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### Reading Assignment 6

5.12: Using the fact that  $S$  is a state function to determine the dependence of  $S$  on  $V$  and  $T$

$$dS = \frac{1}{T} du + \frac{P}{T} dv$$

-  $B/C$ ,  $1/T$ , and  $P/T$  are greater than zero, the entropy of a system increases with the internal energy at constant volume + increases with the volume at constant internal energy.

- We first write the total differential  $dS$  in terms of the partial derivatives with respect to  $V$  and  $T$

$$dS = \left(\frac{\partial S}{\partial T}\right)_V dT + \left(\frac{\partial S}{\partial V}\right)_T dV$$

- To evaluate  $(\partial S / \partial T)_V$  and  $(\partial S / \partial V)_T$  for  $dS$ :

$$dS = \frac{1}{T} \left[ C_V dT + \left(\frac{\partial U}{\partial V}\right)_T dV \right] + \frac{P}{T} dV = \frac{C_V}{T} dT + \frac{1}{T} \left[ P + \left(\frac{\partial U}{\partial V}\right)_T \right] dV$$

$$\left(\frac{\partial S}{\partial T}\right)_V = \frac{C_V}{T} \text{ and } \left(\frac{\partial S}{\partial V}\right)_T = \frac{1}{T} \left[ P + \left(\frac{\partial U}{\partial V}\right)_T \right]$$

- The expressions for  $(\partial S / \partial V)_T$  is not a form that allows for a direct comparison with experiment to be made

$$\left(\frac{\partial}{\partial T} \left(\frac{\partial S}{\partial V}\right)\right)_V = \left(\frac{\partial}{\partial V} \left(\frac{\partial S}{\partial T}\right)_V\right)_T$$

- Taking the mixed second derivative

$$\left(\frac{\partial}{\partial V} \left(\frac{\partial S}{\partial T}\right)_V\right)_T = \frac{1}{T} \left(\frac{\partial}{\partial V} \left(\frac{\partial U}{\partial T}\right)_V\right)_T$$

$$\left(\frac{\partial}{\partial T} \left(\frac{\partial S}{\partial V}\right)_T\right)_V = \frac{1}{T} \left[ \left(\frac{\partial P}{\partial T}\right)_V + \left(\frac{\partial}{\partial T} \left(\frac{\partial U}{\partial V}\right)_T\right)_V \right] - \frac{1}{T^2} \left[ P + \left(\frac{\partial U}{\partial V}\right)_T \right]$$

Simplify

$$P + \left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial P}{\partial T}\right)_V$$

A practical equation is obtained for the dependance of entropy on volume under constant temp conditions:

$$\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V = \frac{\left(\frac{\partial V}{\partial T}\right)_P}{\left(\frac{\partial V}{\partial P}\right)_T} = \frac{\beta}{\alpha}$$

The result of these considerations is that  $dS$  can be expressed in terms of  $dT$  and  $dV$  as:



$$dS = \frac{C_V}{T} dT + \frac{B}{X} dV$$

$$\Delta S = \int_{T_i}^{T_f} \frac{C_V}{T} dT + \int_{V_i}^{V_f} \frac{B}{X} dV$$

### 5.13: The dependence of S on T and P

The total differential  $dS$  is written in the form

$$dS = \left( \frac{\partial S}{\partial T} \right)_P dT + \left( \frac{\partial S}{\partial P} \right)_T dP$$

- starting from the relation  $U = H - PV$  total differential  $dU$ :

$$dU = TdS - PdV = dH - PdV - VdP$$

$$dS = \frac{1}{T} dH - \frac{V}{T} dP$$

$$dH = \left( \frac{\partial H}{\partial T} \right)_P dT + \left( \frac{\partial H}{\partial P} \right)_T dP = C_P dT + \left( \frac{\partial H}{\partial P} \right)_T dP$$

$$dS = \frac{C_P}{T} dT + \frac{1}{T} \left[ \left( \frac{\partial H}{\partial P} \right)_T - V \right] dP = \left( \frac{\partial S}{\partial T} \right)_P dT + \left( \frac{\partial S}{\partial P} \right)_T dP$$

The coefficient of  $dT$  and  $dP$  must be same on both sides:

$$\left( \frac{\partial S}{\partial T} \right)_P = \frac{C_P}{T} \quad \text{and} \quad \left( \frac{\partial S}{\partial P} \right)_T = \frac{1}{T} \left[ \left( \frac{\partial H}{\partial P} \right)_T - V \right]$$

Just as in our evaluation of  $(\partial S / \partial V)_T$  we equate the mixed second partial derivatives of  $(\partial S / \partial T)_P$  and  $(\partial S / \partial P)_T$ :

$$\left( \frac{\partial}{\partial T} \left( \frac{\partial S}{\partial P} \right)_T \right)_P = \left( \frac{\partial}{\partial P} \left( \frac{\partial S}{\partial T} \right)_P \right)_T$$

There mixed partial derivative can be evaluated

$$\left( \frac{\partial}{\partial P} \right)_T \left( \frac{\partial S}{\partial T} \right)_P = \frac{1}{T} \left( \frac{\partial C_P}{\partial P} \right)_T = \frac{1}{T} \left( \frac{\partial}{\partial P} \left( \frac{\partial H}{\partial T} \right)_P \right)_T$$

$$\left( \frac{\partial}{\partial T} \right)_P \left( \frac{\partial S}{\partial P} \right)_T = \frac{1}{T} \left[ \left( \frac{\partial}{\partial T} \left( \frac{\partial H}{\partial P} \right)_T \right)_P - \left( \frac{\partial V}{\partial T} \right)_P \right] - \frac{1}{T^2} \left[ \left( \frac{\partial H}{\partial P} \right)_T - V \right]$$

$$\frac{1}{T} \left( \frac{\partial}{\partial T} \left( \frac{\partial H}{\partial P} \right)_T \right)_P = \frac{1}{T} \left[ \left( \frac{\partial}{\partial T} \left( \frac{\partial}{\partial T} \left( \frac{\partial H}{\partial P} \right)_T \right)_P - \left( \frac{\partial V}{\partial T} \right)_P - \frac{1}{T^2} \left[ \left( \frac{\partial H}{\partial P} \right)_T - V \right] \right]$$

Simplify

$$\left( \frac{\partial H}{\partial P} \right)_T - V = -T \left( \frac{\partial V}{\partial T} \right)_P$$

$$\left( \frac{\partial S}{\partial P} \right)_T = - \left( \frac{\partial V}{\partial T} \right)_P = -V\beta$$

Using these results, the total differential  $dS$  can be written in terms of experimentally accessible parameter as:



$$ds = \frac{C_p}{T} dT - V \beta \Delta P$$

integrating both sides:

$$\Delta s = \int_{T_i}^{T_f} \frac{C_p}{T} dT - \int_{P_i}^{P_f} V \beta \Delta P$$

### 6.1 The Gibbs Energy and the Helmholtz Energy

The fundamental expression governing spontaneity is the Clausius inequality, written in the form.

$$T ds \geq dq$$

The equality is satisfied only for a reversible process b/c  $dq = du - dw$

$$T ds \geq du - dw, \text{ or } -du + dw + T ds \geq 0$$

It's useful to distinguish between expansion work, in which work arises from a volume change in system nonexpansion work

$$-du - P_{\text{external}} dV + dw_{\text{nonexpansion}} + T ds \geq 0$$

For isothermal processes,  $T ds = d(Ts)$

$$-du + T ds \geq -dw_{\text{expansion}} - dw_{\text{nonexpansion}} \text{ or}$$

$$d(u - Ts) \leq dw_{\text{expansion}} + dw_{\text{nonexpansion}}$$

The combination of state function  $u - Ts$ , which has the units of energy, defines a new state function that we call the Helmholtz energy, abbreviated  $A$ .

$$dA = dw_{\text{expansion}} - dw_{\text{nonexpansion}} \leq 0$$

$$dw_{\text{total}} = dw_{\text{expansion}} + dw_{\text{nonexpansion}} \geq dA$$

If nonexpansion work also is not possible in the transformation  $dw_{\text{nonexpansion}} = dw_{\text{expansion}} = 0$  the condition that defines spontaneity and equilibrium becomes

$$dA \leq 0$$

using the relation  $H = U + PV$

$$d(U + PV - Ts) = d(H - Ts) \leq dw_{\text{nonexpansion}}$$

The combination of state functions  $H - Ts$ , which has the units of energy, defines a new state function called the Gibbs energy,  $G$

$$dG - dw_{\text{nonexpansion}} \leq 0$$

$$dG \leq 0$$

Clausius inequality

$$ds - \frac{dq}{T} \geq 0$$

$$ds + ds_{\text{surroundings}} \geq 0$$



Maximum nonexpansion work: can be produced by a chemical transformations for macroscopic changes at constant  $P$  and  $T$  in which no nonexpansion work is possible, the condition for spontaneity is  $\Delta G_R < 0$  where:

$$\Delta G_R = \Delta H_R - T \Delta S_R$$

- The entropic contribution to  $\Delta G_R$  is greater for higher temp
- The chemical transformation is always spontaneous if  $\Delta H_R < 0 + \Delta S_R > 0$
- The chemical transformation is always nonspontaneous if  $\Delta H_R > 0 + \Delta S_R < 0$
- The relative magnitudes of  $\Delta H_R + T \Delta S_R$ , determine if chemical trans. is spontaneous
- If the chemical reaction is not spontaneous, then the reverse process is spontaneous
- If  $\Delta G_R = 0$ , the reaction mixture is at equilibrium, + neither direction of  $\Delta$  is spontaneous
- For macroscopic changes at constant  $V$  +  $T$  in which no nonexpansion work is possible, the condition for spontaneity is  $\Delta A_R < 0$  where

$$\Delta A_R = \Delta U_R - T \Delta S_R$$

b.2: The differential forms of  $U, H, A$ , and  $G$

- Differential forms are essential in calculating how  $U, H, A, G$  vary with state variable such as  $P$  and  $T$ . Starting from the definitions:

$$H = U + PV$$

$$A = U - TS$$

$$G = H - TS = U + PV - TS$$

The following total differentials can be formed:

$$dU = Tds - PdV$$

$$dH = Tds - PdV + PdV + VdP = Tds + VdP$$

$$dA = Tds - PdV - Tds - sdT = -sdT - PdV$$

$$dG = Tds + VdP - Tds - sdT = -sdT + VdP$$

Although other combinations variables can be used, these natural variables are used b/c the differential expression are compact

$$dU = Tds - PdV = \left(\frac{\partial U}{\partial s}\right)_V ds + \left(\frac{\partial U}{\partial V}\right)_s dV$$

$$\left(\frac{\partial U}{\partial s}\right)_V = T \text{ and } \left(\frac{\partial U}{\partial V}\right)_s = -P$$

$$\left(\frac{\partial H}{\partial s}\right)_P = T \text{ and } \left(\frac{\partial H}{\partial P}\right)_s = V$$

$$\left(\frac{\partial A}{\partial T}\right)_V = -S \text{ and } \left(\frac{\partial A}{\partial V}\right)_T = -P$$

$$\left(\frac{\partial G}{\partial T}\right)_P = -S \text{ and } \left(\frac{\partial G}{\partial P}\right)_T = V$$

These expressions states how  $U$ ,  $H$ ,  $A$  and  $G$  vary with their natural variables