ECE 6553: Optimal Control Notes

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Chapter 1

Parameter Optimization

1.1 What is Optimal Control?

Optimal Maximize/minimize cost (subject to constraints): $\min_u g(u)$ With constraints,

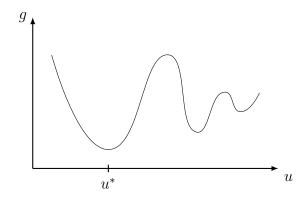
$$\min_{u} g(u)$$
s.t.
$$\begin{cases}
h_1(u) = 0 \\
h_2(u) \le 0
\end{cases}$$

First-order necessary condition (FONC):

$$\frac{\partial g}{\partial u}(u^*) = 0$$

Optimality can be

- $\bullet\,$ local vs global
- max vs min



Control control design: pick u such that specifications are satisfied:

$$\dot{x} = f(x, u), \qquad \dot{x} = Ax + Bu,$$

where $x(t) \in \mathbb{R}^n$ is the state, $u(t) \in \mathbb{R}^m$ is the control, and $f(\cdot)$ is the dynamics. Actually, x and u are signals:

$$x:[0,T]\to\mathbb{R}^n, \qquad u:[0,T]\to\mathbb{R}^m$$

Optimal control find the "best" u!

For "best" to mean anything, we need a cost. The big/deep question is

$$\frac{\partial \text{"cost"}}{\partial u} = 0$$

Example

Suppose we have a car with position p. Its acceleration \ddot{p} is controlled by the gas/brake input u ($\ddot{p} = u$). In order to express the dynamics of the system in the form $\dot{x} = f(x, u)$, we introduce state variables:

$$\begin{array}{c} x_1 = p \\ x_2 = \dot{p} \end{array} \Longrightarrow \begin{cases} \dot{x}_1 = x_2 \\ \dot{x}_2 = u \end{cases}$$

The task is to move the car from its initial position to a stop at a distance c away.

Minimum energy problem

$$\min_{u} \int_{0}^{T} u^{2}(t) dt$$
s.t.
$$\begin{cases}
\dot{x}_{1} = x_{2} \\
\dot{x}_{2} = u
\end{cases}$$

$$x_{1}(0) = 0, x_{2}(0) = 0$$

$$x_{1}(T) = c, x_{2}(T) = 0$$

Minimum time problem

$$\min_{u,T} T = \int_{0}^{T} dt$$
s.t.
$$\begin{cases}
\dot{x}_{1} = x_{2} \\
\dot{x}_{2} = u
\end{cases}$$

$$x_{1}(0) = 0, x_{2}(0) = 0$$

$$x_{1}(T) = c, x_{2}(T) = 0$$

$$u(t) \in [u_{\min}, u_{\max}]$$

The general optimal control problem we will solve will look like

$$\min_{u,T} \int_0^T L(x(t), u(t), t) dt + \Psi(x(T))$$
s.t. $\dot{x}(t) = f(x(t), u(t), t), t \in [0, T]$

$$x(0) = x_0$$

$$x(T) \in S$$

$$u(t) \in \Omega, t \in [0, T]$$

where $\Psi(\cdot)$ is the terminal cost and S is the terminal manifold. This is a so-called **Bolza Problem**.

What tools do we need to solve this?

- 1. optimality conditions $\partial \cos t/\partial u = 0$
- 2. some way of representing the optimal signal $u^*(x,t)$
- 3. some way of actually finding/computing the optimal controllers

1.2 Unconstrained Optimization

Let the decision variable be $u = [u_1, \dots, u_m]^T \in \mathbb{R}^m$. The cost is $g(u) \in C^1$ (C^k means k times continuously differentiable). The problem is

$$\min_{u} g(u), \quad g: \mathbb{R}^m \to \mathbb{R}$$

For u^* to be a minimizer, we need

$$\frac{\partial g}{\partial u}(u^*) = 0$$

Definition. u^* is a (local) minimizer to q if $\exists \delta > 0$ s.t.

$$g(u^*) \le g(u) \quad \forall u \in B_{\delta}(u^*)$$

$$B_{\delta}(u^*) = \{u \mid ||u - u^*|| \le \delta\}$$

Note:

• $\frac{\partial g}{\partial u}(u^*)\delta u \in \mathbb{R}$ and δu is $m \times 1$, so $\frac{\partial g}{\partial u}$ is a $1 \times m$ row vector. For the column vector,

$$\nabla g = \frac{\partial g^{\mathrm{T}}}{\partial u} \in \mathbb{R}^m$$

• $\frac{\partial g}{\partial u} \delta u$ is an inner product

$$\langle \nabla g, \delta u \rangle = \left\langle \frac{\partial g^{\mathrm{T}}}{\partial u}, \delta u \right\rangle$$

• $o(\varepsilon)$ encodes higher-order terms

$$\lim_{\varepsilon \to 0} \frac{o(\varepsilon)}{\varepsilon} = 0 \qquad \text{``faster than linear''}$$

This is opposed to big-O notation:

$$\lim_{\varepsilon \to 0} \frac{\mathcal{O}(\varepsilon)}{\varepsilon} = c$$

• δu has direction and scale so we could write it as

$$\delta u = \varepsilon v, \quad \varepsilon \in \mathbb{R}, \ v \in \mathbb{R}^m$$

Theorem. For u^* to be a minimizer, we need

$$\frac{\partial g}{\partial u}(u^*) = 0$$

or, equivalently,

$$\frac{\partial g}{\partial u}(u^*)v = 0 \quad \forall v \in \mathbb{R}^m$$

Proof. Let u^* be a minimizer. Evaluating the cost g(u) in the ball and using Taylor's expansion,

$$g(u^* + \delta u) = g(u^*) + \frac{\partial g}{\partial u}(u^*)\delta u + o(\|\delta u\|) = g(u^*) + \varepsilon \frac{\partial g}{\partial u}(u^*)v + o(\varepsilon)$$

Assume that $\frac{\partial g}{\partial u} \neq 0$. Then we could pick $v = -\frac{\partial g^{\mathrm{T}}}{\partial u}(u^*)$, i.e.

$$g(u^* + \varepsilon v) = g(u^*) - \varepsilon \left\| \frac{\partial g^{\mathrm{T}}}{\partial u}(u^*) \right\|^2 + o(\varepsilon)$$

Note that the second term is negative per our assumptions. So, for ε sufficiently small, we have

$$g\left(u^* - \varepsilon \frac{\partial g^{\mathrm{T}}}{\partial u}(u^*)\right) < g(u^*)$$

This contradicts u^* being a minimizer. \times (crossed swords)

Definition (Positive definite). $M = M^{T} \succ 0$ if

$$z^{\mathrm{T}}Mz > 0 \quad \forall z \neq 0, \ z \in \mathbb{R}^m$$

 \iff M has real and positive eigenvalues

Theorem. If $g \in C^2$, then a **sufficient** condition for u^* to be a (local) minimizer is

$$1. \ \frac{\partial g}{\partial u}(u^*) = 0$$

2.
$$\frac{\partial^2 g}{\partial u^2}(u^*) \succ 0$$
 (the Hessian is positive definite)

Definition. $g: \mathbb{R}^m \to \mathbb{R}$ is convex if

$$g(\alpha u_1 + (1 - \alpha)u_2) \le \alpha g(u_1) + (1 - \alpha)g(u_2) \quad \forall \alpha \in [0, 1], \ u_1, u_2 \in \mathbb{R}^m$$



Theorem. If $\frac{\partial^2 g}{\partial u^2}(u) \succeq 0 \ \forall u \in \mathbb{R}^m$, then g is convex. \iff for $g \in C^2$)

Example $\min_{u} u^{\mathrm{T}} Q u - b^{\mathrm{T}} u$ where $Q = Q^{\mathrm{T}} \succ 0$ (positive definite matrix)

$$\frac{\partial g}{\partial u} = \frac{\partial}{\partial u} (u^{\mathrm{T}} Q u - b^{\mathrm{T}} u)
= u^{\mathrm{T}} Q^{\mathrm{T}} + u^{\mathrm{T}} Q - b^{\mathrm{T}}
= 2u^{\mathrm{T}} Q - b^{\mathrm{T}}
\frac{\partial^2 g}{\partial u^2} = 2Q
\frac{\partial^2 g}{\partial u^2} = \begin{bmatrix} \frac{\partial^2 g}{\partial u_1^2} & \cdots & \frac{\partial^2 g}{\partial u_1 \partial u_m} \\ \vdots & \ddots & \vdots \\ \frac{\partial^2 g}{\partial u_m \partial u_1} & \cdots & \frac{\partial^2 g}{\partial u_m^2} \end{bmatrix}$$

From $\frac{\partial g}{\partial u} = 2u^{\mathrm{T}}Q - b^{\mathrm{T}} = 0$,

$$u = \frac{1}{2}Q^{-1}b$$

To see whether this is a minimizer, consider the Hessian. Since $Q \succ 0$, it follows that $\frac{\partial^2 g}{\partial u^2}(u^*) \succ 0$ and $u^* = \frac{1}{2}Q^{-1}b$ is a (local) minimizer. Additionally, since $\frac{\partial^2 g}{\partial u^2} \succ 0$, g is convex and u^* is a global minimizer. In fact, since we have strict convexity ($\succ 0$ rather than $\succeq 0$), it is the unique global minimizer.

In optimal control, *local* is typically all we can ask for. In optimization, we can do better! But wait, just because we know $\frac{\partial g}{\partial u} = 0$, it doesn't follow that we can actually find u^* ...

1.3 Numerical Methods

Idea: $u_{k+1} = u_k + \text{step}_k$. What should step_k be? For small step_k = $\gamma_k v_k$,

$$g(u_k \cdot \operatorname{step}_k) = g(u_k) + \frac{\partial g}{\partial u}(u_k) \cdot \operatorname{step}_k + o(\|\operatorname{step}_k\|) = g(u_k) + \gamma_k \frac{\partial g}{\partial u}(u_k)v_k + o(\gamma_k)$$

A perfectly reasonable choice of step direction is

$$v_k = -\frac{\partial g^{\mathrm{T}}}{\partial u}(u_k),$$

known as the steepest descend direction. This produces

$$g\left(u_k - \gamma_k \frac{\partial g}{\partial u}(u_k)\right) = g(u_k) - \gamma_k \left\| \frac{\partial g}{\partial u}(u_k) \right\|^2 + o(\gamma_k)$$

Steepest descent

$$u_{k+1} = u_k - \gamma_k \frac{\partial g^{\mathrm{T}}}{\partial u}(u_k)$$

Note:

• What should γ_k be?

• This method "pretends" that g(u) is linear. If we pretend g(u) is quadratic, we get

$$u_{k+1} = u_k - \left(\frac{\partial^2 g}{\partial u^2}(u_k)\right)^{-1} \frac{\partial g^{\mathrm{T}}}{\partial u}(u_k),$$

i.e. Newton's Method

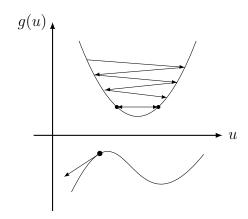
This course: steepest descent

Step-size selection?

• Choice 1: $\gamma_k = \gamma$ "small" $\forall k$; will get close to a minimizer if u_0 is close enough and γ small enough

Problems:

- You may not converge! (but you'll get close)
- You may go off to infinity (diverge)



• Choice 2: Reduce γ_k as a function of k; will get close to a minimizer if u_0 is close enough

Problem: slow

Theorem. If u_0 is close enough to u^* and γ_k satisfies

$$-\sum_{k=0}^{\infty} \gamma_k = \infty$$
$$-\sum_{k=0}^{\infty} \gamma_k^2 < \infty$$

e.g. $\gamma_k = c/k$, then $u_k \to u^*$ as $k \to \infty$.

• Choice 3: **Armijo step-size:** Take as big a step as possible, but no larger Pick $\alpha \in (0,1)$, $\beta \in (0,1)$. Let *i* be the smallest non-negative integer such that

$$g\left(u_k - \beta^i \frac{\partial g^{\mathrm{T}}}{\partial u}(u_k)\right) - g(u_k) < -\alpha\beta^i \left\| \frac{\partial g}{\partial u}(u_k) \right\|^2$$
$$u_{k+1} = u_k - \beta^i \frac{\partial g^{\mathrm{T}}}{\partial u}(u_k)$$

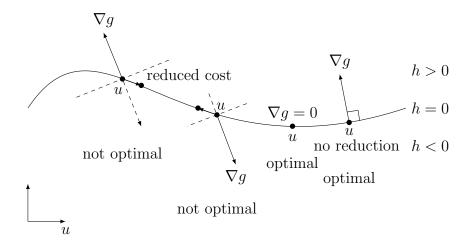
This will get to a minimizer blazingly fast if u_0 is close enough.

1.4 Constrained Optimization

Equality constraints:

$$\min_{u \in \mathcal{R}^m} g(u)$$
s.t. $h(u) = \mathbf{0}$

Consider $u \in \mathbb{R}^2$, $h: \mathbb{R}^2 \to \mathbb{R}$



So u is (locally) optimal if $\nabla g \parallel$ (is parallel to) the normal vector to tangent plane to h.

Fact: (HW# 1)

 $\nabla h \perp Th$ (tangent plane to h)



We need $\nabla g \parallel \nabla h$ at u^* for optimality, i.e.

$$\frac{\partial g}{\partial u}(u^*) = \alpha \frac{\partial h}{\partial u}(u^*), \text{ for some } \alpha \in \mathbb{R}$$

or $(\lambda = -\alpha)$,

$$\frac{\partial g}{\partial u}(u^*) + \lambda \frac{\partial h}{\partial u}(u^*) = 0$$

or

$$\frac{\partial}{\partial u} (g(u^*) + \lambda h(u^*)) = 0$$
, for some $\lambda \in \mathbb{R}$

More generally,

$$\min_{u \in \mathcal{R}^m} g(u)$$

s.t. $h(u) = \mathbf{0}, \quad h : \mathbb{R}^m \to \mathbb{R}^k$

Note that $h(u) = [h_1(u), ..., h_k(u)]^{T}$.

We need $\frac{\partial g}{\partial u}(u^*)$ to be a linear combination of $\frac{\partial h_i}{\partial u}(u^*)$, $i=1,\ldots,k$, for exactly the same reasons, i.e.

$$\frac{\partial g}{\partial u}(u^*) = \sum_{i=1}^k \alpha_i \frac{\partial h_i}{\partial u}(u^*)$$

or $(\lambda = -[\alpha_1, \dots, \alpha_k]^T)$

$$\frac{\partial g}{\partial u}(u^*) + \lambda^{\mathrm{T}} \frac{\partial h}{\partial u}(u^*) = 0$$

or

$$\frac{\partial}{\partial u} \big(g(u^*) + \lambda^{\mathrm{T}} h(u^*) \big) = 0, \quad \text{for some } \lambda \in \mathbb{R}^k$$

Theorem. If u^* is a minimizer to

$$\min_{u \in \mathcal{R}^m} g(u)$$
s.t. $h(u) = \mathbf{0}, \quad h : \mathbb{R}^m \to \mathbb{R}^k$

then $\exists \lambda \in \mathbb{R}^k \ s.t.$

$$\begin{cases} \frac{\partial L}{\partial u}(u^*, \lambda) = 0\\ \frac{\partial L}{\partial \lambda}(u^*, \lambda) = 0 \end{cases}$$

where the Lagrangian L is given by

$$L(u, \lambda) = g(u) + \lambda^T h(u)$$

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Note:

- λ are the Lagrange multipliers
- $\frac{\partial L}{\partial \lambda} = 0$ is fancy speak for $h(u^*) = 0$

Example

$$\min_{u \in \mathbb{R}^m} \frac{1}{2} ||u||^2$$

s.t. $Au = b$

where A is $k \times m$, $k \leq m$. Assume $(AA^{T})^{-1}$ exists (constraints are linearly independent, none of the constraints are "duplicates", all the constraints are essential).

$$L = \frac{1}{2}u^{\mathrm{T}}u + \lambda^{\mathrm{T}}(Au - b)$$
$$\frac{\partial L}{\partial u} = u^{\mathrm{T}} + \lambda^{\mathrm{T}}A = 0$$
$$u^* = -A^{\mathrm{T}}\lambda$$

Using the equality constraint,

$$Au^* = b$$

$$-AA^{T}\lambda = b$$

$$\lambda = -(AA^{T})^{-1}b$$

$$u^* = A^{T}(AA^{T})^{-1}b$$

Example

$$\min \ u_1 u_2 + u_2 u_3 + u_1 u_3$$
s.t. $u_1 + u_2 + u_3 = 3$

$$L = u_1 u_2 + u_2 u_3 + u_1 u_3 + \lambda (u_1 + u_2 + u_3 - 3)$$

$$\begin{cases} \frac{\partial L}{\partial u_1} = u_2 + u_3 + \lambda = 0 \\ \frac{\partial L}{\partial u_2} = u_1 + u_3 + \lambda = 0 \\ \frac{\partial L}{\partial u_3} = u_2 + u_1 + \lambda = 0 \\ \frac{\partial L}{\partial \lambda} = u_1 + u_2 + u_3 = 3 \end{cases} \implies \begin{cases} u_1^* = 1 \\ u_2^* = 1 \\ u_3^* = 1 \\ \lambda = -2 \end{cases}$$
 optimal solution

Note: This was actually the worst we can do—maximize! Even weirder: no local minimizer!

1.4.1 Equality Constraints

$$\min_{u \in \mathcal{R}^m} g(u)$$

s.t. $h(u) = \mathbf{0}, \quad h : \mathbb{R}^m \to \mathbb{R}^k$

Theorem. If u^* is a minimizer/maximizer then $\exists \lambda \in \mathbb{R}^k$ s.t.

$$\frac{\partial L}{\partial u}(u^*, \lambda) = 0$$

$$\frac{\partial L}{\partial \lambda}(u^*, \lambda) = 0 \qquad (\iff h(u^*) = 0)$$

where $L(u, \lambda) = g(u) + \lambda^T h(u)$.

Example [Entropy Maximization]

Given $S = \{x_1, \ldots, x_n\}$ and a distribution over S such that it takes the value x_j with probability p_j . The entropy is

$$E(p) = \sum_{j=1}^{n} (-p_j \ln p_j).$$

The mean is

$$m = \sum_{j=1}^{n} p_j x_j.$$

Problem: Given m, find p such that E is maximized.

$$\min_{p} - \sum_{j=1}^{n} p_{j} \ln p_{j}$$
s.t.
$$\sum_{j=1}^{n} p_{j} x_{j} = m$$

$$\sum_{j=1}^{n} p_{j} = 1$$

$$p_{j} \ge 0, \ j = 1, \dots, n \quad \text{(ignore this...)}$$

$$L = -\sum p_j \ln p_j + \lambda_1 \left[\sum p_j x_j - m \right] + \lambda_2 \left[\sum p_j - 1 \right]$$

$$\frac{\partial L}{\partial p_j} = -\ln p_j - 1 + \lambda_1 x_j + \lambda_2 = 0$$

$$p_j = e^{\lambda_2 - 1 + \lambda_1 x_j}, \quad j = 1, \dots, n \qquad (p_j \ge 0 \text{ so we're ok with ignoring that})$$

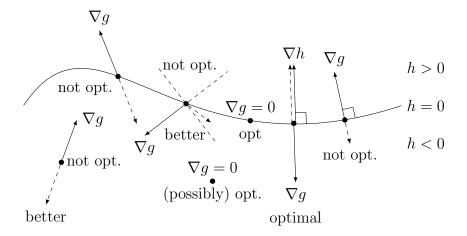
$$\sum e^{\lambda_2 - 1 + \lambda_1 x_j} x_j = m$$
 $n + 2$ equations and
$$\sum e^{\lambda_2 - 1 + \lambda_1 x_j} = 1$$
 $n + 2$ unknowns...

No analytical solution, but numerically "solvable"

1.4.2 Inequality Constraints

$$\min_{u \in \mathcal{R}^m} g(u)$$

s.t. $h(u) \le \mathbf{0}, \quad h : \mathbb{R}^m \to \mathbb{R}^k$



We need:

- if $h(u^*) < 0$ then $\frac{\partial g}{\partial u}(u^*) = 0$
- if $h(u^*) = 0$ then we need either

$$\frac{\partial g}{\partial u}(u^*) = 0$$

$$\frac{\partial g}{\partial u}(u^*) = -\lambda \frac{\partial h}{\partial u}(u^*) \quad \text{for } \lambda > 0$$

or

Or, even better,

$$\frac{\partial}{\partial u}(g(u^*) + \lambda h(u^*)) = 0 \text{ for } \lambda \ge 0,$$

where $\lambda h(u^*) = 0$. $(h < 0 \rightarrow \lambda = 0, h = 0 \rightarrow \lambda \ge 0)$

In general, if $u \in \mathbb{R}^m$ and $h: \mathbb{R}^m \to \mathbb{R}^k$, we have that u^* , if optimal, has to satisfy

$$\frac{\partial}{\partial u}L(u^*,\lambda) = 0$$
$$h(u^*) \le \mathbf{0}$$
$$\lambda^{\mathrm{T}}h(u^*) = 0$$
$$\lambda \ge \mathbf{0}$$

where the Lagrangian is $L(u, \lambda) = g(u) + \lambda^{T} h(u)$. Note that if we're maximizing, the same holds except we need $\lambda \leq 0$.

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Example

min
$$2u_1^2 + 2u_1u_2 + u_2^2 - 10u_1 - 10u_2$$

s.t.
$$\begin{cases} u_1^2 + u_2^2 \le 5\\ 3u_1 + u_2 \le 6 \end{cases}$$

$$L = 2u_1^2 + 2u_1u_2 + u_2^2 - 10u_1 - 10u_2 + \lambda_1(u_1^2 + u_2^2 - 5) + \lambda_2(3u_1 + u_2 - 6)$$

FONC:

i)
$$\partial L/\partial u_1 = 4u_1 + 2u_2 - 10 + 2\lambda_1 u_1 + 3\lambda_2$$

ii)
$$\partial L/\partial u_2 = 2u_1 + 2u_2 - 10 + 2\lambda_1 u_2 + \lambda_2$$

iii)
$$u_1^2 + u_2^2 \le 5$$

iv)
$$3u_1 + u_2 \le 6$$

v)
$$\lambda_1(u_1^2 + u_2^2 - 5) = 0$$

vi)
$$\lambda_2(3u_1 + u_2 - 6) = 0$$

vii)
$$\lambda_1 \geq 0$$

viii)
$$\lambda_2 \geq 0$$

To solve, assume different constraints are active/inactive:

1. Both constraints are inactive $(u_1^2 + u_2^2 < 5, 3u_1 + u_2 < 6) \Longrightarrow \lambda_1 = \lambda_2 = 0$

$$\begin{cases} 4u_1 + 2u_2 - 10 = 0 \\ 2u_1 + 2u_2 - 10 = 0 \end{cases} \implies \begin{cases} u_1 = 0 \\ u_2 = 5 \end{cases}$$

Note: iii) $0^2 + 5^2 \nleq 5$

Not feasible

2. Assume constraint 1 is active and constraint 2 is inactive $(u_1^2 + u_2^2 = 5, \lambda_2 = 0)$

$$\begin{cases} 4u_1 + 2u_2 - 10 + 2\lambda_1 u_1 = 0 \\ 2u_1 + 2u_2 - 10 + 2\lambda_1 u_2 = 0 \\ u_1^2 + u_2^2 = 5 \end{cases} \implies \begin{cases} u_1 = 1 \\ u_2 = 2 \\ \lambda_1 = 1 \end{cases}$$

This is a local minimizer

- 3. Assume constraint 2 is active and constraint 1 is inactive
- 4. Assume both constraints are active

Kuhn-Tucker Conditions (KKT conditions, Karush-Kuhn-Tucker)

Problem:

$$\min_{u \in \mathbb{R}^m} g(u)$$
s.t.
$$\begin{cases}
h_1(u) = 0, & h_1 : \mathbb{R}^m \to \mathbb{R}^p \\
h_2(u) \le 0, & h_2 : \mathbb{R}^m \to \mathbb{R}^k
\end{cases}$$
(1.1)

Theorem. Let u^* be feasible $(h_1 = 0, h_2 \le 0)$. If u^* is a minimizer to (1.1) than there exists vectors $\lambda \in \mathbb{R}^p$, $\mu \in \mathbb{R}^k$ with $\mu \ge \mathbf{0}$ such that

$$\begin{cases} \frac{\partial g}{\partial u}(u^*) + \lambda^T \frac{\partial h_1}{\partial u}(u^*) + \mu^T \frac{\partial h_2}{\partial u}(u^*) = 0\\ \mu^T h_2(u^*) = 0 \end{cases}$$

Looking ahead: $\min \operatorname{cost}(u(\cdot))$ s.t. $\dot{x} = f(x, u)$ (dynamics), where u is a function. Note the equality constraint.

Question: How do we go from $u \in \mathbb{R}^m$ to $u \in \mathcal{U}$ (function space)?

Note: Function space is a set of functions of a given kind from a set X to a set Y

- 1. linear function
- 2. square-integrable functions: $L_2[0,T]: \int_0^T \|u(t)\|^2 dt < \infty$
- 3. $C^{\infty}(\mathbb{R})$

What would ∂ "cost" $/\partial u$ mean?

Chapter 2

Calculus of Variations

2.1 Directional Derivatives

Recall: To minimize g(u), let u^* be a candidate minimizer and pitch a perturbation on u^* of εv , where ε is the scale and v is the direction. Taking Taylor's expansion at the perturbation produces



$$g(u^* + \varepsilon v) = g(u^*) + \varepsilon \frac{\partial g}{\partial u}(u^*)v + o(\varepsilon)$$

FONC: $\frac{\partial g}{\partial u}(u^*) = 0$

Note: $\frac{\partial g}{\partial u}(u^*)v$ tells us how much g(u) increases/decreases in the direction of v.

Definition. The directional (Gateaux) derivative is given by

$$\delta g(u;v) = \lim_{\varepsilon \to 0} \frac{g(u + \varepsilon v) - g(u)}{\varepsilon}$$

Example

$$g(u) = \frac{1}{2}u_1^2 - u_1 + 2u_2, \quad g: \mathbb{R}^2 \to \mathbb{R}$$

Let's consider $e_1 = [1 \ 0]^T$, $e_2 = [0 \ 1]^T$. What is $\delta g(u; e_i)$, i = 1, 2?

$$\delta g(u; v) = \lim_{\varepsilon \to 0} \frac{g(u + \varepsilon v) - g(u)}{\varepsilon}$$

$$= \lim_{\varepsilon \to 0} \frac{g(u) + \varepsilon \frac{\partial g}{\partial u}(u)v + o(\varepsilon) - g(u)}{\varepsilon}$$

$$= \frac{\partial g}{\partial u}(u)v$$

$$\frac{\partial g}{\partial u}(u) = [u_1 - 1 \ 2]$$

$$\delta g(u; e_1) = [u_1 - 1 \ 2]e_1 = u_1 - 1$$

$$\delta g(u; e_2) = [u_1 - 1 \ 2]e_2 = 2$$

But the beauty of directional derivatives is that they generalize beyond vectors, $u \in \mathbb{R}^m$, to function spaces (\mathcal{U}) or other "objects" like matrices.

Example $M \in \mathbb{R}^{n \times n}$, $F(M) = M^2$ What is $\frac{\partial F}{\partial M}$? (ponder at home...) We can easily compute $\delta F(M; N)$!

$$\begin{split} F(M+\varepsilon N) &= (M+\varepsilon N)(M+\varepsilon N) = M^2 + \varepsilon MN + \varepsilon NM + \varepsilon^2 N^2 \\ \delta F(M;N) &= \lim_{\varepsilon \to 0} \frac{F(M+\varepsilon N) - F(M)}{\varepsilon} \\ &= \lim_{\varepsilon \to 0} \frac{\varepsilon MN + \varepsilon NM + \varepsilon^2 N^2}{\varepsilon} = MN + NM \end{split}$$

Infinite Dimensional Optimization Let $u \in \mathcal{U}$ (function space) and let J(u) be the cost:

$$\min_{u \in \mathcal{U}} J(u)$$

Theorem. If $u^* \in \mathcal{U}$ is a (local) minimizer then

$$\delta J(u^*;v) = 0, \quad \forall v \in \mathcal{U}$$

Example Find minimizer u^* to

$$J(u) = \int_0^T L(u(t)) \, \mathrm{d}t$$

$$\begin{split} J(u+\varepsilon v) - J(u) &= \int_0^T L(u(t)+\varepsilon v(t)) \, \mathrm{d}t - \int_0^T L(u(t)) \, \mathrm{d}t, \quad u,v \in \mathcal{U} \\ &= \int_0^T \left[L(u(t)) + \varepsilon \frac{\partial L}{\partial u}(u(t)) v(t) + o(\varepsilon) - L(u(t)) \right] \mathrm{d}t \\ \delta J(u^*;v) &= \lim_{\varepsilon \to 0} \frac{J(u+\varepsilon v) - J(u)}{\varepsilon} \\ &= \lim_{\varepsilon \to 0} \frac{\int_0^T \varepsilon \frac{\partial L}{\partial u}(u(t)) v(t) \, \mathrm{d}t + o(\varepsilon)}{\varepsilon} \\ &= \int_0^T \frac{\partial L}{\partial u}(u(t)) v(t) \, \mathrm{d}t \end{split}$$

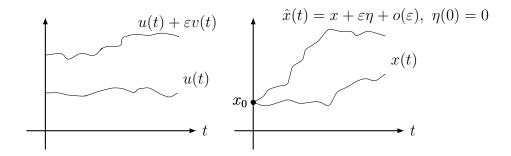


Figure 2.1: Variation in u causes a variation in x.

 u^* optimizer:

$$\delta J(u^*; v) = \int_0^T \frac{\partial L}{\partial u}(u(t))v(t) dt = 0 \quad \forall v \in \mathcal{U}$$

$$\updownarrow$$

$$\frac{\partial L}{\partial u}(u(t)) = 0 \quad \forall t \in [0, T]$$

But, we want optimal control! We want our cost to look like

$$\int_0^T L(x(t), u(t)) dt$$
$$\dot{x} = f(x, u)$$

2.2 Calculus of Variations

What happens to x(t) when u(t) changes to $u(t) + \varepsilon v(t)$? Let the system be given by

$$\begin{cases} \dot{x} = f(x, u) \\ x(0) = x_0 \end{cases}$$

After perturbation of u, the new system is

$$\begin{cases} \dot{\hat{x}} = f(\hat{x}, u + \varepsilon v) \\ x(0) = x_0 \end{cases}$$

Consider

$$\tilde{x} = x + \varepsilon \eta,$$

where

$$\dot{x} = f(x, u),$$
 $x(0) = x_0$
 $\dot{\eta} = \frac{\partial f}{\partial x}(x, u)\eta + \frac{\partial f}{\partial u}(x, u)v,$ $\eta(0) = 0$

Theorem. If f is continuously differentiable in x and u then

$$\hat{x}(t) = \tilde{x}(t) + o(\varepsilon)$$

Proof.

i) Initial conditions:

$$\hat{x}(0) = x_0$$

$$\tilde{x}(0) = x(0) + \varepsilon \eta(0) = x_0$$

ii) Dynamics:

$$\begin{split} \dot{\hat{x}} &= f(\hat{x}, u + \varepsilon v) \\ \dot{\hat{x}} &= \dot{x} + \varepsilon \dot{\eta} = f(x, u) + \varepsilon \frac{\partial f}{\partial x}(x, u) \eta + \varepsilon \frac{\partial f}{\partial u}(x, u) v \\ &= f(x + \varepsilon \eta, u + \varepsilon v) + o(\varepsilon) \\ &= f(\tilde{x}, u + \varepsilon v) + o(\varepsilon) \end{split}$$

We can see that the dynamics of $\hat{x}(t)$ are equal to those of $\tilde{x}(t)$ plus higher order terms:

$$\dot{\tilde{x}} = f(\tilde{x}, u + \varepsilon v) + o(\varepsilon)$$
$$\dot{\hat{x}} = f(\hat{x}, u + \varepsilon v)$$

Therefore, if our perturbation is small enough, we can model $\hat{x}(t)$ as $\tilde{x}(t)$.

Note: Taylor expansion with two elements is

$$h(w + \varepsilon v, z + \varepsilon y) = h(w, z + \varepsilon y) + \frac{\partial h}{\partial w}(w, z + \varepsilon y)\varepsilon v + o(\varepsilon)$$

$$= \left\{h(w, z) + \frac{\partial h}{\partial z}(w, z)\varepsilon y + o(\varepsilon)\right\}$$

$$+ \left\{\frac{\partial h}{\partial w}(w, z)\varepsilon v + \underbrace{\frac{\partial^2 h}{\partial z\partial w}\varepsilon v \odot \varepsilon y}_{o(\varepsilon)} + o(\varepsilon)\right\}$$

$$= h(w, z) + \frac{\partial h}{\partial z}\varepsilon y + \frac{\partial h}{\partial w}\varepsilon v + o(\varepsilon)$$

Last class:

1. $u \in \mathcal{U}$ (space of functions), $J : \mathcal{U} \to \mathbb{R}$ (cost).

FONC: If u^* is optimal, then

$$\delta J(u; \nu) = 0 \quad \forall \nu \in \mathcal{U},$$

where the directional derivative is given by

$$\delta J(u;\nu) = \lim_{\varepsilon \to 0} \frac{J(u+\varepsilon\nu) - J(u)}{\varepsilon}.$$

2. If

$$\begin{cases} \dot{x} = f(x, u) \\ x(0) = x_0 \end{cases}$$

then a variation in u:

$$u \longmapsto u + \varepsilon \nu$$

results in a variation in x:

$$x \longmapsto x + \varepsilon \eta + o(\varepsilon)$$

See Figure 2.1. Note $\eta(0) = 0$.

2.2.1 An (Almost) Optimal Control Problem

Let $\dot{x} = f(x)$, $x(0) = x_0$. Note we get to pick the initial condition!

Problem

$$\min_{x_0 \in \mathbb{R}^m} J(x_0) = \int_0^T L(x(t)) dt$$
 s.t.
$$\begin{cases} \dot{x}(t) = f(x(t)) & \text{the } constraint! \text{ (equality)} \\ x(0) = x_0 \end{cases}$$

Note every constraint needs a Lagrange multiplier. We have infinitely many constraints:

$$\dot{x}(t) = f(x(t)) \quad \forall t \in [0, T]$$

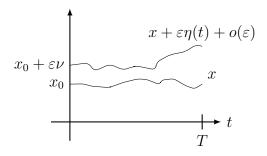
We need $\lambda(t)$ as a function of t. Also, the sum in the Lagrangian has to become an integral. The continuous-time Lagrangian thus becomes

$$\tilde{J}(x_0, \lambda) = \int_0^T \left[L(x(t)) + \lambda^{\mathrm{T}}(t) (f(x(t)) - \dot{x}(t)) \right] dt$$

The task is to perturb x_0 as $x_0 \longmapsto x_0 + \varepsilon \nu$, $\nu \in \mathbb{R}^m$ and compute

$$\delta \tilde{J}(x_0; \nu) = \lim_{\varepsilon \to 0} \frac{\tilde{J}(x_0 + \varepsilon \nu) - \tilde{J}(x_0)}{\varepsilon}$$

and make this equal to $0 \ \forall \nu \in \mathbb{R}^m$. The variation in x is



Note:

 x_0 decision variable

 ν variation in x_0

x(t) trajectory starting at x_0

 $\eta(t)$ change in trajectory resulting from ν -variation in x_0

 $\lambda(t)$ time-varying Lagrange multiplier

$$\begin{split} \tilde{J}(x_0 + \varepsilon \nu) &= \int_0^T \left\{ L(x(t)) + \lambda^{\mathrm{T}}(t) [f(x(t) + \varepsilon \eta(t)) - \dot{x}(t) - \varepsilon \dot{\eta}(t)] \right\} \mathrm{d}t + o(\varepsilon) \\ &= \int_0^T \left[L(x) + \varepsilon \frac{\partial L}{\partial x}(x) \eta + \lambda^{\mathrm{T}} \left(f(x) + \varepsilon \frac{\partial f}{\partial x}(x) \eta - \dot{x} - \varepsilon \dot{\eta} \right) \right] \mathrm{d}t + o(\varepsilon) \\ \tilde{J}(x_0 + \varepsilon \nu) - \tilde{J}(x_0) &= \int_0^T \left[\varepsilon \frac{\partial L}{\partial x}(x) \eta + \lambda^{\mathrm{T}} \left(\varepsilon \frac{\partial f}{\partial x} \eta - \varepsilon \dot{\eta} \right) \right] \mathrm{d}t + o(\varepsilon) \\ \delta \tilde{J}(x_0; \nu) &= \int_0^T \left[\frac{\partial L}{\partial x}(x) \eta + \lambda^{\mathrm{T}} \left(\frac{\partial f}{\partial x} \eta - \dot{\eta} \right) \right] \mathrm{d}t \end{split}$$

A powerful idea: we want $\delta \tilde{J}(x_0; \nu) = 0 \ \forall \nu$. Somehow get this in the form

$$\int_0^T \left(\operatorname{stuff}(t) \right) \eta(t) \, \mathrm{d}t = 0$$

We can pick $stuff(t) = 0 \ \forall t \in [0, T].$

In $\delta \tilde{J}(x_0; \nu)$ we have $\dot{\eta}$ (problem!). We can solve this using integration by parts.

$$\int_0^T \lambda^{\mathrm{T}} \dot{\eta} \, \mathrm{d}t = \lambda^{\mathrm{T}}(T) \eta(T) - \lambda^{\mathrm{T}}(0) \eta(0) - \int_0^T \dot{\lambda}^{\mathrm{T}} \eta \, \mathrm{d}t$$

Hence,

$$\delta \tilde{J}(x_0; \nu) = \int_0^T \underbrace{\left(\frac{\partial L}{\partial x} + \lambda^{\mathrm{T}} \frac{\partial f}{\partial x} + \dot{\lambda}^{\mathrm{T}}\right)}_{\mathrm{pick} = 0} \eta \, \mathrm{d}t - \underbrace{\lambda^{\mathrm{T}}(T)}_{\mathrm{pick} = 0} \eta(T) + \lambda^{\mathrm{T}}(0) \underbrace{\eta(0)}_{\nu}$$

We are free to pick λ freely if it gives $\delta \tilde{J} = 0$.

Pick:
$$\begin{cases} \dot{\lambda}(t) = -\frac{\partial L^{\mathrm{T}}}{\partial x}(x(t)) - \frac{\partial f^{\mathrm{T}}}{\partial x}(x(t))\lambda(t) \\ \lambda(T) = 0 \end{cases}$$
 backwards diff. eq

Under this choice of λ we get

$$\delta \tilde{J}(x_0; \nu) = \lambda^{\mathrm{T}}(0)\nu$$

This is linear in ν so the FONC is $\lambda(0) = 0$.

Moreover, we really have a "normal" optimization problem

$$\min_{x_0 \in \mathbb{R}^m} \tilde{J}(x_0)$$
$$\delta \tilde{J}(x_0; \nu) = \frac{\partial \tilde{J}}{\partial x_0}(x_0)\nu$$

which means that

$$\frac{\partial \tilde{J}}{\partial x_0} = \lambda^{\mathrm{T}}(0)$$

If x_0^* minimizes

$$\text{s.t. } \begin{cases} \int_0^T L(x(t)) \, \mathrm{d}t \\ \\ \dot{x}(t) = f(x(t)) \\ x(0) = x_0^* \end{cases}$$

then

$$\lambda(0) = \mathbf{0}$$

where $\lambda(t)$ satisfies

$$\begin{cases} \dot{\lambda}(t) = -\frac{\partial L^{\mathrm{T}}}{\partial x}(x(t)) - \frac{\partial f^{\mathrm{T}}}{\partial x}(x(t))\lambda(t) \\ \lambda(T) = 0 \end{cases}$$

So what? We actually have a two-point boundary value problem.

$$\dot{x} = f(x) \qquad \qquad \dot{\lambda} = -\frac{\partial L^{\mathrm{T}}}{\partial x} - \frac{\partial f^{\mathrm{T}}}{\partial x} \lambda$$

$$x(0) = x_0 \qquad \qquad \lambda(T) = 0$$

$$x_0 \qquad \qquad \lambda(0) \qquad \qquad \lambda$$

We want to find x_0 that gives f(x) such that after solving backwards for $\lambda(t)$, we find that

$$\lambda(0) = \frac{\partial \tilde{J}^{\mathrm{T}}}{\partial x_0} = 0.$$

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This leads to the following:

```
Pick x_{0,0}

k=1

repeat

Simulate x(t) from x_{0,k} over [0,T]

Simulate \lambda(t) from \lambda(T)=0 backwards using x(t)

Update x_{0,k} as x_{0,k+1}=x_{0,k}-\gamma\lambda(0) \Rightarrow \lambda(0) is the gradient k:=k+1

until \lambda(0)=0
```

An algorithm

Example: optinit.m

$$\dot{x} = Ax, \quad L = x^{\mathrm{T}}Qx - q, \quad Q = Q^{\mathrm{T}} \succ 0$$

$$\dot{\lambda} = -2Qx - A^{\mathrm{T}}\lambda$$

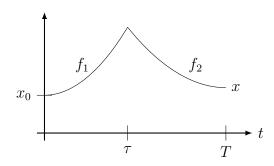
$$\lambda(0) = 0$$

2.2.2 Optimal Timing Control

When to switch between modes?

$$\dot{x} = \begin{cases} f_1(x) & \text{if } t \in [0, \tau) \\ f_2(x) & \text{if } t \in [\tau, T] \end{cases}$$

$$x(0) = x_0 \tag{2.1}$$

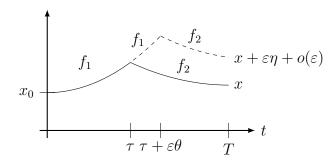


$$\min_{\tau} \int_{0}^{T} L(x(t)) dt = J(\tau)$$
 s.t. (2.1) holds

Step 1: Augment cost with constraint

$$\tilde{J} = \int_0^{\tau} \left[L(x) + \lambda^{\mathrm{T}} (f_1(x) - \dot{x}) \right] dt + \int_{\tau}^{T} \left[L(x) + \lambda^{\mathrm{T}} (f_2(x) - \dot{x}) \right] dt$$

Step 2: Variation $\tau \longmapsto \tau + \varepsilon \theta$



Step 3: Compute $\delta \tilde{J}(\tau;\theta)$

$$\tilde{J}(\tau + \varepsilon \theta) = \int_0^{\tau + \varepsilon \theta} \left\{ L(x + \varepsilon \eta) + \lambda^{\mathrm{T}} [f_1(x + \varepsilon \eta) - \dot{x} - \varepsilon \dot{\eta}] \right\} dt
+ \int_{\tau + \varepsilon \theta}^T \left\{ L(x + \varepsilon \eta) + \lambda^{\mathrm{T}} [f_2(x + \varepsilon \eta) - \dot{x} - \varepsilon \dot{\eta}] \right\} dt + o(\varepsilon)$$

Note that $\eta = \dot{\eta} = 0$ on $[0, \tau)$.

$$\tilde{J}(\tau + \varepsilon\theta) = \int_{0}^{\tau} \left\{ L(x) + \lambda^{\mathrm{T}} [f_{1}(x) - \dot{x}] \right\} dt
+ \int_{\tau}^{\tau + \varepsilon\theta} \left\{ \underbrace{L(x + \varepsilon\eta)}_{L(x) + \varepsilon \frac{\partial L}{\partial x} \eta} + \lambda^{\mathrm{T}} \underbrace{\left[f_{1}(x + \varepsilon\eta) - \dot{x} - \varepsilon \dot{\eta} \right]}_{f_{1}(x) + \varepsilon \frac{\partial f_{1}}{\partial x} \eta} + \int_{\tau + \varepsilon\theta}^{T} \left\{ \underbrace{L(x + \varepsilon\eta)}_{L(x) + \varepsilon \frac{\partial L}{\partial x} \eta} + \lambda^{\mathrm{T}} \underbrace{\left[f_{2}(x + \varepsilon\eta) - \dot{x} - \varepsilon \dot{\eta} \right]}_{f_{2}(x) + \varepsilon \frac{\partial f_{2}}{\partial x} \eta} + \varepsilon \dot{\eta} \right\} dt + o(\varepsilon)$$

$$\delta \tilde{J}(\tau;\theta) = \lim_{\varepsilon \to 0} \frac{\tilde{J}(\tau + \varepsilon\theta) - \tilde{J}(\tau)}{\varepsilon}$$

$$\tilde{J}(\tau + \varepsilon\theta) - \tilde{J}(\tau) = \int_{0}^{\tau} 0 \cdot dt + \underbrace{\int_{\tau}^{\tau + \varepsilon\theta} \left[\varepsilon \frac{\partial L}{\partial x} \eta + \lambda^{\mathrm{T}} \left(f_{1}(x) + \varepsilon \frac{\partial f_{1}}{\partial x} \eta - f_{2}(x) - \varepsilon \dot{\eta} \right) \right] dt}_{I_{1}}$$

$$+ \underbrace{\int_{\tau + \varepsilon\theta}^{T} \left[\varepsilon \frac{\partial L}{\partial x} \eta + \lambda^{\mathrm{T}} \left(\varepsilon \frac{\partial f_{2}}{\partial x} \eta - \varepsilon \dot{\eta} \right) \right] dt}_{I_{2}} + o(\varepsilon)$$

Theorem (Mean-value theorem).

$$\int_{t_1}^{t_2} h(t) dt = (t_2 - t_1)h(\xi) \quad \text{for some } \xi \in [t_1, t_2]$$

The first integral is

$$I_{1} = \int_{\tau}^{\tau + \varepsilon \theta} \left\{ \varepsilon \frac{\partial L}{\partial x} \eta + \lambda^{\mathrm{T}} \left[f_{1}(x) + \varepsilon \frac{\partial f_{1}}{\partial x} \eta - \varepsilon \dot{\eta} - f_{2}(x) \right] \right\} dt$$
$$= \varepsilon \theta \left\{ \lambda^{\mathrm{T}}(\xi) \left[f_{1}(x(\xi)) - f_{2}(x(\xi)) \right] \right\} + o(\varepsilon)$$

Note that as $\varepsilon \to 0$, $\xi \to \tau$. Using integration by parts, the second integral is

$$\begin{split} \int_{\tau}^{T} \lambda^{\mathrm{T}} \dot{\eta} \, \mathrm{d}t &= \lambda^{\mathrm{T}}(T) \eta(T) - \lambda^{\mathrm{T}}(\tau) \underbrace{\eta(\tau)}_{=0} - \int_{\tau}^{T} \dot{\lambda}^{\mathrm{T}} \eta \, \mathrm{d}t \\ I_{2} &= \int_{\tau}^{T} \left[\varepsilon \frac{\partial L}{\partial x} \eta + \lambda^{\mathrm{T}} \left(\varepsilon \frac{\partial f_{2}}{\partial x} \eta - \varepsilon \dot{\eta} \right) \right] \mathrm{d}t - \underbrace{\int_{\tau}^{\tau + \varepsilon \theta} \left[\varepsilon \frac{\partial L}{\partial x} \eta + \lambda^{\mathrm{T}} \left(\varepsilon \frac{\partial f_{2}}{\partial x} \eta - \varepsilon \dot{\eta} \right) \right] \mathrm{d}t}_{o(\varepsilon)} \\ &= \varepsilon \int_{\tau}^{T} \left[\frac{\partial L}{\partial x} + \lambda^{\mathrm{T}} \frac{\partial f_{2}}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right] \eta \, \mathrm{d}t - \varepsilon \lambda^{\mathrm{T}}(T) \eta(T) + o(\varepsilon) \end{split}$$

Hence,

$$\delta \tilde{J}(\tau;\theta) = \lim_{\varepsilon \to 0} \frac{\tilde{J}(\tau + \varepsilon\theta) - \tilde{J}(\tau)}{\varepsilon}$$
$$= \theta \lambda^{\mathrm{T}}(\tau) \left[f_1(x(\tau)) - f_2(x(\tau)) \right] + \int_{\tau}^{T} \left[\frac{\partial L}{\partial x} + \lambda^{\mathrm{T}} \frac{\partial f_2}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right] \eta \, \mathrm{d}t - \lambda^{\mathrm{T}}(T) \eta(T)$$

Step 4: Select the costate $\lambda(t)$. The key idea is to get rid of any term that has η in it, i.e.

$$\dot{\lambda} = -\frac{\partial L^{\mathrm{T}}}{\partial x} - \frac{\partial f_2^{\mathrm{T}}}{\partial x} \lambda \quad \text{on } [\tau, T]$$
$$\lambda(T) = 0$$

Step 5: With this choice of $\lambda(t)$, we have

$$\delta \tilde{J}(\tau;\theta) = \theta \lambda^{\mathrm{T}}(\tau) \Big[f_1(x(\tau)) - f_2(x(\tau)) \Big] = \frac{\partial \tilde{J}}{\partial \tau} \theta.$$

Therefore,

$$\frac{\partial \tilde{J}}{\partial \tau} = \lambda^{\mathrm{T}}(\tau) \left[f_1(x(\tau)) - f_2(x(\tau)) \right] = 0 \quad \text{(for optimality)}$$

Algorithm

```
Pick \tau_0
k = 0
repeat

Simulate x forward in time from x(0) = x_0

Simulate \lambda backwards from \lambda(T) = 0

Update \tau_k as \tau_{k+1} = \tau_k - \gamma \lambda^{\mathrm{T}}(\tau_k) \left[ f_1(x(\tau_k)) - f_2(x(\tau_k)) \right]
k := k + 1

until \|\lambda^{\mathrm{T}}(f_1 - f_2)\| < \varepsilon
```

Where are we going? Come up with general principles for $\min_{u \in \mathcal{U}} J(u)$:

- Costate equations
- Optimality conditions
- Algorithms
- Applications

Chapter 3

The Maximum Principle

3.1 The Bolza Problem

Up until now, we have optimized with respect to finite-dimensional parameters. Today, we will minimize with respect to $u \in \mathcal{U}$.

$$\min_{u \in \mathcal{U}} J(u) = \int_0^T L(x(t), u(t), t) \, \mathrm{d}t + \underbrace{\Psi(x(T))}_{\substack{\text{terminal cost} \\ (\text{parking cost})}}$$
 s.t.
$$\dot{x}(t) = f(x(t), u(t), t)$$

$$x(0) = x_0$$

Assume that f and L are C^1 in x, u and piecewise continuous in t. Then, a small change in u causes small changes in f and L. The variation: $u \mapsto u + \varepsilon v$, $\varepsilon \in \mathbb{R}$, $v \in \mathcal{U}$. See Figure 2.1.

$$\begin{split} \tilde{J}(u) &= \int_0^T \left[L(x,u,t) + \lambda^{\mathrm{T}} (f(x,u,t) - \dot{x}) \right] \mathrm{d}t + \Psi(x(T)) \\ \tilde{J}(u + \varepsilon v) &= \int_0^T \left[L(x + \varepsilon \eta, u + \varepsilon v, t) + \lambda^{\mathrm{T}} (f(x + \varepsilon \eta, u + \varepsilon v, t) - \dot{x} - \varepsilon \dot{\eta}) \right] \mathrm{d}t \\ &+ \Psi(x(T) + \varepsilon \eta(T)) + o(\varepsilon) \\ \tilde{J}(u + \varepsilon v) - \tilde{J}(u) &= \int_0^T \left[L(x + \varepsilon \eta, u + \varepsilon v, t) - L(x, u, t) \right. \\ &+ \lambda^{\mathrm{T}} \left(f(x + \varepsilon \eta, u + \varepsilon v, t) - f(x, u, t) - \dot{x} - \varepsilon \dot{\eta} + \dot{x} \right) \right] \mathrm{d}t \\ &+ \Psi(x(T) + \varepsilon \eta(T)) - \Psi(x(T)) + o(\varepsilon) \\ &= \int_0^T \left[\frac{\partial L}{\partial x} \varepsilon \eta + \frac{\partial L}{\partial u} \varepsilon v + \lambda^{\mathrm{T}} \left(\frac{\partial f}{\partial x} \varepsilon \eta + \frac{\partial f}{\partial u} \varepsilon v - \varepsilon \dot{\eta} \right) \right] \mathrm{d}t \\ &+ \frac{\partial \Psi}{\partial x} (x(T)) \varepsilon \eta(T) + o(\varepsilon) \end{split}$$

(See Taylor expansion with respect to two variables.)

$$\delta \tilde{J}(u;v) = \int_0^T \left(\frac{\partial L}{\partial u} + \lambda^{\mathrm{T}} \frac{\partial f}{\partial u} \right) v \, \mathrm{d}t + \int_0^T \left[\left(\frac{\partial L}{\partial x} + \lambda^{\mathrm{T}} \frac{\partial f}{\partial x} \right) \eta - \lambda^{\mathrm{T}} \dot{\eta} \right] \mathrm{d}t + \frac{\partial \Psi}{\partial x} (x(T)) \eta(T)$$

Integrating by parts,

$$\begin{split} \int_0^T \lambda^\mathrm{T} \dot{\eta} \, \mathrm{d}t &= \lambda^\mathrm{T}(T) \eta(T) - \lambda^\mathrm{T}(0) \eta(0) - \int_0^T \dot{\lambda}^\mathrm{T} \eta \, \mathrm{d}t \\ &= \lambda^\mathrm{T}(T) \eta(T) - \int_0^T \dot{\lambda}^\mathrm{T} \eta \, \mathrm{d}t \\ \delta \tilde{J}(u;v) &= \int_0^T \left(\frac{\partial L}{\partial u} + \lambda^\mathrm{T} \frac{\partial f}{\partial u} \right) v \, \mathrm{d}t + \int_0^T \left(\frac{\partial L}{\partial x} + \lambda^\mathrm{T} \frac{\partial f}{\partial x} + \dot{\lambda}^\mathrm{T} \right) \eta \, \mathrm{d}t \\ &+ \left(\frac{\partial \Psi}{\partial x} (x(T)) - \lambda^\mathrm{T}(T) \right) \eta(T) \end{split}$$

For optimality, we need the directional derivative to be zero for every $v \in \mathcal{U}$, where v represents the direction of the derivative. Therefore, the term $(\frac{\partial L}{\partial u} + \lambda^{\mathrm{T}} \frac{\partial f}{\partial u})$ in the first integral has to be identically zero. Thus, we need

$$\begin{cases} \frac{\partial L}{\partial u} + \lambda^{\mathrm{T}} \frac{\partial f}{\partial u} = 0, & \forall t \in [0, T] \\ \frac{\partial L}{\partial x} + \lambda^{\mathrm{T}} \frac{\partial f}{\partial x} + \dot{\lambda}^{\mathrm{T}} = 0, & \forall t \in [0, T] \\ \frac{\partial \Psi}{\partial x} (x(T)) - \lambda^{\mathrm{T}} (T) = 0 \end{cases}$$

Definition. Let the *Hamiltonian* $H(x, u, t, \lambda)$ be given by

$$H(x, u, t, \lambda) = L(x, u, t) + \lambda^{\mathrm{T}} f(x, u, t)$$

Theorem. For u to solve the Bolza problem, it has to satisfy

$$\frac{\partial H}{\partial u}(x, u, t, \lambda) = 0,$$

where the costate satisfies

$$\begin{cases} \dot{\lambda} = -\frac{\partial H^T}{\partial x}(x, u, t, \lambda) \\ \lambda(T) = \frac{\partial \Psi^T}{\partial x}(x(T)) \end{cases}$$

Example

$$\min_{u} \int_{0}^{1} \frac{1}{2} u^{2}(t) dt + \frac{1}{2} x^{2}(1)$$
s.t.
$$\begin{cases} \dot{x} = u, & x, u \in \mathbb{R} \\ x(0) = 1 \end{cases}$$

$$H = \frac{1}{2}u^2 + \lambda u$$

$$\frac{\partial H}{\partial u} = u + \lambda = 0 \Longrightarrow u = -\lambda$$

$$\dot{\lambda} = -\frac{\partial H}{\partial x} = 0 \Longrightarrow \lambda(t) = c$$

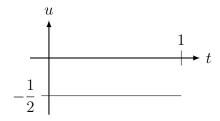
$$\lambda(T) = c = \frac{\partial \Psi}{\partial x}(x(1)) = x(1)$$

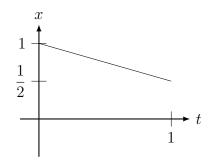
$$\dot{x} = u = -c \Longrightarrow x(t) = -ct + x(0) = -ct + 1$$

$$x(1) = -c + 1$$

$$\lambda(1) = c = x(1) = -c + 1 \Longrightarrow c = \frac{1}{2}$$

$$u^* = -\frac{1}{2}$$





We really used five different equations to solve this!

i)
$$\frac{\partial H}{\partial u} = 0$$

ii)
$$\dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x}$$

iii)
$$\lambda(T) = \frac{\partial \Psi^{\mathrm{T}}}{\partial x}(x(T))$$

iv)
$$\dot{x} = f(x, u, t)$$

v)
$$x(0) = x_0$$

There is a sixth condition that is pretty useful if L and f do not depend on t (L(x, u), f(x, u)). This is called a *conservative system*. Then, along optimal trajectories (equations i-v are satisfied), the total time derivative of the Hamiltonian is

$$\frac{\mathrm{d}}{\mathrm{d}t}H = \underbrace{\frac{\partial H}{\partial t}}_{H(x,u,\lambda)} + \underbrace{\frac{\partial H}{\partial x}}_{-\dot{\lambda}^{\mathrm{T}}} \dot{x} + \underbrace{\frac{\partial H}{\partial u}}_{u \text{ is optimal}} \dot{u} + \underbrace{\frac{\partial H}{\partial \lambda}}_{f^{\mathrm{T}} = \dot{x}^{\mathrm{T}}} \dot{\lambda} = -\dot{\lambda}^{\mathrm{T}} \dot{x} + \dot{x}^{\mathrm{T}} \dot{\lambda} = 0$$

Therefore, for conservative systems,

vi) H is constant along optimal trajectories. (Hamilton's Principle in analytical mechanics)

Back to the example,

$$H = \frac{1}{2}u^2 + \lambda u = \frac{1}{2}c^2 - c^2 = -\frac{1}{2}c^2 = -\frac{1}{8}$$

The Hamiltonian

$$H(x, u, t, \lambda) = L(x, u, t) + \lambda^{\mathrm{T}} f(x, u, t)$$

lets us write the Lagrangian as

$$\tilde{J}(u) = \int_0^T \left[L + \lambda^{\mathrm{T}} (f - \dot{x}) \right] \mathrm{d}t + \Psi = \int_0^T \left(H - \lambda^{\mathrm{T}} \dot{x} \right) \mathrm{d}t + \Psi$$

The optimality conditions are

$$\frac{\partial H}{\partial u} = 0, (3.1)$$

where

$$\begin{cases} \dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x} \\ \lambda(T) = \frac{\partial \Psi}{\partial x}(x(T)) \end{cases}$$
(3.2)

Example Hamilton's Principle

Let q be the generalized coordinates (positions and angles). Then, $\dot{q} = u$ are generalized velocities, which we assume we can control. Let $T(q, u) = u^{T}M(q)u$, $M \succ 0$, be the kinetic energy and V(q) be the potential energy.

For conservative systems, the following quantity is minimized:

$$\int_{0}^{T} \underbrace{\left[T(q, u) - V(q)\right]}_{L(q, u) = \text{Lagrange's "action function}} dt$$

The Hamiltonian is

$$H(q, u, \lambda) = L(q, u) + \lambda^{\mathrm{T}} f(q, u) = L(q, u) + \lambda^{\mathrm{T}} u$$

In mechanics, λ is called a generalized momentum, satisfying

$$\dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial q} = -\frac{\partial L^{\mathrm{T}}}{\partial q} + 0$$

$$0 = \frac{\partial H}{\partial u} = \frac{\partial L}{\partial u} + \lambda^{\mathrm{T}} \Longrightarrow \lambda = -\frac{\partial L^{\mathrm{T}}}{\partial u}$$

$$\dot{\lambda} = -\frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L^{\mathrm{T}}}{\partial u} = -\frac{\partial L^{\mathrm{T}}}{\partial q}$$

This produces the Euler-Lagrange Equation:

$$\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} = 0$$

Recall, along optimal trajectories

$$\frac{\mathrm{d}H}{\mathrm{d}t} = \underbrace{\frac{\partial H}{\partial t}}_{=0 \text{ if } L \text{ and } f \text{ do not depend explicitly on } t}_{=0 \text{ depend explicitly on } t} + \underbrace{\frac{\partial H}{\partial u}}_{=0} \dot{u} \underbrace{\frac{\partial H}{\partial \lambda}}_{f^{\mathrm{T}} = \dot{x}^{\mathrm{T}}} \dot{\lambda} = -\dot{\lambda}^{\mathrm{T}} \dot{x} + \dot{x}^{\mathrm{T}} \dot{\lambda} = 0$$

Therefore, along optimal trajectories, the Hamiltonian is constant! We had

$$\begin{split} H &= L + \lambda^{\mathrm{T}} u \\ \frac{\partial H}{\partial u} &= \lambda^{\mathrm{T}} + \frac{\partial L}{\partial u} = 0 \end{split}$$

Along optimal trajectories,

$$H = L - \frac{\partial L}{\partial u}u$$

Recall, L(q, u) = T(q, u) - V(q).

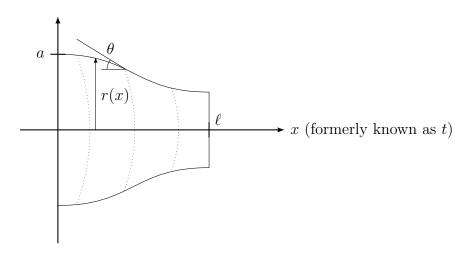
$$\frac{\partial L}{\partial u} = \frac{\partial T}{\partial u} - 0$$
$$T(q, u) = u^{\mathrm{T}} M(q) u$$
$$\frac{\partial T}{\partial u} = 2u^{\mathrm{T}} M$$

$$H = T - V - 2u^{T}Mu = -(V + u^{T}Mu) = -(V + T)$$

$$u^{T}Mu$$

Therefore, the total energy (kinetic plus potential energy) remains constant for conservative systems.

Example minimum drag nose shape (Newton 1686)



The drag is

$$D = -2\pi q \int_{r=0}^{\ell} C_p(\theta) r \, \mathrm{d}r,$$

where q is a pressure constant and $C_p(\theta) = 2\sin^2\theta$ is Newton's pressure formula.

Geometry tells us

$$\frac{\mathrm{d}r}{\mathrm{d}x} = -\tan\theta = -u$$

Choose the control as $\tan \theta$. Manipulating the drag,

$$\frac{D}{4\pi q} = \int_0^\ell \frac{ru^3}{1+u^2} \, \mathrm{d}x + \frac{1}{2}r(\ell)^2$$

The optimal control problem is

$$\min_{u} \int_{0}^{\ell} \frac{ru^3}{1+u^2} dx + \frac{1}{2}r(\ell)^2$$

s.t.
$$\frac{dr}{dx} = -u$$

This is in the standard form with the following changes of variables:

$$\ell \longleftarrow T$$

$$x \longleftarrow t$$

$$r \longleftarrow x$$

Refer to (3.1) and (3.2) for the following steps.

$$H = \frac{ru^3}{1+u^2} - \lambda u$$

$$\frac{\partial H}{\partial u} = \frac{3ru^2(1+u^2) - ru^3 \cdot 2u}{(1+u^2)^2} - \lambda$$

$$= \frac{ru^4 + 3ru^2}{(1+u^2)^2} - \lambda = 0$$

$$\lambda = \frac{ru^2(u^2+3)}{(1+u^2)^2}$$

$$\frac{d\lambda}{dx} = -\frac{\partial H}{\partial r} = -\frac{u^3}{1+u^2}$$

$$\lambda(\ell) = r(\ell)$$
(3.3)

Right now, we know

$$\begin{cases} \frac{\mathrm{d}r}{\mathrm{d}x} = -u\\ r(0) = a\\ \frac{\mathrm{d}\lambda}{\mathrm{d}x} = -\frac{u^3}{1+u^2}\\ \lambda(\ell) = r(\ell) \end{cases}$$

We need to remove u and get a function of r and λ instead. However, it is difficult to solve (3.3). Maybe H = const. gives us something nicer?

$$H = \frac{ru^3}{1+u^2} - \lambda u$$

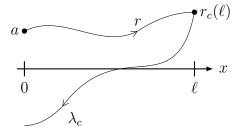
$$= \frac{ru^3}{1+u^2} - \frac{ru^2(u^2+3)}{(1+u^2)^2} u$$

$$= -\frac{2ru^3}{(1+u^2)^2} = c$$

Assume we can find u = G(r, c), either numerically or some other way. So, now we have

$$\begin{cases} \frac{\mathrm{d}r}{\mathrm{d}x} = -G(r,c) \\ r(0) = a \end{cases}$$
$$\begin{cases} \frac{\mathrm{d}\lambda}{\mathrm{d}x} = -\frac{G^3(r,c)}{1 + G^2(r,c)} \\ \lambda(\ell) = r(\ell) \end{cases}$$

We do not know c, but we can guess c and simulate r forward in "time" (x) from r(0) = a. Then, we simulate λ backwards from $r(\ell)$.



Problem: we can do this for any c. Which c is it? Last 15 minutes was a dead end! Back to $u = F(r, \lambda)$. Assume we have F (numerically).

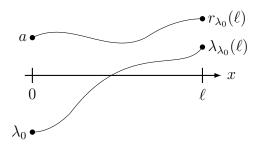
$$\frac{\mathrm{d}r}{\mathrm{d}x} = -F(r,\lambda)$$

$$r(0) = a$$

$$\frac{\mathrm{d}\lambda}{\mathrm{d}x} = -\frac{F^3(r,\lambda)}{1 + F^2(r,\lambda)}$$

$$\lambda(\ell) = r(\ell)$$

The mistake before was that the simulation forward from a depends on λ .



Therefore, we should guess λ_0 and simulate both r and λ to get $r_{\lambda_0}(\ell)$ and $\lambda_{\lambda_0}(\ell)$. We need

$$r_{\lambda_0}(\ell) = \lambda_{\lambda_0}$$

for optimality. To do this, we need numerics.

Terminal Constraints

Let $x = [x_1, \dots, x_n]^T \in \mathbb{R}^n$ and solve

$$\min_{u \in \mathcal{U}} \int_0^T L(x, u, t) dt + \Psi(x(T))$$
s.t.
$$\dot{x} = f(x, u, t)$$

$$x(0) = x_0$$

$$x_i(T) = x_{iT} \quad \text{given for } i \in \mathcal{T} \subset \{1, \dots, n\}$$

First, we augment the cost:

$$\begin{split} \tilde{J}(u) &= \int_0^T \left[L + \lambda^{\mathrm{T}} (f - \dot{x}) \right] \mathrm{d}t + \Psi \\ &= \int_0^T (H - \lambda^{\mathrm{T}} \dot{x}) \, \mathrm{d}t + \Psi \\ \tilde{J}(u + \varepsilon v) - \tilde{J}(u) &= \int_0^T \left(\varepsilon \frac{\partial H}{\partial u} v + \varepsilon \frac{\partial H}{\partial x} \eta - \varepsilon \lambda^{\mathrm{T}} \dot{\eta} \right) \mathrm{d}t + \varepsilon \frac{\partial \Psi}{\partial x} (x(T)) \eta(T) + o(\varepsilon) \\ \delta \tilde{J}(u; v) &= \int_0^T \left(\frac{\partial H}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right) \eta \, \mathrm{d}t + \int_0^T \frac{\partial H}{\partial u} v \, \mathrm{d}t \\ &+ \lambda^{\mathrm{T}} (0) \eta(0) - \lambda^{\mathrm{T}} (T) \eta(T) + \frac{\partial \Psi}{\partial x} (x(T)) \eta(T) \end{split}$$

As always,

$$\dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x}$$

$$\frac{\partial H}{\partial u} = 0 \quad (\text{FONC})$$

Additionally,

$$\eta(0) = 0$$
 $\eta_i(T) = 0 \quad \text{for } i \in \mathcal{T}$

Note that if $x(T) = x_T$ is given, then $x(T) = x(T) + \varepsilon \eta(T) + o(\varepsilon)$, so $\eta(T) = 0$. Here, we have $x_i(T) = x_{iT}$ fixed for $i \in \mathcal{T}$ so $\eta_i(T) = 0$ for $i \in \mathcal{T}$.

For optimality, we want

$$\left[-\lambda^{\mathrm{T}}(T) + \frac{\partial \Psi}{\partial x}(x(T)) \right] \eta(T) = 0 \quad \text{for all } admissible \text{ variations}$$

$$\left[\frac{\partial \Psi}{\partial x_1} - \lambda_1, \cdots, \frac{\partial \Psi}{\partial x_n} - \lambda_n \right] \begin{bmatrix} \eta_1(T) \\ \vdots \\ \eta_n(T) \end{bmatrix} = 0$$

Hence, we need

$$\lambda_j(T) = \frac{\partial \Psi}{\partial x_j}(x(T))$$
 if $j \notin \mathcal{T}$
 $\lambda_i(T) = \text{free}$ if $i \in \mathcal{T}$

So we have

$$\begin{cases} \dot{x} = f \\ \dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x}, \end{cases}$$

an ODE with 2n variables. We need 2n boundary conditions for this ODE to be well-posed.

So we have n + q + (n - q) = 2n boundary conditions.

We could even fix some but not all of x(0), i.e.

$$x_i(0) = x_{i0}$$
 if $i \in \mathcal{I}$
 $x_j(0) = \text{free}$ if $j \notin \mathcal{I}$

Recall,

$$\delta \tilde{J}(u;v) = \int_0^T \left(\frac{\partial H}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right) \eta \, \mathrm{d}t + \int_0^T \frac{\partial H}{\partial u} v \, \mathrm{d}t + \lambda^{\mathrm{T}}(0) \eta(0) + \left[\lambda^{\mathrm{T}}(T) - \frac{\partial \Psi}{\partial x}(x(T)) \right] \eta(T)$$

For $x_i(0) = x_{i0}$ fixed, we have $\eta_i(0) = 0$ and $\lambda_i(0)$ free. For $x_j(0)$ free, we have $\eta_j(0)$ free and $\lambda_j(0) = 0$.

To ponder, what if $J = \int L dt + \Psi(x(T)) + \Theta(x(0))$?

To summarize, the minimizer to

$$\min_{u \in \mathcal{U}} \int_0^T L(x, u, t) dt + \Psi(x(T))$$
s.t.
$$\dot{x} = f(x, u, t)$$

$$x_i(0) = x_{i0}, \quad i \in \mathcal{I}$$

$$x_j(T) = x_{jT} \quad j \in \mathcal{T}$$

has to satisfy

$$\begin{split} \frac{\partial H}{\partial u} &= 0\\ \dot{\lambda} &= -\frac{\partial H^{\mathrm{T}}}{\partial x}\\ \lambda_i(0) &= 0, \quad i \not\in \mathcal{I}\\ \lambda_j(T) &= \frac{\partial \Psi}{\partial x_j}(x(T)), \quad j \not\in \mathcal{T} \end{split}$$

Example

$$\frac{d}{dt} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix} = f(x_1, x_2, x_3, x_4)$$

$$x_1(0) = 1, x_3(0) = 7, x_4(0) = 0, x_1(1) = 2$$

$$\mathcal{I} = \{1, 3, 4\}, \mathcal{T} = \{1\}$$

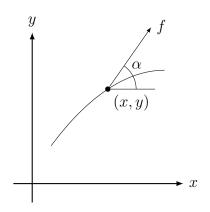
$$\min \int_0^1 L(x, u) dt + (x_2^2(1) - x_3^2(1) + 7x_1(1) + 14)$$

Note there are 4 boundary conditions on x so there must be 4 boundary conditions on λ :

$$\begin{array}{lll} \lambda_1(0) \ \text{free/unspecified} & \lambda_1(1) \ \text{free} \\ \lambda_2(0) = 0 & \lambda_2(1) = 2x_2(1) \\ \lambda_3(0) \ \text{free} & \lambda_3(1) = -2x_3(1) \\ \lambda_4(0) \ \text{free} & \lambda_4(1) = 0 \end{array}$$

Example

A force f acts on a particle at position (x, y) (mass = 1).



$$\begin{split} \dot{x} &= v_x \\ \dot{y} &= v_y \\ \dot{v}_x &= |f| \cos \alpha \\ \dot{v}_y &= |f| \sin \alpha \\ \alpha &= \text{control variable} \end{split}$$

Assume we only care about where the particle ends up (to be specified later), i.e. L=0.

$$H = \begin{bmatrix} \lambda_x & \lambda_y & \lambda_{v_x} & \lambda_{v_y} \end{bmatrix} \begin{bmatrix} v_x \\ v_y \\ |f| \cos \alpha \\ |f| \sin \alpha \end{bmatrix}$$

$$\dot{\lambda}_x = -\frac{\partial H}{\partial x} = 0 \qquad \Longrightarrow \qquad \lambda_x(t) = c_1$$

$$\dot{\lambda}_y = -\frac{\partial H}{\partial y} = 0 \qquad \Longrightarrow \qquad \lambda_y(t) = c_2$$

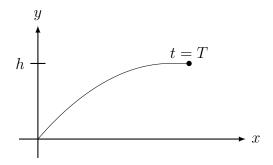
$$\dot{\lambda}_{v_x} = -\frac{\partial H}{\partial v_x} = -\lambda_x \qquad \Longrightarrow \qquad \lambda_{v_x}(t) = -c_1 t + c_3$$

$$\dot{\lambda}_{v_y} = -\frac{\partial H}{\partial v_y} = -\lambda_y \qquad \Longrightarrow \qquad \lambda_{v_y}(t) = -c_2 t + c_4$$

Moreover,

$$\frac{\partial H}{\partial \alpha} = -\lambda_{v_x} |f| \sin \alpha + \lambda_{v_y} |f| \cos \alpha = 0$$
$$\tan \alpha = \frac{\lambda_{v_y}}{\lambda_{v_x}} = \frac{-c_2 t + c_4}{-c_1 t + c_3}$$

We want to drive the particle from $[0,0,0,0]^T$ to a path parallel to the x-axis with y(T) = h.



Choose $\Psi = -v_x$,

$$y(T) = h$$
 $v_y(T) = 0$
 $x(T)$ free $v_x(T)$ free, but costs
 $\lambda_i(0)$ free
 $\lambda_y(T)$ free $\lambda_{v_y}(T)$ free
 $\lambda_x(T) = 0$ $\lambda_{v_x}(T) = -1$

$$c_1 = \lambda_x(t) = 0$$

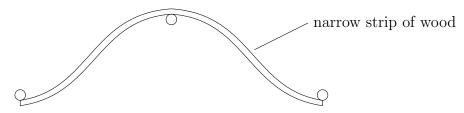
$$\Longrightarrow \lambda_{v_x} = -c_1 t + c_3 = c_3 = -1$$

$$\Longrightarrow \tan \alpha = -\frac{-c_2 t + c_4}{-1} = c_2 t + c_4$$

How do we find c_2 and c_4 ? Plug into \dot{x} and $\dot{\lambda}$ and try to satisfy the remaining boundary conditions. (This is hard=numerics.)

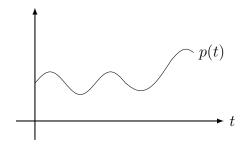
3.2 Splines

From ship building. Splines are used a lot in path-planning, e.g. cubic splines.



But, they are solutions to optimal control problems.

Let p(t) be a curve we'd like to shape.



We want to minimize the "energy" put into the curve, a.k.a acceleration. Let $x_1 = p$ and $x_2 = \dot{p}$, so

$$\begin{cases} \dot{x}_1 = x_2 \\ \dot{x}_2 = u \end{cases}$$

3.2.1 Minimum-Energy

$$\min_{u \in \mathcal{U}} \frac{1}{2} \int_{0}^{T} u^{2}(t) dt + \text{Boundary conditions on } x$$

$$H = L + \lambda^{T} f = \frac{1}{2} u^{2} + \lambda_{1} x_{2} + \lambda_{2} u$$

$$\frac{\partial H}{\partial u} = u + \lambda_{2} = 0 \Longrightarrow u = -\lambda_{2}$$

$$\dot{\lambda}_{1} = -\frac{\partial H}{\partial x_{1}} = 0 \Longrightarrow \lambda_{1} = c_{1}$$

$$\dot{\lambda}_{2} = -\frac{\partial H}{\partial x_{2}} = -\lambda_{1} \Longrightarrow \lambda_{2} = -c_{1} t + c_{2}$$

$$u = -\lambda_{2} = c_{1} t - c_{2}$$

$$\dot{x}_{2} = u = c_{1} t - c_{2} \Longrightarrow x_{2} = c_{1} \frac{t^{2}}{2} - c_{2} t + c_{3}$$

$$\dot{x}_{1} = x_{2} = c_{1} \frac{t^{2}}{2} - c_{2} t + c_{3}$$

$$\Longrightarrow x_{1} = \frac{c_{1}}{6} t^{3} - \frac{c_{2}}{2} t^{2} + c_{3} t + c_{4}$$

p(t) is a cubic polynomial!

What about boundary conditions?

Let T = 1, p(0) given, p(1) given, $\dot{p}(0) = 0$, $\dot{p}(1) = 0$, e.g. p(0) = 0, p(1) = 1. Since the boundary conditions for x are all specified, those for the costate are free.

$$\begin{array}{l}
x_1(0) = 0 \\
x_2(0) = 0 \\
x_1(1) = 1 \\
x_2(1) = 0
\end{array}
\Longrightarrow
\begin{cases}
\lambda_1(0) \\
\lambda_2(0) \\
\lambda_1(1)
\end{cases}$$
 free/unspecified $\lambda_2(1)$

$$x_2(0) = c_3 = 0 x_1(1) = \frac{2c_2}{6} - \frac{c_2}{2} = 1$$

$$x_1(0) = c_4 = 0 c_2 = -6$$

$$x_2(1) = \frac{c_1}{2} - c_2 + \underbrace{c_3}_{0} = 0 c_1 = -12$$

$$c_1 = 2c_2$$

$$\implies p(t) = -2t^3 + 3t^2$$
$$u(t) = -12t + 6$$

Or, what if $\dot{p}(0)$, $\dot{p}(1)$ are not specified?

$$x_1(0) = 0$$

$$x_2(0) \text{ unspec.}$$

$$x_1(1) = 1$$

$$x_2(1) \text{ unspec.}$$

$$\Rightarrow \begin{cases} \lambda_1(0) \text{ unspec.} \\ \lambda_2(0) = 0 \\ \lambda_1(1) \text{ unspec.} \\ \lambda_2(1) = 0 \end{cases}$$

$$\lambda_2(0) = c_2 = 0
\lambda_2(1) = -c_1 + c_2 = 0$$

$$\begin{vmatrix}
x_1(0) = c_4 = 0 \\
x_1(1) = c_3 = 1
\end{vmatrix} \implies p(t) = t$$

What did we do?

Case 1:

$$\begin{bmatrix} 0 & 0 & 0 & 1 \\ 1/6 & -1/2 & 1 & 1 \\ 0 & 0 & 1 & 0 \\ 1/2 & -1 & 1 & 0 \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} x_1(0) \\ x_1(1) \\ x_2(0) \\ x_2(1) \end{bmatrix}$$

Case 2:

$$\begin{bmatrix} 0 & 0 & 0 & 1 \\ 1/6 & -1/2 & 1 & 1 \\ 0 & 1 & 0 & 0 \\ -1 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} x_1(0) \\ x_1(1) \\ \lambda_2(0) \\ \lambda_2(1) \end{bmatrix}$$

3.2.2 Generalized Splines

We had $\dot{x} = Ax + Bu$ with

$$A = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}.$$

This A is nilpotent $(A^k = 0 \text{ for some } k \in \mathbb{Z}^+)$. This means e^{At} is a polynomial in t. (This e^{At} is cubic.)

In general, e^{At} is a mix of polynomials, exponentials, and trignometric terms. The eigenvalues of A determine the form of x(t).

$$\dot{x} = Ax$$
 $\Longrightarrow x(t) = e^{At}x(0)$
 $\dot{x} = Ax + Bu \Longrightarrow x(t) = e^{At}x(0) + \int_0^t e^{A(t-\tau)}Bu(\tau) d\tau$

The general problem to solve is

$$\min_{u \in \mathcal{U}} \int_0^T \frac{1}{2} ||u||^2 dt$$

s.t. $\dot{x} = Ax + Bu$
+ Boundary conditions

$$H = \frac{1}{2} ||u||^2 + \lambda^{\mathrm{T}} (Ax + Bu)$$
$$\frac{\partial H}{\partial u} = u^{\mathrm{T}} + \lambda^{\mathrm{T}} B = 0$$
$$\Rightarrow u = -B^{\mathrm{T}} \lambda$$
$$\dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x} = -A^{\mathrm{T}} \lambda$$

We have the Hamiltonian Dynamics:

$$\begin{bmatrix} \dot{x} \\ \dot{\lambda} \end{bmatrix} = \underbrace{\begin{bmatrix} A & -BB^{\mathrm{T}} \\ 0 & -A^{\mathrm{T}} \end{bmatrix}}_{M} \begin{bmatrix} x \\ \lambda \end{bmatrix}$$

Where we used $\dot{x} = Ax + Bu = Ax - BB^{T}\lambda$. Then,

$$\begin{bmatrix} x(t) \\ \lambda(t) \end{bmatrix} = e^{Mt} \begin{bmatrix} x(0) \\ \lambda(0) \end{bmatrix}$$

Suppose we want to drive from $x(0) = x_0$ to $x(T) = x_T$.

$$\begin{bmatrix} x_T \\ \lambda(T) \end{bmatrix} = e^{MT} \begin{bmatrix} x_0 \\ \lambda(0) \end{bmatrix} = \begin{bmatrix} N_{xx} & N_{x\lambda} \\ N_{\lambda x} & N_{\lambda\lambda} \end{bmatrix} \begin{bmatrix} x_0 \\ \lambda(0) \end{bmatrix}$$
$$x_T = N_{xx}x_0 + N_{x\lambda}\lambda(0)$$

 $N_{x\lambda}$ is invertible if (A, B) is completely controllable. Assume it is.

$$\lambda(0) = N_{x\lambda}^{-1}(x_T - N_{xx}x_0)$$

$$\Longrightarrow \begin{bmatrix} x(t) \\ \lambda(t) \end{bmatrix} = e^{Mt} \begin{bmatrix} x_0 \\ N_{x\lambda}^{-1}(x_T - N_{xx}x_0) \end{bmatrix}$$

$$\Longrightarrow u(t) = -B^{T}\lambda(t)$$

This is the optimal trajectory, but there is no feedback. We will consider closed-loop systems after the midterm.

As a preview, we need to find λ as a function of x. For example, $u = -R^{-1}B^{T}Px$ minimizes $u^{T}Ru$, so $\lambda = Px$ where P is the solution to the Riccati equation.

3.3 Numerical Methods

Optimal control boils down to solving two sets of differential equations:

$$\dot{x} = f(x, u) \qquad \frac{\partial H}{\partial u}(x, u, \lambda) = 0$$

$$\dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x}(x, u, \lambda) \qquad u = F(x, \lambda)$$

$$\Longrightarrow \begin{cases} \dot{x} = f(x, F(x, \lambda)) \\ \dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x}(x, F(x, \lambda), \lambda) \end{cases}$$

The equations are functions of x and λ . They are completely determined by the boundary conditions on x(0), x(T), $\lambda(0)$, $\lambda(T)$. This is known as the *Boundary Value Problem*. This is solved using *test shooting*:

- 1. Guess initial conditions
- 2. Simulate forward in time
- 3. Update the guess (cleverly...)

Exmaple: Bolza problem

$$\min_{u \in \mathcal{U}} \int_0^T L(x, u) \, \mathrm{d}t + \Psi(x(T))$$
s.t.
$$\begin{cases} \dot{x} = f(x, u) \\ x(0) = x_0 \end{cases}$$

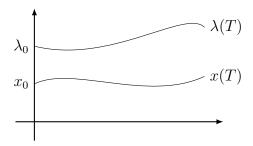
$$H(x, u, \lambda) = L(x, u) + \lambda^{\mathrm{T}} f(x, u)$$

$$u^*(x, \lambda) \text{ satisfies } \frac{\partial H}{\partial u} = 0$$

The optimal control satisfies

$$\begin{cases} x = f(x, u^*(x, \lambda)) \\ x(0) = x_0 \\ \lambda = -\frac{\partial H^{\mathrm{T}}}{\partial x}(x, u^*(x, \lambda), \lambda) \\ \lambda(T) = \frac{\partial \Psi}{\partial x}(x(T)) \end{cases}$$

Algorithm Guess λ_0 and solve for x(t), $\lambda(t)$.



Let's define a cost:

$$\left\| \lambda(T) - \frac{\partial \Psi^{\mathrm{T}}}{\partial x}(x(T)) \right\|^2 = g(\lambda_0)$$

Update λ_0 through

$$\lambda_0 \coloneqq \lambda_0 - \gamma \frac{\partial g^{\mathrm{T}}}{\partial \lambda_0} (\lambda_0)$$

any choice of step size works

Repeat

Problem: What is $\partial g/\partial \lambda_0$? We estimate $\partial g/\partial \lambda_0$ numerically. This is where "test shooting" comes into play.

Let e_i be the *i*th unit vector, i = 1, ..., n:

$$e_{1} = \begin{bmatrix} 1\\0\\\vdots\\0 \end{bmatrix}, e_{2} = \begin{bmatrix} 0\\1\\\vdots\\0 \end{bmatrix}, \dots, e_{n} = \begin{bmatrix} 0\\0\\\vdots\\1 \end{bmatrix}$$
$$\frac{\partial g}{\partial \lambda_{0}} = \left(\frac{\partial g}{\partial \lambda_{0,1}}, \frac{\partial g}{\partial \lambda_{0,2}}, \dots, \frac{\partial g}{\partial \lambda_{o,n}}\right)$$

The *i*th component of $\partial g/\partial \lambda_0$ is given by the directional derivative

$$\frac{\partial g}{\partial \lambda_{0,i}} = \frac{\partial g}{\partial \lambda_0} \cdot e_i = \delta g(\lambda_0; e_i) = \lim_{\varepsilon \to 0} \frac{g(\lambda_0 + \varepsilon e_i) - g(\lambda_0)}{\varepsilon}$$

So, if $x \in \mathbb{R}^n$ (and thus so is λ_0), we have to do this n times (with a small ε) and get the full derivative $\partial g/\partial \lambda_0$.

Given
$$\lambda_0$$
, $g(\lambda_0)$
for $i=1$ to n do
Compute $g(\lambda_0 + \varepsilon e_i)$
 $dg_i = \frac{1}{\varepsilon} [g(\lambda_0 + \varepsilon e_i) - g(\lambda_0)]$
end for
 $\frac{\partial g}{\partial \lambda_0} = [dg_1, \dots, dg_n]$

Algorithm

Example LQ

$$\min_{u} \frac{1}{2} \int_{0}^{1} (x^{\mathrm{T}}Qx + u^{\mathrm{T}}Ru) \, \mathrm{d}t + \frac{1}{2}x^{\mathrm{T}}(1)Sx(1)$$
s.t.
$$\begin{cases} \dot{x} = Ax + Bu \\ x(0) = x_{0} \end{cases}$$

$$Q, R, S \succ 0$$

$$H = \frac{1}{2}x^{\mathrm{T}}Qx + \frac{1}{2}u^{\mathrm{T}}Ru + \lambda^{\mathrm{T}}(Ax + Bu)$$

$$\frac{\partial H}{\partial u} = u^{\mathrm{T}}R + \lambda^{\mathrm{T}}B = 0$$

$$u^* = -R^{-1}B^{\mathrm{T}}\lambda$$

$$\dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x} = -Qx - A^{\mathrm{T}}\lambda$$

$$\lambda(1) = \frac{\partial \Psi^{\mathrm{T}}}{\partial x}(x(1)) = Sx(1)$$

So putting it all together,

$$\dot{x} = Ax - BR^{-1}B^{T}\lambda$$
 $x(0) = x_0$
 $\dot{\lambda} = -Qx - A^{T}\lambda$ $\lambda(1) = Sx(1)$

Example Newton's nose shape problem (revisited, see previous)

$$\min_{u} \int_{0}^{\ell} \frac{ru^3}{1+u^2} dx + \frac{1}{2}r(\ell)^2$$

s.t.
$$\frac{dr}{dx} = -u \qquad r(0) = a$$

$$H = \frac{ru^{3}}{1 + u^{2}} + \lambda(-u)$$
$$\frac{\partial H}{\partial u} = \frac{ru^{2}(3 + u^{2})}{(1 + u^{2})^{2}} - \lambda = 0$$

We solve the above numerically to get $u^*(r, \lambda)$.

$$\frac{\partial \lambda}{\partial x} = -\frac{\partial H}{\partial r} = -\frac{u^3}{1 + u^2}$$
$$\lambda(\ell) = r(\ell)$$

So, we have

$$\frac{\mathrm{d}r}{\mathrm{d}x} = -u \qquad r(0) = a \qquad u = F(x, \lambda)$$

$$\frac{\mathrm{d}\lambda}{\mathrm{d}x} = -\frac{u^3}{1+u^2} \qquad \lambda(\ell) = r(\ell)$$

Example Fixed terminal constraints (revisited, see previous)

$$\begin{aligned} \min_{\alpha} -v_x(T) & \alpha = \text{control} \\ \text{s.t.} & \dot{x} = v_x & x(0) = 0 \\ & \dot{y} = v_y & y(0) = 0 \\ & \dot{v}_x = |f| \cos \alpha & v_x(0) = 0 \\ & \dot{v}_y = |f| \sin \alpha & v_y(0) = 0 \\ & y(T) = h \\ & v_y(T) = 0 \end{aligned}$$

$$H = -v_x(T) + \begin{bmatrix} \lambda_x & \lambda_y & \lambda_{v_x} & \lambda_{v_y} \end{bmatrix} \begin{bmatrix} v_x \\ v_y \\ |f| \cos \alpha \\ |f| \sin \alpha \end{bmatrix}$$

$$\frac{\mathrm{d}H}{\mathrm{d}\alpha} = 0 \Rightarrow \tan \alpha = \frac{\lambda_{v_y}}{\lambda_{v_x}}$$

$$\dot{\lambda}_x = 0$$

$$\dot{\lambda}_y = 0$$

$$\dot{\lambda}_{v_x} = -\lambda_x$$

$$\dot{\lambda}_{v_y} = -\lambda_y$$

$$\lambda(0) \text{ unspecified}$$

$$\lambda_x(T) = \frac{\partial \Psi^T}{\partial x}(x(T)) = 0$$

$$\lambda_y(T) \text{ unspecified}$$

$$\lambda_{v_x}(T) = \frac{\partial \Psi^T}{\partial v_x}(v_x(T)) = -1$$

$$\lambda_{v_y}(T) \text{ unspecified}$$

Again, we guess λ_0 and solve forward in time. But, we have terminal constraints on y and v_y as well.

$$g(\lambda_0) = \frac{1}{2} \left[(y(T) - h)^2 + (v_y(T))^2 + (\lambda_x(T))^2 + (\lambda_{v_x} + 1)^2 \right]$$

3.4 Terminal Manifolds

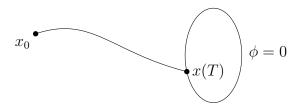
We can solve

$$\begin{split} \min_{u \in \mathcal{U}} \int_0^T L(x, u, t) \, \mathrm{d}t + \Psi(x(T)) \\ \text{s.t. } \dot{x} &= f(x, u, t) \end{split}$$

with all sorts of boundary conditions on x:

- $x(0) = x_0, x(T)$ free (typical)
- $x_i(0) = x_{i0}, i \in \mathcal{I} \text{ and } x_j(T) = x_{jT}, j \in \mathcal{T}$

But what if we want x(T) to belong to a set?



Problem

$$\begin{split} \min_{u \in \mathcal{U}} \int_0^T L(x, u, t) \, \mathrm{d}t + \Psi(x(T)) \\ \text{s.t.} \ \ \dot{x} &= f(x, u, t), \quad x \in \mathbb{R}^n \\ x(0) &= x_0 \\ \phi(x(T)) &= 0, \quad \phi : \mathbb{R}^n \to \mathbb{R}^q, \ q \leq n \end{split}$$

The augmented cost is

$$\tilde{J} = \int_0^T [H(x, u, t, \lambda) - \lambda^T \dot{x}] dt + \Psi(x(T)) + \underset{\text{q-dimensional Lagrange multiplier}}{\psi^T} \phi(x(T))$$

Let $\Phi(x(T), \nu) = \Psi(x(T)) + \nu^{\mathrm{T}} \phi(x(T))$. Then,

$$\tilde{J} = \int_0^T \left(\frac{\partial H}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right) \eta \, \mathrm{d}t + \Phi(x(T), \nu)$$

We know how to solve this! With $u \mapsto u + \varepsilon v$, $x \mapsto x + \varepsilon \eta + o(\varepsilon)$,

$$\delta \tilde{J} = \int_{0}^{T} \left(\frac{\partial H}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right) \eta \, \mathrm{d}t + \int_{0}^{T} \frac{\partial H}{\partial u} v \, \mathrm{d}t + \frac{\partial \Phi}{\partial x} (x(T), \nu) \eta(T)$$

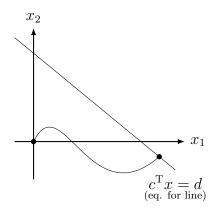
$$- \lambda^{\mathrm{T}}(T) \eta(T) + \lambda^{\mathrm{T}}(0) \underbrace{\eta(0)}_{=0}$$

$$\begin{cases} \frac{\partial H}{\partial u} = 0 \\ \dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x} \\ \lambda(T) = \frac{\partial \Phi^{\mathrm{T}}}{\partial x} (x(T), \dot{\nu}) \\ \phi(x(T)) = 0 & \leftarrow q \text{ new equations} \end{cases}$$

$$\Longrightarrow u^{*}$$

Spling to line

$$\min_{u} \frac{1}{2} \int_{0}^{1} u^{2}(t) dt$$
s.t.
$$\begin{cases}
\dot{x}_{1} = x_{2} \\
\dot{x}_{2} = u \\
x_{1}(0) = 0, \ x_{2}(0) = 0 \\
c_{1}x_{1}(1) + c_{2}x_{2}(1) = d
\end{cases}$$



$$H = \frac{1}{2}u^2 + \lambda_1 x_2 + \lambda_2 u$$

$$\frac{\partial H}{\partial u} = u + \lambda_2 \Longrightarrow u = -\lambda_2$$

$$\dot{\lambda}_1 = -\frac{\partial H}{\partial x_1} = 0 \Longrightarrow \lambda_1 = k_1$$

$$\dot{\lambda}_2 = -\frac{\partial H}{\partial x_2} = -\lambda_1 \Longrightarrow \lambda_2 = -k_1 t + k_2$$

$$\phi(x(1)) = c_1 x_1(1) + c_2 x_2(1) - d$$

$$\Psi = 0 \Longrightarrow \Phi = \nu(c_1 x_1(1) + c_2 x_2(1) - d)$$

$$\lambda_1(1) = \frac{\partial \Phi}{\partial x_1} = \nu c_1$$

$$\lambda_2(1) = \frac{\partial \Phi}{\partial x_2} = \nu c_2$$

So,

$$\lambda_{1}(1) = \nu c_{1} = k_{1}$$

$$\lambda_{2}(1) = \nu c_{2} = -k_{1} + k_{2}$$

$$k_{2} = \nu (c_{1} + c_{2})$$

$$\dot{x}_{2} = u = -\lambda_{2} = k_{1}t - k_{2}$$

$$x_{2} = \frac{k_{1}}{2}t^{2} - k_{2}t + 0$$

$$\dot{x}_{1} = x_{2}$$

$$x_{1} = \frac{k_{1}}{6}t^{3} - \frac{k_{2}}{2}t^{2} + 0$$

Substituting k_1 and k_2 into $c_1x_1(1) + c_2x_2(1) = d$,

$$\nu \left(-\frac{c_1^2}{3} - c_1 c_2 - c_2^2 \right) = d$$

$$\nu = -\frac{d}{c_1^2/3 + c_1 c_2 + c_2^2}$$

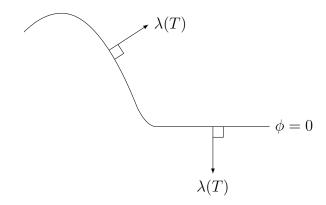
And finally

$$u = k_1 t - k_2 = \frac{d}{c_1^2/3 + c_1 c_2 + c_2^2} (c_1 + c_2 - c_1 t)$$

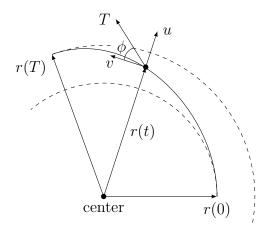
As a final observation if $\Psi = 0$ then

$$\lambda(T) = \nu^{\mathrm{T}} \frac{\partial \phi}{\partial x}(x(T)),$$

which means $\lambda(T)$ is orthogonal to the tangent plane to $\phi(x(T))$.



Example Maximum orbit transform (e.g. Hidden Figures)



r = radial distance from spacecraft to

center

u = radial velocity

v =tangential velocity

m =mass of spacecraft

 $\dot{m} = -$ fuel consumption rate

 $\phi = \text{thrust angle (control input)}$

T = thrust

$$\begin{split} \max_{\phi} r(T) &\iff \min_{\phi} - r(T) \\ \text{s.t.} & \begin{cases} \dot{r} = u \\ \dot{u} = \frac{v^2}{r} - \frac{g}{r^2} + \frac{T \sin \phi}{m_0 - |\dot{m}|t} \\ \dot{v} = -\frac{uv}{r} + \frac{T \cos \phi}{m_0 - |\dot{m}|t} \\ r(0) = r_0 \\ u(0) = 0 \\ v(0) = \sqrt{\frac{g}{r_0}} \\ u(T) = 0 = \phi_1 \\ v(T) = \sqrt{\frac{g}{r(T)}} = \phi_2 \end{split}$$

$$H = \lambda_r u + \lambda_u \left(\frac{v^2}{r} - \frac{g}{r^2} + \frac{T \sin \phi}{m_0 - |\dot{m}|t} \right) + \lambda_v \left(-\frac{uv}{r} + \frac{T \cos \phi}{m_0 - |\dot{m}|t} \right)$$

$$\Phi = \underbrace{\nu_1 u(T) + \nu_2 \left(v(T) - \sqrt{\frac{g}{r(T)}} \right)}_{\nu^T \phi} \underbrace{-r(T)}_{\Psi}$$

$$\frac{\partial H}{\partial \phi} = \frac{\lambda_u T \cos \phi - \lambda_v T \sin \phi}{m_0 - |\dot{m}|t} = 0$$

$$\Rightarrow \tan \phi = \frac{\lambda_u}{\lambda_v}$$

$$\dot{\lambda}_r = -\frac{\partial H}{\partial r} = -\lambda_u \left(-\frac{v^2}{r^2} + \frac{2g}{r^3} \right) - \lambda_v \cdot \frac{uv}{r^2}$$

$$\dot{\lambda}_u = -\frac{\partial H}{\partial u} = -\lambda_r + \lambda_v \cdot \frac{v}{r}$$

$$\dot{\lambda}_v = -\frac{\partial H}{\partial v} = -\lambda_u \cdot \frac{2v}{r} + \lambda_v \cdot \frac{u}{r}$$

$$\left\{ \lambda_r(T) = \frac{\partial \Phi}{\partial r} = -1 + \frac{\nu_2 \sqrt{g}}{2(r(T))^{3/2}} \right.$$

$$\lambda_u(T) = \frac{\partial \Phi}{\partial u} = \nu_1$$

$$\lambda_v(T) = \frac{\partial \Phi}{\partial v} = \nu_2$$

$$u(T) = 0$$

$$v(T) = \sqrt{\frac{g}{r(T)}} \end{split}$$

This needs numerics to solve.

3.4.1 Terminal manifold with inequality constraints

$$\min_{u} \int_{0}^{T} L \, \mathrm{d}t + \Psi$$

$$\dot{x} = f(x, u)$$

$$\phi(x(T)) \le 0$$

$$\phi < 0$$

$$\phi > 0$$

$$\phi = 0$$

Repeat process: $\tilde{J} = \int (H - \lambda^T \dot{x}) dt + \Psi + \nu^T \phi$. The optimality conditions are

$$\begin{cases} \frac{\partial H}{\partial u} = 0 \\ \dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x} \\ \lambda(T) = \frac{\partial \Psi^{\mathrm{T}}}{\partial x} (x(T)) + \nu^{\mathrm{T}} \frac{\partial \phi^{\mathrm{T}}}{\partial x} (x(T)) \\ \nu \ge 0 \\ \phi(x(T)) \le 0 \\ \nu^{\mathrm{T}} \phi(x(T)) = 0 \quad (\mathrm{KKT}) \end{cases}$$

3.4.2 Initial manifold

$$\min_{x_0,u} \int L + \Psi(x(T)) + \Theta(x(0))$$

s.t. $\dot{x} = f(x,u)$
$$\phi(x(T)) = 0$$

$$\xi(x(0)) = 0$$

$$\begin{split} \tilde{J} &= \int (H - \lambda^{\mathrm{T}} \dot{x}) \, \mathrm{d}t + \Psi(x(T)) + \Theta(x(0)) + \nu_{\phi}^{\mathrm{T}} \phi(x(T)) + \nu_{\xi}^{\mathrm{T}} \xi(x(0)) \\ \delta \tilde{J} &= \int \left[\left(\frac{\partial H}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right) \eta + \frac{\partial H}{\partial u} v \right] \mathrm{d}t + \left[\frac{\partial \Psi}{\partial x} (x(T)) + \nu_{\phi}^{\mathrm{T}} \frac{\partial \phi}{\partial x} (x(T)) - \lambda^{\mathrm{T}} (T) \right] \eta(T) \\ &+ \left[\frac{\partial \Theta}{\partial x} (x(0)) + \nu_{\xi}^{\mathrm{T}} \frac{\partial \xi}{\partial x} (x(0)) + \lambda^{\mathrm{T}} (0) \right] \eta(0) \end{split}$$

The optimality conditions are

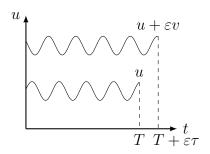
$$\begin{cases} \frac{\partial H}{\partial u} = 0 \\ \dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x} \\ \lambda(T) = -\frac{\partial \Psi^{\mathrm{T}}}{\partial x}(x(T)) - \nu_{\phi}^{\mathrm{T}} \frac{\partial \phi^{\mathrm{T}}}{\partial x}(x(T)) \\ \lambda(0) = -\frac{\partial \Theta^{\mathrm{T}}}{\partial x}(x(0)) - \nu_{\xi}^{\mathrm{T}} \frac{\partial \xi^{\mathrm{T}}}{\partial x}(x(0)) \end{cases}$$

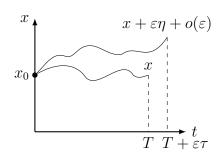
3.4.3 Unspecified Terminal Times

For example, instead of driving to the moon using minimum fuel, we want to get there as soon as possible:

$$\min_{u,T} \int_0^T L(x, u, t) dt + \Psi(x(T), T).$$

The variations are $u \mapsto u + \varepsilon v$ and $T \mapsto T + \varepsilon \tau$





$$\begin{split} \tilde{J}(u,T) &= \int_0^T [L(x,u,t) + \lambda^{\mathrm{T}}(f(x) - \dot{x})] \, \mathrm{d}t + \Psi(x(T),T) \\ &= \int_0^T [H - \lambda^{\mathrm{T}} \dot{x}] \, \mathrm{d}t + \Psi \\ \tilde{J}(u + \varepsilon v, T + \varepsilon \tau) &= \int_0^T [H(x + \varepsilon \eta, u + \varepsilon v, t, \lambda) - \lambda^{\mathrm{T}} (\dot{x} + \varepsilon \dot{\eta})] \, \mathrm{d}t \\ &+ \int_T^{T + \varepsilon \tau} [H(x + \varepsilon \eta, u + \varepsilon v, t, \lambda) - \lambda^{\mathrm{T}} (\dot{x} + \varepsilon \dot{\eta})] \, \mathrm{d}t \\ &+ \Psi(x(T + \varepsilon \tau) + \varepsilon \eta (T + \varepsilon \tau), T + \varepsilon \tau) \end{split}$$

$$\tilde{J}(u+\varepsilon v,T+\varepsilon\tau) - \tilde{J}(u,T) = \varepsilon \int_{0}^{T} \left(\frac{\partial H}{\partial x} + \dot{\lambda}^{T}\right) \eta \, dt + \varepsilon \int_{0}^{T} \frac{\partial H}{\partial u} v \, dt
- \varepsilon \lambda^{T}(T) \eta(T) + \varepsilon \lambda^{T}(0) \eta(0) + o(\varepsilon)
+ \underbrace{\int_{T}^{T+\varepsilon\tau} \left[H(x+\varepsilon\eta,u+\varepsilon v,t,\lambda) - \lambda^{T}(\dot{x}+\varepsilon\dot{\eta})\right] dt}_{(I)}
+ \underbrace{\Psi(x(T+\varepsilon\tau) + \varepsilon\eta(T+\varepsilon\tau),T+\varepsilon\tau) - \Psi(x(T),T)}_{(II)}$$

For term I, use the mean value theorem to get rid of terms inside the integral that have a ε before them:

$$\int_{T}^{T+\varepsilon\tau} [L + \lambda^{T} (f - \dot{x} - \varepsilon \dot{\eta})] dt$$

$$= \int_{T}^{T+\varepsilon\tau} \left[L(x, u, t) + \varepsilon \frac{\partial L}{\partial x} \eta + \varepsilon \frac{\partial L}{\partial u} v + \lambda^{T} \left(f + \varepsilon \frac{\partial f}{\partial x} \eta + \varepsilon \frac{\partial f}{\partial u} v - \dot{x} - \varepsilon \dot{\eta} \right) \right] dt + o(\varepsilon)$$

$$= \varepsilon \tau \left[L + \lambda^{T} (f - \dot{x}) \right] \Big|_{t=\xi} + o(\varepsilon) = \varepsilon \tau L \Big|_{t=\xi} + o(\varepsilon)$$

$$= \varepsilon \tau L(x(\xi), u(\xi), \xi) + o(\varepsilon), \quad \xi \in [T, T + \varepsilon \xi] \tag{3.5}$$

Note that as $\varepsilon \to 0$, $\xi \to T$.

For term II, we further split it into two parts:

$$\Psi(x+\varepsilon\eta,T+\varepsilon\tau) - \Psi(x,T) = \underbrace{\Psi(x,T+\varepsilon\tau)}_{\text{(II.a)}} + \underbrace{\varepsilon\frac{\partial\Psi}{\partial x}(x,T+\varepsilon\tau)\eta(T+\varepsilon\tau)}_{\text{(II.b)}} - \Psi(x,T)$$

$$(\text{II.a}) \Longrightarrow \Psi(x, T + \varepsilon \tau) = \Psi(x(T), T + \varepsilon \tau) + \varepsilon \frac{\partial \Psi}{\partial x}(x(T), T + \varepsilon \tau)\dot{x}(T)\tau + o(\varepsilon)$$

$$= \Psi(x(T), T) + \varepsilon \frac{\partial \Psi}{\partial x}(x(T), T)\dot{x}(T)\tau + \varepsilon \frac{\partial \Psi}{\partial T}(x(T), T)\tau + o(\varepsilon)$$

$$(\text{II.b}) \Longrightarrow \varepsilon \frac{\partial \Psi}{\partial x}(x, T + \varepsilon \tau)\eta(T + \varepsilon \tau)$$

$$= \varepsilon \left[\frac{\partial \Psi}{\partial x}(x(T), T) + \varepsilon \frac{\partial^2 \Psi}{\partial x^2}\dot{x}\tau + \varepsilon \frac{\partial^2 \Psi}{\partial T\partial x}\tau + o(\varepsilon) \right]$$

$$\times \left[\eta(T) + \varepsilon \dot{\eta}(T)\tau + o(\varepsilon) \right]$$

$$= \varepsilon \frac{\partial \Psi}{\partial x}(x(T), T)\eta(T) + o(\varepsilon)$$

$$(\text{II}) \Longrightarrow \Psi(x + \varepsilon \eta, T + \varepsilon \tau) - \Psi(x, T)$$

$$= \varepsilon \frac{\partial \Psi}{\partial x}(x(T), T)[\dot{x}(T)\tau + \eta(T)] + \varepsilon \frac{\partial \Psi}{\partial T}(x(T), T)\tau + o(\varepsilon) \quad (3.6)$$

Substituting (3.5) and (3.6) into (3.4) and taking the directional derivative,

$$\delta \tilde{J} = \int_0^T \left(\frac{\partial H}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right) \eta \, \mathrm{d}t + \int_0^T \frac{\partial H}{\partial u} v \, \mathrm{d}t + \lambda^{\mathrm{T}}(0) \eta(0)$$
$$+ \left[L + \frac{\partial \Psi}{\partial T} + \frac{\partial \Psi}{\partial x} f \right] \tau \bigg|_{t=T} + \left(\frac{\partial \Psi}{\partial x} - \lambda^{\mathrm{T}} \right) \eta \bigg|_{t=T}$$

So we have a mix of old and new:

old:
$$\frac{\partial H}{\partial u} = 0$$

 $\dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x}$
 $\lambda(T) = \frac{\partial \Psi}{\partial x}\Big|_{T}$
new: $L + \frac{\partial \Psi}{\partial T} + \lambda^{\mathrm{T}} f\Big|_{T} = 0$

This last condition is known as the *Transversality condition*.

Example Pure minimum time question

$$\min_{u,T} \int_0^T dt$$

$$\dot{x} = f(x, u)$$

$$x(0) = x_0$$

$$x(T) = x_T$$

$$H = L + \lambda^T f = 1 + \lambda^T f$$

The transversality condition is

$$\begin{split} L + \frac{\partial \Psi}{\partial T} + \lambda^{\mathrm{T}} f \bigg|_{T} &= 0 \\ \lambda^{\mathrm{T}} f \big|_{T} &= -1 \\ H(T) &= 1 + \lambda^{\mathrm{T}} f \big|_{T} = 1 - 1 = 0 \end{split}$$

But this is a conservative system, so H is a constant. Therefore,

$$H(t) = 0 \quad \forall t \in [0, T]$$

Example Zermelo's problem: sail from A to B as quickly as possible in the presence of known winds and currents.

$$v = \text{known}$$

$$\phi = \text{steering angle (input)}$$

$$A \bullet \qquad \qquad \bullet B$$

$$\text{wind, current}$$

The dynamics are

$$\dot{x} = v\cos\phi + c_1(x, y)
\dot{y} = v\sin\phi + c_2(x, y) \qquad \lambda = \begin{bmatrix} \lambda_x \\ \lambda_y \end{bmatrix}$$

For minimum time, L=1.

$$H = 1 + \lambda_x (v \cos \phi + c_1) + \lambda_y (v \sin \phi + c_2)$$
$$0 = \frac{\partial H}{\partial \phi} = -v \lambda_x \sin \phi + v \lambda_y \cos \phi$$
$$\phi = \tan^{-1} \left(\frac{\lambda_y}{\lambda_x}\right)$$

Since this is a conservative system and $\partial \Psi/\partial T=0$, then H(t)=H(T)=0.

$$-1 = \lambda_x (v \cos \phi + c_1) + \lambda_y (v \sin \phi + c_2)$$

$$\lambda_x = -\frac{\cos \phi}{v + c_1 \cos \phi + c_2 \sin \phi}$$

$$\lambda_y = -\frac{\sin \phi}{v + c_1 \cos \phi + c_2 \sin \phi}$$

$$\dot{\lambda} = -\frac{\partial H^{\mathrm{T}}}{\partial x}$$

$$\dot{\lambda}_x = -\lambda_x \frac{\partial c_1}{\partial x} - \lambda_y \frac{\partial c_2}{\partial x}$$

$$\dot{\lambda}_y = -\lambda_x \frac{\partial c_1}{\partial y} - \lambda_y \frac{\partial c_2}{\partial y}$$

$$\dot{\phi} = \sin^2 \phi \frac{\partial c_2}{\partial x} + \sin \phi \cos \phi \left(\frac{\partial c_1}{\partial x} - \frac{\partial c_2}{\partial y}\right) - \cos^2 \phi \frac{\partial c_1}{\partial y}$$

This is an ODE that completely determines ϕ if we just had ϕ_0 .

Example We want to drive a car and stop at a stop sign as quickly as possible. Assume that the stop sign is at the origin, and our control is the acceleration $(\ddot{x} = u)$.

$$\min_{u,T} \int_0^T dt$$
s.t.
$$\begin{cases}
\dot{x}_1 = x_2, & x(0) = x_0 \\
\dot{x}_2 = u, & x(T) = 0
\end{cases}$$

Recall the transversality condition:

$$H + \frac{\partial \Psi}{\partial T} \bigg|_{t=T} = 0.$$

For minimum-time problems, L=1 and $\Psi=0$, so $\lambda^{\mathrm{T}} f|_{t=T}=-1$.

$$H = 1 + \lambda_1 x_2 + \lambda_2 u$$

$$\lambda_1(T) \underbrace{x_2(T)}_{=0 \text{ (rest)}} + \lambda_2(T) u(T) = -1$$

$$\underbrace{\lambda_2(T) u(T)}_{\partial u} = -1$$

$$\underbrace{\frac{\partial H}{\partial u}}_{\partial u} = \underbrace{\lambda_2 = 0}_{,}$$

i.e. $0 \cdot u(T) = -1$? This problem is ill-posed; we need to go infinitely fast...

Idea 1: Constrain u. We don't know how to do this.

Idea 2: Pay for gas. This is a design choice.

For the second idea,

$$\min_{u,T} \int_0^T \frac{1}{2} u^2(t) dt$$
s.t.
$$\begin{cases}
\dot{x}_1 = x_2, & x(0) = x_0 \\
\dot{x}_2 = u, & x(T) = 0
\end{cases}$$

$$H = \frac{1}{2} u^2 + \lambda_1 x_2 + \lambda_2 u$$

$$\frac{1}{2} u^2(T) + \lambda_1(T) x_2(T) + \lambda_2(T) u(T) = 0$$

$$\frac{1}{2} u^2(T) + \lambda_2(T) u(T) = 0$$

$$\frac{\partial H}{\partial u} = u + \lambda_2 = 0 \Longrightarrow u = -\lambda_2$$

$$\frac{1}{2} \lambda_2^2(T) - \lambda_2^2(T) = 0$$

$$\lambda_2(T) = 0$$

$$\dot{\lambda}_1 = -\frac{\partial H}{\partial x_1} = 0 \Longrightarrow \lambda_1 = c$$

$$\dot{\lambda}_2 = -\frac{\partial H}{\partial x_2} = -\lambda_1 \Longrightarrow \lambda_2 = -ct + d$$

$$\lambda_2(T) = -cT + d = 0 \Longrightarrow T = \frac{d}{c}$$

$$\dot{x}_2 = u = -\lambda_2 = ct - d$$

$$x_2 = c\frac{t^2}{2} - dt + x_{2,0}$$

$$\dot{x}_1 = x_2 \Longrightarrow x_1 = c\frac{t^3}{6} - d\frac{t^2}{2} + x_{2,0}t + x_{1,0}$$

$$\begin{cases} x_1(T) = c\frac{T^3}{6} - d\frac{T^2}{2} + x_{2,0}T + x_{1,0} = 0 \\ x_2(T) = c\frac{T^2}{2} - dT + x_{2,0} = 0 \end{cases}$$

$$T = \frac{d}{c}$$

$$d = \sqrt{-\frac{4}{3}\frac{x_{2,0}^3}{x_{1,0}}}$$

$$T = \frac{d}{c}$$

$$u = ct - d$$

Fine, but we really want to get there as quickly as possible! We have to constrain u, e.g. $u(t) \in [-1, 1], \forall t \in [0, T]$. How do we deal with the constraints on u?

3.5 Hamilton's Minor "Mistake"

$$\min_{u \in \mathcal{U}_{\text{constr.}}} \int_0^T L(x, u, t) \, \mathrm{d}t + \Psi(x(T))$$
s.t. $\dot{x} = f(x, u, t)$

$$x(0) = x_0$$

$$(u(t) \in U)$$

Augment the cost:

$$\tilde{J}(u) = \int_0^T \left(H(x, u, t, \lambda) - \lambda^{\mathrm{T}} \dot{x} \right) \mathrm{d}t + \Psi(x(T))$$

Vary $u \mapsto u + \varepsilon v$ s.t. $u + \varepsilon v \in \mathcal{U}_{\text{constr.}} \Rightarrow x \mapsto x + \varepsilon \eta + o(\varepsilon)$:

$$\tilde{J}(u+\varepsilon v) = \int_0^T \left(H(x+\varepsilon \eta, u+\varepsilon v, t, \lambda) - \lambda^{\mathrm{T}} \dot{x} - \lambda^{\mathrm{T}} \varepsilon \dot{\eta} \right) \mathrm{d}t + \Psi(x(T) + \varepsilon \eta(T)) + o(\varepsilon)$$

Instead of computing $\delta \tilde{J}(u;v)$, let's check $\Delta \tilde{J} = \tilde{J}(u+\varepsilon v) - \tilde{J}(u)$. If $\Delta \tilde{J} \geq 0 \ \forall v \ \text{s.t.} u + \varepsilon v \in \mathcal{U}_{\text{constr.}}$ for ε small enough, then u is a local minimum!

$$\Delta \tilde{J} = \int_0^T \left[H(x + \varepsilon \eta, u + \varepsilon v, t, \lambda) - H(x, u, t, \lambda) - \lambda^{\mathrm{T}} (\dot{x} + \varepsilon \dot{\eta} - \dot{x}) \right] dt + \Psi(x(T) + \varepsilon \eta(T)) - \Psi(x(T)) + o(\varepsilon)$$

Only Taylor expanding w.r.t. x:

$$\begin{split} \Delta \tilde{J} &= \int_0^T \left[\varepsilon \frac{\partial H}{\partial x}(x,u,t,\lambda) \eta - \varepsilon \lambda^{\mathrm{T}} \dot{\eta} \right] \mathrm{d}t + \int_0^T \left[H(x,u+\varepsilon v,t,\lambda) - H(x,u,t,\lambda) \right] \mathrm{d}t \\ &+ \varepsilon \frac{\partial \Psi}{\partial x}(x(T)) \eta(T) + o(\varepsilon) \\ &= \varepsilon \int_0^T \left(\frac{\partial H}{\partial x} + \dot{\lambda}^{\mathrm{T}} \right) \eta \, \mathrm{d}t + \varepsilon \lambda^{\mathrm{T}}(0) \eta(0) - \varepsilon \lambda^{\mathrm{T}}(T) \eta(T) + \varepsilon \frac{\partial \Psi}{\partial x}(x(T)) \eta(T) \\ &+ \int_0^T \left[H(x,u+\varepsilon v,t,\lambda) - H(x,u,t,\lambda) \right] \mathrm{d}t + o(\varepsilon) \end{split}$$

With $\dot{\lambda} = -\partial H^{\mathrm{T}}/\partial x$ and $\lambda(T) = \partial \Psi(x(T))/\partial x$,

$$\Delta \tilde{J} = \int_0^T \left[H(x, u + \varepsilon v, t, \lambda) - H(x, u, t, \lambda) \right] dt + o(\varepsilon)$$

Here, Hamilton did Taylor's expansion and set $\partial H/\partial u = 0$. Instead, Pontryagin desired $\Delta \tilde{J} \geq 0 \ \forall v | u + \varepsilon v \in \mathcal{U}_{\text{constr.}}, \ \varepsilon$ small enough, i.e. we need

$$H(x, u^* + \varepsilon v, t, \lambda) \ge H(x, u^*, t, \lambda)$$

 $\forall t \in [0, t], \forall v | u + \varepsilon v \in \mathcal{U}_{\text{constr.}}, \varepsilon \text{ small enough. That is, we need}$

$$u^* = \arg\min_{u} H(x, u, t, \lambda)$$

In summary,

Hamilton:
$$\frac{\partial H}{\partial u} = 0$$

Pontryagin: $\min_{u} H$

Theorem (Pontryagin's Maximum Principle (PMP)). Consider the problem:

$$\min_{u,T} \int_0^T L(x, u, t) dt + \Psi(x(T), T)$$

$$s.t. \quad \dot{x} = f(x, u, t)$$

$$u(t) \in U(x, t), \quad \forall t \in [0, T]$$

$$x_i(0) = x_{i0}, \qquad i \in \mathcal{I}$$

$$x_j(T) = x_{jT}, \qquad j \in \mathcal{T}$$

The necessary condition for optimality is

$$H = L + \lambda^{T} f$$

$$\dot{\lambda} = -\frac{\partial H^{T}}{\partial x}$$

$$\lambda_{j}(0) = 0, \quad j \notin \mathcal{I}$$

$$\lambda_{i}(T) = \frac{\partial \Psi}{\partial x_{i}}(x(T)), \quad i \notin \mathcal{T}$$

$$H + \frac{\partial \Psi}{\partial T}\Big|_{t=T} = 0$$

$$u^{*}(x, t, \lambda) = \underset{u \in U(x, t)}{\operatorname{arg min}} H(x, u, t, \lambda)$$

We have two paths to solve optimality problems: we always start with the Hamiltonian, find the costate dynamics and boundary conditions, and apply the transversality condition; then, we can either apply calculus of variations (COV) or Pontryagin's Maximum Principle (PMP). COV only works for unconstrained problems, while with PMP we can deal with constraints.

3.6 Bang-Bang Control

Return to the car problem:

$$\min_{u,T} \int_0^T dt$$
s.t.
$$\begin{cases}
\dot{x}_1 = x_2, & x_1(0) = x_{1,0}, & x_1(T) = 0 \\
\dot{x}_2 = u, & x_2(0) = x_{2,0}, & x_2(T) = 0 \\
u(t) \in [-1, 1] & \forall t \in [0, T]
\end{cases}$$

How do we minimize H w.r.t. u?

$$H = 1 + \lambda_1 x_2 + \lambda_2 u$$

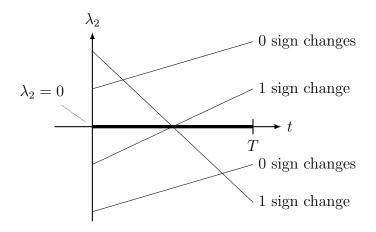
Clearly, we minimize H by letting

$$u = \begin{cases} -1, & \lambda_2 > 0 \\ +1, & \lambda_2 < 0 = -\operatorname{sign}(\lambda_2) \\ ??, & \lambda_2 = 0 \end{cases}$$

Therefore, the optimal u switches between -1 and +1 (bang-bang control).

$$\begin{split} \dot{\lambda}_1 &= -\frac{\partial H}{\partial x_1} = 0 \Longrightarrow \lambda_1 = c \\ \dot{\lambda}_2 &= -\frac{\partial H}{\partial x_2} = -\lambda_1 \Longrightarrow \lambda_2 = -ct + d \end{split}$$

Notice that $\lambda_2(t)$ is a line, so it has at most one sign change. Thus, u also changes sign (from ± 1 to ∓ 1) at most one time.



Let's solve this for all x_0 !

i) Assume $\lambda_2 > 0 \ \forall t \in [0, T], \ \therefore u = -1 \quad \forall t \in [0, T]$

$$\dot{x}_2 = -1 \Longrightarrow x_2 = -t + k_1
x_2(T) = 0 = -T + k_1 \Longrightarrow k_1 = T
x_2(t) = T - t \Longrightarrow x_2 > 0, \ t \in [0, T)
\dot{x}_1 = x_2 = T - t \Longrightarrow x_1 = -\frac{t^2}{2} + Tt + k_2
x_1(T) = 0 = -\frac{T^2}{2} + T^2 + k_2 \Longrightarrow k_2 = -\frac{T^2}{2}
x_1(t) = -\frac{t^2}{2} + Tt - \frac{T^2}{2} = -\frac{(T - t)^2}{2} \quad (< 0, \ t \in [0, T))
= -\frac{x_2^2(t)}{2}$$

Let's consider the curve

$$\begin{bmatrix} -x_2^2/2 \\ x_2 \end{bmatrix}$$

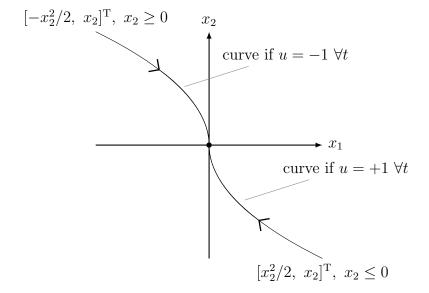
for $x_2 \ge 0$. If x_0 lies on this curve, use u = -1 and drive to the origin.

ii) Assume $u = +1 \quad \forall t \in [0, T]$

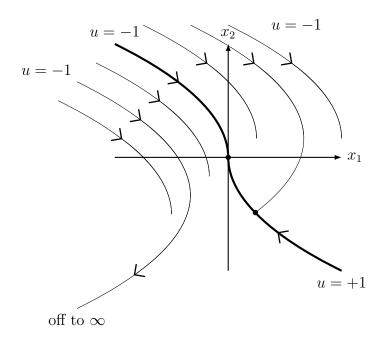
$$x_2 = t - T \quad (\le 0 \text{ on } [0, T])$$

 $x_1 = \frac{x_2^2}{2} \quad (\ge 0)$

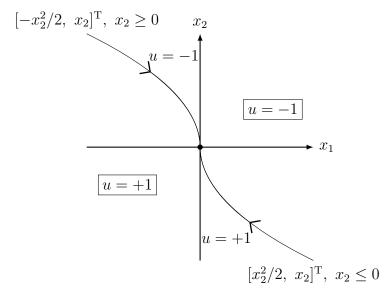
For this curve, use u = +1.



What happens when we do not start on the curves? We start with a certain u depending on x_0 and perform a single switch of u when we encounter one of the initial curves that travel to the origin. Note that for the case $\lambda_2 = 0 \forall t$, we start at the stop sign at rest, so the control does not matter.



The optimal solution is given by the following *switching curve*.



Note 1: Bang-bang control typically involves

- a) finding the number of switches
- b) find the switching surfaces

Note 2: This is a feedback law! (u depends on x!!)

3.6.1 Linear Systems (scalar input)

$$\min_{u,T} \int_0^T dt$$
s.t. $\dot{x} = Ax + Bu$

$$x(0) = x_0, \quad x(T) = 0$$

$$u \in [-1, 1]$$

$$H = 1 + \lambda^T (Ax + Bu)$$

$$u = -\operatorname{sign}(\lambda^T B) \quad \text{(bang-bang)}$$

Aside...

$$\dot{x} = f(x) + g(x)u$$
 (control affine)
 $H = 1 + \lambda^{T} f + \lambda^{T} g u$
 $u = -\operatorname{sign}(\lambda^{T} g(x))$ (bang-bang)

Back to linear...

$$\dot{\lambda} = -\frac{\partial H^{T}}{\partial x} = -A^{T}\lambda$$
$$\lambda(t) = e^{-A^{T}t}\lambda_{0}$$
$$u(t) = -\operatorname{sign}\left(\lambda_{0}^{T}e^{-At}B\right)$$

How do we find λ_0 ?

$$\dot{x} = Ax + Bu, \quad x(0) = x_0$$

$$x(T) = e^{AT}x_0 + \int_0^T e^{A(T-\tau)}Bu(\tau) d\tau$$

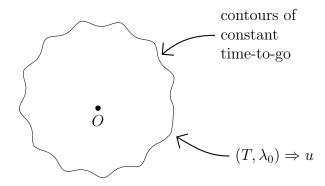
$$x(T) = 0 = e^{AT}x_0 - \int_0^T e^{A(T-t)}B \operatorname{sign}\left(\lambda_0^T e^{-At}B\right) dt$$
(3.7)

Problem 1: Given x_0 , figure out λ_0 from (3.7). Then, $u = -\operatorname{sign}(\lambda_0^{\mathrm{T}} e^{-At}B)$. This has to be done numerically in general (not super simple...).

Problem 2: Find all x_0 s from which it takes the same amount of time to get to x(T) = 0.

$$e^{At}x_0 = \int_0^T e^{A(T-t)}B\operatorname{sign}\left(\lambda_0^{\mathrm{T}}e^{-At}B\right)\mathrm{d}t$$
$$x_0 = \int_0^T e^{-At}B\operatorname{sign}\left(\lambda_0^{\mathrm{T}}e^{-At}B\right)\mathrm{d}t$$

Fix T. By varying λ_0 , we will get the x_0 s that take time T to go to x(T) = 0 optimally.



So by solving problem 2, we find λ_0 associated with all x_0 , i.e. we have "solved" problem 1 as well.