

# Real dimension in the Newtonian simulation of disk-like pressure-free systems

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## Abstract

Abstract...

## 1 Brute force: field line intersection density gradient

Regarding the holographic principle, where  $n$  is the gravitational field line count, and  $A_s$  is the Schwarzschild black hole event horizon area:

$$n = \frac{A_s k c^3}{4G\hbar \log 2}, \quad (1)$$

the Schwarzschild radius is:

$$r_s = \sqrt{\frac{A_s}{4\pi}} = \sqrt{\frac{nG\hbar \log 2}{kc^3\pi}}, \quad (2)$$

and the mass is:

$$M = \frac{c^2 r_s}{2G} = \sqrt{\frac{nc\hbar \log 2}{4Gk\pi}}. \quad (3)$$

Where  $R$  is some far distance from the centre of the gravitating body (e.g,  $R \gg r_s$ ),  $\beta$  is the get intersecting line length function, and  $\epsilon$  is some small value (e.g  $10^{-5}$ ), the gradient is:

$$\gamma = \frac{\beta(R + \epsilon) - \beta(R)}{\epsilon}. \quad (4)$$

The gradient strength is:

$$g = -\gamma\pi = \frac{n}{2R^3}. \quad (5)$$

The Newtonian acceleration  $a_{Newton}$  is:

$$a_{Newton} = \frac{v_{Newton}^2}{R} = \sqrt{\frac{gGc\hbar \log 2}{2R^2k\pi}}. \quad (6)$$

The Newtonian acceleration  $a_{flat}$  for a flat rotation curve is:

$$a_{flat} = \frac{v_{flat}^2}{R} = \frac{gRc\hbar \log 2}{2k\pi M}. \quad (7)$$

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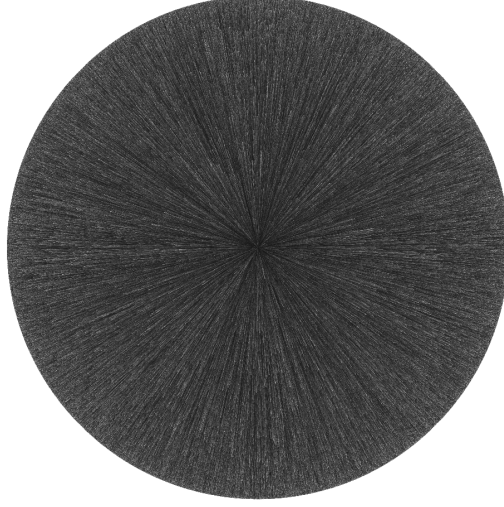


Figure 1: Where  $D = 3$ . The normals are isotropic.

The ratio of the acceleration is

$$\frac{a_{flat}}{a_{Newton}} = R^d, \quad (8)$$

where  $d = 3 - D$  stands for disk-like, and the dimension of the gravitation field is:

$$D = 3 - \frac{\log \frac{a_{flat}}{a_{Newton}}}{\log R} = 3 - \frac{\log \frac{v_{flat}^2}{v_{Newton}^2}}{\log R}. \quad (9)$$

```
#include <cmath>
#include <iostream>
using namespace std;

const double G = 6.67430e-11;
const double c = 299792458;
const double c2 = c * c;
const double c3 = c * c * c;
const double pi = 4.0 * atan(1.0);
const double h = 6.62607015e-34;
const double hbar = h / (2.0 * pi);
const double k = 1.380649e-23;

int main(void)
{
    double M = 1e41;

    double r_s = 2 * G * M / c2;
    double A_s = 4 * pi * r_s * r_s;
    double n = A_s * k * c3 / (4 * G * hbar * log(2.0));

    double R = 3e20;
    double g = n / (2 * R * R * R);

    double a_Newton = sqrt((g * G * c * hbar * log(2.0)) / (2 * R * R * k * pi));
```

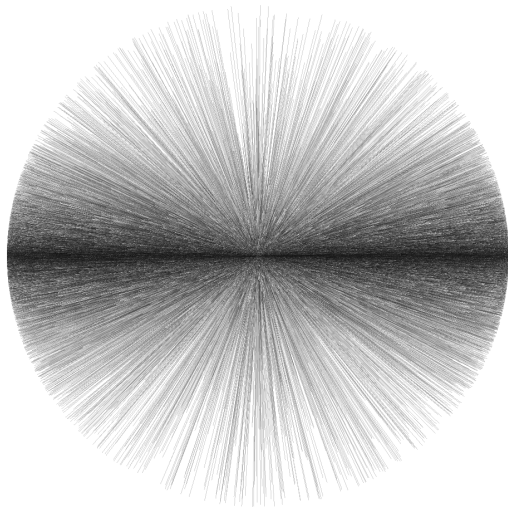


Figure 2: Where  $D = 2.1$ . The normals are increasingly anisotropic.

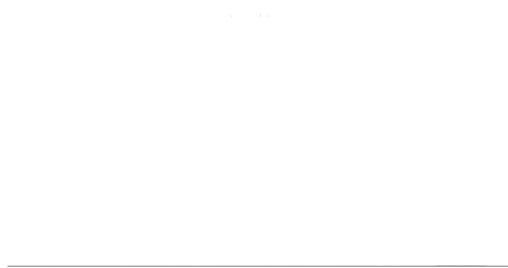


Figure 3: Where  $D = 2.001$ . The normals are anisotropic.



Figure 4: Where  $D = 3$ . This figure shows an axis-aligned bounding box and an isotropic emitter, looking from above. The bounding box is filled with the intersecting line segments.

```
double a_flat = pow(220000, 2.0) / R;  
  
double v_Newton = sqrt(a_Newton * R);  
double v_flat = 220000;  
  
double D = 3.0 - log(a_flat / a_Newton) / log(R);  
double D_ = 3.0 - log(pow(v_flat, 2.0) / pow(v_Newton, 2.0)) / log(R);  
  
cout << D << endl;  
cout << D_ << endl;  
  
return 0;  
}
```

## References

- [1] Misner et al. Gravitation. (1970)
- [2] 't Hooft. Dimensional reduction in quantum gravity. (1993)
- [3] Susskind. The World as a Hologram. (1994)