))}/\hbar$. This equals $\sqrt{2mE}/\hbar$ inside the well and $\sqrt{2m(E - V_0)}/\hbar$

outside. k is therefore real everywhere, so $\psi(x)$ is an oscillatory function both

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E

Energy in units of V_0

Figure 3

gap between the energies grows as n increases. Note that the energies in the case of a string

are also proportional to n^2

2

, because although $\omega = ck \propto n$, the energy is proportional to ω^2

i am number 1

i am number 2

i am number 3

i am number 4

i am number 5

i am number 6

i am number 7

i am number 8

i am number 9

i am number 10

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