

Scintillation loss

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Scintillation

- Index-of-refraction turbulence in the atmosphere induces temporal and spatial fluctuations in the intensity of the received signal
- Aim to find the fade statistics of the intensity falling below a specific threshold value
- Scintillation is expressed with the scintillation index, which gives the normalised variance of the optical intensity at a point in the detector plane:

$$\sigma_I^2 = \frac{\langle I^2 \rangle - \langle I \rangle^2}{\langle I \rangle^2} = \frac{\langle I^2 \rangle}{\langle I \rangle^2} - 1$$

Uplink vs Downlink

- Differs because of where the atmospheric turbulence layer is located
- For downlink, the turbulence layer is located only near the receiver, whereas for uplink, it is located near the transmitter
- Thus, beam wander affects the uplink case but it can be considered negligible for the downlink case

Weak vs strong fluctuation theory

- Weak fluctuation theory based on Rytov approximation is sufficient for most of the calculations required provided the zenith angle is sufficiently small (less than 60 degrees)
- Strong fluctuation theory applies for larger-diameter beams in the presence of beam wander or with larger zenith angles
- Weak or strong fluctuation theory will determine what the scintillation index looks like
- Used weak fluctuation theory for downlink

Probability Distribution of Intensity

Probability density function of intensity will help us find the fractional fade time.

Lognormal vs gamma-gamma distribution:

- Under weak turbulence ($SI < 1$), PDF can be said to be lognormal but values start disagreeing at tail ends, which is where we calculate fade probabilities.
- For both weak and strong turbulence, the gamma-gamma distribution holds and is more accurate, but it's more difficult to implement with the computer and is not accurate when aperture averaging occurs.
- Lognormal distribution was chosen for downlink.

Downlink model

Atmospheric Turbulence Strength, C_n^2

H-V 5/7 model:

$$C_n^2(h) = 0.00594(w/27)^2(10^{-5}h)^{10} \exp(-h/1000) \\ + 2.7 \times 10^{-16} \exp(-h/1500) + A \exp(-h/100),$$

w = rms wind speed

A = near ground level of C_n^2

h = height

For H-V 5/7 model, $w = 21\text{ms}^{-1}$ and $A = 1.7 \times 10^{-14} \text{ m}^{-2/3}$

Scintillation Index

Under weak fluctuation theory, scintillation index for downlink can be approximated by the Rytov Variance:

$$\sigma_R^2 = 2.25k^{7/6} \sec^{11/6}(\zeta) \int_{h_0}^H C_n^2(h)(h - h_0)^{5/6} dh.$$

$k = 2\pi / \text{wavelength}$

$h_0 = \text{ground height}$

$H = \text{height of satellite}$

$\zeta = \text{Zenith Angle}$

Aperture averaging

- Whether or not aperture averaging occurs depends on the correlation width p_c , which identifies the maximum receiver aperture size that will act like a “point receiver”.
- Aperture sizes larger than p_c will experience aperture averaging
- For downlink the correlation width is normally 7-10cm. Thus, aperture averaging is taken into account for our aperture
- The scintillation index is related to the aperture averaging effect through:

$$A_f(D) = \frac{\sigma_P^2(D)}{\sigma_I^2}$$

Aperture averaging

$$A_f = \left[1 + 1.062 \left(\frac{D}{2\rho_I} \right)^2 \right]^{-7/6}$$

$$\rho_I \approx 1.5 \cdot \sqrt{\frac{L \cdot \lambda}{2\pi}}$$

$$L \approx H_d \frac{\varepsilon/90^\circ}{\left(\varepsilon/90^\circ \right)^2 + \left(\varepsilon_{\max}/90^\circ \right)^2}$$

D = diameter of aperture

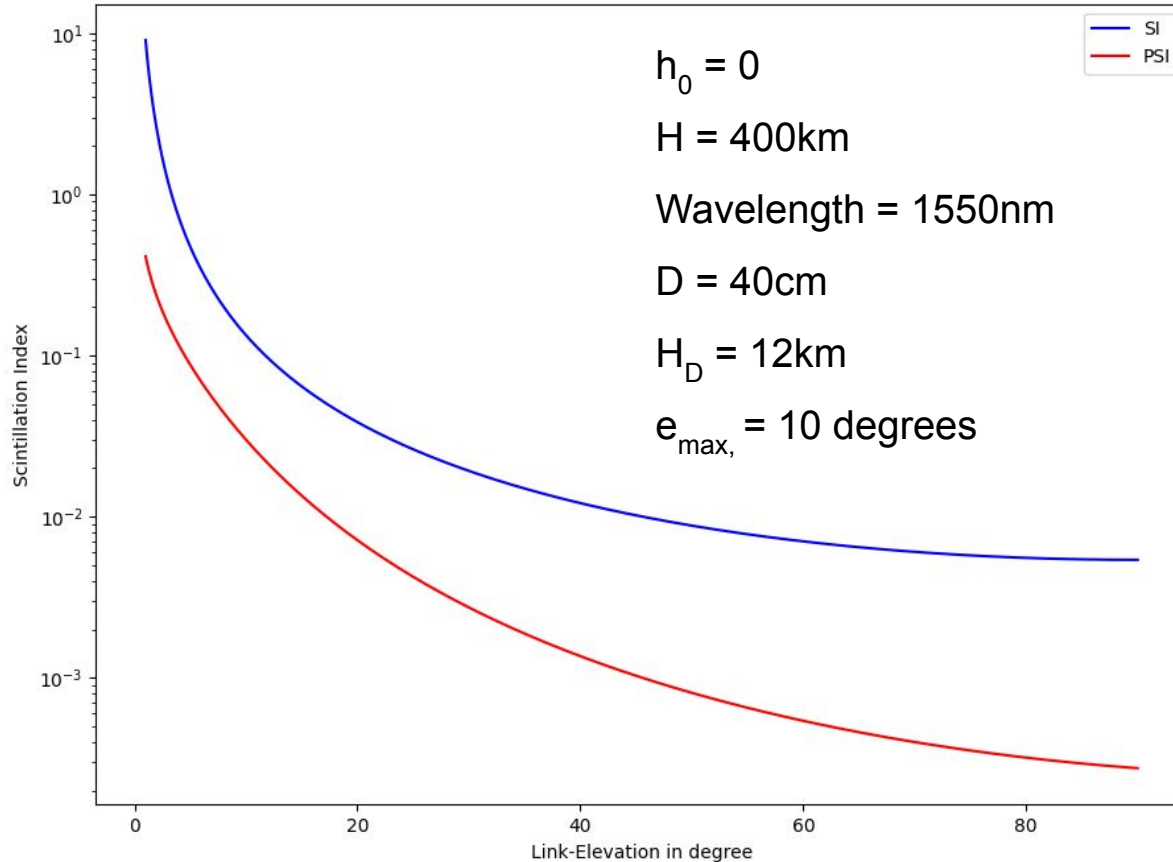
ρ_I = intensity structure size parameter

L = distance from a dominant turbulent layer

H_d = height of dominant turbulent layer
(assumed to be tropopause)

ε_{\max} = maximum-size elevation

Graph of scintillation index against link elevation

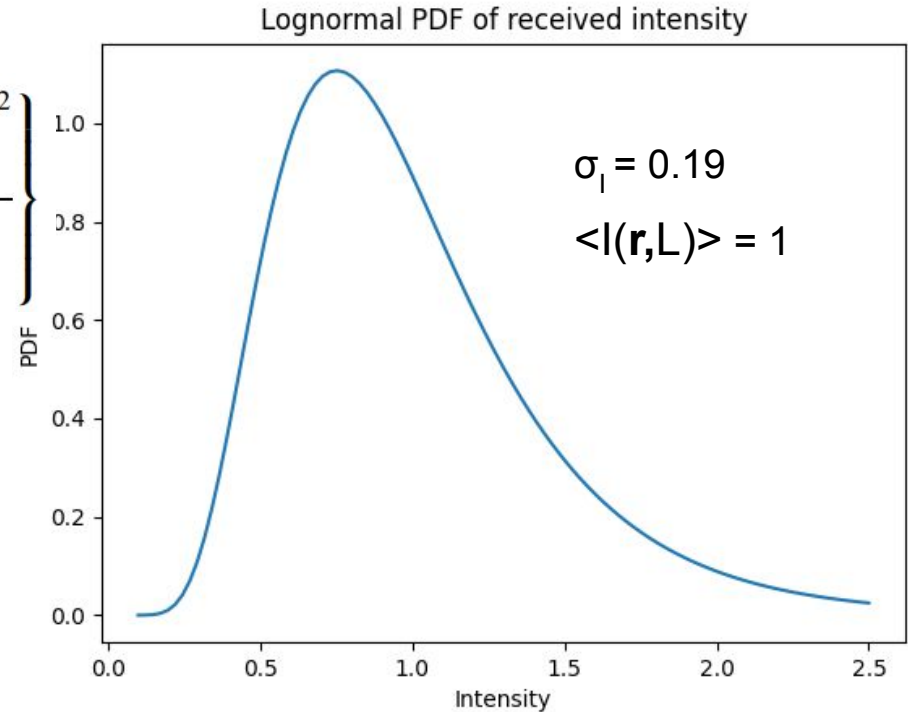


Probability density function of received intensity

$$p_I(I) = \frac{1}{I\sigma_I(\mathbf{r}, L)\sqrt{2\pi}} \exp \left\{ -\frac{\left[\ln\left(\frac{I}{\langle I(\mathbf{r}, L) \rangle}\right) + \frac{1}{2}\sigma_I^2(\mathbf{r}, L) \right]^2}{2\sigma_I^2(\mathbf{r}, L)} \right\}$$

σ_I = scintillation index (SI)

$\langle I(\mathbf{r}, L) \rangle$ = mean intensity

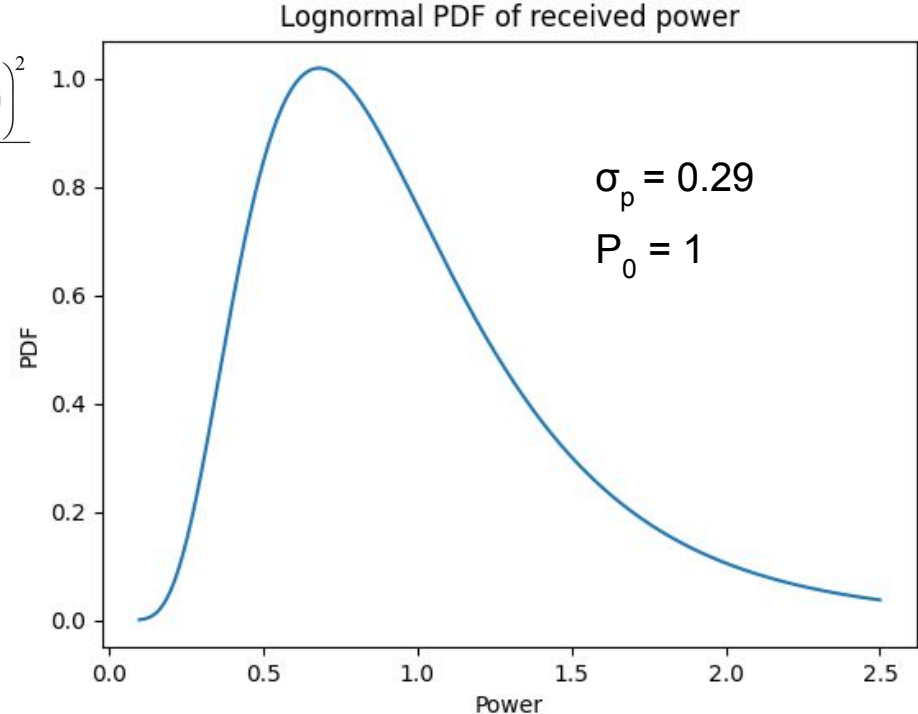


Probability density function of received power

$$p(P, \sigma_P^2) = \frac{1}{P \cdot \sqrt{2\pi \cdot \ln(\sigma_P^2 + 1)}} \cdot e^{-\frac{\left(\ln \frac{P}{P_0} + \frac{1}{2} \ln(\sigma_P^2 + 1)\right)^2}{2 \cdot \ln(\sigma_P^2 + 1)}}$$

σ_p = power scintillation index (PSI)

P_0 = mean power



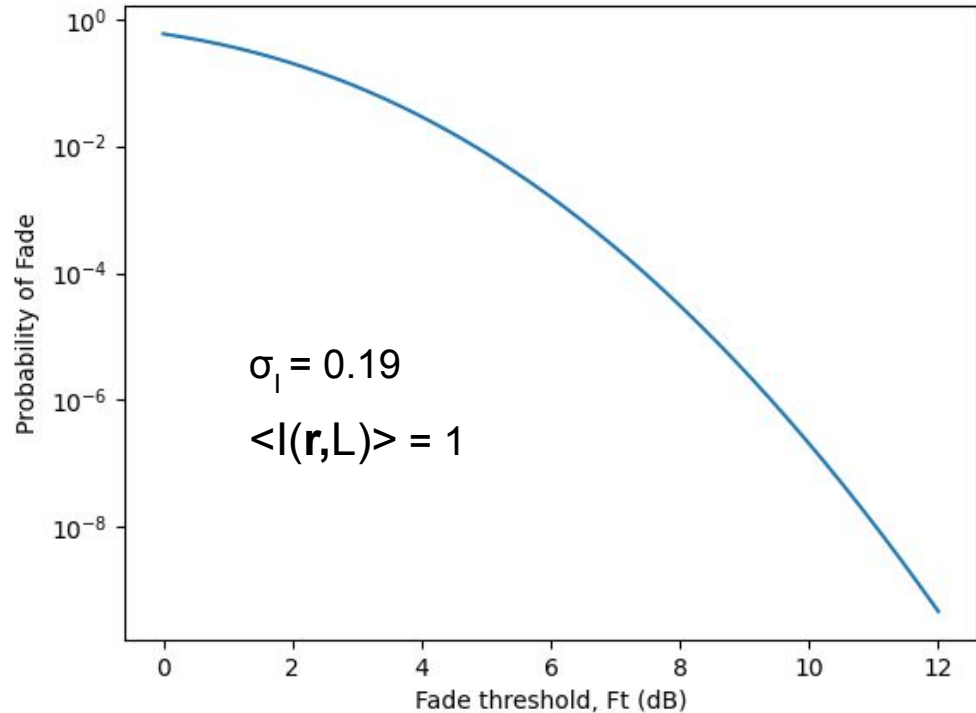
Probability of fade

- Describes the fractional fade time
- Obtained by integrating the PDF up to the intensity threshold, I_T

$$P(I \leq I_T) = \int_0^{I_T} p(I) dI$$

- Plotted against F_T , which expresses the intensity threshold I_T below the mean intensity in the form of decibels:

$$F_T = 10 \log_{10} \left(\frac{\langle I(0, L) \rangle}{I_T} \right). \quad [\text{dB}]$$



Expected number of fades per second

$$\langle n(I_T) \rangle = v_0 \exp \left\{ - \frac{\left[\frac{1}{2} \sigma_I^2(\mathbf{r}, L) + \frac{2r^2}{W_{LT}^2} - 0.23 F_T \right]^2}{2 \sigma_I^2(\mathbf{r}, L)} \right\},$$

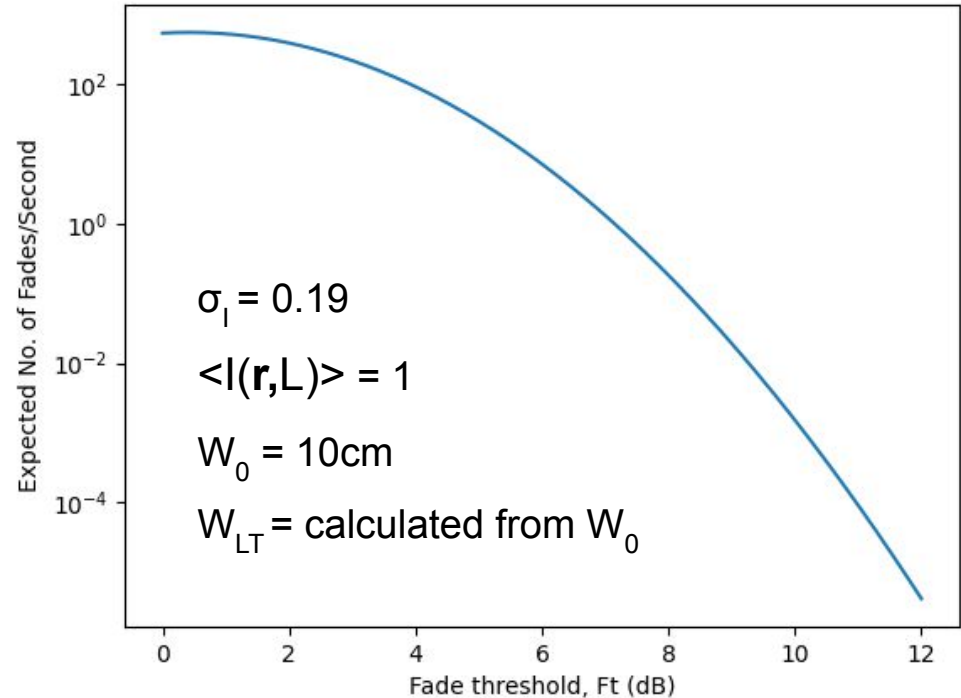
where r is related through

$$\langle I(\mathbf{r}, L) \rangle = \frac{W_0^2}{W_{LT}^2} \exp \left(- \frac{2r^2}{W_{LT}^2} \right)$$

$$v_0 = 550 \text{ Hz}$$

W_0 = beam waist at transmitter

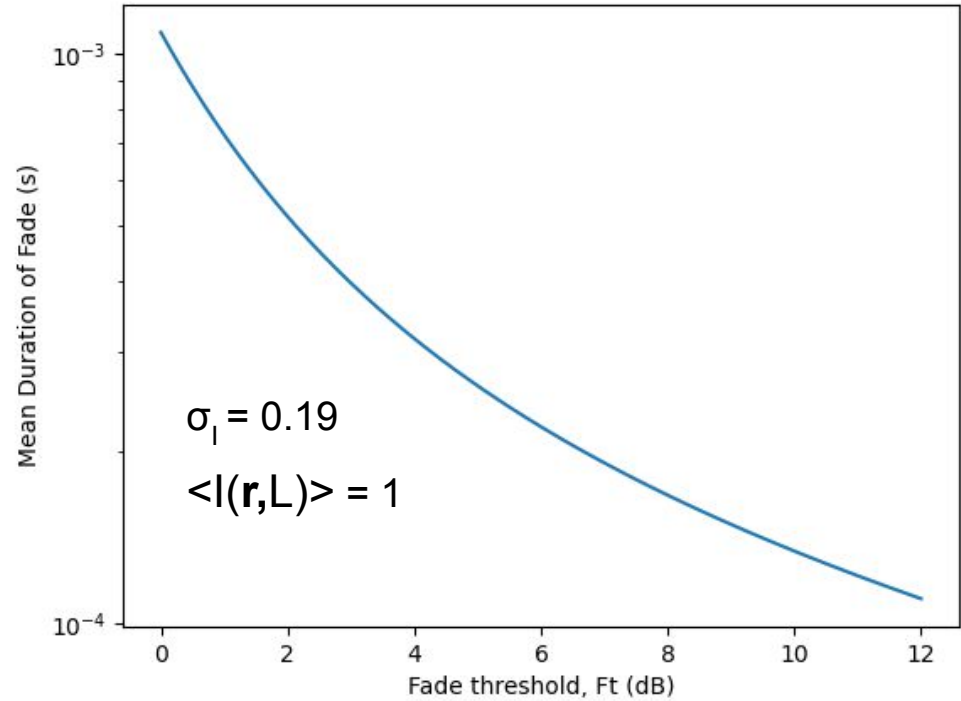
W_{LT} = long-term beam waist



Mean duration of fade

- Obtained by dividing the probability of fade with the expected number of fades per second

$$\langle t(F_T) \rangle = \frac{P_I(I \leq I_T)}{\langle n(I_T) \rangle},$$

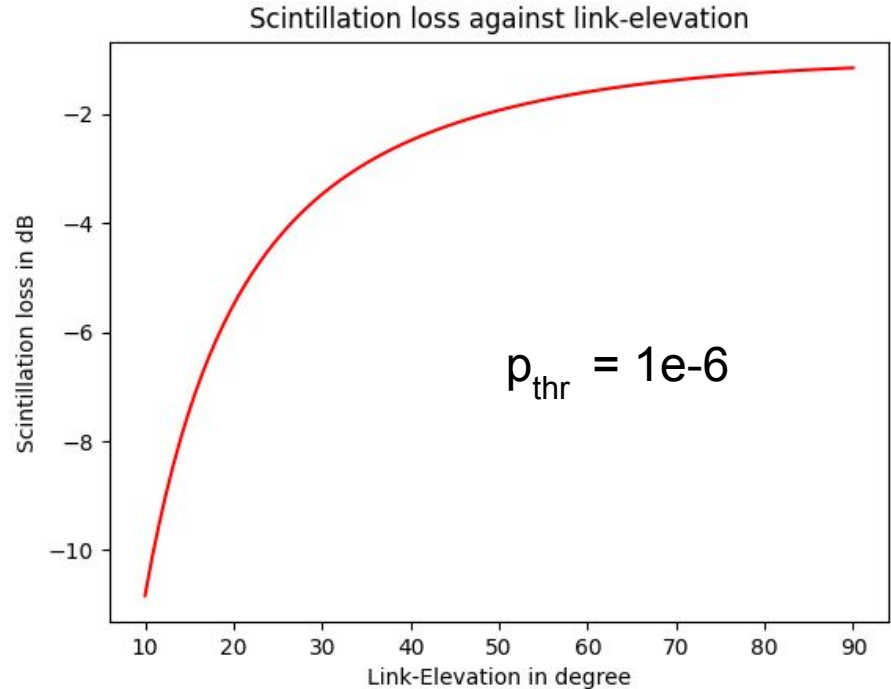


Loss expressed in dB

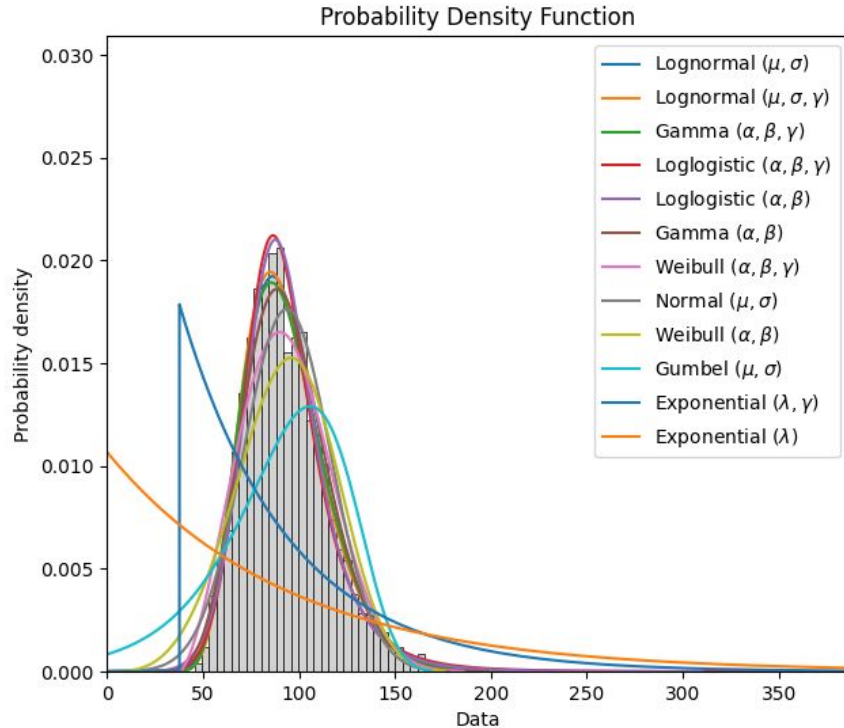
a_{sci} = loss associated with fade time

$$a_{\text{sci}} = -4.343 \left[\operatorname{erfinv}(2p_{\text{thr}} - 1) \sqrt{2 \ln(\sigma_P^2 + 1)} - \frac{1}{2} \ln(\sigma_P^2 + 1) \right]$$

p_{thr} = loss fraction = share where the received power falls below a given threshold from mean power (fade)



Measurements from Archenar Star



Wavelength: 780nm

Elevation angle: 30 degrees

Scintillation index calculated from
measurements: $0.12141/0.00269$

Scintillation index calculated from
model:

Uplink

- Calculation for scintillation index takes into account beam wander and comprises more steps:
 - a. Atmospheric coherence width, r_o
 - b. Pointing error variance, σ_{pe}^2
 - c. Split into tracked and untracked beams
 - d. Calculate off-axis and on-axis components separately
- Aperture averaging is negligible for uplink

Atmospheric coherence width

$$r_0 = \left[0.42 \sec(\zeta) k^2 \int_{h_0}^H C_n^2(h) dh \right]^{-3/5}$$

Pointing error variance

$$\sigma_{pe}^2 = 0.54 (H - h_0)^2 \sec^2(\zeta) \left(\frac{\pi}{kW_0} \right)^2 \left(\frac{2W_0}{r_0} \right)^{5/3} \\ \times \left[1 - \left(\frac{C_r^2 W_0^2 / r_0^2}{1 + C_r^2 W_0^2 / r_0^2} \right)^{1/6} \right]$$

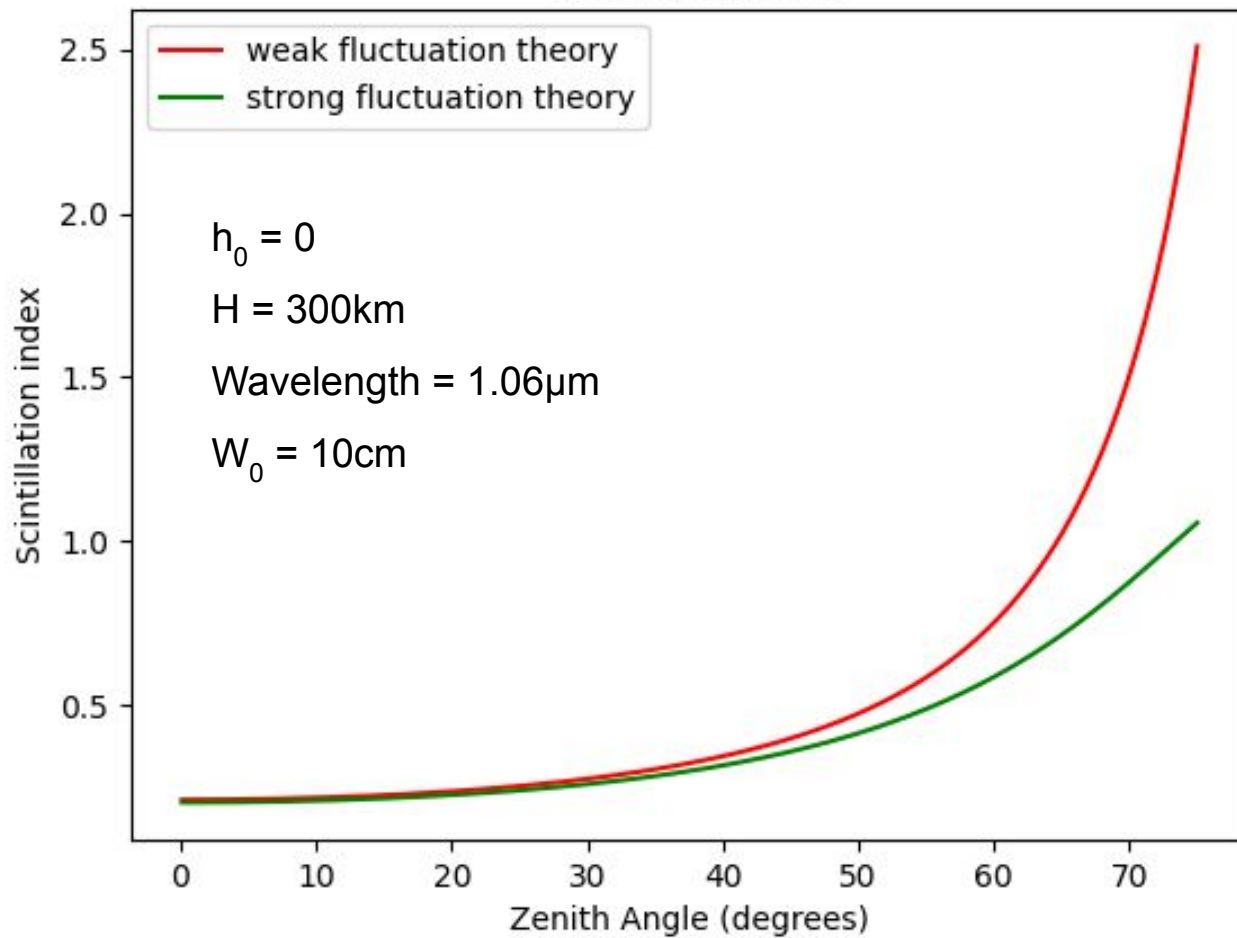
SI of a spherical wave under weak fluctuation theory

$$\sigma_{Bu}^2 = 2.25k^{7/6}(H - h_0)^{5/6} \sec^{11/6}(\zeta) \int_{h_0}^H C_n^2(h) \left(1 - \frac{h - h_0}{H - h_0}\right)^{5/6} \left(\frac{h - h_0}{H - h_0}\right)^{5/6} dh,$$

Longitudinal component of SI of a tracked beam under strong fluctuation theory

$$\sigma_{I,l}^2(L)_{\text{tracked}} = \exp \left[\frac{0.49\sigma_{Bu}^2}{\left(1 + (1 + \Theta)0.56\sigma_{Bu}^{12/5}\right)^{7/6}} + \frac{0.51\sigma_{Bu}^2}{\left(1 + 0.69\sigma_{Bu}^{12/5}\right)^{5/6}} \right] - 1,$$

on-axis, tracked



Gamma-gamma distribution

- More accurate for uplink

$$p_I(I) = \frac{2(\alpha\beta)^{(\alpha+\beta)/2}}{\Gamma(\alpha)\Gamma(\beta)I} \left(\frac{I}{\langle I(\mathbf{r}, L) \rangle} \right)^{(\alpha+\beta)/2} K_{\alpha-\beta} \left(2\sqrt{\frac{\alpha\beta I}{\langle I(\mathbf{r}, L) \rangle}} \right),$$

- α and β are parameters for the function representing large-scale and small-scale fluctuations respectively, where $\alpha = 1/\sigma_x^2$ and $\beta = 1/\sigma_y^2$

$$\sigma_X^2 = 5.95(H - h_0)^2 \sec^2(\zeta) \left(\frac{2W_0}{r_0} \right)^{5/3} \left(\frac{\alpha_{pe}}{W} \right)^2 \\ + \exp \left[\frac{0.49\sigma_{Bu}^2}{\left(1 + (1 + \Theta)0.56\sigma_{Bu}^{12/5} \right)^{7/6}} \right] - 1,$$

$$\sigma_Y^2 = \exp \left[\frac{0.51\sigma_{Bu}^2}{\left(1 + 0.69\sigma_{Bu}^{12/5} \right)^{5/6}} \right] - 1.$$