Natalie Williams The Born-Oppenheimer Approximation - 4.2.3 This approximation makes it possible to separate the motion of the nuclei Cneutrons & protons) from the motion of the electrons. The chicleons are 4800x larger than electrons, so this approximation tales into account the nucleus moving slower than electrons Electronic Schrödinger equation - 5 45 + 5 Est + 5 Es The Secular Equation - 4.3.2 We want to want the best one-election wavefunction for a molecular System. We first want to select N basis functions, then find N2 values for Hij and Siji. Form the the secular determinant to find IV nods Ej of the equation. Then lastly, for all of the N valves of E; , solve the set of linear equations to Rnothe basis set cooppicients ais for the orbital over 12 box integral

A. 4.5.1. Hartree-product Wave Functions

The approximation talked close saccester a goldeller one-electron model of a Hamiltonian. The In a multiplied to each other to resemble a wavefunction for a multiplied with electron system.

(4.35) (YHP= Y, Y2... YN

In short, the energy of the Ymp is the sum of the one electron values. Does not take into account attraction/repulsion forces accordage

4.5.2 The Hartee Flamilitoniana

This approximation uses the ideas of the Hartreproduct wave Rinction but takes into account electron-electron interactions. Equation (4.39) is a more better approximation than (4.35) in the text.

E = = = = = 5 [ 1 vil2 1 vil2 dridg

1 y 2 is the propability density · i +; are electrons

\* 8 4.5.3 Election Spin & Antisymmetry

This approximation deals with involving spin eigenfunctions.

It uses the spin quantum number, and also involver with tells us no two electrons uses the same quantum number.

In different These principles provides the basis for what we know orbits with the Pauli Exp exclusion principle. So when taking into spin acrount spin, we use (4.43) to satisfy the Paulie Expectation principle.

State and the spin electron 1 electron 2

State and 12 [N(1) a (1) y (2) a (2) - Y2 (2) a (2) Y6(1) a)

COLPHA 15

EXECUTION 1

EXECUTION 2

EXECUTION 1

EXECUTION 1

EXECUTION 1

EXECUTION 2

EXECUTION 1

EXECUTION 2

EXECUTION 1

EXECUTION 2

EXECUTION 2

EXECUTION 3

EXECUTI

## & 4.5.4 Slater Determinants

The State Determinant takes into account to the wavefulling when you have electrons the different spins in an energy orbital. In (4.50) differs the from (4.43) for this orbital. In (4.50) differs the from (4.43) for this or (eason. White you solve by the wave function, you integrate each wave function, you are able to solve for the loss of variable kab, the Coulomb repulsion. We find this to be less than for the approximation above because electrons of the Same spin repel each other more than electrons of different spin.

Ad+ B different spins

## X 4.5.5 The Hartree-Fock Self-consistent Field Method

This approximation takes into account the interactional interaction of etections, including the exchange effects of Coulomb repulsion. It's main drawback is its one-election nature and other election interactions. It is still flawed in its approximation but it provides a stepping store until we have a better approximation to solve the schröding equation.

Flow Chart of HF SCF procedure

thatia good

The electron

The electron  $f_{i}^{2} = -\frac{1}{2} \nabla_{i}^{2} - \frac{Z_{i}}{Z_{i}} + V_{i}^{2} + \sum_{k=1}^{2} \frac{Z_{k}}{C_{ik}} + V_{i}^{2} + \sum_{k=$