Linear Systems Solution of LDS

Sivakumar Balasubramanian

Department of Bioengineering Christian Medical College, Bagayam Vellore 632002

State space representation of LDS

▶ A state space representation of a LTI system takes the following form,

$$\dot{\mathbf{x}}(t) = \mathbf{A}\mathbf{x}(t) + \mathbf{B}\mathbf{u}(t)$$

$$\mathbf{y}\left(t\right) = \mathbf{C}\mathbf{x}\left(t\right) + \mathbf{D}\mathbf{u}\left(t\right)$$

 Obtaining the solution to the above equations can be posed as the following prolbem,

Determine:
$$\mathbf{x}(t)$$
, $\mathbf{y}(t)$ $\forall t \geq 0$ Given: $\mathbf{u}(t)$, $\forall t \geq 0$ and $\mathbf{x}(0^{-})$

- We first solve the state equation to obtain $\mathbf{x}(t)$, which is then used to obtain $\mathbf{y}(t)$.
- ▶ Because the system is linear, we can separate the solution into zero-input and zero-state solutions.

Zero-input solution for $\mathbf{x}(t)$

Zero-Input Solution: We will start by assuming $\mathbf{u}\left(t\right) = \mathbf{0}$.

$$\dot{\mathbf{x}}\left(t\right) = \mathbf{A}\mathbf{x}\left(t\right)$$

- For the scalar case, $\dot{x}(t) = ax(t)$, we know the solution to be the following, $x(t) = e^{at}x(0^-)$.
- A similar approach works for the vector case. Let us assume to the solution to zero-input state equation to be of the form, $\mathbf{x}\left(t\right)=e^{t\mathbf{A}}\mathbf{x}\left(0^{-}\right)$.

Zero-input solution for $\mathbf{x}(t)$

ightharpoonup What is $e^{t\mathbf{A}}$? Functions of matrices are often defined to have properties consistent with that of their scalar counterparts.

$$e^{t\mathbf{A}} = \mathbf{I} + t\mathbf{A} + \frac{1}{2!}t^2\mathbf{A}^2 + \dots = \sum_{k=0}^{\infty} \frac{1}{k!}t^k\mathbf{A}^k \implies \frac{d}{dt}e^{t\mathbf{A}} = \mathbf{A}e^{t\mathbf{A}}$$

▶ Thus, $\mathbf{x}(t) = e^{t\mathbf{A}}\mathbf{x}(0^{-})$ is the zero-input solution.

 \blacktriangleright How do we evaluate $e^{t\mathbf{A}}$? We do not need to evaluate the infinite series.

Cayley-Hamilton Theorem

Every square matrix $\mathbf{A} \in \mathbb{R}^{n \times n}$ satisfies it own characteristic equation $p(\lambda) = 0$, i.e. $p(\mathbf{A}) = \mathbf{0}$. $p(\mathbf{A}) = \mathbf{A}^n + a_1 \mathbf{A}^{n-1} + a_2 \mathbf{A}^{n-2} + \ldots + a_n \mathbf{I} = \mathbf{0}$

$$p(\mathbf{A}) = \mathbf{A}^n + a_1 \mathbf{A}^{n-1} + a_2 \mathbf{A}^{n-2} + \dots + a_n \mathbf{I} = \mathbf{0}$$

▶ Consider a analytic function, $f(x) = \sum_{k=0}^{\infty} a_k x^k$, and a matrix $\mathbf{A} \in \mathbb{R}^{n \times n}$ with characteristic polynomial p(x). Then,

$$f(x) = q(x) p(x) + r(x)$$

where, q(x) and r(x) are the quotient and reminder polynomials, respectively, and $r(x) = c_0 + c_1 x + c_2 x^2 + \ldots + c_{n-1} x^{n-1}$. Since $p(\mathbf{A}) = \mathbf{0}$, we have $f(\mathbf{A}) = r(\mathbf{A}) = \sum_{k=0}^{n-1} c_k \mathbf{A}^k$. Determining c_k s will allow us to calulate $f(\mathbf{A})$.

- ▶ If **A** has n distinct eigenvalues, c_k s solved through the n equations, $f(\lambda_i) = q(\lambda_i) p(\lambda_i) + r(\lambda_i) = r(\lambda_i)$.
- For repeated eigenvalues, we will need the following,

$$\left. \frac{d^{m-1}}{dx^{m-1}} f\left(x\right) \right|_{x=\lambda_{i}} = \left. \frac{d^{m-1}}{dx^{m-1}} r\left(x\right) \right|_{x=\lambda_{i}}$$

- ightharpoonup For a diagonalizable matrix, $e^{t\mathbf{A}} = \mathbf{X}e^{t\mathbf{\Lambda}}\mathbf{X}^{-1}$, with $e^{t\mathbf{\Lambda}} = \operatorname{diag}\left(e^{\lambda_1 t}\dots e^{\lambda_n t}\right)$.

Evaluate $e^{t\mathbf{A}}$ for the following \mathbf{A} : (a) $\begin{bmatrix} 1 & 2 \\ 0 & 2 \end{bmatrix}$; (b) $\begin{bmatrix} 1 & 2 \\ 0 & 1 \end{bmatrix}$; (c) $\begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$

Evaluate $e^{t\mathbf{J}}$ for \mathbf{J} with GM=1: (a) AM=2; (b) AM=3; (c) AM=n.

Laplace transform approach to zero-input response $\mathbf{x}\left(t\right)$

► Taking the Unilateral Laplace transform of the zero-input state equation,

$$s\mathbf{x}_{\mathcal{L}}\left(s\right) - \mathbf{x}\left(0^{-}\right) = \mathbf{A}\mathbf{x}_{\mathcal{L}}\left(s\right)$$
where, $\mathbf{x}_{\mathcal{L}}\left(s\right) = \begin{bmatrix} X_{1}\left(s\right) & X_{2}\left(s\right) \dots X_{n}\left(s\right) \end{bmatrix}^{T}$, where $x_{i}\left(t\right) \overset{\mathcal{L}}{\longleftrightarrow} X_{i}\left(s\right)$.
$$\left(s\mathbf{I} - \mathbf{A}\right)\mathbf{x}_{\mathcal{L}}\left(s\right) = \mathbf{x}\left(0^{-}\right) \implies \mathbf{x}_{\mathcal{L}}\left(s\right) = \left(s\mathbf{I} - \mathbf{A}\right)^{-1}\mathbf{x}\left(0^{-}\right)$$

$$\implies \mathbf{x}\left(t\right) = \mathcal{L}^{-1}\left\{\left(s\mathbf{I} - \mathbf{A}\right)^{-1}\right\}\mathbf{x}\left(0^{-}\right)$$

Using the analogy from the scalar case, we could guess that $\mathbf{x}(t) = e^{t\mathbf{A}}\mathbf{x}(0^{-})$. One can obtain the same solution by first finding the $(s\mathbf{I} - \mathbf{A})^{-1}$ and taking the inverse Laplace of each entry of this matrix.

Laplace transform approach to zero-input response $\mathbf{x}\left(t\right)$

Find
$$\mathbf{x}(t)$$
 for $t \ge 0$: $\dot{\mathbf{x}}(t) = \begin{bmatrix} 1 & 2 \\ 0 & 2 \end{bmatrix} \mathbf{x}(t)$.

Properties of $e^{t\mathbf{A}}$

▶ The columns of $e^{t\mathbf{A}} = \begin{bmatrix} \mathbf{a}_1\left(t\right) & \mathbf{a}_2\left(t\right) & \dots \mathbf{a}_n\left(t\right) \end{bmatrix}$ represent the solutions to different initial conditions, i.e. $\mathbf{x}\left(t\right) = \mathbf{a}_i\left(t\right) = e^{t\mathbf{A}}\mathbf{e}_i$.

$$\mathbf{x}(t) = e^{t\mathbf{A}}\mathbf{x}(0^{-}) = e^{t\mathbf{A}}\sum_{i=1}^{n} x_i(0^{-})\mathbf{e}_i = \sum_{i=1}^{n} x_i(0^{-})\mathbf{a}_i$$

If we know the response of a system to a set of n linearly independent initial conditions. Let $\mathbf{X}(t)$ represent the matrix whose columns are the solutions to the different initial conditions, then for any given intial condition $\mathbf{x}(0^-)$, we have the solution,

$$\mathbf{x}(t) = \mathbf{X}(t) \left(\mathbf{X}(0^{-}) \right)^{-1} \mathbf{x}(0^{-})$$

Properties of $e^{t\mathbf{A}}$

For any arbitrary initial time τ , instead of 0, we can still use the exponential formula to find out the solution,

$$\mathbf{x}\left(t\right) = e^{(t-\tau)\mathbf{A}}\mathbf{x}\left(\tau\right)$$

 $ightharpoonup e^{t{f A}}$ is called the *state transition matrix*, which takes the state at any given time to its value t seconds forward in time.

Consider the case where, **A** is diagonalizable. Let $\{\lambda_i, \mathbf{v}_i\}_{i=1}^n$ be the eigenpairs of **A**. Then, $\mathbf{A} = \mathbf{V} \mathbf{\Lambda} \mathbf{V}^{-1}$, and we could write the zero-input state equation as,

$$\dot{\mathbf{x}}\left(t\right) = \mathbf{V}\mathbf{\Lambda}\mathbf{V}^{-1}\mathbf{x}\left(t\right) \xrightarrow{\tilde{\mathbf{x}}\left(t\right) = \mathbf{V}^{-1}\mathbf{x}\left(t\right)} \dot{\tilde{\mathbf{x}}}\left(t\right) = \mathbf{\Lambda}\tilde{\mathbf{x}}\left(t\right)$$

The set of coupled first order differential equations are decoupled by this transforamtion.

▶ The individual states of $\tilde{\mathbf{x}}(t)$ evolve independently of each other.

$$\tilde{\mathbf{x}}(t) = e^{t\mathbf{\Lambda}}\tilde{\mathbf{x}}(0^{-}) = \begin{bmatrix} e^{\lambda_{1}t} & & & \\ & e^{\lambda_{2}t} & & \\ & & \ddots & \\ & & & e^{\lambda_{n}t} \end{bmatrix} \tilde{\mathbf{x}}(0^{-})$$

▶ An arbitrary initial state $\mathbf{x}(0^-) = \sum_{i=1}^n \alpha_i \mathbf{v}_i$ evolves as follows,

$$\mathbf{x}(t) = \alpha_1 e^{\lambda_1 t} \mathbf{v}_1 + \alpha_2 e^{\lambda_2 t} \mathbf{v}_2 + \ldots + \alpha_n e^{\lambda_n t} \mathbf{v}_n$$

When **A** is not diagonalizable: $\mathbf{A} = \mathbf{VJV}^{-1} \implies \dot{\tilde{\mathbf{x}}}(t) = \mathbf{J}\tilde{\mathbf{x}}(t)$.

$$\tilde{\mathbf{x}}(t) = e^{t\mathbf{J}}\tilde{\mathbf{x}}(0^{-}) = \begin{bmatrix} e^{t\mathbf{J}_{1}} & & \\ & \ddots & \\ & & e^{t\mathbf{J}_{k}} \end{bmatrix} \tilde{\mathbf{x}}(0^{-})$$

Consider
$$\mathbf{A} = \mathbf{V} \begin{bmatrix} \lambda & 1 \\ 0 & \lambda \end{bmatrix} \mathbf{V}^{-1} \implies \dot{\tilde{\mathbf{x}}}(t) = \begin{bmatrix} \lambda & 1 \\ 0 & \lambda \end{bmatrix} \tilde{\mathbf{x}}(t)$$

$$\dot{\tilde{x}}_1(t) = \lambda \tilde{x}_1(t) + \tilde{x}_2(t)$$

$$\dot{\tilde{x}}_2(t) = \lambda \tilde{x}_2(t)$$

We do not have complete decoupling as in the case where A was diagonalizable.

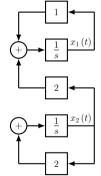
The exponential of a Jordan block has terms $e^{\lambda t}$, $te^{\lambda t}$, $te^{\lambda t}$, ...

$$e^{t\mathbf{J}_{1}} = e^{\lambda_{1}t} \begin{bmatrix} 1 & t & \frac{t^{2}}{2!} & \dots & \frac{t^{n}}{n!} \\ 0 & 1 & t & \dots & \frac{t^{n-1}}{(n-1)!} \\ 0 & 0 & 1 & \dots & \frac{t^{n-2}}{(n-2)!} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & 1 \end{bmatrix}$$

Thus,

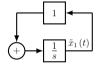
$$\tilde{x}_1(t) = \tilde{x}_1(0^-)e^{\lambda t} + \tilde{x}_2(0^-)te^{\lambda t}$$
$$\tilde{x}_2(t) = \tilde{x}_2(0^-)e^{\lambda t}$$

$$\dot{\mathbf{x}}\left(t\right) = \begin{bmatrix} 1 & 2 \\ 0 & 2 \end{bmatrix} \mathbf{x}\left(t\right)$$



When ${f A}$ is diagonalizable.

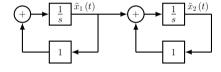
$$\dot{\tilde{\mathbf{x}}}(t) = \begin{bmatrix} 1 & 0 \\ 0 & 2 \end{bmatrix} \tilde{\mathbf{x}}(t)$$





When ${f A}$ is not-diagonalizable.

$$\dot{\tilde{\mathbf{x}}}(t) = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix} \tilde{\mathbf{x}}(t)$$



A Jordan block results in series of simple (scalar) first order blocks, where the output of a block acts as the input to another.

$$\dot{\mathbf{x}}(t) = \mathbf{A}\mathbf{x}(t) \implies \mathbf{x}(t) = \sum_{i=1}^{n} \alpha_i e^{\lambda_i t} \mathbf{v}_i$$

- The eigenvalues $\{\lambda_i\}_{i=1}^n$ of the system matrix $\mathbf A$ characterize the "natural" behavior of the system. These are called the *modes of the system*.
- ▶ The modes are exclusively expressed when the system starts in some specific set of states. When the system starts in an arbitrary state, the response contains a linear mixtute of these modes.

$$\dot{\mathbf{x}}(t) = \mathbf{A}\mathbf{x}(t) \implies \mathbf{x}(t) = \sum_{i=1}^{n} \alpha_i e^{\lambda_i t} \mathbf{v}_i$$

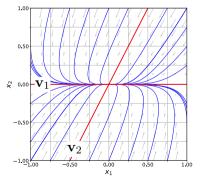
- ▶ **Dominant mode**: Determines the long-term behavior of the system. In the case of continuous-time systems, this would be the eigenvalue with the largest real part.
- If λ_i is a dominant mode $\Longrightarrow \left|\alpha_i e^{\lambda_i t}\right| \gg \left|\alpha_j e^{\lambda_j t}\right|, \forall j \neq i$ and t > T. This implies that after some time, the response almost only has that particular mode,

$$\mathbf{x}(t) \approx \alpha_i e^{\lambda_i t} \mathbf{v}_i, \ \forall t > T$$

▶ **Subdominant mode**: These are the other modes of the system, and these essentially determine how fast the system moves to the dominant mode.

Consider the system,
$$\dot{\mathbf{x}}\left(t\right) = \begin{bmatrix} -1 & 2 \\ 0 & -5 \end{bmatrix} \mathbf{x}\left(t\right)$$
.

Modes:
$$\begin{cases} \lambda_1 = -1, & \mathbf{v}_1 = \begin{bmatrix} 1 & 0 \end{bmatrix}^T \\ \lambda_2 = -5, & \mathbf{v}_2 = \begin{bmatrix} 1 & 2 \end{bmatrix}^T \end{cases}$$



Consider a system with modes: $(-1, \mathbf{v}_1)$, $(-1, \mathbf{v}_2)$, $(-3, \mathbf{v}_3)$, and $(-10, \mathbf{v}_4)$. What are the dominant modes? How does any arbitrary state evolve?

Describe the state equation of a mass M in free space. What are its modes?

Zero-solution for $\mathbf{x}(t)$

Let us now assume that the LTI system is relaxed when the input the applied to the system, i.e. $\mathbf{x}(0^-) = \mathbf{0}$. The effect of the input $\mathbf{u}(t)$ can be obtained by working in the Laplace domain,

$$\dot{\mathbf{x}}(t) = \mathbf{A}\mathbf{x}(t) + \mathbf{B}\mathbf{u}(t) \implies \mathbf{x}_{\mathcal{L}}(s) = (s\mathbf{I} - \mathbf{A})^{-1}\mathbf{B}\mathbf{u}_{\mathcal{L}}(s)$$

Taking the inverse Laplace transform, we get,

$$\mathbf{x}(t) = \int_{0}^{\infty} e^{(t-\tau)\mathbf{A}} \mathbf{B} \mathbf{u}(\tau) d\tau$$

Zero-solution for $\mathbf{x}(t)$

What do the columns of $e^{t\mathbf{A}}\mathbf{B}$ represent? What about the row of $e^{t\mathbf{A}}\mathbf{B}$? What about the ij^{th} element of $e^{t\mathbf{A}}\mathbf{B}$?

Complete solution for $\mathbf{x}(t)$ and $\mathbf{y}(t)$

▶ The complete solution for the state equations is given by the following,

$$\mathbf{x}(t) = e^{t\mathbf{A}}\mathbf{x}\left(0^{-}\right) + \int_{0}^{\infty} e^{(t-\tau)\mathbf{A}}\mathbf{B}\mathbf{u}\left(\tau\right)d\tau$$

▶ The output of the system is given by,

$$\mathbf{y}(t) = \mathbf{C}e^{t\mathbf{A}}\mathbf{x}\left(0^{-}\right) + \int_{0}^{\infty} \mathbf{C}e^{(t-\tau)\mathbf{A}}\mathbf{B}\mathbf{u}\left(\tau\right)d\tau + \mathbf{D}\mathbf{u}\left(t\right) = \mathbf{C}e^{t\mathbf{A}}\mathbf{x}\left(0^{-}\right) + \int_{0}^{\infty} \mathbf{G}\left(t-\tau\right)u$$

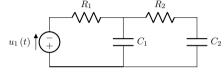
where, $\mathbf{G}(t) = \mathbf{C}e^{t\mathbf{A}}\mathbf{B} + \mathbf{D}\delta(t)$ is the *impulse response matrix* of the system.

▶ The transfer function of the system is given by: $\mathbf{H}\left(s\right) = \mathbf{C}\left(s\mathbf{I} - \mathbf{A}\right)^{-1}\mathbf{B} + \mathbf{D}$.

Find the impulse response matrix for
$$\mathbf{A} = \begin{bmatrix} 1 & 2 \\ 0 & 1 \end{bmatrix}$$
, $\mathbf{B} = \begin{bmatrix} 1 & -0.5 \\ 1 & 1 \end{bmatrix}$, $\mathbf{C} = \begin{bmatrix} 1 & 0 \end{bmatrix}$, and $\mathbf{D} = 0$.

Complete solution for $\mathbf{x}(t)$ and $\mathbf{y}(t)$

Find the expression for $\mathbf{y}\left(t\right)=\begin{bmatrix}v_{C_{1}}\left(t\right) & v_{R_{2}}\left(t\right)\end{bmatrix}^{T}$ for the following system, such that $v_{C_{1}}\left(0^{-}\right)=1V$, $v_{C_{2}}\left(0^{-}\right)=-0.5V$, $u_{1}\left(t\right)=1\left(t\right)V$, and $R=1k\Omega,C=1mF$.



System equations:

$$\mathbf{x}[k+1] = \mathbf{A}\mathbf{x}[k] + \mathbf{B}\mathbf{u}[k]$$

 $\mathbf{y}[k] = \mathbf{C}\mathbf{x}[k] + \mathbf{D}\mathbf{u}[k]$

Zero-input solution:

$$\mathbf{x}\left[k\right] = \mathbf{A}^k \mathbf{x}\left[0\right]$$

Zero-state solution:

$$\mathbf{x}\left[k\right] = \sum_{l=0}^{k-1} \mathbf{A}^{k-l-1} \mathbf{B} \mathbf{u}\left[l\right]$$

Complete solution:

$$\mathbf{x}[k] = \mathbf{A}^{k}\mathbf{x}[0] + \sum_{l=0}^{k-1} \mathbf{A}^{k-l-1}\mathbf{B}\mathbf{u}[l]$$

▶ \mathbf{A}^k is the state transition matrix and $\mathbf{G}[k] = \mathbf{A}^{k-1}\mathbf{B}$ is the impulse response matrix.

► We can approach this problem through the z-transform. Taking the unilateral z-transform of the state equation,

$$z\mathbf{X}_{\mathcal{Z}}(z) - z\mathbf{x}(0) = \mathbf{A}\mathbf{X}_{\mathcal{Z}}(z) + \mathbf{B}\mathbf{U}_{\mathcal{Z}}(z)$$

$$\mathbf{X}_{\mathcal{Z}}(z) = z (z\mathbf{I} - \mathbf{A})^{-1} \mathbf{x} [0] + (z\mathbf{I} - \mathbf{A})^{-1} \mathbf{B} \mathbf{U}_{\mathcal{Z}}(z)$$

The inverse z-tansform of this leads us to,

$$\mathbf{x}[k] = \mathbf{A}^{k}\mathbf{x}[0] + \sum_{l=1}^{k-1} \mathbf{A}^{k-l-1}\mathbf{B}\mathbf{u}[l]$$

Output:

$$\mathbf{y}\left[k\right] = \mathbf{C}\mathbf{A}^{k}\mathbf{x}\left[0\right] + \sum_{l=0}^{k-1}\mathbf{C}\mathbf{A}^{k-l-1}\mathbf{B}\mathbf{u}\left[l\right] + \mathbf{D}\mathbf{u}\left[k\right]$$

▶ The transfer function of the system is, $\mathbf{H}\left(z\right) = \mathbf{C}\left(z\mathbf{I} - \mathbf{A}\right)^{-1}\mathbf{B} + \mathbf{D}$

When A is diagonalizable, then we have

$$\mathbf{x}\left[k+1\right] = \mathbf{V}\mathbf{\Lambda}\mathbf{V}^{-1}\mathbf{x}\left[k\right] \implies \tilde{\mathbf{x}}\left[k+1\right] = \mathbf{\Lambda}\tilde{\mathbf{x}}\left[k\right]$$
 where, $\tilde{\mathbf{x}}\left[k\right] = \mathbf{V}^{-1}\mathbf{x}\left[k\right]$.

$$ilde{\mathbf{x}}\left[k
ight] = \mathbf{\Lambda}^k ilde{\mathbf{x}}\left[0
ight] = egin{bmatrix} \lambda_1^k & & & & \ & \lambda_2^k & & \ & & \ddots & \ & & & \lambda_n^k \end{bmatrix} ilde{\mathbf{x}}\left[0
ight]$$

An arbitrary initial state $\mathbf{x}[0] = \sum_{i=1}^{n} \alpha_i \mathbf{v}_i$ evolves as follows,

$$\mathbf{x}[k] = \alpha_1 \lambda_1^k \mathbf{v}_1 + \alpha_2 \lambda_2^k \mathbf{v}_2 + \ldots + \alpha_n \lambda_n^k \mathbf{v}_n$$

When A is not diagonalizable, $A = VJV^{-1}$

$$\tilde{\mathbf{x}}[k+1] = \mathbf{J}\tilde{\mathbf{x}}[k+1]$$

$$\tilde{\mathbf{x}}[k] = \mathbf{J}^k\tilde{\mathbf{x}}[0] = \begin{bmatrix} \mathbf{J}_1^k & & \\ & \ddots & \\ & & \mathbf{J}_l^k \end{bmatrix} \tilde{\mathbf{x}}[0]$$

Consider
$$\mathbf{A} = \mathbf{V} \begin{bmatrix} \lambda & 1 \\ 0 & \lambda \end{bmatrix} \mathbf{V}^{-1} \implies \tilde{\mathbf{x}} [k+1] = \begin{bmatrix} \lambda & 1 \\ 0 & \lambda \end{bmatrix} \tilde{\mathbf{x}} [k]$$

$$\tilde{x}_1 [k+1] = \lambda \tilde{x}_1 [k] + \tilde{x}_2 [k]$$

$$\tilde{x}_2 [k+1] = \lambda \tilde{x}_2 [k]$$

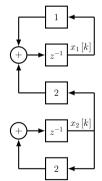
$$\mathbf{J}^k = \lambda^k \begin{bmatrix} 1 & \frac{k!\lambda^{-1}}{(k-1)!1!} & \frac{k!\lambda^{-2}}{(k-2)!2!} & \dots & \frac{k!\lambda^{-(n-1)}}{(k-n+1)!(n-1)!} \\ 0 & 1 & \frac{k!\lambda^{-1}}{(k-1)!1!} & \dots & \frac{k!\lambda^{-(n-2)}}{(k-n+2)!(n-2)!} \\ 0 & 0 & 1 & \dots & \frac{k!\lambda^{-(n-3)}}{(k-n+3)!(n-3)!} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & 1 \end{bmatrix}$$

Thus,

$$\tilde{x}_1[k] = \tilde{x}_1[0] \lambda^k + \tilde{x}_2[0] k \lambda^k$$

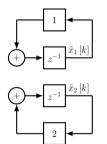
$$\tilde{x}_2[k] = \tilde{x}_2[0] \lambda^k$$

$$\mathbf{x}\left[k+1\right] = \begin{bmatrix} 1 & 2\\ 0 & 2 \end{bmatrix} \mathbf{x}\left[k\right]$$



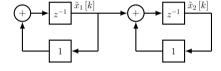
When ${f A}$ is diagonalizable.

$$\tilde{\mathbf{x}}[k+1] = \begin{bmatrix} 1 & 0 \\ 0 & 2 \end{bmatrix} \tilde{\mathbf{x}}[k]$$



When ${f A}$ is not-diagonalizable.

$$\tilde{\mathbf{x}}[k+1] = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix} \tilde{\mathbf{x}}[k]$$



A Jordan block results in series of simple (scalar) first order blocks, where the output of a block acts as the input to another.

Modes of a discrete-time system

$$\mathbf{x}[k] = \alpha_1 \lambda_1^k \mathbf{v}_1 + \alpha_2 \lambda_2^k \mathbf{v}_2 + \ldots + \alpha_n \lambda_n^k \mathbf{v}_n$$

- The eigenvalues $\{\lambda_i\}_{i=1}^n$ of the system matrix $\mathbf A$ characterize the "natural" behavior of the system. These are called the *modes of the system*.
- ▶ **Dominant mode**: Determines the long-term behavior of the system. In the case of discrete-time systems, this would be the eigenvalue with the largest magnitude.
- ▶ If λ_i is a dominant mode $\implies \left|\alpha_i\lambda_i^k\right| \gg \left|\alpha_j\lambda_j^k\right|, \forall j\neq i \text{ and } k>N.$ This implies that after some time, the response almost only has that particular mode,

$$\mathbf{x}[k] \approx \alpha_i \lambda_i^k \mathbf{v}_i, \ \forall t > T$$

▶ **Subdominant mode**: These are the other modes of the system, and these essentially determine how fast the system moves to the dominant mode.