

# An axiomatic review of anisotropic quantum gravity

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## Abstract

In Newton's and Einstein's theory, all mass gravitates in an *isotropic* (spherical) manner. In this paper, we will consider *anisotropic* gravitating processes. We discuss dark matter, as well as dark energy and the possibility of a final, 5th interaction.

## 1 Questions

In this paper we will address the following three questions:

1. The hierarchy problem: why is gravitation so weak?
2. The dark matter problem: why is gravitation stronger than that predicted by general relativity in pressure-free dusts such as the Galactic disc?
3. The dark energy problem: why is the Universe undergoing accelerating expansion?

## 2 Axioms

Here we provide a list of 11 axioms regarding relativity:

1. Speed causes kinematic time dilation.
2. The gravitational field causes gravitational time dilation.
3. Physical processes are interruptible, and are indeed interrupted when undergoing time dilation.
4. Processes undergoing heavy kinematic time dilation are deflected twice as much as in Newtonian gravitation – for neutrinos, there is (practically) no internal process occurring to resist the gravitational attraction.

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Table 1: Table of interactions, including a 5th interaction.

Type	Inherent spatial dimension $D$	Communication spatial dimension
Gravitation (isotropic)	3	4
Gravitation (oblate)	2	3
Gravitation (prolate)	1	2
Weak	0	1
Electromagnetism	1	0
Strong	2	1
5th interaction	3	2

5. Physical processes are computations. This is not merely an analogy, it's reality.
6. Computations can be optimized, so there is time contraction and length dilation to consider.
7. The densest process for any given mass  $M$  or entropy  $S$  is a black hole – they are the most optimal of computations.
8. The gravitational field is quantized into gravitons.
9. There is no gravitational shadow – the relaying of gravitons is the cause of gravitational time dilation.
10. Gravitationally-bound, pressure-free dusts, such as the Galactic disc, have fractional dimension. As the dimension reduces, the strength of the gravitation increases – the dimension reduces as the shape of the Milky Way goes from spherical to disc-like with distance from the Galactic centre.
11. The self-optimization of the Universal process over time leads to length dilation, in the form of expansion – the antithesis of attractive gravitation.

### 3 Results

We have constructed Table 1 by first taking into account the inherent 3-D nature of isotropic gravitation, and its 4-D communications (e.g. *the* Wilson hypervolume). Next, we extrapolate all the way down, to where the strong interaction is 2-D, with 1-D communications (e.g. some Wilson lines). Finally, the possibility of a 5th interaction follows suit, in order to bring balance to the interactions in terms of their inherent spatial dimension.

Note that, unlike with the many Wilson lines per process, there is only one Wilson hypervolume, shared by all processes. This means that gravitational interactions are *connectionless* – isotropic gravitation is *broadcast*; there is no specific recipient (e.g. everyone is a target). On the other hand, the strong interactions are more directed, and *connected* – strong interaction is *unicast* or *multicast*; there is a specific recipient (e.g. not everyone is a target). For instance, the transition from broadcast transmission to directed transmission occurs as dark matter is factored in (e.g. where  $D < 3$ ). See Figs. 1 and 2 for a simplified analytical solution for anisotropic gravitation in the Milky Way. Connectedness is an attribute of the non-gravitational interactions (e.g. weak, electromagnetic, strong, and 5th interaction).

## 4 Conclusion

To answer the questions from the first section:

1. Why is gravitation so weak? Because it's often isotropic due to isotropic internal pressure.
2. Why is gravitation stronger than that predicted by general relativity in pressure-free dusts such as the Galactic disc? Fractional dimension and anisotropic gravitation.
3. Why is the Universe undergoing accelerating expansion? Accelerating Universal computational self-optimization.

## 5 Further questions

- Is a 5th interaction the same tetrahedral process that is predicted by (Wilson) loop quantum gravity? If so, then are superstring theory and loop quantum gravity fundamentally compatible?
- Do gravitons undergo Shapiro time delay? If so, then are graviton condensates naturally cold?
- Is a photon a gas of constituent particles related to a 5th interaction? If so, are all particles made up of constituent particles related to a 5th interaction?
- Can a human-scale object become pressure-free internally? If so, then does the gravitation become anisotropic as predicted in this paper?

## 6 Numerical solution for anisotropic gravitation

A copy of the C++ code for Figs. 3 through 9 can be obtained from:

[https://github.com/sjhalayka/ellipsoid\\_emitter](https://github.com/sjhalayka/ellipsoid_emitter)

This C++ code calculates an anisotropic gravitational field, and allows the coder to view the results using OpenGL or gnuplot. For instance, Fig. 3 shows the intersection count (e.g. field strength) as the dimension of the gravitational field reduces from  $D = 3.0$  to  $D = 2.001$ . Figs. 4 through 6 show an example gravitational field that goes from  $D = 3.0$  to  $D = 2.5$  to  $D = 2.001$ . Finally, Figs. 7 through 9 show that there are fractal patterns to be found when analyzing the data – here we define the fractionality of  $D$  as it goes from  $D = 3.0$  to  $D = 2.001$  as

$$\alpha = 1 - 2 \left( \frac{1}{2} - (D \text{ fmod } 1.0) \right), \quad (1)$$

where fmod is the floating-point (e.g. real) modulo operator.

## References

- [1] Halayka. Is the anisotropic interaction of luminous matter responsible for the extrinsic gravitation usually attributed to exotic dark matter? (2008)
- [2] Halayka. A note on anisotropic quantum gravity. (2023)

```

#include <iostream>
using namespace std;

int main(void)
{
    const double c = 299792458; // Speed of light in vacuum
    const double G = 6.674e-11; // Gravitational constant
    const double M = 1e41; // Galactic bulge mass

    const double start_distance = 6.1495e19; // Galactic bulge radius
    const double end_distance = 1e21; // just past the solar radius

    double v = sqrt(G * M / start_distance); // Speed with dark matter

    const size_t resolution = 10000;

    const double step_size = (end_distance - start_distance) / (resolution - 1);

    for (double r = start_distance; r <= end_distance; r += step_size)
    {
        const double v_N = sqrt(G * M / r); // Speed without dark matter

        if (v < v_N)
            v = v_N;

        const double D = 3.0 - log(v / v_N) / log(c);

        cout << r << " " << D << endl;
    }

    return 0;
}

```

Figure 1: C++ code for Galactic orbit. Here  $D$  represents dimension. Note that where  $D < 3$  that  $D$  is the lower bound, and that the value of  $D$  will actually be a little bit higher since the strength of the gravitation field increases *and* falloff decreases as  $D$  reduces. In other words, the code does not take this small change in falloff into account, in order to keep things as simple as possible while still highlighting the role of fractional dimension.

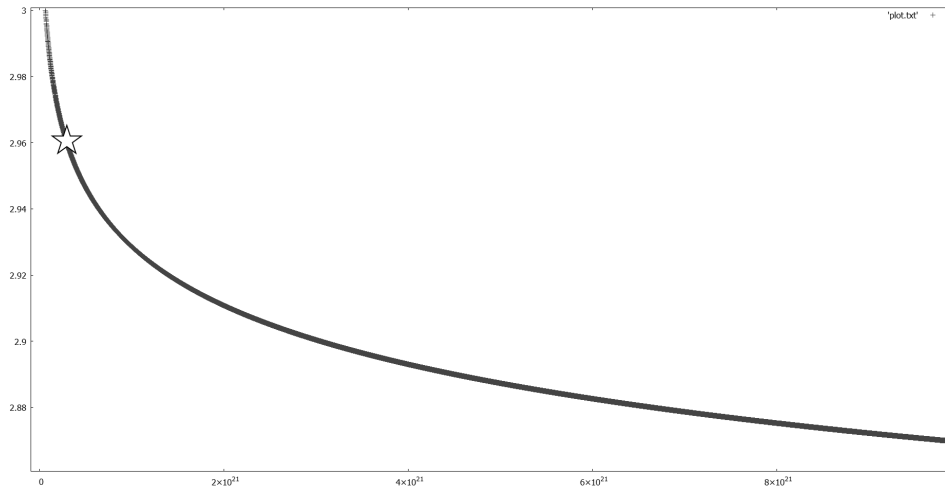


Figure 2: Output from the C++ code in Fig 1. Here the  $x$  axis is the distance from the Galactic centre  $r$ , and the  $y$  axis is dimension  $D$ . Note that the lower bound is  $D = 2.96$  at  $r = 3 \times 10^{20}$  metres from the Galactic centre (e.g. the Solar orbit radius).

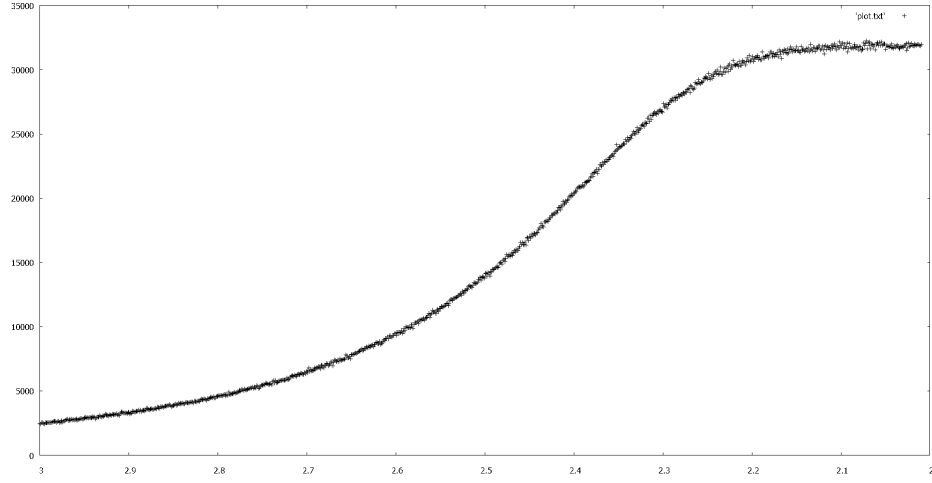


Figure 3: Here the  $x$  axis is dimension  $D$ , and the  $y$  axis is intersection count (e.g. field strength). It forms a nice sigmoid curve. Although not shown here, it is important to also note that falloff decreases from  $1/r^2$  to  $1/r^1$  as  $D$  reduces from 3 to 2 – see Figs. 4 through 6.

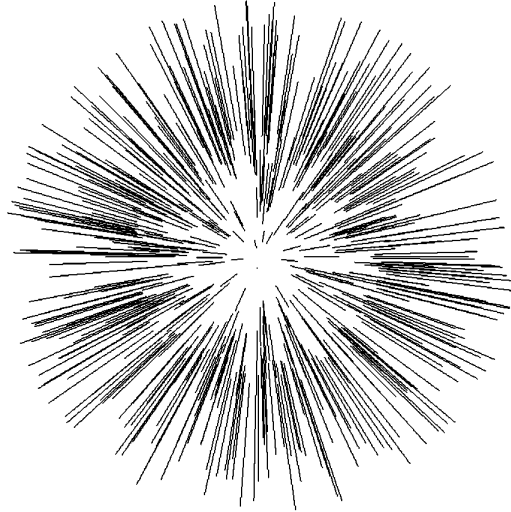


Figure 4: Example of an isotropic emitter. Here  $D = 3.0$ . The emitter is spherical. The positions are placed pseudorandomly on a 2-sphere, and the normals are calculated using the same sphere. The strength of the gravitational field is  $G$ , and the falloff is related to  $1/r^2$  – in other words, standard Newtonian gravitation.



Figure 5: Example of an anisotropic emitter. Here  $D = 2.5$ . The emitter is ellipsoidal. The positions are placed on an oblate (e.g. flattened) ellipsoid, and the normals are calculated by using the dual prolate (e.g. tall) ellipsoid. The strength of the gravitational field is like  $Gc^{0.5}$ , and the falloff is related to  $1/r^{1.5}$ .



Figure 6: Example of an anisotropic emitter. Here  $D = 2.001$ . The emitter is practically circular. This figure shows the circle edge-on, which is why it looks like a straight line. The strength of the gravitational field is like  $Gc$ , and the falloff is related to  $1/r^1$ .

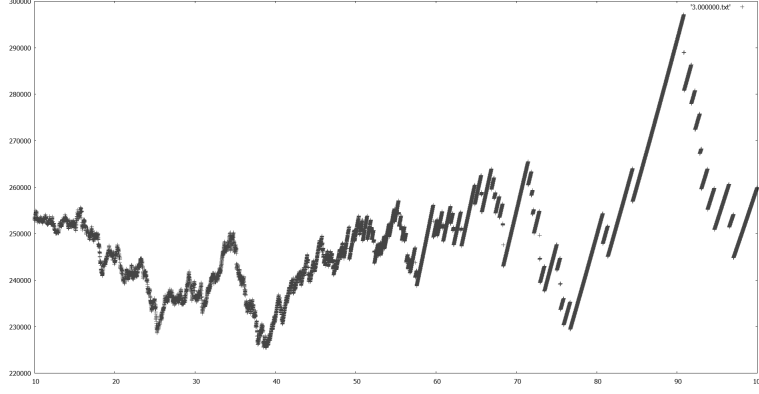


Figure 7: Here  $D = 3.0$ . Fractionality is  $\alpha \sim 0$ . The  $x$  axis is distance  $r$ , and the  $y$  axis is intersection count times inverse falloff  $r^2$ . Note that the output appears to be fractal in nature – non-differentiable anyway.

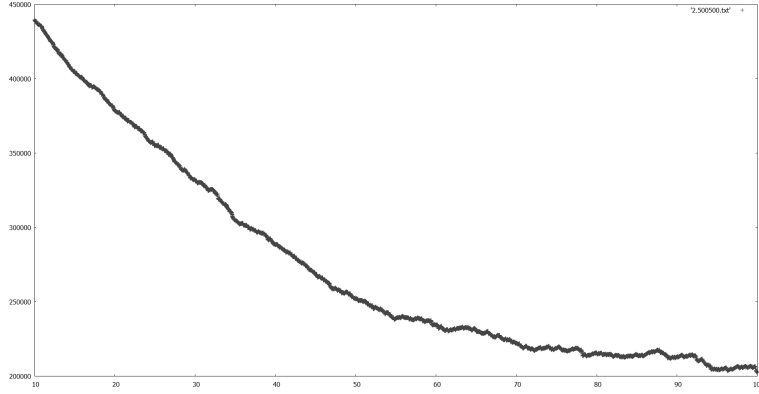


Figure 8: Here  $D = 2.5$ . Fractionality is  $\alpha \sim 1$ . The  $x$  axis is distance  $r$ , and the  $y$  axis is intersection count times inverse falloff  $r^{1.5}$ . Note that the output appears to be differentiable.

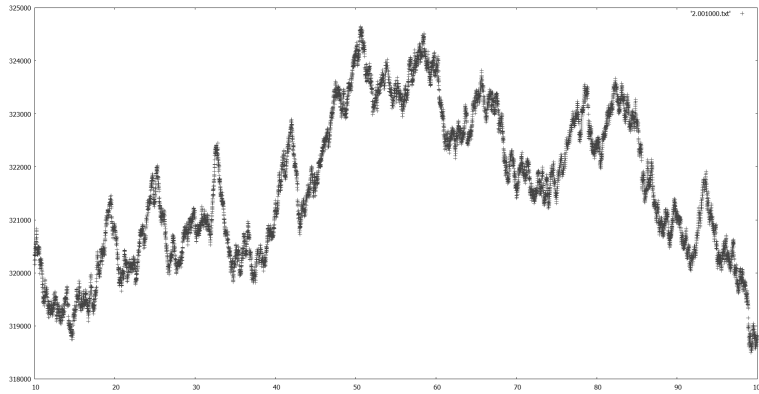


Figure 9: Here  $D = 2.001$ . Fractionality is  $\alpha \sim 0$ . The  $x$  axis is distance  $r$ , and the  $y$  axis is intersection count times inverse falloff  $r^1$ . Note that the output appears to be fractal in nature – non-differentiable anyway. And so, it appears that the greater the fractionality, the more differentiable the result.