Comparison of Nonlinear Filtering Methods for Battery State of Charge Estimation

A Thesis

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Abstract

In battery management systems, the State of Charge, obtained from voltage and current measurements, is a figure of merit. The estimation of State of Charge is challenging due to the nonlinear behavior of the battery, measurement noise, and the trade-off between accuracy and computational complexity in selecting the sample rate. Additionally, for the purposes of system design, an electrical-circuit battery model is useful, which presents additional filtering difficulties when compared to presently-used analytical models. This thesis investigated the performance of various nonlinear filters for estimating the State of Charge using an electrical-circuit battery model.

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Chapter 1

Introduction

Batteries, particularly rechargeable ones, are used extensively in daily life. They provide the energy for such electrical systems as communication, automotive, and renewable power systems. In order to design for and operate these systems, an accurate battery model and a means of simulating the model efficiently is needed. For example, modern battery charge and health management schemes use high-fidelity battery models to track the state of charge (SOC) and state of health (SOH); this information is then used to predict and optimize the runtime of the battery. However, widely-used chemical batteries have nonlinear capacitive effects, which require the use of a nonlinear filter for accurate prediction of its states in the presence of noise. This thesis explores one possible solutions to this problem by choosing an appropriate battery model and testing the accuracy of various nonlinear filters in determining the SOC through simulation.

1.1 Electrical Characteristics of Rechargeable Batteries

A high-fidelity battery model has to accurately reproduce the various characteristics of a battery. Most models keep track of the total capacity and SOC in order to predict remaining runtime. More accurate models include nonlinear effects, such as the rate-capacity effect and the recovery effect, along with self-discharge and the effects of ambient temperature. The dynamic electrical attributes, such as the current-voltage (i-v) characteristics and transient responses, can also be modeled. The remainder of this section defines these characteristics.

The capacity of a battery is the amount of electric charge it can store, measured in the SI unit Ampere-hours (Ah). Commonly, for rechargeable battery specifications, the subunit milliampere-hour (mAh) is used. Related is the available capacity, which is the amount of charge available for use. Due to the electrochemical nature of batteries, a battery's available capacity decreases as the rate of discharge increases, known as the rate-capacity effect. Therefore, the capacity for a battery is typically stated for a given discharge rate. Related to this is the recovery effect, so called because when a battery is allowed to rest during an idle period, the battery "recovers" available capacity previously lost during discharge due to the rate-capacity effect. Thus, a battery discharged at a high rate until seemingly fully discharged, when allowed to rest, regains a portion of its lost capacity.

Both the rate-capacity effect and the recovery effect can be explained by the electrochemical nature of the battery. During discharge, the concentration of the active material around the electrode is depleted, and the active materials in the depletion region move towards the electrode to reduce the concentration gradient [1]. Because the speed at which the concentration gradient is equalized is limited, the

faster the rate of discharge, the less the active material is replenished, resulting in a decrease in the available capacity. Likewise, when the battery is allowed to rest, the active material gradient has additional time to equalize, and the available capacity is increased.

Closely related to the capacity is the SOC. This thesis defines it as the ratio between the remaining capacity and the maximum capacity, with both capacities measured using the amount of active material within the battery. This definition then denotes the proportion of remaining chemical energy rather the available energy and is unaffected by the rate-capacity and recovery effects. Note that a fully charged battery has an SOC of unity and a fully discharged battery has an SOC of zero, regardless of the available capacity. Additionally, it is convenient to establish the relationship between the SOC of the battery and its open-circuit voltage V_{OC}, which is useful for simulation of the i-v characteristics and transient responses. V_{OC} can be thought as the limit of the measured battery voltage after recovery.

Other more minor effects that are usually incorporated into models are self-discharge, the effect of ambient temperature, and aging. Self-discharge refers to an idle battery decreasing its SOC over time due to internal chemical reactions. It is dependent on the type of battery, SOC, ambient temperatures, and other factors. The ambient temperature has effects on the internal resistance of the battery and the self-discharge rate. Commonly, the battery is designed to operate with a narrow range of temperatures. Below the operating temperature range, the internal resistance increases, decreasing the capacity. Above the operating range, the internal resistance decreases, not only increasing the capacity but also the self-discharge rate; thus, the actual capacity is lowered due to the increased self-discharge. Aging refers to the decrease in battery performance measures, such as capacity, self-discharge, and internal resistance, over time due to unwanted chemical reactions. In practice,

aging is indicated by the SOH, defined as the ratio between the current maximum capacity and that of a new battery. The SOH threshold at which the battery performance is considered too degraded varies by application.

1.2 Battery Models

This thesis studied the predition of the SOC of a battery given knowledge of the resistive load on the battery as well as noisy measurements of the voltage across its terminals. In order to do so for a general load profile, incorporation of the rate-capacity and recovery effects as well as the transient i-v characteristics is desirable. Furthermore it is useful to have a model easily tunable for different battery types. This section reviews the characteristics of major types of battery models. Battery models can be divided into five categories, namely electrochemical, computational intelligence, analytical, stochastic, and electrical-circuit. The remainder of this section reviews each type and determines the most suitable battery model for this study.

1.2.1 Electrochemical

Electrochemical models are describe the chemical processes that place in the battery in great detail. These are generally the most accurate, but they require in-depth knowledge of the chemical processes to create and impose large computational costs [2]. One of the most widely known electrochemical models was developed by Doyle, Fuller, and Newman for lithium and lithium-ion batteries using noninvasive voltage-current cycling experiments [3–5]. It consists of six coupled, nonlinear differential equatons that capture lithium diffusion dynamics and charge transfer kinetics. The model is able to predict i-v response and provides a design guide for thermody-

namics, kinetics, and transport across electrodes. A implementation of their model in Fortran, called Dualfoil, is available for free online.[1] The program needs more than 60 parameters along with the load profile in order to compute the battery properties. Setting the parameters requires detailed knowledge of the battery, but the result of the program is highly accurate. Other battery models are often compared to it rather than to experimental results.

1.2.2 Computational Intelligence

Computational intelligence is a brance of computer science interested in problems that require the intelligence of humans and animals to solve. One of the earliest definitions by Bezdek states that computational intelligent systems use pattern recognition on low-level, numerical data and do not use knowledge as with artificial intelligence [7, 8]. Methods such as neural networks, fuzzy systems, and evolutionary computation are commonly classified as computational intelligence. Battery models using such methods as neural networks [9, 10], support vector machines [11], and hybrid neural-fuzzy models [12] have been studied. These models learn the nonlinear relationships between battery properties, such as SOC, current, voltage, and temperature, through a computationally costly training process. However, once trained, they incur a much lower cost and can achieve comparable accuracy to electrochemical models.

1.2.3 Analytical

Analytical models are simplified electrochemical models that trade off accuracy for simplicity. One of the simplest such models is Peukert's law for lead-acid batteries,

[1] J. Newman, *Fortran programs for the simulation of electrochemical systems*, `http://www.cchem.berkeley.edu/jsngrp/fortran.html`, 1998.

which states that for a one-ampere discharge rate [13]

$$C_p = I^k t, \tag{1.1}$$

where C_p is the capacity at a one-ampere discharge rate in Ah, I is the discharge current in A, t is the time to discharge the battery in hours, and $k \geq 1$ is the dimensionless Peukert constant, typically between 1.1 and 1.3 for a lead-acid battery. The constant k only equals unity for an ideal accumulator, so for real batteries, k is always greater than unity. Thus, for a given increase in the discharge current, the discharge time decreases by a proportionally greater amount. Therefore, the effective, or available, capacity $C \times t$ is reduced. Peukert's law can be extended to some other battery chemistries, such as lithium-ion. Note that Peukert's law only models the rate-capacity effect and not the recovery effect. More complicated models, such as the kinetic battery model and the diffusion model, are able to describe both effects.

The kinetic battery model (KiBaM), initially created for large lead-acid batteries, describes the battery as a kinetic process, using two charge wells for the bound and available charges connected by a valve whose flow rate is proportional to the height difference between the wells [14]. The change of charge in the wells is given by

$$\begin{cases} \dfrac{dy_1}{dt} = -I + k\left(h_2 - h_1\right) \\ \dfrac{dy_2}{dt} = -k\left(h_2 - h_1\right), \end{cases} \tag{1.2}$$

where y_1, y_2 are the charges, h_1, h_2 are the heights of the wells, the parameter k controls the rate of charge flow between the wells, and I is the applied load. The flow rate of the valve should be lower than the typical discharge rate of the battery.

During discharge from the available-charge well, the bound charges flow through the valve to equalize the heights of the two wells. It can be seen that for slower discharge rates, more charge flows through the valve and the effective capacity increases. Likewise, during idle periods for the battery, the available charge increases.

Related to the KiBaM is the diffusion model, which describes the movement of the ions in the electrolyte of a lithium-ion battery [15]. Like in the kinetic battery model, the difference in the concentration of adjacent ions along the length of the battery determines the diffusion rate of the ions. The available charges are those ions directly touching the electrode of the battery. It can be seen that the KiBaM is a first-order approximation of the diffusion model [2], since the individual ions in the diffusion model are replaced by two charge wells in the KiBaM.

1.2.4 Stochastic

Stochastic models describe the discharging and the recovery effect as stochastic processes. The first models were developed by Chiasserini and Rao and based on discrete-time Markov chains [16]. They studied two models of a battery in a communication device that transmitted packets. The simpler model described the battery as a discrete-time Markov chain with $N + 1$ states, numbered from 0 to N and corresponding to the number of charge units available in the battery. Transmitting one packet requires one charge unit of energy. Thus, in continuous transmission, N packets can be sent. At every time step, a charge unit is either consumed with probability $a_1 = q$ or recovered with probability $a_0 = 1 - q$. The battery is considered empty when the 0 state is reached or when a theoretical maximum of T charge units have been consumed. The second model is an extension of the first, allowing for more than one charge unit to be consumed in a time step, modeling more bursty

usage. Additionally, the battery has a non-zero probability of staying in the same charge state, indicating no consumption or recovery during a time step. Chiasserini and Rao extended their model further in following papers by adding state and phase dependence [1, 17, 18]. The state number is the number of charge units, and the phase number is the number of consumed charge units. Having fewer charge units decreases the probability of recovery, while having more consumed charge units increases the probability of recover. Using these models, one can model different loads by setting the transitions probabilities. However, the order of the transitions is uncontrollable, so it is impossible to model fixed load patterns and compute their impact on battery life.

Chiasserini and Rao mainly investigated the gain G in transmitted packets using a pulsed discharge relative to using a constant discharge, defined as $G = m/N$, where m is the mean number of transmitted packets. The gain increases when the load decreases, due to an increase in the recovery probability. Additionally, the gain increases for lower discharge demand rates and higher current densities. These load profiles result in discharge currents close to the specified limits of the battery, causing the available capacity to decrease overly quickly. Therefore, the recovery effect is especially strong for these cases during pulsed discharge, greatly increasing the gain. Chiasserini and Rao compared the computation of the gain parameter for different current densities and demand rates using the stochastic model to that of the electrochemical model of Doyle et al. They found an average deviation of 1% and a maximum deviation of 4%. This shows that the stochastic model accurately describes battery behavior during pulsed discharge. However, this model is only able to compute relative lifetimes.

In 2005, Rao et al. [19] proposed a stochastic battery model based on the Kinetic Battery Model (KiBaM) of Manwell and McGowan. This stochastic KiBaM was for

a nickel-metal hydride (NiMH) battery. The differential equations governing the original KiBaM were modified to include an extra factor h_2 governing the flow of charge between the wells. This changes Eq. (1.2) into

$$\begin{cases} \dfrac{dy_1}{dt} = -I + k_s h_2 \left(h_2 - h_1\right) \\ \dfrac{dy_2}{dt} = -k_s h_2 \left(h_2 - h_1\right), \end{cases} \tag{1.3}$$

This change causes the recovery effect to weaken as the remaining charge decreases. The stochastic model was also modified to allow the possibility of no recovery during idle periods. The stochastic KiBaM describes the battery using a discrete-time, transient Markov process. The states are labeled with the parameters (i, j, t), with i and j representing the discrete charge levels of the available and bound charge wells and t representing the length of the current idle period. Like the stochastic model of Chiasserini and Rao, it is impossible to fully model a real-life discharge pattern using the stochastic KiBaM. Rao et al. compared the results of their model with experimental results using an AAA NiMH battery. Two sets of experiments were conducted, the first with varying frequency of the load and a 50% duty cycle and the second with varying off-time and a constant on-time. Their model accurately predicted the lifetime and delivered charge from the battery, with a maximum error of 2.65%.

1.2.5 Electrical-Circuit

Electrical-circuit models for batteries developed from the discovery of capacitative effects at the electrode-electrolyte interface. Helmholtz first proposed the existence of a double layer of charge at the interface in 1879. In 1899, Warburg proposed a series resistance and capacitance circuit model with an infinitely low current density.

The Warburg capacitance C_W named after him varies inversely with the square root of the frequency [20]. In 1947, Randles proposed a model consisting of a double-layer polarization capacitance C_p in parallel with the series combination of a resistor R and a capacitance C [21]. In 1994, Kovacs improved Randles circuit with the addition of Warburg impedence Z_W replacing the capacitance C and the solution resistance R_s in series with the original Randles circuit [22]. In addition, he renamed C_p to the double layer capacitance C_{dl} and R to the charge-transfer resistance R_{ct}. These proposals came from a desire to represnt impedance spectra created using electrochemical impedance spectroscopy (EIS). The various elements in the models represent the different processes within a battery, which have different time constants. While these attempts model the impedance and, thus, account for the nonlinear rate-capacity and recovery effects, they do not consider the capacity and self-discharge of the battery.

In 1993, Hageman created simplified electrical-circuit models using PSpice for nickel-cadmium (NiCd), lead-acid, and alkaline batteries [23]. The circuits shared the common elements of *i*) a capacitor that represents the battery capacity, *ii*) a discharge rate normalizer that determines the additional capacity loss at high discharge rates, *iii*) a circuit that discharges the battery, *iv*) a lookup table of battery voltage versus SOC, and *v*) a resistor that represents the battery's internal resistance [23, 24]. In addition, battery models for NiCd batteries simulated the thermal effects under high discharge rates. The main lookup table is formed by discharging a battery at a low rate at a constant current (20 to 200 hours). At high discharge rates, the discharge rate normalizer reduces the battery voltage below the value from looking up the SOC in the table. This normalizer is implemented using additional lookup tables. These circuit models are much simpler than electrochemical models, but they are also less accurate with an approximate error of 10%. Furthermore,

creation of the lookup tables requires considerable data. These circuit-based models are mainly concerned with modeling the remaining discharge time and are referred to as runtime-based models.

In 2006, Chen and Rincón-Mora proposed a combination of a runtime-based model and a impedance-based model consisting of a series resistor and two parallel resistor-capacitor networks [25]. The elements of the impedance part of the model had parameters that depended on the SOC. Additionally, the runtime model included a resistance that modeled the self-discharge rate. Their proposed model has the advantage of accurate prediction of the SOC using the runtime-based portion while also modeling nonlinear transient effects, such as the rate-capacity and recover effects, with the impedance-based portion. Furthermore, the battery data can be collected using EIS measurements, which requires neither detailed knowledge of the battery chemistry nor lengthy, low-rate discharge experiments.

1.2.6 Evaluation

Of the model types, only some are fit for use with filtering algorithms. The computational-intelligence and stochastic models do not model the dyanmics of the battery system, so they cannot be used. The electrochemical models, the related analytical models, and the electrical-circuit models describe the system dynamics and the nonlinear rate-capacity and recovery effects. However, only the electrical-circuit model has the advantage of modeling the internal impedance of the battery, which is useful in the design of battery systems. Of the electrical-circuit models, the proposal by Chen and Rincón-Mora is best suited for the purposes of this thesis since it is the only one discussed by this paper that describes both the capacity and the transient effects. Therefore, their proposal is used for the simulation of the battery and comparing

the performance of different filters.

1.3 Nonlinear Filtering Methods

Filtering refers to the methodology for estimating the state of a time-varying system that is indirectly observed through noisy measurements. Specifically, the current state is estimated from the current and previous measurements. The state of a system is a group of dynamic variables that evolve through time, and its evolution through time is governed by a dynamic system, perturbed by process noise. The measurements are functions of the state and measurement noise. A battery can be modeled as a time-varying system, with state variables that describe such states as the SOC and the SOH. The measurements are typically the voltage and the current. Note that this thesis uses a resistive load instead of the current, because this is more realistic from a usage standpoint. Additionally, the SOH is not considered by this thesis. The dynamic system for the battery is described in more detail in the following chapter.

For linear systems, the optimal filtering solution with respect to the minimum mean squared error (MMSE) is given by the least squares solution, meaning the optimal least squares solution equals the posterior mean. A closed form solution to the discrete-time linearing filtering problem is given by the Kalman filter [26], which is a linear MMSSE (LMMSE) filter. Under the assumption that the noises are Gaussian, the posterior distribution is also Gaussian and numerical approximations are unnecessary. The Kalman filter has a prediction phase and an update phase. The prediction phase estimates the current state using the state estimate from the previous time step. The update phase refines the state estimate using the current measurement.

For nonlinear systems, optimal filtering solutions are generally intractable, so various numerical approximation methods have been developed, mainly classified in three groups: function approximation, moment approximation, and stochastic model approximation [27]. Function approximation techniques approximate the nonlinear dynamic and measurement processes, commonly using the Taylor series expansion. A moment approximation technique uses a representative set of sample points and calculates integrals as weighted averages. Stochastic model approximation simplifies the original nonlinear stochastic system to a linear system so that linear filtering results are applicable.

Note that the choice of discretization method tends to depend on the filter, especially for highly nonlinear systems. Therefore, the simulations used a continuous-discrete (mixed) time system, with the form

$$\dot{\mathbf{x}}(t) = \mathbf{f}(\mathbf{x}(t), \mathbf{u}(t), \mathbf{w}(t)) \tag{1.4}$$

$$\mathbf{z}_k = \mathbf{h}(\mathbf{x}(t_k), \mathbf{u}(t_k), \mathbf{v}(t_k)), \tag{1.5}$$

where $\mathbf{x}$, $\mathbf{u}$, $\mathbf{z}$, $\mathbf{w}$, and $\mathbf{v}$ are the vectors of the states, the inputs, the measurements, the process noises, and the measurement noises, respectively. For simplicity of simulation, the input $\mathbf{u}(t)$ is assumed to be constant over each time step so that $\mathbf{u}(t) = \mathbf{u}(t_k)$ for $t_k \leq t < t_{k+1}$. For the prediction phase, the continuous-time dynamics were solved numerically, with each filter using different approximation techniques. The state error covariances are also approximated numerically. Then, the updates were performed based on the update of the linear Kalman filter.

Furthermore, note that the system is extremely stiff for some inputs $\mathbf{u}$, which can be seen from the stiffness ratio based on the ratio of the largest eigenvalue of F to its smallest eigenvalue, where F is the jacobian of the dynamics $\mathbf{f}(\mathbf{x}, \mathbf{u})$; when the

stiffness ratio is much greater than unity, the system is stiff [28, 29]. From EIS studies of batteries, the major chemical processes have widely differing time constants; low frequency mass transport effects like diffusion are on the order of 10^{-6} to 10^{0} Hz, middle frequency effects caused by charge transfer and the electrochemical double layer are on the order of 10^{0} to 10^{3} Hz, and the high frequency conductance and skin effects are on the order of 10^{3} to 10^{4} Hz [30]. Therefore, the approximate stiffness ratio is $10^{10} \gg 1$, and the system is stiff. As a result of the stiffness, any numerical integration method needs to be A-stable, i.e. the method converges for all systems whose eigenvalues have negative real parts. For example, simulation results show that the fourth-order Runge-Kutta method diverges even at step sizes $< 10^{-2}$ seconds.

Särkkä and Solin state that a linearized discretion approach, in which the continuous-time system is first discretized and then approximated as Gaussian, tends to work better than a discretized linearization approach, in which the system is first approximated as a Gaussian process and then discretized [31]. This thesis follows this guideline and performs the prediction using linearized approximations of a discretization of the continuous-time dynamics. The specifics of the linearized discretion approach for each filter along with their general implementation are discussed in the remainder of this chapter.

1.3.1 Function Approximation

One of the most popular nonlinear filters is the extended Kalman filter (EKF), which approximates the nonlinear state and measurement equations, typically using Taylor series expansion. The standard Kalman filter formulas are used for the update. For the discrete-time EKF, Taylor series expansion can be directly used on the system

dynamics to linearize them. For the mixed-time system of this study, a linearized discretization approach proposed by Mazzoni [32] is used, in which the dynamics are first discretized and then approximated using Taylor series expansion. This approach has the advantage of A-stability. The discretization is performed using the trapezoidal approximation (Heun's method) of Eq. (1.4). For convenience, denote the value of a quantity at time t_k using the subscript k and assume that $\delta = t_{k+1} - t_k$ is the time step. Additionally, a subscript of $k|k-1$ indicates the value at time t_k given the information at t_{k-1}. Then, the approximation produces

$$\mathbf{x}_{k+1} \approx \mathbf{x}_k + \frac{1}{2}\big(\mathbf{f}(\mathbf{x}_k) + \mathbf{f}(\mathbf{x}_{k+1})\big)\delta. \tag{1.6}$$

The vector field $\mathbf{f}$ at $\mathbf{x}_{k+1}$ is approximated by first-order Taylor expansion around $\mathbf{x}_k$, giving

$$\mathbf{x}_{k+1} \approx \mathbf{x}_k + \mathbf{f}(\mathbf{x}_k)\delta + \frac{1}{2}F(\mathbf{x}_k)\left(\mathbf{x}_{k+1} - \mathbf{x}_k\right)\delta, \tag{1.7}$$

where $F(\mathbf{x}_k)$ is the Jacobian of $\mathbf{f}$ at $\mathbf{x}_k$. Solving for $\mathbf{x}_{k+1}$ yields

$$\mathbf{x}_{k+1} \approx \mathbf{x}_k + \left(I - F(\mathbf{x}_k)\frac{\delta}{2}\right)^{-1}\mathbf{f}(\mathbf{x}_k)\delta, \tag{1.8}$$

with the identity matrix I. This Taylor-Heun scheme uses linear Taylor expansion of f rather than then Euler prediction of the standard Heun scheme. This numerical integration scheme is convergent with order $\mathcal{O}(\delta^2)$ and A-stable [32]. For the state error covariance ODE given by

$$\dot{P} = F(\mathbf{x})P + PF^\top(\mathbf{x}) + L(\mathbf{x})Q(\mathbf{x})L^\top(\mathbf{x}), \tag{1.9}$$

where Q is the covariance of the error variables and L is the Jacobian of $\mathbf{f}$ with respect to the state error variables $\mathbf{w}$, the numerical integration method should have the key features of being consistent with the same order as the Taylor-Heun scheme, able to process nonautonomous differential equations, and A-stable. Mazzoni proposed using a modified Gauss-Legendre formula with an implicit increment rule, producing

$$P_{k+1} \approx P_k + M_\tau \left(F(\mathbf{x}_\tau)P_k + P_k F^\top(\mathbf{x}_\tau) + L(\mathbf{x}_\tau)Q(\mathbf{x}_\tau)L^\top(\mathbf{x}_\tau) \right) M_\tau^\top \delta, \tag{1.10}$$

with

$$M_\tau = \left(I - F(\mathbf{x}_\tau)\frac{\delta}{2} \right)^{-1} \quad \text{and} \quad \tau = t_k + \frac{\delta}{2}. \tag{1.11}$$

The value of $\mathbf{x}_\tau$ can be interpolated from $\mathbf{x}_k$ and $\mathbf{x}_{k+1}$ with a precision of $\mathcal{O}(\delta^3)$ using series expansion, producing

$$\mathbf{x}_\tau \approx \frac{1}{2} \left(\mathbf{x}_k + \mathbf{x}_{k+1} - F(\mathbf{x}_k) f(\mathbf{x}_k) \frac{\delta^2}{4} \right). \tag{1.12}$$

This modified Gauss-Legendre approximation is A-stable, consistent with order $\mathcal{O}(\delta^2)$, and ensures the positive definiteness of the resultant error covariance matrix [32]. This numerical integration can be repeated multiple times with a smaller step size should greater accuracy be desired. The update equations for the EKF come from the LMMSE filter and are

$$K_k = P_{k|k-1} H_k^\top \left(H_k P_{k|k-1} H_k^\top + M(\mathbf{x}_k) R(\mathbf{x}_k) M^\top(\mathbf{x}_k) \right)^{-1} \tag{1.13}$$

$$\mathbf{x}_{k|k} = \mathbf{x}_{k|k-1} + K_k \big(\mathbf{z}_k - \mathbf{h}(\mathbf{x}_{k|k-1}) \big) \tag{1.14}$$

$$P_{k|k} = (I - K_k H_k) P_{k|k-1}, \tag{1.15}$$

where R is the covariance of the measurement error variables, M is the Jacobian of $\mathbf{f}$ with respect to the measurement error variables $\mathbf{v}$, and H_k is the Jacobian of $\mathbf{f}$ at $\mathbf{x}_{k|k-1}$.

1.3.2 Moment Approximation

In moment approximation, the mean, covariance, and possibly higher-order moments are approximated directly, rather than approximating the nonlinear integrations of expectations. Of the numerical moment approximation methods, two of the best known are the unscented Kalman filter (UKF) and the cubature Kalman filter (CKF). Note that the CKF is a special case of the UKF, but they are considered separately by this study.

In order to handle the mixed-time system, the Itô-Taylor expansion with a strong order of 1.5 (IT-1.5) proposed by Arasaratnam et al. is used to discretize the stochastic differential equations (SDE) for numerical integration [33]. In the context of Gaussian filtering, the weak order of convergence is more important, because the moments of the distributions and not the functionals of the path are of importance. Thus, the IT-1.5 expansion is equivalent to the weak Itô-Taylor expansion of order 2.0 [34], which was proposed as a SDE discretization method for target tracking applications by Morelande and Gordon [35]. For a SDE of the form

$$d\mathbf{x}(t) = \mathbf{f}(\mathbf{x}(t), t)dt + \sqrt{Q}d\beta(t), \tag{1.16}$$

Arasaratnam et al. gives the IT-1.5 expansion as [33]

$$\mathbf{x}_{k+1} = \mathbf{x}_k + \delta\mathbf{f}(\mathbf{x}_k) + \frac{\delta^2}{2}\big(\mathbb{L}_0\mathbf{f}(\mathbf{x}_k)\big) + \sqrt{Q}\mathbf{w} + \big(\mathbb{L}\mathbf{f}(\mathbf{x}_k)\big)\mathbf{y}, \tag{1.17}$$

where

$$\mathbb{L}_0 = \frac{\partial}{\partial t} + \sum_{i=1}^{n} \mathbf{f}_i \frac{\partial}{\partial \mathbf{x}_i} + \frac{1}{2} \sum_{j,p,q=1}^{n} \sqrt{Q}_{p,j} \sqrt{Q}_{q,j} \frac{\partial^2}{\partial \mathbf{x}_p \partial \mathbf{x}_q} \tag{1.18}$$

$$\mathbb{L}_j = \sum_{i=1}^{n} \sqrt{Q}_{i,j} \frac{\partial}{\partial \mathbf{x}_i} \tag{1.19}$$

and $(\mathbf{w}, \mathbf{y})$ are a pair of correlated Gaussian random variables generated from a pair of independent standard Gaussian random variables. In the absense of noise, the discretization is

$$\mathbf{x}_{k+1} = \mathbf{f}_d(\mathbf{x}_k) = \mathbf{x}_k + \delta \mathbf{f}(\mathbf{x}_k) + \frac{\delta^2}{2} \big(\mathbb{L}_0 \mathbf{f}(\mathbf{x}_k) \big). \tag{1.20}$$

It can be shown that the expected value of the state is given by

$$\mathbf{x}_{k+1} = \int_{\mathcal{R}^n} \mathbf{f}_d(\mathbf{x}_k) \mathcal{N}(\mathbf{x}_k; \hat{\mathbf{x}}_{k|k}, P_{k|k}) \, \mathrm{d}\mathbf{x}_k. \tag{1.21}$$

The propagation of the covariance matrix is given by

$$\begin{aligned} P_{k|k} = & \int_{\mathcal{R}^n} \mathbf{f}_d(\mathbf{x}_k) \mathbf{f}_d^\top(\mathbf{x}_k) \mathcal{N}(\mathbf{x}_k; \hat{\mathbf{x}}_{k|k}, P_{k|k}) \, \mathrm{d}\mathbf{x}_k + \frac{\delta^3}{3} \big(\mathbb{L}\mathbf{f}(\mathbf{x}_k) \big) \big(\mathbb{L}\mathbf{f}(\mathbf{x}_k) \big)^\top \\ & + \frac{\delta^2}{2} \Big[\sqrt{Q} \big(\mathbb{L}\mathbf{f}(\mathbf{x}_k) \big)^\top + \big(\mathbb{L}\mathbf{f}(\mathbf{x}_k) \big)^\top \sqrt{Q} \Big] - (\hat{\mathbf{x}}_{k|k})(\hat{\mathbf{x}}_{k|k})^\top + \delta Q. \end{aligned} \tag{1.22}$$

The Gaussian integrals can then be approximated using the unscented transform in the UKF and the third-degree cubature rule in the CKF. It can be seen that the IT-1.5 expansion needs to evaluate the first and second order derivatives of the dynamic system, which makes these moment approximation methods no longer derivative-free. There is also a performance hit associated with the evaluations of the derivatives. However, the advantages are greater stability of the filter and impproved numerical

accuracy, which are necessary for the stiff system.

The following two sections discuss the implementation of the UKF and CKF in mixed-time using IT-1.5 for the propagation of the dynamics.

1.3.2.1 Unscented Kalman Filter

The UKF is an efficient, generally derivative-free filtering algorithm that relies on the unscented transformation (UT). The UT is useful for forming the Gaussian approximation to the joint distribution of random variables x and y for $x \sim \mathcal{N}(m, P)$ and $y = g(x)$, where $x \in \mathbb{R}^n$, $y \in \mathbb{R}^m$, and $g : \mathbb{R}^n \mapsto \mathbb{R}^m$ is a nonlinear function. Then, the first and second moments corresponding to the mean and covariance can be easily found. Specifically, suppose the Gaussian approximation of the joint probability density of x and y has the form

$$\begin{bmatrix} x \\ y \end{bmatrix} = \mathcal{N}\left(\begin{bmatrix} m \\ \mu_U \end{bmatrix}, \begin{bmatrix} P & C_U \\ C_U^\top & S_U \end{bmatrix} \right). \tag{1.23}$$

Then, the UT picks $2n + 1$ sample points $\{x_i\}$, commonly known as sigma points, along with the same number of weights $\{w_i\}$, as follows [36]. First, the sigma points are chosen from the columns of the matrix $\sqrt{(n+\lambda)P}$, giving

$$x^{(0)} = m_x \tag{1.24}$$

$$x^{(i)} = m_x + \left[\sqrt{(n+\lambda)P}\right]_i, \quad i = 1, \ldots, n \tag{1.25}$$

$$x^{(i)} = m_x - \left[\sqrt{(n+\lambda)P}\right]_{i-n}, \quad i = n+1, \ldots, 2n \tag{1.26}$$

with the weights

$$W_0^{(m)} = \frac{\lambda}{n+\lambda} \tag{1.27}$$

$$W_0^{(c)} = \frac{\lambda}{n+\lambda} + (1 - \alpha^2 + beta) \tag{1.28}$$

$$W_i^{(m)} = W_i^{(c)} = \frac{1}{2(n+\lambda)}, \quad i = 1, \ldots, 2n. \tag{1.29}$$

The parameter λ is defined as

$$\lambda = \alpha^2(n + \kappa) - n, \tag{1.30}$$

and the constants α, β, and κ are parameters of the method. For the UKF, α is a small positive number, e.g. 10^{-3}, $\beta = 2$ is ideal for a Gaussian distribution, and κ is typically 0. Each sigma point is transformed by

$$y^{(i)} = g(x^{(i)}, \quad i = 0, \ldots, 2n. \tag{1.31}$$

Then, moments are approximated by

$$\mu_U = \sum_{i=0}^{2n} W_i^{(m)} y^{(i)} \tag{1.32}$$

$$S_U = \sum_{i=0}^{2n} W_i^{(c)} (y^{(i)} - \mu_U)(y^{(i)} - \mu_U)^\top \tag{1.33}$$

$$C_U = \sum_{i=0}^{2n} W_i^{(c)} (x^{(i)} - m)(y^{(i)} - \mu_U)^\top. \tag{1.34}$$

The square root of the positive definite matrix P is defined as a matrix A such that $P = AA^\top$. Note that A is not unique. For performance reasons, the Cholesky

factorization is typically used. The UT algorithm is denoted as

$$[\mu_U, S_U, C_U] = \mathrm{UT}(g, m, P). \tag{1.35}$$

Then, for the discretized system

$$\mathbf{x}_k = \mathbf{f}_d(\mathbf{x}_{k-1}) + \mathbf{q}_{k-1} \tag{1.36}$$

$$\mathbf{y}_k = \mathbf{h}(\mathbf{x}_k) + \mathbf{r}_k, \tag{1.37}$$

where $\mathbf{q}_k \sim \mathcal{N}(0, Q_k)$ and $\mathbf{r}_k \sim \mathcal{N}(0, R_k)$, the prediction and update steps for the UKF can be written as

$$[m_k^-, \tilde{P}_k] = \mathrm{UT}(\mathbf{f}_d, m_{k-1}, P_{k-1}) \tag{1.38}$$

$$P_k^- = \tilde{P}_k + Q_{k-1} \tag{1.39}$$

$$[\mu_k, \tilde{S}_k, C_k] = \mathrm{UT}(\mathbf{h}, m_k^-, P_k^-) \tag{1.40}$$

$$S_k = \tilde{S}_k + R_k \tag{1.41}$$

$$K_k = C_k S_k^{-1} \tag{1.42}$$

$$m_k = m_k^- + K_k(y_k - \mu_k) \tag{1.43}$$

$$P_k = P_k^- - K_k S_k K_k^\top. \tag{1.44}$$

The discretized dynamics are those from the previous section.

1.3.2.2 Cubature Kalman Filter

The CKF uses the cubature rule rather than the UT to approximate the Gaussian integrals. For the CKF, $2n$ sigma points are chosen, giving [33]

$$x^{(i)} = m_x + \left[\sqrt{nP}\right]_i, \quad i = 1, \ldots, n \tag{1.45}$$

$$x^{(i)} = m_x - \left[\sqrt{nP}\right]_{i-n}, \quad i = n+1, \ldots, 2n. \tag{1.46}$$

The weights are uniformly $W_i = 1/2n$. Thus, the moments are approximated by

$$\mu_U = \frac{1}{2n}\sum_{i=0}^{2n} y^{(i)} \tag{1.47}$$

$$S_U = \frac{1}{2n}\sum_{i=0}^{2n} (y^{(i)} - \mu_U)(y^{(i)} - \mu_U)^\top \tag{1.48}$$

$$C_U = \frac{1}{2n}\sum_{i=0}^{2n} (x^{(i)} - m)(y^{(i)} - \mu_U)^\top. \tag{1.49}$$

It can be seen that the cubature rule is a special case of the UT with $\alpha = 1$, $\beta = 0$, and $\kappa = 0$. This choice of parameters results in $W_0 = 0$, eliminating one of the sigma points. The prediction and update steps then proceed as in the UKF.

Compared to the UKF, the CKF is numerically more stable due to its positive weights. While the UKF has some desirable theoretical properties, some of its weights can be negative, causing numerical problems in some cases [31].

1.3.3 Stochastic Model Approximation

In stochastic model approximation, the nonlinear state and measurement functions are statistically linearized to minimize the MSE. Then, the resulting linear system can be filtered using the linear Kalman filter. This study refers to this method as the

statistically linearized filter (SLF). Specifically, given $\mathbf{x} \sim \mathcal{N}(m, P)$, the nonlinear function $\mathbf{f}(\mathbf{x})$ is linearized as [27, 37]

$$\mathbf{f}(\mathbf{x}) \approx \mathbf{b} + A(\mathbf{x} - m), \tag{1.50}$$

where the parameters $\mathbf{b}$ and A are chosen to minimize the error

$$\mathrm{MSE}(\mathbf{b}, A) = E\left[\|\mathbf{f}(\mathbf{x}) - \mathbf{b} - A(\mathbf{x} - m)\|^2\right]. \tag{1.51}$$

Differentiating the MSE expression and setting the derivatives to zero produces the optimal values gives

$$\mathbf{b} = E[\mathbf{f}(\mathbf{x})] \tag{1.52}$$

$$A = E[\mathbf{f}(\mathbf{x})(\mathbf{x} - m)^\top]P^{-1}. \tag{1.53}$$

These values reproduce the mean exactly but the covariance is an approximation. These expectations can be calculated analytically or numerically. Due to the difficulty of finding the analytical forms of the expectations, they are approximated numerically using the third-order cubature rule, giving

$$\mathbf{b} = E[\mathbf{f}_d(\mathbf{x})] \approx \frac{1}{2n}\sum_{i=1}^{2n} \mathbf{f}_d(\mathbf{x}^{(i)}) \tag{1.54}$$

$$A = E[\mathbf{f}(\mathbf{x})(\mathbf{x} - m)^\top]E[(\mathbf{x} - m)(\mathbf{x} - m)^\top]^{-1} = E[F(\mathbf{x})] \approx \frac{1}{2n}\sum_{i=1}^{2n} F(\mathbf{x}^{(i)}), \tag{1.55}$$

where the sigma points come from the columns of $\sqrt{nP}$ as in the CKF described in the previous section. A similar linearization can be created for the nonlinear measurement function. With the given statistically optimal linearization, the resulting

linear system can be filtered using the Kalman filter. For the given mixed-time system, the Taylor-Heun scheme discussed in the section on the EKF is used to discretize the continuous-time dynamics. The resulting filter is similar to the CKF implementation in the estimation of the mean, but the covariance is calculated using the first-order derivatives of the nonlinear function. The schemes for the discretization and the estimation of the expectations were chosen for numerical stability and computational complexity. Thus, this SLF implementation with numerically approximated expectation values is less computationally expensive than the CKF and UKF implementations and more computationally expensive than the EKF implementation. However, as long as the discretization and the estimation of the mean are sufficiently accurate, the SLF should perform best in the MMSE sense. Overall, ignoring the numerical approximation of the expectations, the SLF resembles the EKF expect that the Taylor series expansion is replaced by a first-order Fourier-Hermite series expansion.

Chapter 2

Methodology

2.1 Battery Model

As discussed in the previous section, this thesis considers the electrical-circuit battery model proposed by Chen and Rincón-Mora [25] and shown in Figure 2.1. The left portion of the circuit models the capacity, SOC, and runtime, while the right portion models the transient i-v characteristics. For convenience, the model is designed so that the SOC of the battery equals the voltage V_{SOC}, in volts. The parameters C_{cap} and R_{sd} are assumed constant for a given battery and determine the capacity and self-discharge rate of the battery. The other parameters are all nonlinear functions of V_{SOC} and determine the transient i-v response as well as the open-circuit voltage V_{OC}. From a typical TCL PL-383562 polymer lithium-ion battery, Chen and Rincón-Mora extracted these parameters and fit them to curves, obtaining

$$R_s(V_{\mathrm{SOC}}) = 0.1562e^{-24.37V_{\mathrm{SOC}}} + 0.07446 \tag{2.1}$$

$$R_{ts}(V_{\mathrm{SOC}}) = 0.3208e^{-29.14V_{\mathrm{SOC}}} + 0.04669 \tag{2.2}$$

$$C_{ts}(V_{\mathrm{SOC}}) = -752.9e^{-13.51V_{\mathrm{SOC}}} + 703.6 \tag{2.3}$$

$$R_{tl}(V_{\mathrm{SOC}}) = 6.603e^{-155.2V_{\mathrm{SOC}}} + 0.04984 \tag{2.4}$$

$$C_{tl}(V_{\mathrm{SOC}}) = -6056e^{-27.12V_{\mathrm{SOC}}} + 4475 \tag{2.5}$$

$$V_{\mathrm{OC}}(V_{\mathrm{SOC}}) = -1.031e^{-35V_{\mathrm{SOC}}} + 3.685 + 0.2156V_{\mathrm{SOC}} - 0.1178V_{\mathrm{SOC}}^2 + 0.3201V_{\mathrm{SOC}}^3 \tag{2.6}$$

The resistance and capacitance parameters shown above are approximately constant for SOC > 0.2 and change exponentially for SOC < 0.2. The open-circuit voltage also changes exponentially for SOC < 0.2 but is approximately linear for SOC > 0.2. Note that the capacitances C_{ts} and C_{tl} are negative for SOC values close to zero, which is both unrealistic according to the experimental data collected by Chen and Rinón-Mora and problematic mathematically. To solve this, a lower bound was placed on the V_{SOC} input to the capacitance functions. Thus, for inputs below some threshold value v_T, the capacitances are adjusted to their value at that threshold, producing

$$\hat{C}_{ts}(V_{\mathrm{SOC}}) = \begin{cases} C_{ts}(V_{\mathrm{SOC}}), & V_{\mathrm{SOC}} \geq v_T \\ C_{ts}(v_T), & V_{\mathrm{SOC}} < v_T \end{cases} \tag{2.7}$$

$$\hat{C}_{tl}(V_{\mathrm{SOC}}) = \begin{cases} C_{tl}(V_{\mathrm{SOC}}), & V_{\mathrm{SOC}} \geq v_T \\ C_{tl}(v_T), & V_{\mathrm{SOC}} < v_T \end{cases} \tag{2.8}$$

The threshold v_T was chosen based on the experimental data of Chen and Rinón-Mora, specifically so that the threshold capacitance values are approximately equal to the lowest such values measured by them. A threshold of $v_T = 0.015$ V accomplishes this goal.

This study used the nonlinear parameters given by Chen and Rincón-Mora for

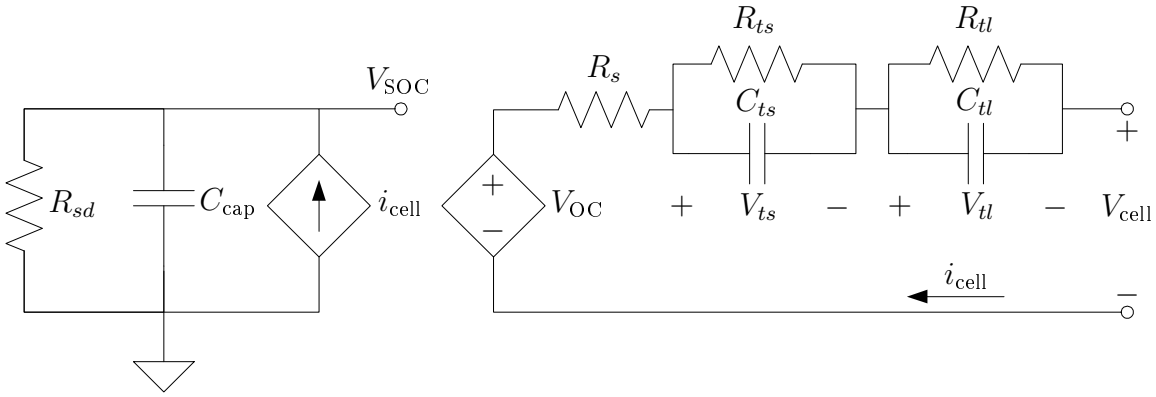

Figure 2.1: Electrical-circuit battery model.

the implementation of a battery using their battery model in Matlab. In addition, the thresholding defined in Eqs. (2.7) and (2.8) was used with $v_T = 0.015$ V. The other, constant parameters were chosen to produce a capacity of 1 Ah and a self-discharge rate of 4% per month. To do so, the capacitance $C_{\rm cap}$ is calculated to hold the desired capacity when $V_{\rm SOC} = 1$ V, and then the resistance R_{sd} is set to produce the desired self-discharge rate. For a given capacity of $C^\dagger$ in Ah, $C_{\rm cap}$ needs to be

$$C_{\rm cap} = \frac{Q}{V_{\rm SOC}} = \frac{C^\dagger}{1\ \text{V}} = 3600 C^\dagger\ [\text{F}]. \tag{2.9}$$

Then, the resistance R_{sd} is chosen so that the time constant $\tau = RC$ results in the desired drop of $\xi = 0.04$ over $T = 1$ month as follows

$$V(t) = V_0 e^{-T/\tau} = V_0(1-\xi) \tag{2.10}$$

$$\tau = -T/\ln(1-\xi) = -2592000/\ln 0.96\ [\text{s}]. \tag{2.11}$$

Then, $R_{sd} = \tau / C_{\rm cap}$. Thus, the parameters are $C_{\rm cap} = 3600$ F and $R_{sd} = 17.6376\ \text{k}\Omega$.

In order to simulate the use of the modeled battery, discharging and charging loads were implemented, as shown in Figures 2.2. For discharging, a resistive load

R_L is placed across the battery terminals, creating a discharge rate of $i_{\text{cell}} = V_{\text{cell}}/R_L$. For charging, a negative resistance $-R_L$, where $R_L > 0$, is used, creating a charging current of $-i_{\text{cell}} = V_{\text{cell}}/R_L$. Thus, any arbitrary charging or discharging current can be set by choosing the appropriate resistance R_L. Furthermore, an open circuit can be simulated by choosing R_L sufficiently large so that $i_{\text{cell}} \approx 0$. Additional consideration has to be taken to produce constant current and constant voltage charging conditions for standard charging procedure. Typically, the specific battery modeled by the given parameters is charged at a rate of $C_5/5$ until a terminal voltage of 4.2 V is reached, where $C_5/5$ is the discharge rate at which a full battery is completely discharged in 5 hours [38]. Then, the battery is charged at a constant voltage of 4.2 V until the charging current is below $C_5/20$. The constant current condition can be met by varying R_L so that V_{cell}/R_L stays constant, while the constant voltage condition is met by varying R_L so that $i_{\text{cell}}R_L$ stays constant.

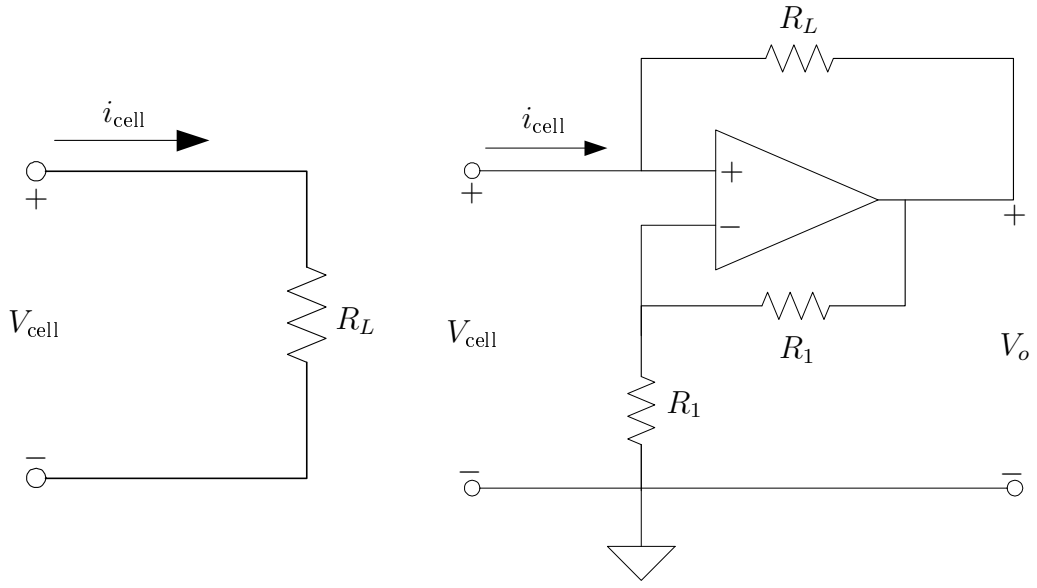

Figure 2.2: Loads to (a) discharge and (b) charge the battery.

This use of the load R_L to control the current i_{cell} suggests that it is the input to the system. Moreover, the outputs of the system are V_{cell} and i_{cell}. However, since

knowledge of one of them along with R_L allows for the calculation of the other, the two outputs have a known relationship between them. Therefore, only one of the outputs is necessary to fully define the input-output relationship of the system. In this study, the voltage V_{cell} was chosen as the output.

For ease of numerical simulation, it is useful to find the state-space system for the circuit. The state-space representation is derived using the physical variable definition, in which the state variables are chosen to represent the voltages across the capacitors. Choosing $x_1 = V_{\text{SOC}}$, $x_2 = V_{ts}$, and $x_3 = V_{tl}$ achieves this goal and results in the state-space representation

$$\dot{x}_1 = -\frac{x_1}{R_{sd}C_{\text{cap}}} - \frac{V_{\text{OC}}(x_1) - x_2 - x_3}{(R_s(x_1) + R_L)C_{\text{cap}}} + f_{w,1}(\mathbf{x}, R_L, \mathbf{w}) \tag{2.12}$$

$$\dot{x}_2 = -\frac{x_2}{R_{ts}(x_1)C_{ts}(x_1)} + \frac{V_{\text{OC}}(x_1) - x_2 - x_3}{(R_s(x_1) + R_L)C_{ts}(x_1)} + f_{w,2}(\mathbf{x}, R_L, \mathbf{w}) \tag{2.13}$$

$$\dot{x}_3 = -\frac{x_3}{R_{tl}(x_1)C_{tl}(x_1)} + \frac{V_{\text{OC}}(x_1) - x_2 - x_3}{(R_s(x_1) + R_L)C_{tl}(x_1)} + f_{w,3}(\mathbf{x}, R_L, \mathbf{w}) \tag{2.14}$$

$$V_{\text{cell}} = \frac{V_{\text{OC}}(x_1) - x_2 - x_3}{1 + R_s(x_1)/R_L} + f_v(\mathbf{x}, R_L, \mathbf{v}), \tag{2.15}$$

where R_L is the input to the system, V_{cell} is the output, f_w is the process noise function, f_v is the measurement noise function, and the nonlinear parameters depending on x_1 are given by Eqs. (2.1) through (2.6) along with the thresholding defined in Eqs. (2.7) and (2.8). It is obvious from this formulation that the system is nonlinear to both the input and the states. In order to establish the noise expressions, the types of noise present in the battery have to first be determined.

This thesis assumes that the process and measurement noises in this system are due to thermal noise in the resistances for the internal impedance of the battery R_s, R_{ts}, and R_{tl}, and for the load R_L. This is motivated by measurements of the voltage noise in batteries conducted by Boggs et al. that showed the measured noise is mainly

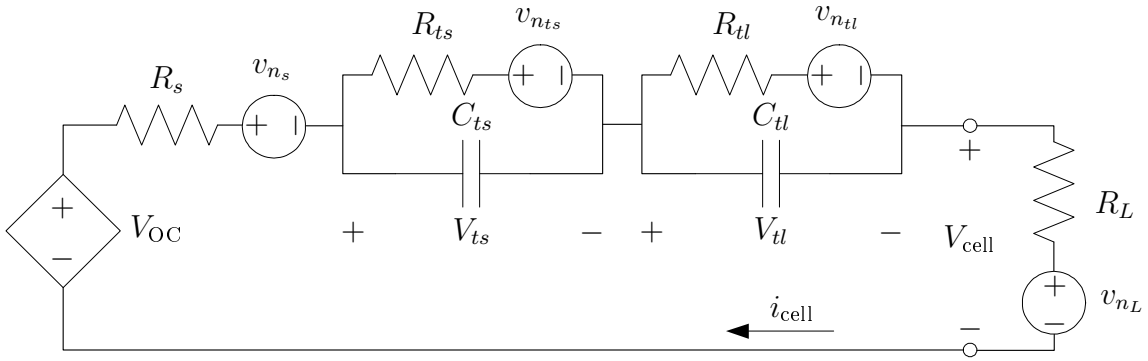

Figure 2.3: Modeling of thermal noise in resistances as voltage sources in series with the resistances.

due to thermal noise; the correlation between the battery terminals suppresses shot noise [39]. This thermal noise is assumed to be Gaussian white noise with a power spectral density (PSD) of [40]

$$S_n(\omega) \cong 2kT \text{ watts per Hz} \qquad \text{for} \qquad |\omega| \ll 2\pi kT/h, \tag{2.16}$$

where T is the temperature of the conducting medium in Kelvin, k is the Boltzmann's constant, and h is the Planck's constant. Figure 2.3 shows that the thermal noise due to the resistances is modeled as voltage sources in series with the resistances, with PSDs of $S_v(\omega) = 2kTR$ for a corresponding resistance R. Using this definition, the noise functions are given by

$$f_{w,1} = \frac{v_{n_s} + v_{n_L}}{(R_s(x_1) + R_L)C_{\text{cap}}} \tag{2.17}$$

$$f_{w,2} = \frac{v_{n_{ts}}}{R_{ts}(x_1)C_{ts}(x_1)} - \frac{v_{n_s} + v_{n_L}}{(R_s(x_1) + R_L)C_{ts}(x_1)} \tag{2.18}$$

$$f_{w,3} = \frac{v_{n_{tl}}}{R_{tl}(x_1)C_{tl}(x_1)} - \frac{v_{n_s} + v_{n_L}}{(R_s(x_1) + R_L)C_{tl}(x_1)} \tag{2.19}$$

$$f_v = -\frac{v_{n_s} + v_{n_L}}{1 + R_s(x_1)/R_L}. \tag{2.20}$$

It can be seen that the resistances change over time, which causes the covariance of the sources v_n to also change. For the purposes of modeling, it is useful to define noise variables that have constant covariance. Using the square root of the power supplied by the noise sources as the noise variables accomplishes this goal and produces the variables $w_1 = v_{n_s}/\sqrt{R_s}$, $w_2 = v_{n_{ts}}/\sqrt{R_{ts}}$, $w_3 = v_{n_{tl}}/\sqrt{R_{tl}}$, and $w_4 = v_{n_L}/\sqrt{|R_L|}$ along with $v_1 = v_{n_s}/\sqrt{R_s}$ and $v_2 = v_{n_L}/\sqrt{|R_L|}$, which all have constant covariances of $2kT$. Note the use of the absolute value of R_L in the definition of w_4 and v_2, since R_L can become negative. In the case of $R_L < 0$, their covariances remain at $2kT$ while the sign of v_{n_L} is negated, which is implemented in the system using the signum function, defined as

$$\operatorname{sgn}(x) := \begin{cases} -1 & \text{if } x < 0 \\ 0 & \text{if } x = 0 \\ 1 & \text{if } x > 0. \end{cases} \tag{2.21}$$

Then, the state space representation of the system becomes

$$\dot{x}_1 = -\frac{x_1}{R_{sd}C_{\text{cap}}} - \frac{V_{\text{OC}}(x_1) - x_2 - x_3 - \sqrt{R_s(x_1)}w_1 - \operatorname{sgn}(R_L)\sqrt{|R_L|}w_4}{(R_s(x_1) + R_L)C_{\text{cap}}} \tag{2.22}$$

$$\dot{x}_2 = -\frac{x_2 - \sqrt{R_{ts}(x_1)}w_2}{R_{ts}(x_1)C_{ts}(x_1)} + \frac{V_{\text{OC}}(x_1) - x_2 - x_3 - \sqrt{R_s(x_1)}w_1 - \operatorname{sgn}(R_L)\sqrt{R_L}w_4}{(R_s(x_1) + R_L)C_{ts}(x_1)} \tag{2.23}$$

$$\dot{x}_3 = -\frac{x_3 - \sqrt{R_{tl}(x_1)}w_3}{R_{tl}(x_1)C_{tl}(x_1)} + \frac{V_{\text{OC}}(x_1) - x_2 - x_3 - \sqrt{R_s(x_1)}w_1 - \operatorname{sgn}(R_L)\sqrt{R_L}w_4}{(R_s(x_1) + R_L)C_{tl}(x_1)} \tag{2.24}$$

$$V_{\text{cell}} = \frac{V_{\text{OC}}(x_1) - x_2 - x_3 - \sqrt{R_s(x_1)}v_1 - \operatorname{sgn}(R_L)\sqrt{R_L}v_2}{1 + R_s(x_1)/R_L}. \tag{2.25}$$

It can be seen that the system can be written in the form

$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}, \mathbf{u}, \mathbf{w}) \tag{2.26}$$

$$y = h(\mathbf{x}, \mathbf{u}, \mathbf{v}). \tag{2.27}$$

It is useful to find the derivatives of $\mathbf{f}$ and h with respect to $\mathbf{x}$, $\mathbf{w}$, and $\mathbf{v}$ for use with the filters. This thesis defines the Jacobians $F = (\partial\mathbf{f}/\partial\mathbf{x})$, $H = (\partial h/\partial\mathbf{x})$, $L = (\partial\mathbf{f}/\partial\mathbf{w})$, and $M = (\partial h/\partial\mathbf{v})$. Their equations are

$$F = \begin{bmatrix} \frac{-1}{R_{sd}C_{\text{cap}}} + \frac{(Voc - x_2 - x_3)R_s' - (R_s + R_L)V_{\text{OC}}'}{(R_s + R_L)^2 C_{\text{cap}}} & & \\ \frac{(R_{ts}C_{ts})'x_2}{R_{ts}C_{ts}} + \frac{(R_s + R_L)C_{ts}V_{\text{OC}}' - (Voc - x_2 - x_3)\left[(R_s + R_L)C_{ts}\right]'}{\left[(R_s + R_L)C_{ts}\right]^2} & & \\ \frac{(R_{tl}C_{tl})'x_3}{R_{tl}C_{tl}} + \frac{(R_s + R_L)C_{tl}V_{\text{OC}}' - (Voc - x_2 - x_3)\left[(R_s + R_L)C_{tl}\right]'}{\left[(R_s + R_L)C_{tl}\right]^2} & & \\ \cdots & \frac{1}{(R_s + R_L)C_{\text{cap}}} & \frac{1}{(R_s + R_L)C_{\text{cap}}} \\ \cdots & \frac{-1}{R_{ts}C_{ts}} + \frac{-1}{(R_s + R_L)C_{ts}} & \frac{-1}{(R_s + R_L)C_{ts}} \\ \cdots & \frac{-1}{(R_s + R_L)C_{tl}} & \frac{-1}{R_{tl}C_{tl}} + \frac{-1}{(R_s + R_L)C_{tl}} \end{bmatrix} \tag{2.28}$$

$$H = \begin{bmatrix} (1 + R_s/R_L)V_{\text{OC}}' - (V_{\text{OC}} - x_2 - x_3)R_s'/R_L}{(1 + R_s/R_L)^2} & \frac{-1}{1 + R_s/R_L} & \frac{-1}{1 + R_s/R_L} \end{bmatrix} \tag{2.29}$$

$$L = \begin{bmatrix} \dfrac{\sqrt{R_s}}{(R_s+R_L)C_{\text{cap}}} & 0 & 0 & \dfrac{\operatorname{sgn} R_L\sqrt{|R_L|}}{(R_s+R_L)C_{\text{cap}}} \\ \dfrac{-\sqrt{R_s}}{(R_s+R_L)C_{\text{cap}}} & \dfrac{\sqrt{R_{ts}}}{R_{ts}C_{ts}} & 0 & \dfrac{-\operatorname{sgn} R_L\sqrt{|R_L|}}{(R_s+R_L)C_{\text{cap}}} \\ \dfrac{-\sqrt{R_s}}{(R_s+R_L)C_{\text{cap}}} & 0 & \dfrac{\sqrt{R_{tl}}}{R_{tl}C_{tl}} & \dfrac{-\operatorname{sgn} R_L\sqrt{|R_L|}}{(R_s+R_L)C_{\text{cap}}} \end{bmatrix} \tag{2.30}$$

$$M = \begin{bmatrix} \dfrac{-\sqrt{R_s}}{(R_s+R_L)C_{\text{cap}}} & \dfrac{-\operatorname{sgn} R_L\sqrt{|R_L|}}{(R_s+R_L)C_{\text{cap}}} \end{bmatrix}, \tag{2.31}$$

where $()'$ indicates derivation with respect to x_1 and the dependence on x_1 has been omitted due to space constraints. Furthermore, it is useful to find the Hessian of $\mathbf{f}$ with respect to $\mathbf{x}$. Due to symmetry and $\partial^2 f_k/\partial x_i \partial x_j = 0$ for $i,j = 2,3$, only the first column of each tensor component of the Hessian is given. The resultant Hessian is

$$\frac{\partial^2 f_1}{\partial x_i \partial x_1} = \begin{bmatrix} \dfrac{(R_s+R_L)\big[(V_{\text{OC}}-x_2-x_3)R_s'' - (R_s+R_L)V_{\text{OC}}''\big] - 2R_s'\big[(V_{\text{OC}}-x_2-x_3)R_s' - (R_s+R_L)V_{\text{OC}}'\big]}{(R_s+R_L)^3 C_{\text{cap}}} \\ \dfrac{-R_s'}{(R_s+R_L)^2 C_{\text{cap}}} \\ \dfrac{-R_s'}{(R_s+R_L)^2 C_{\text{cap}}} \end{bmatrix} \tag{2.32}$$

$$\frac{\partial^2 f_2}{\partial x_i \partial x_1} = \begin{bmatrix} \dfrac{\big\{R_{ts}C_{ts}(R_{ts}C_{ts})'' - \big[(R_{ts}C_{ts})'\big]^2\big\}x_2}{(R_{ts}C_{ts})^2} + \dfrac{(R_s+R_L)C_{ts}\big\{(R_s+R_L)C_{ts}V_{\text{OC}}'' - (V_{\text{OC}}-x_2-x_3)\big[(R_s+R_L)C_{ts}\big]''\big\} + 2\big[(R_s+R_L)C_{ts}\big]'\big\{(R_s+R_L)C_{ts}V_{\text{OC}}' - (Voc-x_2-x_3)\big[(R_s+R_L)C_{ts}\big]'\big\}}{\big[(R_s+R_L)C_{ts}\big]^3} \\ \dfrac{(R_{ts}C_{ts})'}{R_{ts}C_{ts}} + \dfrac{\big[(R_s+R_L)C_{ts}\big]'}{\big[(R_s+R_L)C_{ts}\big]^2} \\ \dfrac{\big[(R_s+R_L)C_{ts}\big]'}{\big[(R_s+R_L)C_{ts}\big]^2} \end{bmatrix} \tag{2.33}$$

$$\frac{\partial^2 f_3}{\partial x_i \partial x_1} = \begin{bmatrix} \dfrac{\{R_{tl}C_{tl}(R_{tl}C_{tl})'' - \left[(R_{tl}C_{tl})'\right]^2\}x_3}{(R_{tl}C_{tl})^2} + \dfrac{(R_s+R_L)C_{tl}\{(R_s+R_L)C_{tl}V''_{\rm OC} - (V_{\rm OC}-x_2-x_3)\left[(R_s+R_L)C_{tl}\right]''\} + 2\left[(R_s+R_L)C_{tl}\right]'\{(R_s+R_L)C_{tl}V'_{\rm OC} - (Voc-x_2-x_3)\left[(R_s+R_L)C_{tl}\right]'\}}{\left[(R_s+R_L)C_{tl}\right]^3} \\ \dfrac{\left[(R_s+R_L)C_{tl}\right]'}{\left[(R_s+R_L)C_{tl}\right]^2} \\ \dfrac{(R_{tl}C_{tl})'}{R_{tl}C_{tl}} + \dfrac{\left[(R_s+R_L)C_{tl}\right]'}{\left[(R_s+R_L)C_{tl}\right]^2} \end{bmatrix} \tag{2.34}$$

Additionally, the covariances of the error variables are

$$Q = 2kTI_4 \tag{2.35}$$

$$R = 2kTI_2, \tag{2.36}$$

where I_n is a $n \times n$ identity matrix.

2.2 Simulation Setup

In order to simulate the stochastic system, the covariances of the noises as well as their generation and simulation methods have to be determined. This thesis assumed that a standard temperature of $T = 290$ Kelvin. Therefore, the noise variables have covariances of $\sigma^2 = 2kT = 8.0078 \times 10^{-21}$ W/Hz. Furthermore, to better differentiate the performance of the filters, additional measurement noise was introduced assuming an oscilloscope was used to perform the measurement. Specifically, consider the Tektronix TBS1022 oscilloscope with a DC gain accuracy of $\pm 3\%$ of the full range. Based on numerical simulations, an approximate range of 4 V peak-to-peak is necessary to fully capture the range of possible $V_{\rm cell}$ values, resulting in a measurement inaccuracy of ± 0.12 V. This noise is assumed to be white Gaussian noise, and its range lies within three standard deviations. Thus, an

additional measurement noise variable v_3 is introduced with a variance of

$$\sigma_{v3}^2 = \left(\frac{0.12 \text{ V}}{3}\right)^2 = 0.0016 \text{ V}. \tag{2.37}$$

Then, the new measurement covariance matrix is

$$R = \text{diag}(2kT, 2kT, 0.0016), \tag{2.38}$$

and the corresponding Jacobian $M = (\partial h/\partial \mathbf{v})$ is

$$M = \begin{bmatrix} \dfrac{-\sqrt{R_s}}{(R_s + R_L)C_{\text{cap}}} & \dfrac{-\operatorname{sgn} R_L \sqrt{|R_L|}}{(R_s + R_L)C_{\text{cap}}} & 1 \end{bmatrix}. \tag{2.39}$$

In order to numerically simulate the effect of white noise on a system, the simulation time step must be sufficiently smaller than the time constant of the fastest battery process. This is approximated as the product of the constant terms of the functions for R_{ts} and C_{ts}, halved to satisfy Nyquist conditions. Then, the simulation time step is taken to be 1/100 of the calculated maximum time step, as suggested by Matlab documentation. Thus, the simulations used a time step of

$$\delta_{\text{sim}} = \frac{R_{ts,\text{const.}} C_{ts,\text{const.}}}{200} = \frac{0.04669 \times 703.6}{200} = 0.164255 \text{ s}. \tag{2.40}$$

Simulation showed that the resulting time step is sufficiently small to capture the effects of the white noise, i.e. further reducing the step size had negligible effect. With the chosen step size, the noise is simulated as band-limited white Gaussian noise with a correlation time equal to the step size. At each time step, the noise values for the process noise sources are generated using random number generators producing normally-distributed numbers with means of zero and variances equal

to the diagonals of the covariance matrix divided by the correlation time. The scaling of the variance by the correlation time ensures the response of the system to the approximate white noise has the same covariance as it would have to actual white noise. Note that the measurement noise is bandlimited not by the system but by the measurement device, in this case an oscilloscope. Thus, the calculated variance already takes into account the measurement bandwidth of the oscilloscope and scaling is unnecessary. For reproducibility, the random number generators were seeded with predictable numbers. This was done in Matlab by first seeding the main random number generator with a seed of 0. Then, for each Monte Carlo trial, five positive integers were generated, for the five noise sources, with the integers uniformly distributed between 1 and $2^{32} - 1$. These integers were used to seed the random number generators in a Simulink model. The use of the Simulink model allows for the random number generators to be easily seeds with different integers and does not affect the predictable sequence in Matlab.

Then, the battery was simulated using a Simulink model with the fourth-order Runge-Kutta method and an initial condition of $x_0 = [1, 0, 0]^\top$. A total of 100 Monte Carlo runs of the battery were performed with seed values for the noise sources generated using the method mentioned in the previously. Figure 2.4 shows the input load on the battery system, where the values off the graph are idle periods, simulated using a very large input of $R_L = 10^{10}\ \Omega$ so that the battery current is approximately zero. It can be seen that the input is piecewise constant. This input was chosen to to test the performance of the filters by gradually increasing the strength of the nonlinear rate-capacity and recovery effects. Note that care was taken to ensure the SOC remained within the range of zero to one, so discharge and charge times could not always be equal. Intially, the battery is idle for 10 minutes to allow the estimated covariances of the filters to converge. Then, the battery is

discharged and charged at $R_L = 20\ \Omega$ for 290 and 270 minutes, respectively. The resulting low current from this load is approximately the testing current $C_5/5$ used to determine battery capacity, as discussed in Section 2.1. In order to increase the strength of the nonlinear effects, the absolute value of the current was increased by decreasing the input to $R_L = 10\ \Omega$ and discharging and charging for 150 minutes, each. Then, the input was further reduced to $R_L = 5\ \Omega$ and discharged and charged for 80 minutes and 70 minutes, respectively. Next, the battery was discharged and charged at $R_L = 4\ \Omega$ for 60 minutes, each. Following were discharge and charge periods at $R_L = 2\ \Omega$ for 35, 25, 25, 20, 25, 25, 20, and 15 minutes. The high current results in very strong rate-capacity effects. Finally, the battery was rested for 70 minutes, discharged at $R_L = 2\ \Omega$ for 25 minutes, rested for 75 minutes, and charged at $R_L = 10\ \Omega$ for 135 minutes. The two resting periods should show the strongest recovery effect. The total input time is 1635 minutes. The SOC and the noisy measurement of V_{cell} resulting from the given input is shown one Monte Carlo trial in Figures 2.5 and 2.6, respectively.

The Monte Carlo trials used varying sampling periods, calculated as some multiple K of the simulation step size. For example, for a desired sampling period of 300 seconds, the actual sampling period is

$$T_s = K\delta_{\text{sim}} = 1826 \times 0.164255\ \text{s} = 299.93\ \text{s}, \tag{2.41}$$

where the factor K is chosen to minimize the difference between the desired and actual periods. For convenience, the sampling period will refer to the actual sampling period calculated in this manner. Additionally, to increase the accuracy of the discretization in the prediction steps, the sampling period is divided by a positive integer M and the prediction step is performed M times prior to each update step.

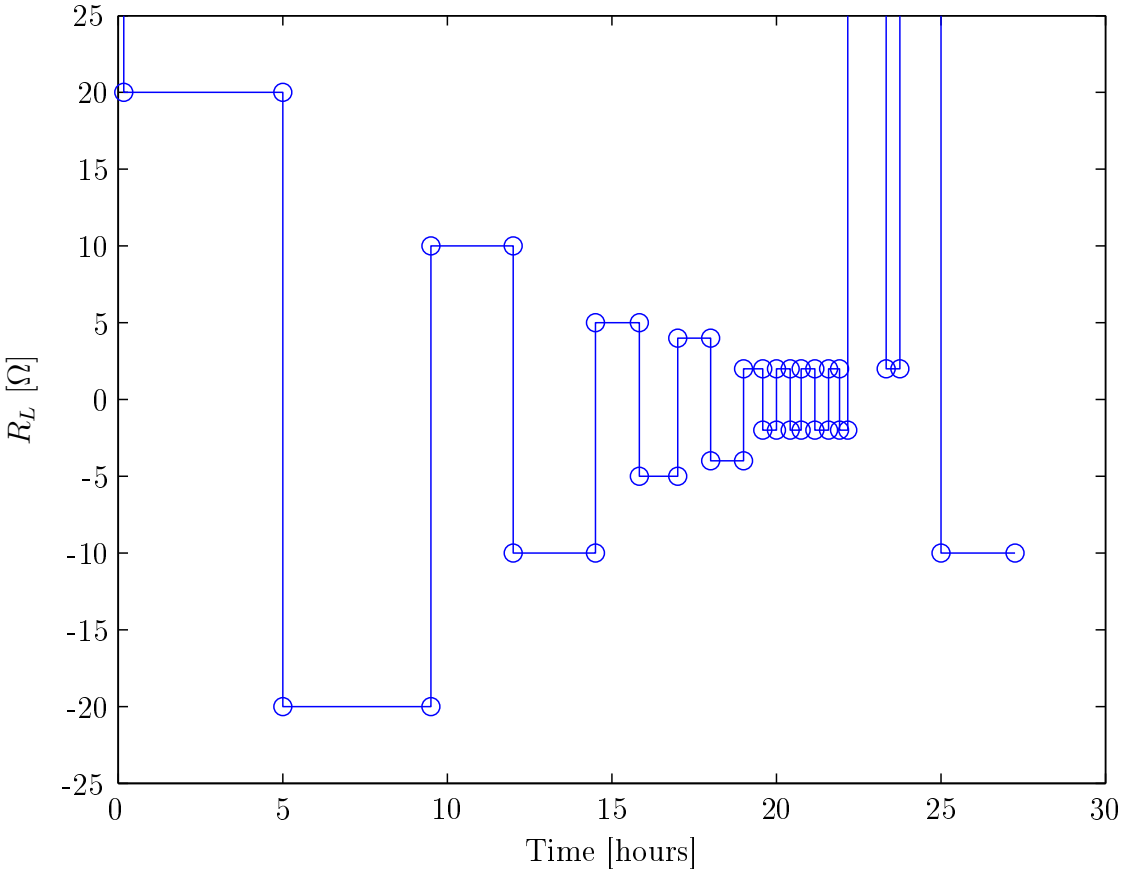

Figure 2.4: Input load R_L on the battery.

Thus, the prediction and update portions of the filters can be calculated at different rates. The next chapter discusses the simulation results for a range of sampling periods and refinement steps M.

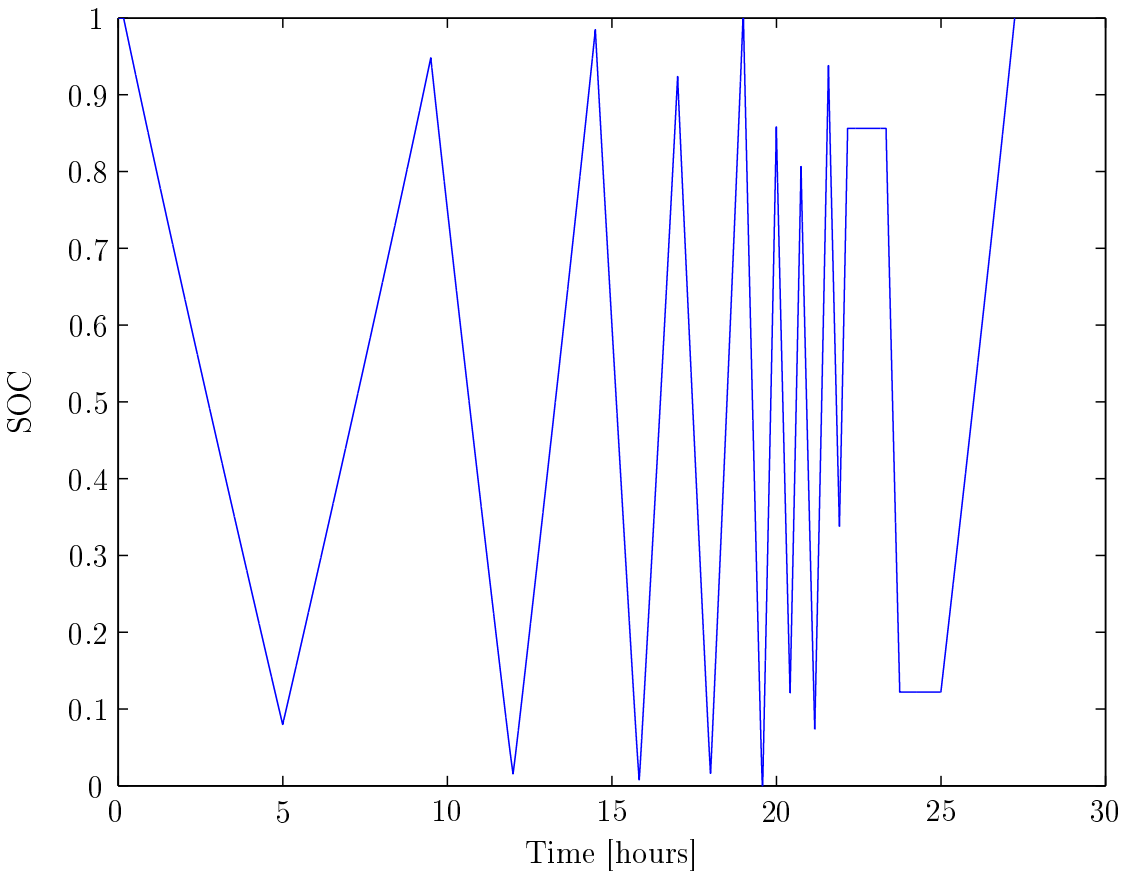

Figure 2.5: True SOC for one run due to input load.

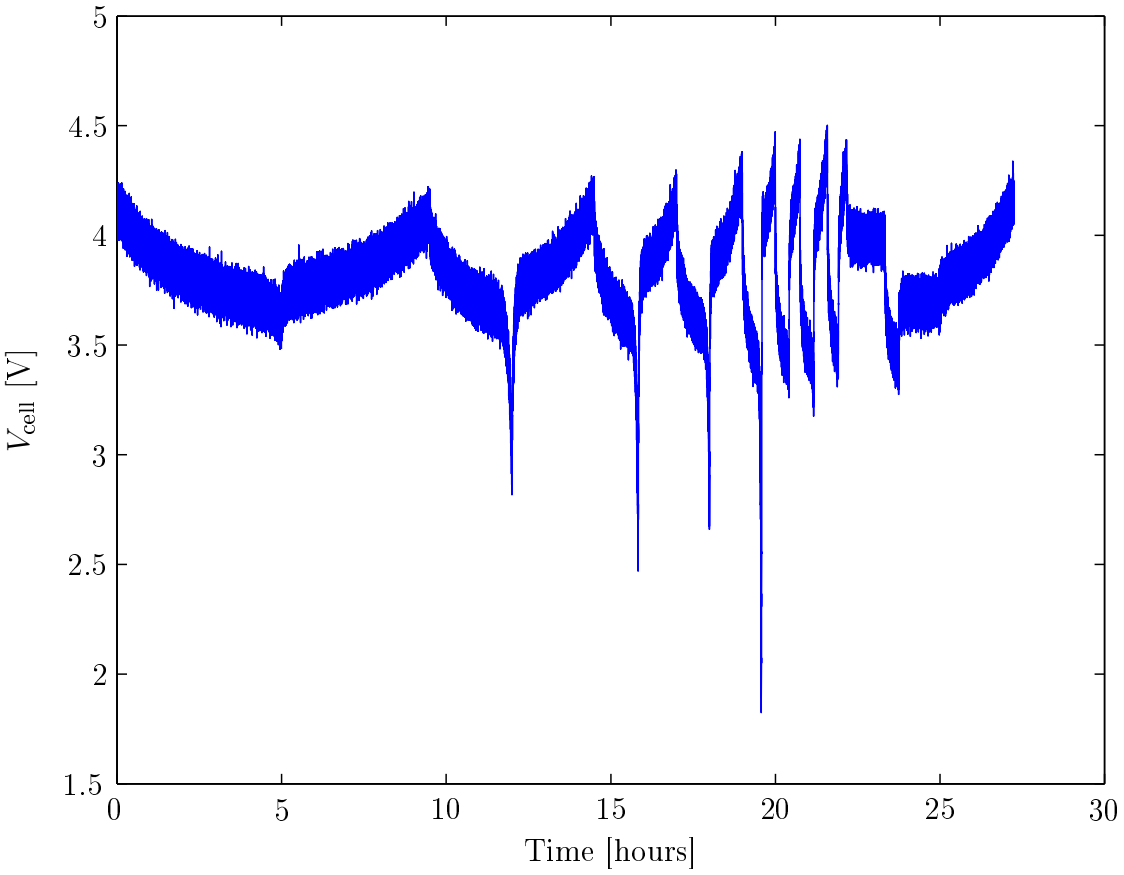

Figure 2.6: Noisy measurement V_{cell} for one run.

Chapter 3

Results

For a sampling period of $T_s = 90,300$ s, refinement steps of $M = 1, 2, \ldots, 2^5$ were used. The filtering was performed assuming an initial state of $x_0 = [1, 0, 0]^\top$ and an initial covariance matrix of $P_0 = 10^{-6} I_3$, where the initial state was chosen to match the actual state and the initial covariance was experimentally tuned for fast convergence. For the 100 Monte Carlo trials and the four filters, the number of divergences were counted, where a divergence is considered an absolute error in the estimated SOC greater than 0.1 V or any failure in the filtering process, such as due to a non-invertible matrix or a non-positive definite covariance matrix. If a failure occurs, no attempt was made to keep filtering the system, and the remainder of the SOC values are assumed to be zero. Additionally, the mean RMSE (MRMSE) is given, where the squared error for each filter is averaged over the trials and the square root of the result is averaged time.

Table 3.1: Divergences and MRMSE for various filters at $T_s = 90$ s.

Filter	$M = 1$		2		4		8		16		32	
	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE
EKF												
CKF												
UKF												
SLF												

Table 3.2: Divergences and MRMSE for various filters at $T_s = 150$ s.

Filter	$M = 1$		2		4		8		16		32	
	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE
EKF	100	0.3625	99	0.5410	99	0.3687	94	0.4450				
CKF	100	5.527	100	0.2774	100	0.2768	100	7.851				
UKF	100	0.3352	100	0.2973	100	0.2733	100	7840				
SLF	100	0.3093	100	0.2716	100	0.3185	100	0.1589				

Table 3.3: Divergences and MRMSE for various filters at $T_s = 300$ s.

Filter	$M = 1$		2		4		8		16		32	
	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE	Divs.	MRMSE
EKF	100	0.8511	100	0.3949	100	0.4471	100	19.01	97	0.2745	100	
CKF	100	513.9	100	0.4047	100	0.2769	100	0.2769	100	15.54	100	
UKF	100	258.1	100	4.182	100	0.2861	100	0.2750	100	13.79	100	
SLF	100	0.2437	100	0.2751	100	0.2737	100	1.144	100	0.1786	100	

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