

On the CCN (de)activation nonlinearities

Sylwester Arabas

Jagiellonian University

introduction

Arabas & Shima 2017

Nonlin. Processes Geophys., 24, 535–542, 2017

<https://doi.org/10.5194/npg-24-535-2017>

© Author(s) 2017. This work is distributed under
the Creative Commons Attribution 3.0 License.



Nonlinear Processes
in Geophysics



On the CCN (de)activation nonlinearities

Sylwester Arabas^{1,2} and Shin-ichiro Shima³

¹Institute of Geophysics, Faculty of Physics, University of Warsaw, Warsaw, Poland

²Chatham Financial Corporation Europe, Cracow, Poland

³Graduate School of Simulation Studies, University of Hyogo, Kobe, Japan

Correspondence to: Sylwester Arabas (sarabas@chathamfinancial.eu) and Shin-ichiro Shima (s_shima@sim.u-hyogo.ac.jp)

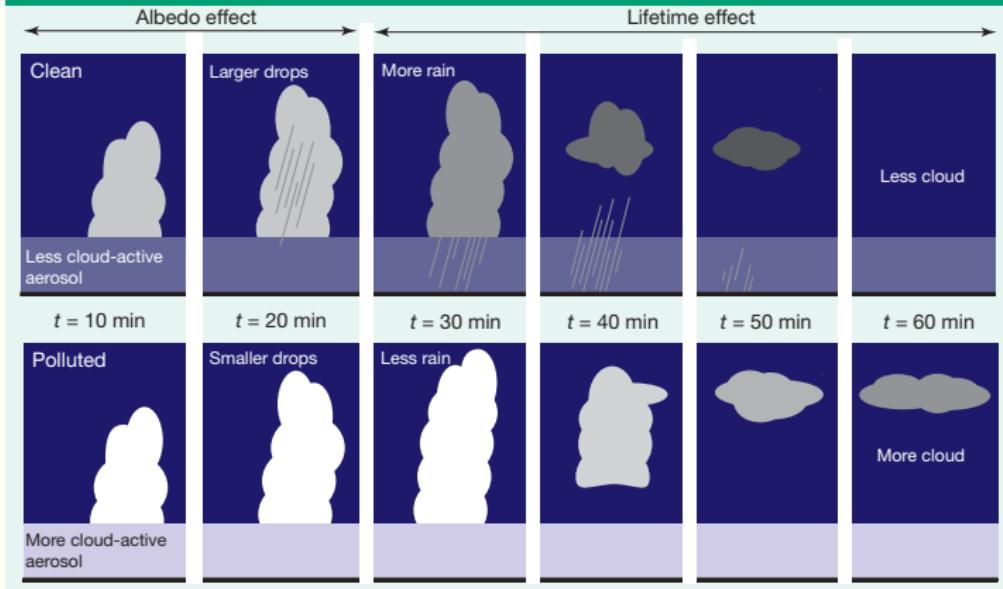
Received: 9 September 2016 – Discussion started: 4 October 2016

Revised: 23 May 2017 – Accepted: 24 July 2017 – Published: 5 September 2017

one-slide aerosol-cloud (micro-macro) interaction primer

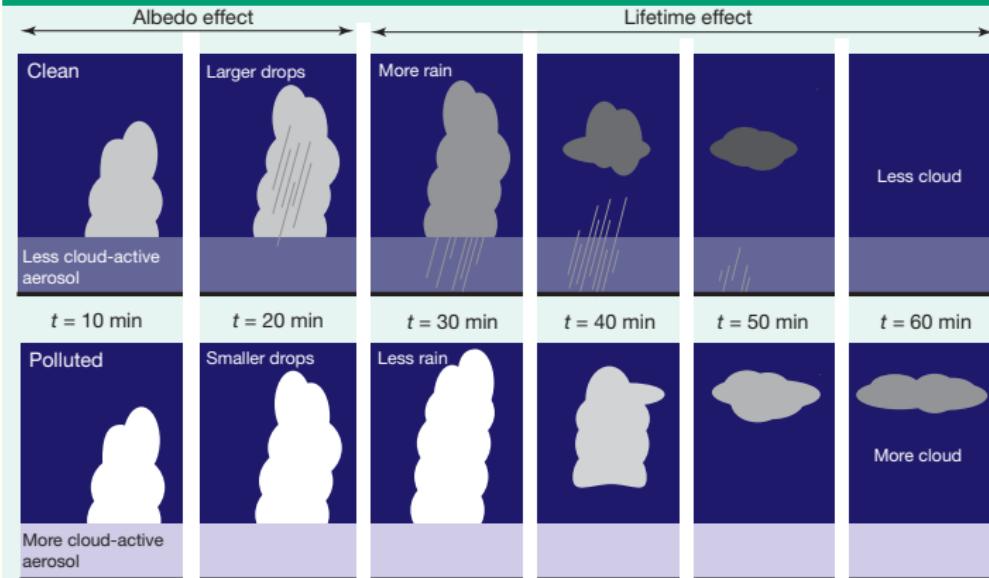
one-slide aerosol-cloud (micro-macro) interaction primer

Stevens and Feingold, 2009 (Nature)



one-slide aerosol-cloud (micro-macro) interaction primer

Stevens and Feingold, 2009 (Nature)

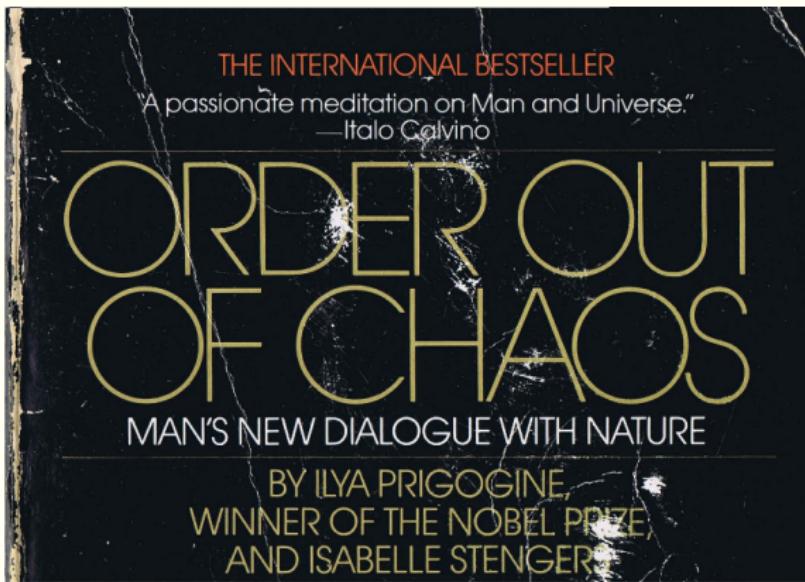


Stevens and Boucher, 2012 (Nature)

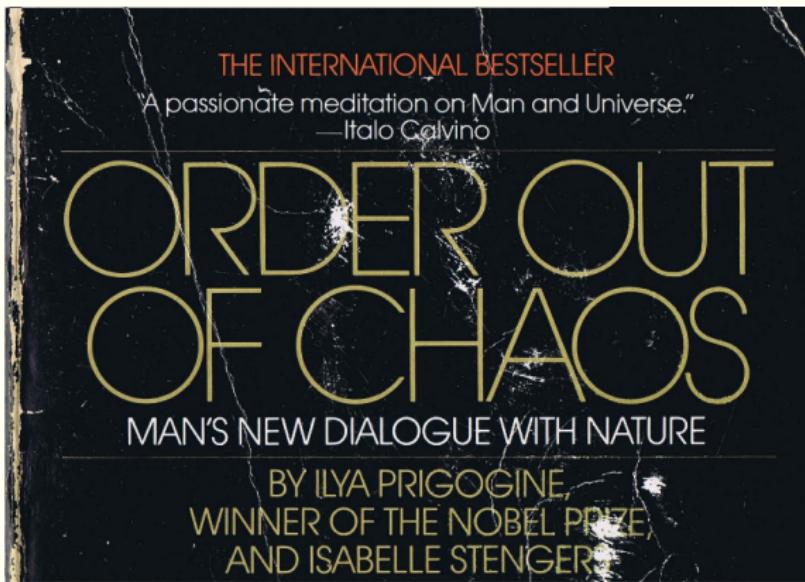
"there is something captivating about the idea that fine particulate matter, suspended almost invisibly in the atmosphere, holds the key to some of the greatest mysteries of climate science"

... others captivated by micro-macro interactions

... others captivated by micro-macro interactions



... others captivated by micro-macro interactions



Prigogine and Stengers 1984

"Much of this book has centered around the relation between the microscopic and the macroscopic. One of the most important problems in evolutionary theory is the eventual feedback between macroscopic structures and microscopic events: macroscopic structures emerging from microscopic events would in turn lead to a modification of the microscopic mechanisms."

Prigogine and Stengers 1984 @ BUW

"Z chaosu ku porządkowi"

Szukaj

[Wyszukiwanie zaawansowane](#)

Logowanie

Wpisz numer karty bibliotecznej i hasło dostępu

Nr karty bibliotecznej

Hasło

Zaloguj

Aktualne wyszukiwanie: "Z chaosu ku porządkowi" 

Wyniki 1 do 1 z 1.

Uporządkuj według Relewancji

Dodaj stronę do Schowka Dodaj wszystko do schowka

1.  **Z chaosu ku porządkowi : nowy dialog człowieka z przyrodą**
Isabelle Stengers ; przeł. Katarzyna Lipszyc ; przedm. opatrzyła Barbara Baranowska.

Prigogine, Ilya (1917-2003).

Klasyfikacja WD Q175 .P8822165 1990

Adres wyd. Warszawa : Państwowy Instytut Wydawniczy, 1990.

Opis fiz. 355, [1] s. : il. ; 20 cm.

Seria Biblioteka Myśli Współczesnej

Dostępne egzemplarze: 8 z 9

BUW Magazyn (1)
BUW Wolny Dostęp (2)
BUW wd Księgozbiór prof. Żurowskiego (1)
Inst. Stosowanych Nauk Społecznych (1)
WYCOFANE (0 z 1)
Wydział Pedagogiczny (2)
Wydział Psychologii - Wypożyczalnia (1)

Pokaż mniej...

Ogranicz wyniki wyszukiwania

Dodatkowe terminy

Dodaj

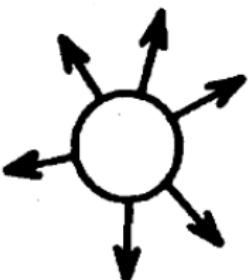
Lokalizacja

BUW Magazyn (1)
 BUW Wolny Dostęp (1)

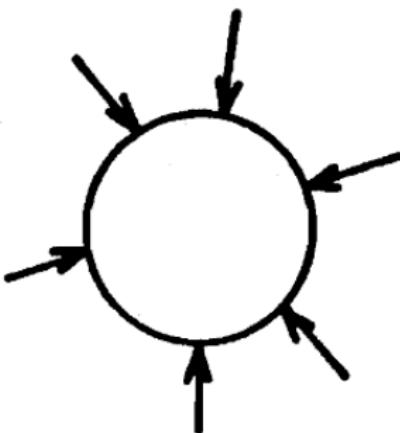
regime-transition (bifurcation) example from P&S 1984

regime-transition (bifurcation) example from P&S 1984

ORDER OUT OF CHAOS 188



(a)



(b)

Figure 19. Nucleation of a liquid droplet in a supersaturated vapor. (a) droplet smaller than the critical size; (b) droplet larger than the critical size. The existence of the threshold has been experimentally verified for dissipative structures.

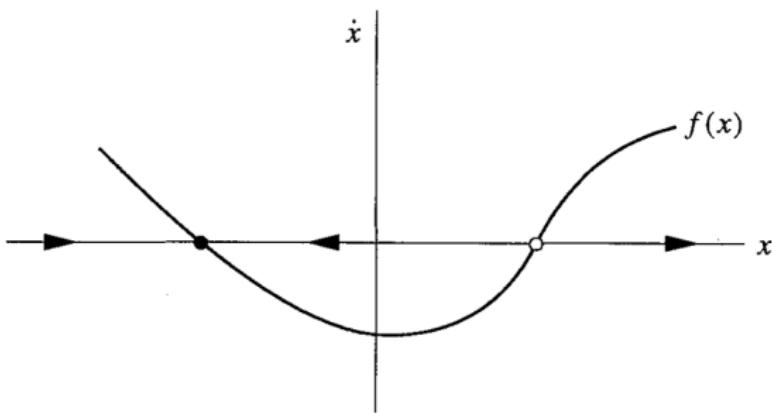
two-slide bifurcation analysis primer (1/2)

two-slide bifurcation analysis primer (1/2)

Strogatz 2014 (sect. 2.2): fixed points and stability

graphical (qualitative) analysis
of a non-linear one-dimensional dynamical system:

$$\dot{x} = f(x)$$



NONLINEAR
With Applications to Physics,
Biology, Chemistry, and Engineering
DYNAMICS
AND CHAOS



Steven H. Strogatz

SECOND EDITION

two-slide bifurcation analysis primer (2/2)

Strogatz 2014 (sect. 3.1): saddle-node bifurcation

prototypical example of saddle-node bifurcation:

$$\dot{x} = r + x^2$$

r : parameter (distinct regimes if positive, negative or zero)

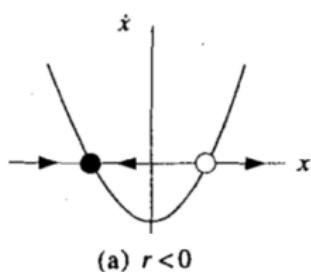
two-slide bifurcation analysis primer (2/2)

Strogatz 2014 (sect. 3.1): saddle-node bifurcation

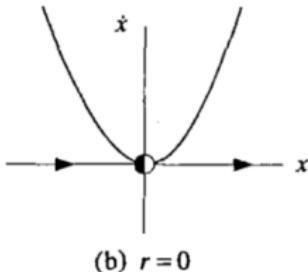
prototypical example of saddle-node bifurcation:

$$\dot{x} = r + x^2$$

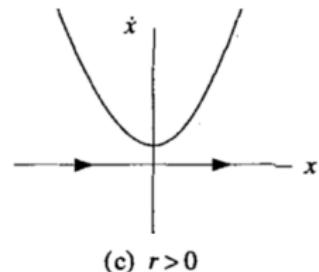
r : parameter (distinct regimes if positive, negative or zero)



(a) $r < 0$



(b) $r = 0$



(c) $r > 0$

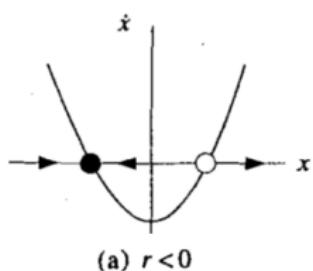
two-slide bifurcation analysis primer (2/2)

Strogatz 2014 (sect. 3.1): saddle-node bifurcation

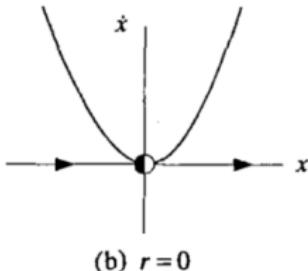
prototypical example of saddle-node bifurcation:

$$\dot{x} = r + x^2$$

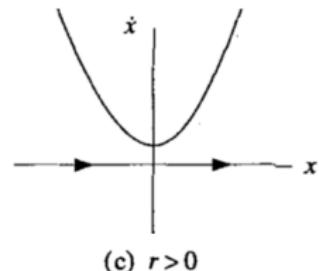
r : parameter (distinct regimes if positive, negative or zero)



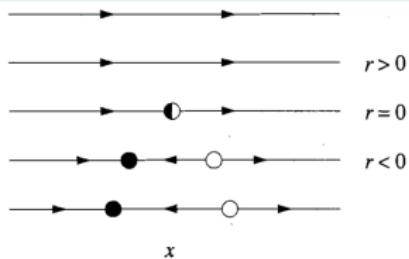
(a) $r < 0$



(b) $r = 0$



(c) $r > 0$



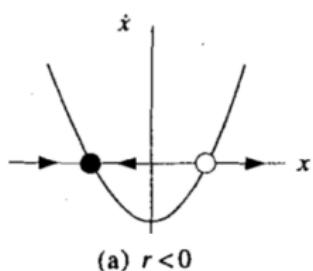
two-slide bifurcation analysis primer (2/2)

Strogatz 2014 (sect. 3.1): saddle-node bifurcation

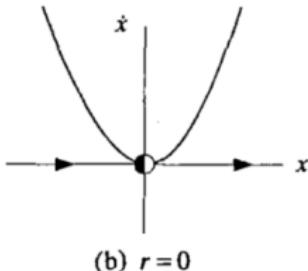
prototypical example of saddle-node bifurcation:

$$\dot{x} = r + x^2$$

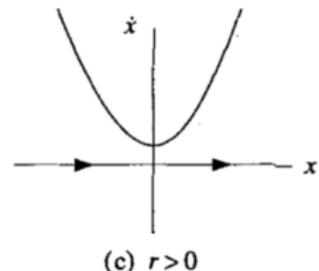
r : parameter (distinct regimes if positive, negative or zero)



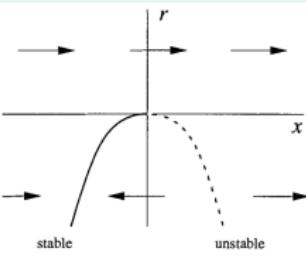
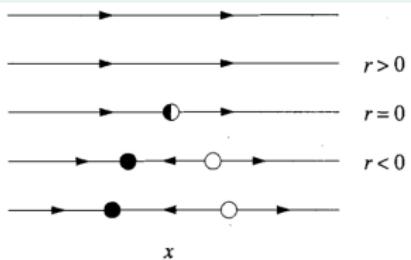
(a) $r < 0$



(b) $r = 0$



(c) $r > 0$



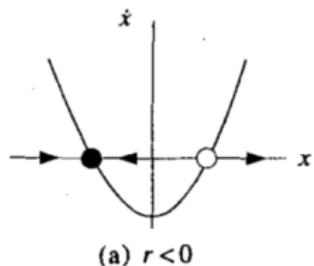
two-slide bifurcation analysis primer (2/2)

Strogatz 2014 (sect. 3.1): saddle-node bifurcation

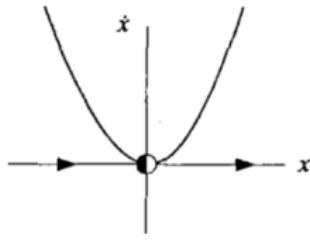
prototypical example of saddle-node bifurcation:

$$\dot{x} = r + x^2$$

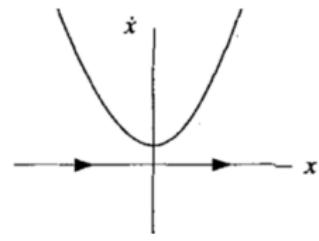
r : parameter (distinct regimes if positive, negative or zero)



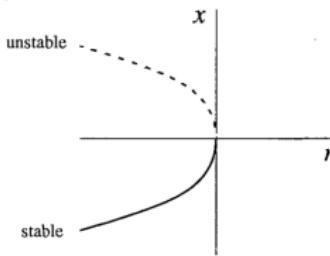
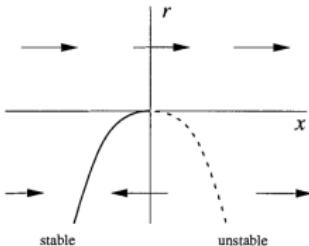
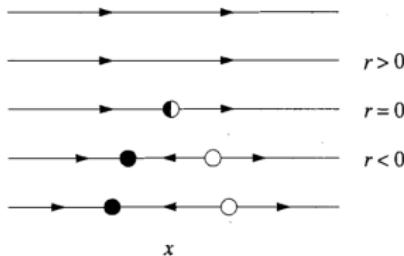
(a) $r < 0$



(b) $r = 0$

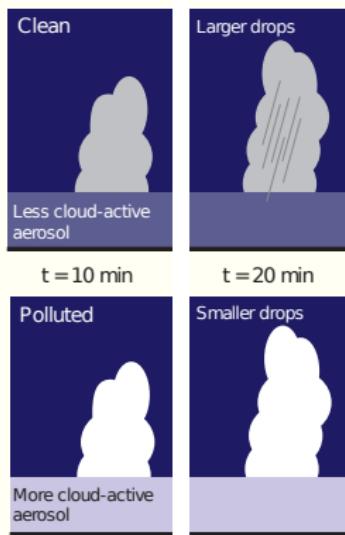


(c) $r > 0$

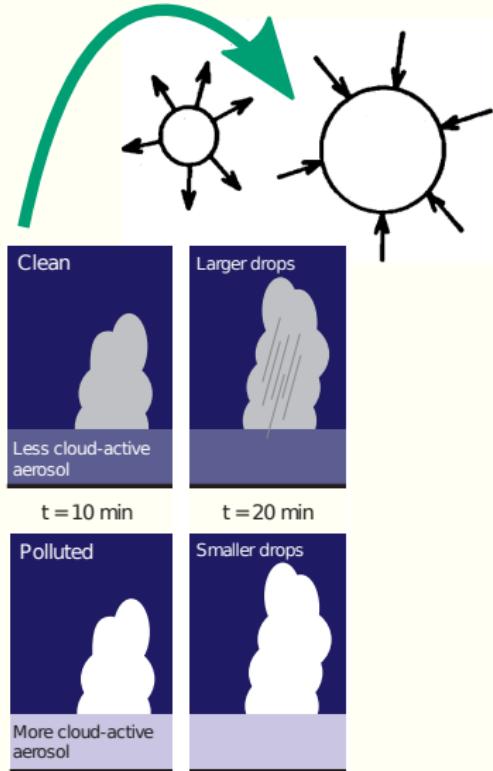


connecting the dots ...

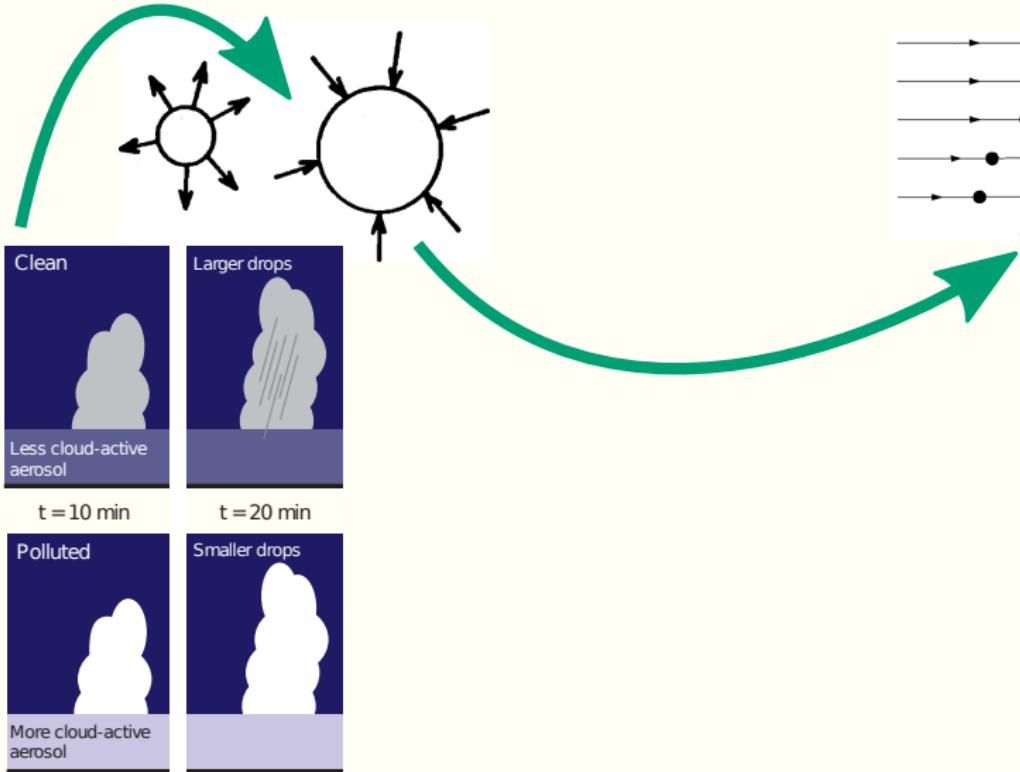
connecting the dots ...



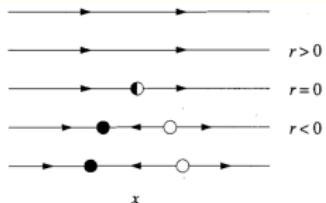
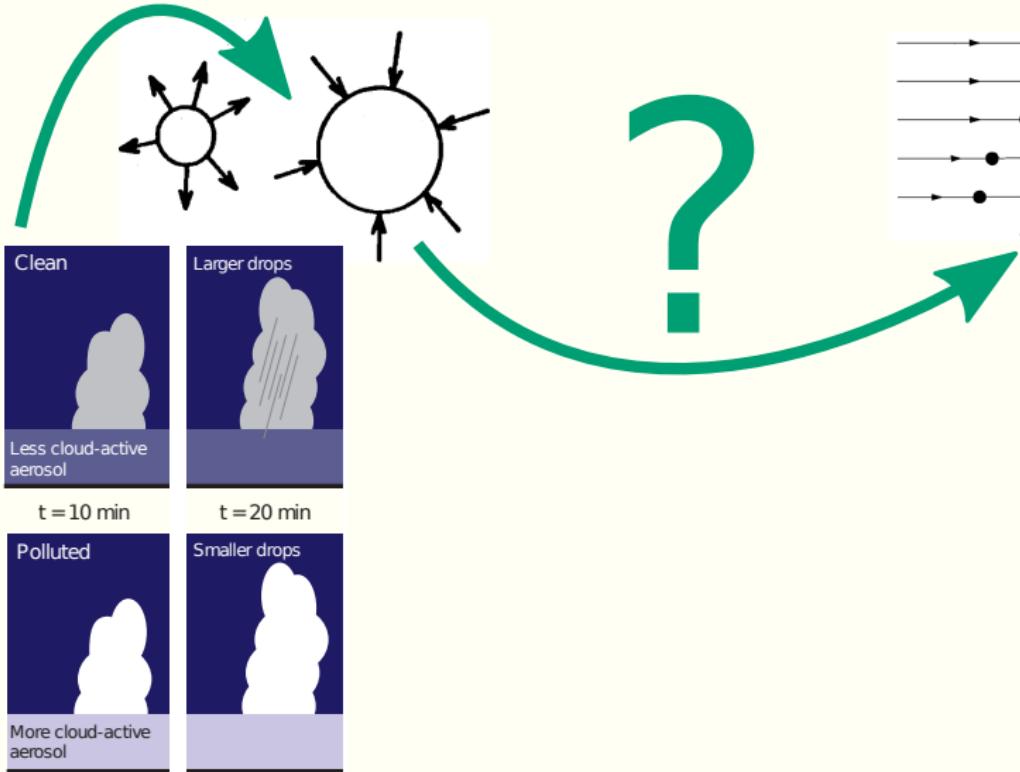
connecting the dots ...



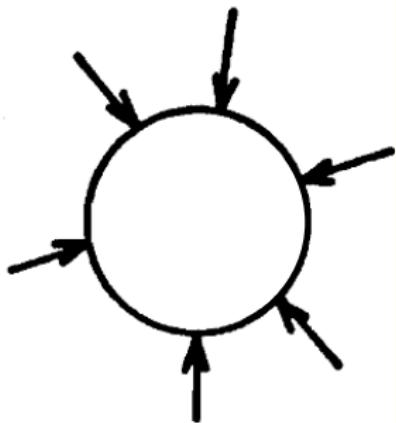
connecting the dots ...



connecting the dots ...

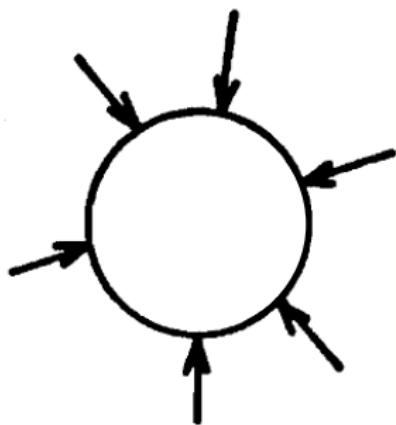


droplet growth laws in a nutshell: mass and heat diffusion



droplet growth laws in a nutshell: mass and heat diffusion

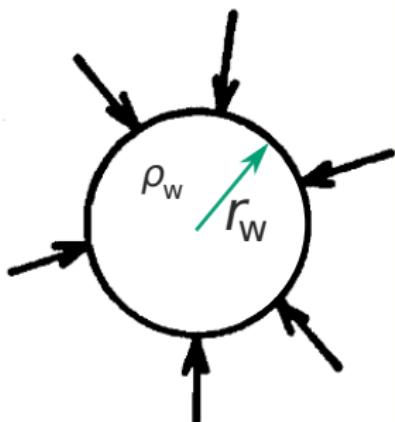
Fick's and Fourier's laws combined
spherical geometry



$$\dot{r}_w = \frac{1}{r_w} \frac{D_{\text{eff}}}{\rho_w} (\rho_v - \rho_o)$$

droplet growth laws in a nutshell: mass and heat diffusion

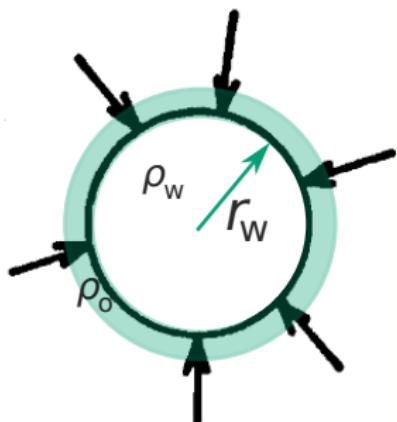
Fick's and Fourier's laws combined
spherical geometry



$$\dot{r}_w = \frac{1}{r_w} \frac{D_{\text{eff}}}{\rho_w} (\rho_v - \rho_o)$$

droplet growth laws in a nutshell: mass and heat diffusion

Fick's and Fourier's laws combined
spherical geometry

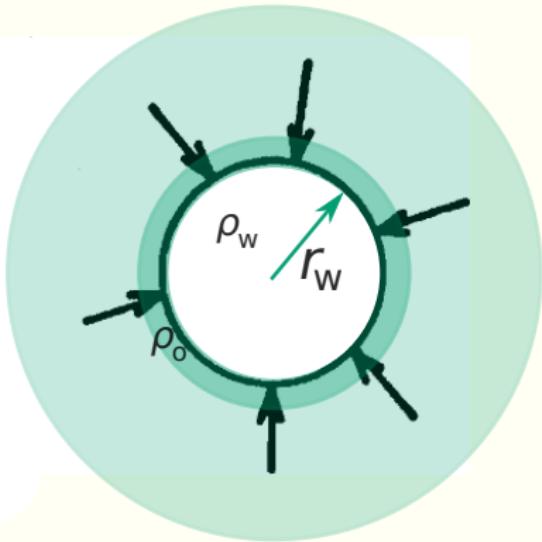


$$\dot{r}_w = \frac{1}{r_w} \frac{D_{\text{eff}}}{\rho_w} (\rho_v - \rho_o)$$

droplet growth laws in a nutshell: mass and heat diffusion

Fick's and Fourier's laws combined
spherical geometry

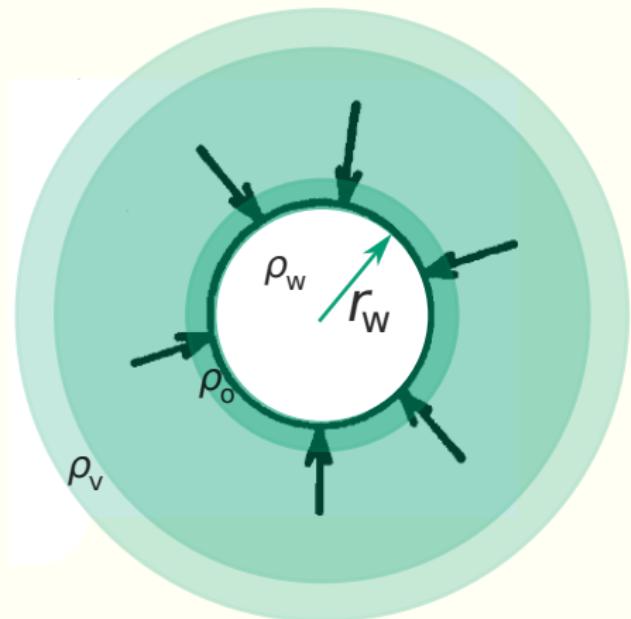
$$\dot{r}_w = \frac{1}{r_w} \frac{D_{\text{eff}}}{\rho_w} (\rho_v - \rho_o)$$



droplet growth laws in a nutshell: mass and heat diffusion

Fick's and Fourier's laws combined
spherical geometry

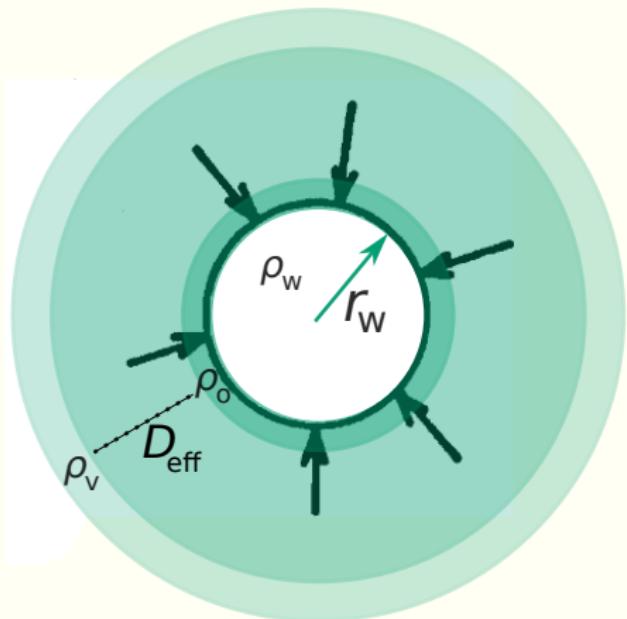
$$\dot{r}_w = \frac{1}{r_w} \frac{D_{\text{eff}}}{\rho_w} (\rho_v - \rho_o)$$



droplet growth laws in a nutshell: mass and heat diffusion

Fick's and Fourier's laws combined
spherical geometry

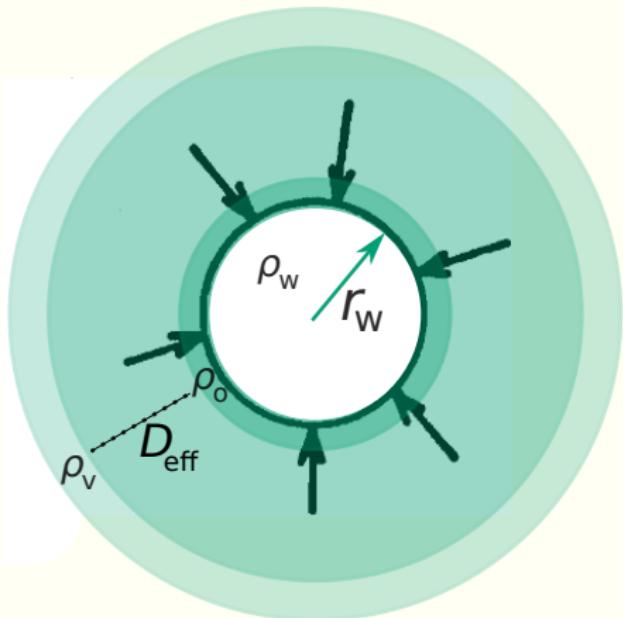
$$\dot{r}_w = \frac{1}{r_w} \frac{D_{\text{eff}}}{\rho_w} (\rho_v - \rho_o)$$



droplet growth laws in a nutshell: mass and heat diffusion

Fick's and Fourier's laws combined
spherical geometry

$$\dot{r}_w = \frac{1}{r_w} \frac{D_{\text{eff}}}{\rho_w} (\rho_v - \rho_o)$$



non-dimensional numbers:

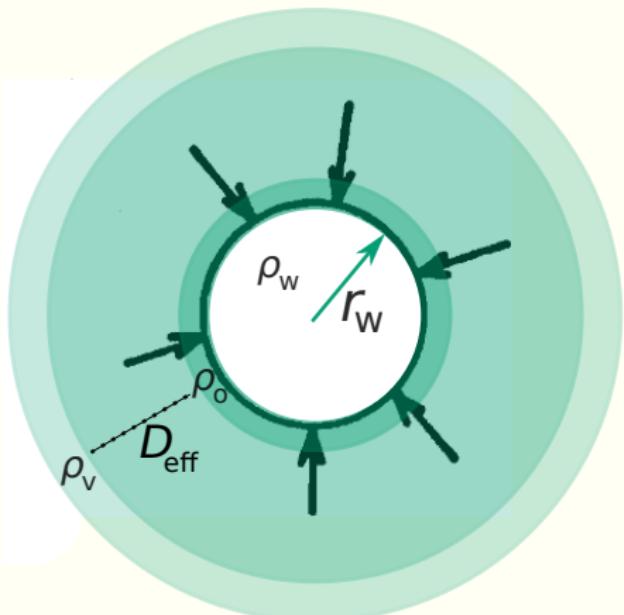
$$\text{RH} = \rho_v / \rho_{vs}$$

$$\text{RH}_{\text{eq}} = \rho_o / \rho_{vs}$$

droplet growth laws in a nutshell: mass and heat diffusion

Fick's and Fourier's laws combined
spherical geometry

$$\dot{r}_w = \frac{1}{r_w} \frac{D_{\text{eff}}}{\rho_w} (\rho_v - \rho_o)$$



non-dimensional numbers:

$$\text{RH} = \rho_v / \rho_{vs}$$

$$\text{RH}_{\text{eq}} = \rho_o / \rho_{vs}$$

$$\dot{r}_w = \frac{1}{r_w} D_{\text{eff}} \frac{\rho_{vs}}{\rho_w} (\text{RH} - \text{RH}_{\text{eq}})$$

droplet growth laws in a nutshell: Köhler curve

$$\dot{r}_w = \frac{1}{r_w} D_{\text{eff}} \frac{\rho_{vs}}{\rho_w} (\text{RH} - \text{RH}_{\text{eq}})$$

droplet growth laws in a nutshell: Köhler curve

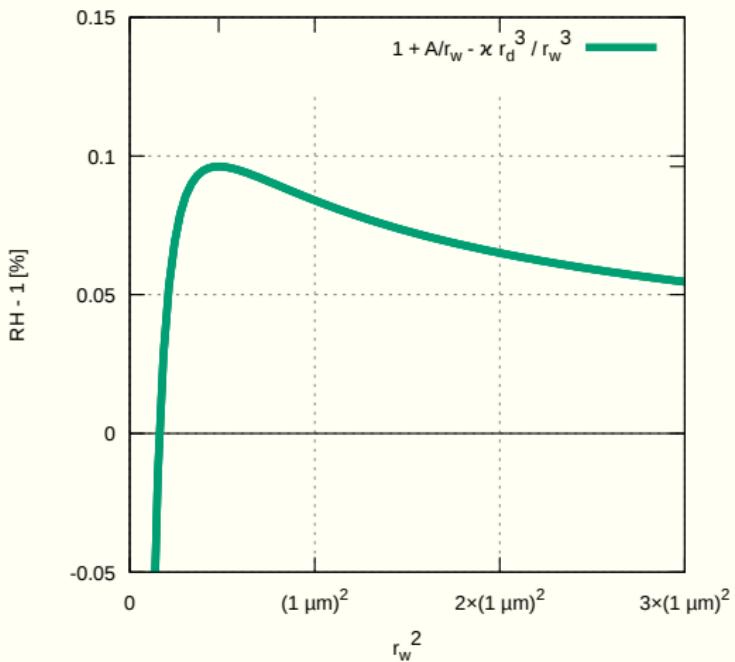
$$\dot{r}_w = \frac{1}{r_w} D_{\text{eff}} \frac{\rho_{vs}}{\rho_w} (\text{RH} - \text{RH}_{\text{eq}})$$
$$\text{RH}_{\text{eq}} = \frac{r_w^3 - r_d^3}{r_w^3 - r_d^3(1 - \kappa)} \exp\left(\frac{A}{r_w}\right)$$
$$\approx 1 + \frac{A}{r_w} - \frac{\kappa r_d^3}{r_w^3}$$

droplet growth laws in a nutshell: Köhler curve

$$\dot{r}_w = \frac{1}{r_w} D_{\text{eff}} \frac{\rho_{\text{vs}}}{\rho_w} (\text{RH} - \text{RH}_{\text{eq}})$$

$$\text{RH}_{\text{eq}} = \frac{r_w^3 - r_d^3}{r_w^3 - r_d^3(1 - \kappa)} \exp\left(\frac{A}{r_w}\right)$$

$$\approx 1 + \frac{A}{r_w} - \frac{\kappa r_d^3}{r_w^3}$$

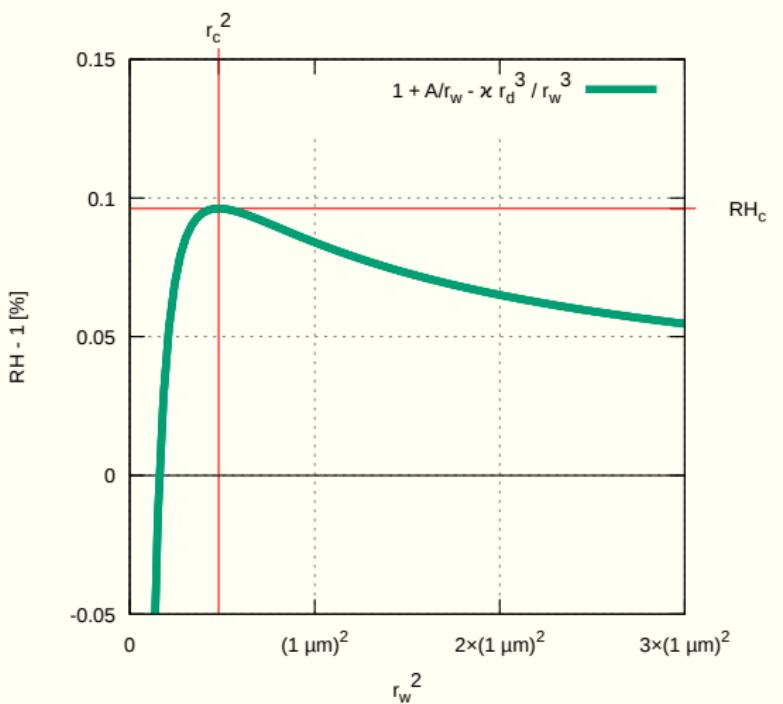


droplet growth laws in a nutshell: Köhler curve

$$\dot{r}_w = \frac{1}{r_w} D_{\text{eff}} \frac{\rho_{\text{vs}}}{\rho_w} (\text{RH} - \text{RH}_{\text{eq}})$$

$$\text{RH}_{\text{eq}} = \frac{r_w^3 - r_d^3}{r_w^3 - r_d^3(1 - \kappa)} \exp\left(\frac{A}{r_w}\right)$$

$$\approx 1 + \frac{A}{r_w} - \frac{\kappa r_d^3}{r_w^3}$$

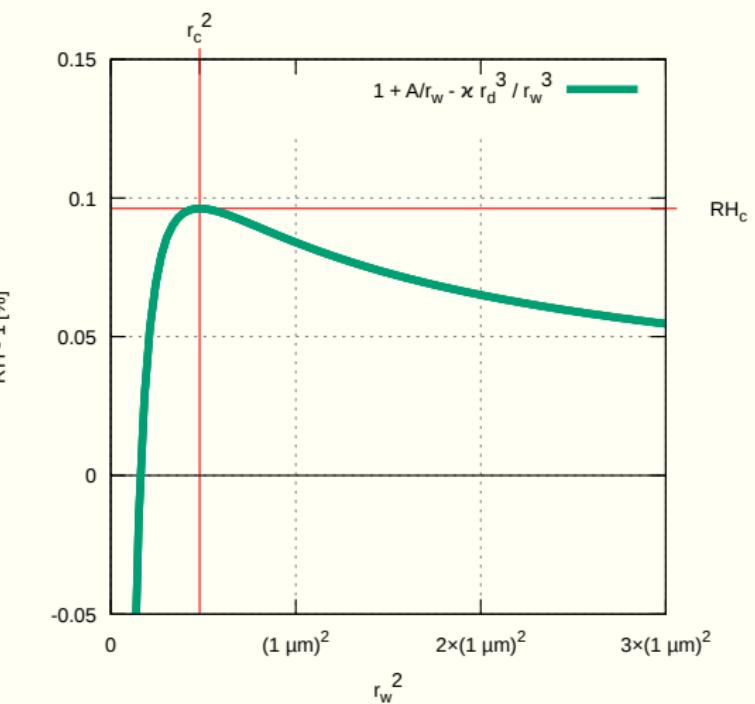


droplet growth laws in a nutshell: Köhler curve

$$\dot{r}_w = \frac{1}{r_w} D_{\text{eff}} \frac{\rho_{\text{vs}}}{\rho_w} (\text{RH} - \text{RH}_{\text{eq}})$$

$$\text{RH}_{\text{eq}} = \frac{r_w^3 - r_d^3}{r_w^3 - r_d^3(1 - \kappa)} \exp\left(\frac{A}{r_w}\right)$$

$$\approx 1 + \frac{A}{r_w} - \frac{\kappa r_d^3}{r_w^3}$$

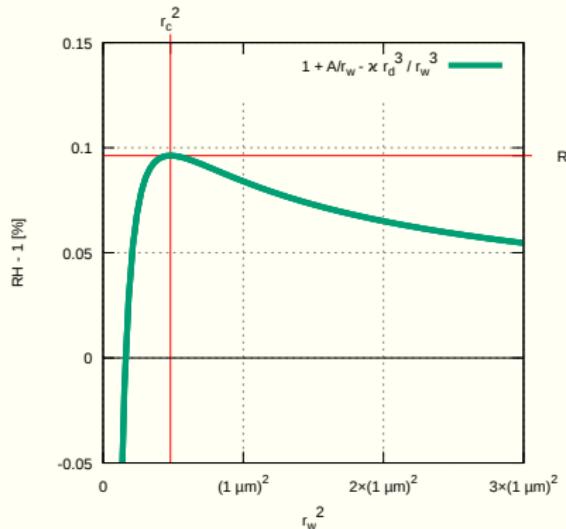


maximum at (r_c, RH_c) :

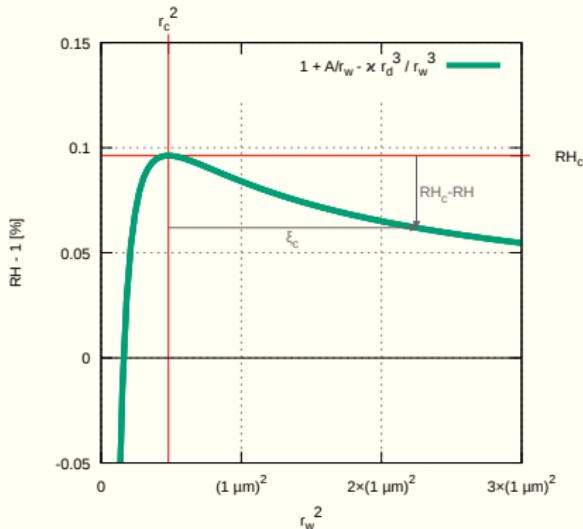
$$r_c = \sqrt{3\kappa r_d^3 / A}$$

$$\text{RH}_c = 1 + \frac{2A}{3r_c}$$

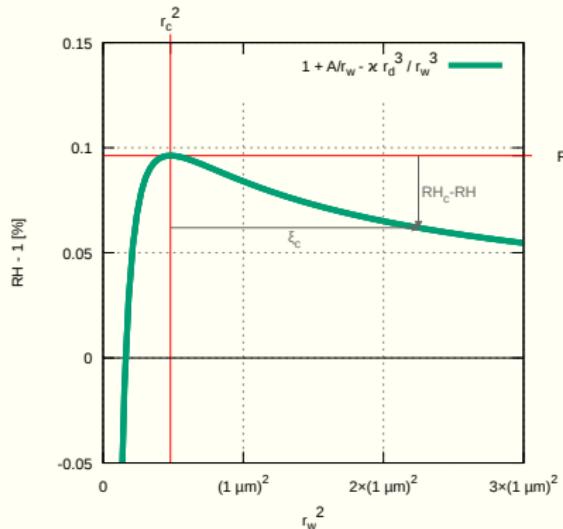
phase portrait of the system: flipped Köhler curve



phase portrait of the system: flipped Köhler curve



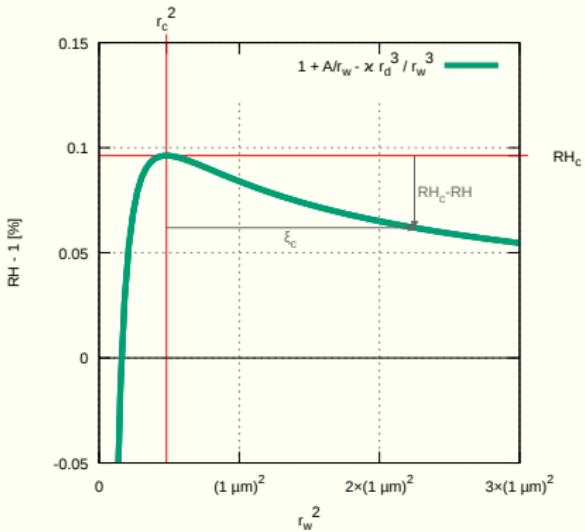
phase portrait of the system: flipped Köhler curve



$$\xi = r_w^2 + C$$

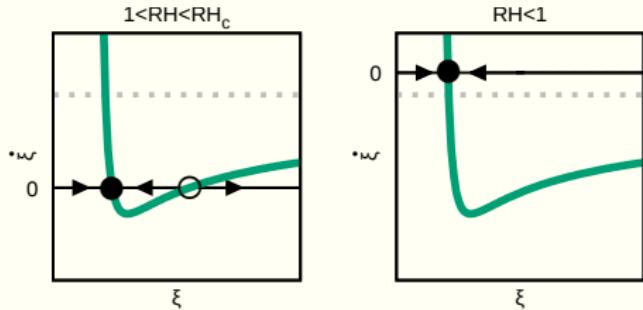
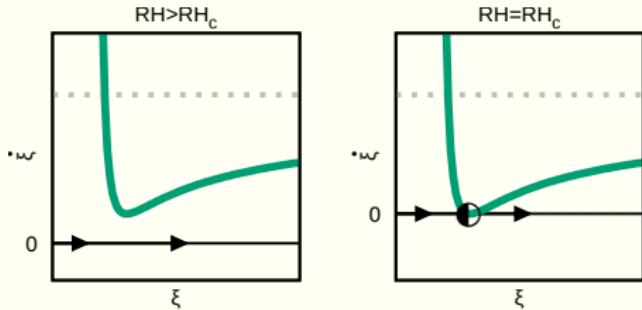
$$\dot{\xi} = 2D_{\text{eff}} \frac{\rho_{\text{vs}}}{\rho_w} (\text{RH} - \text{RH}_{\text{eq}}(\xi))$$

phase portrait of the system: flipped Köhler curve



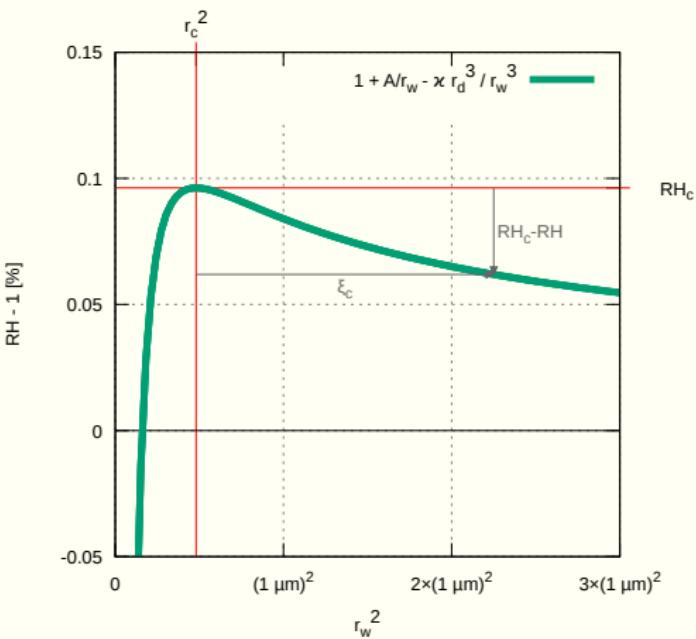
$$\xi = r_w^2 + C$$

$$\dot{\xi} = 2D_{\text{eff}} \frac{\rho_{\text{vs}}}{\rho_w} (\text{RH} - \text{RH}_{\text{eq}}(\xi))$$



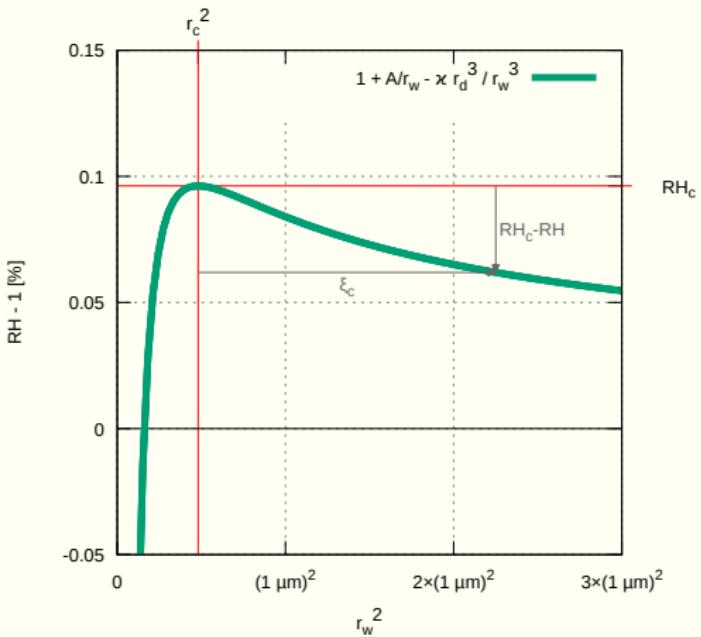
saddle-node bifurcation at Köhler curve maximum

saddle-node bifurcation at Köhler curve maximum



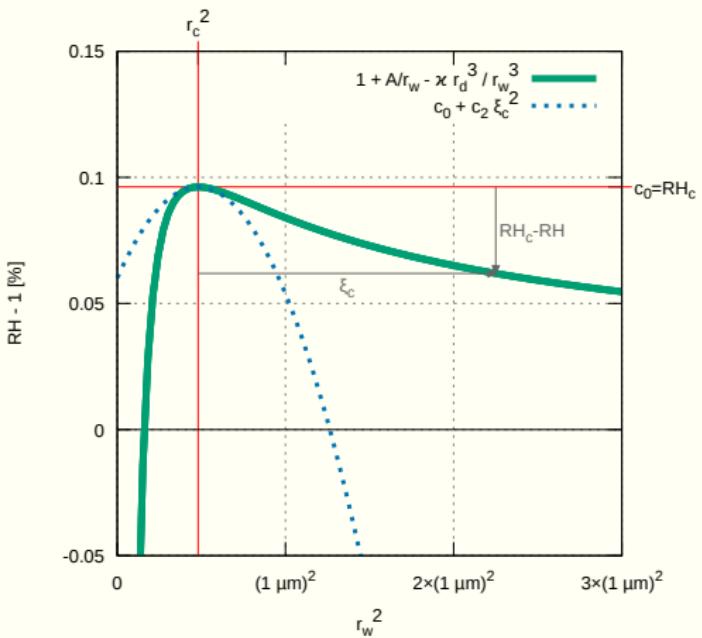
saddle-node bifurcation at Köhler curve maximum

$$RH_{eq}(\xi_c) = c_0 + \cancel{c_1 \xi_c} + c_2 \xi_c^2 + \dots$$



saddle-node bifurcation at Köhler curve maximum

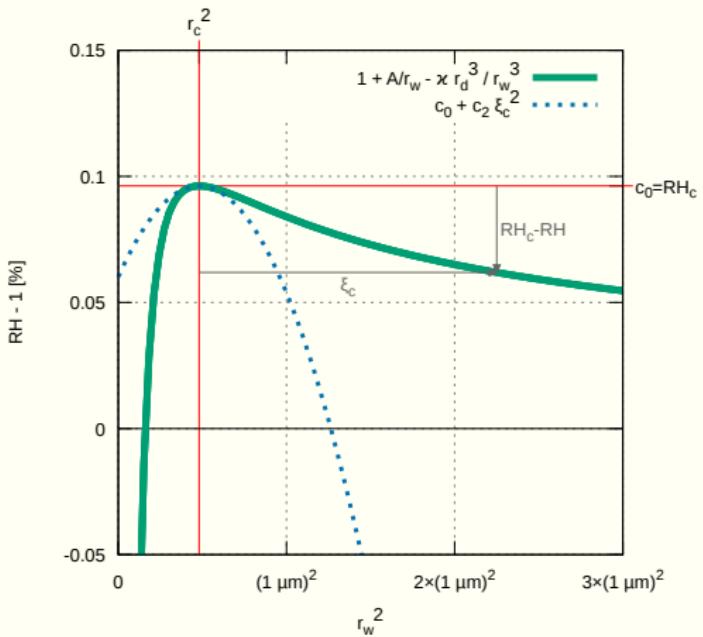
$$RH_{eq}(\xi_c) = c_0 + \cancel{c_1 \xi_c} + c_2 \xi_c^2 + \dots$$



saddle-node bifurcation at Köhler curve maximum

$$RH_{eq}(\xi_c) = c_0 + c_1 \xi_c + c_2 \xi_c^2 + \dots$$

$$\dot{\xi}_c \Big|_{\xi_c \rightarrow 0} \sim \frac{RH - RH_c}{A/(4r_c^5)} + \xi_c^2$$

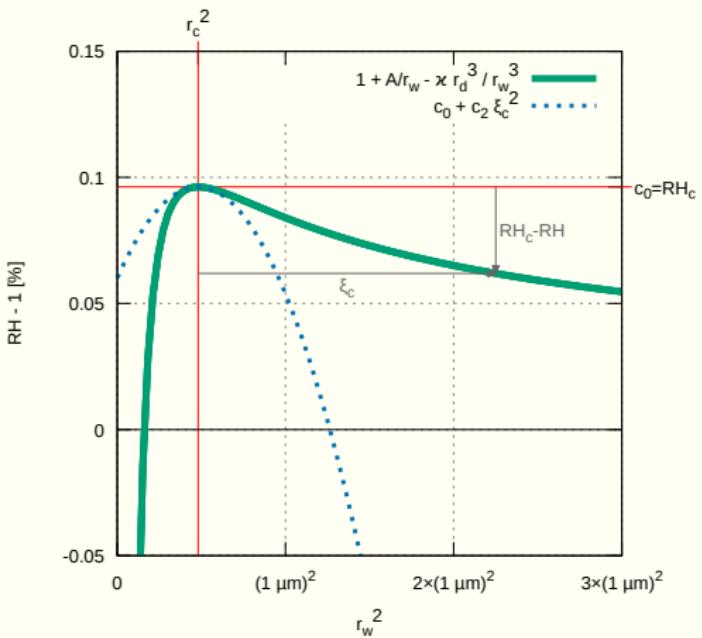


saddle-node bifurcation at Köhler curve maximum

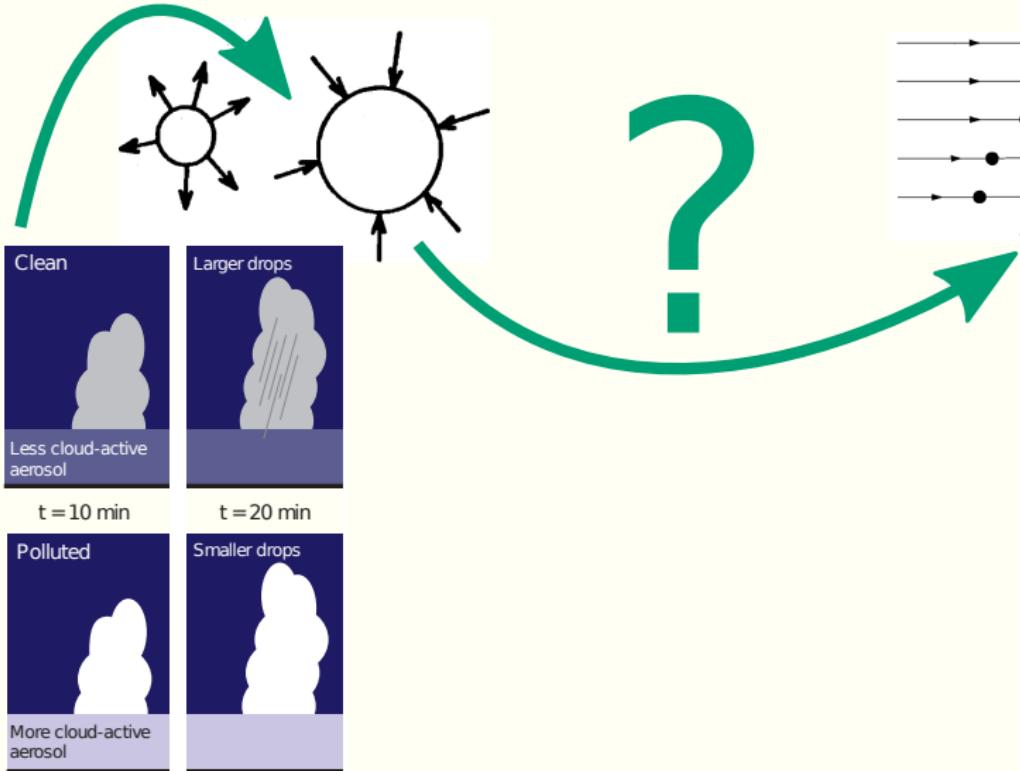
$$RH_{eq}(\xi_c) = c_0 + c_1 \xi_c + c_2 \xi_c^2 + \dots$$

$$\dot{\xi}_c \Big|_{\xi_c \rightarrow 0} \sim \frac{RH - RH_c}{A/(4r_c^5)} + \xi_c^2$$

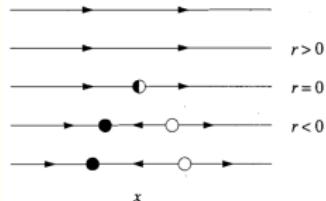
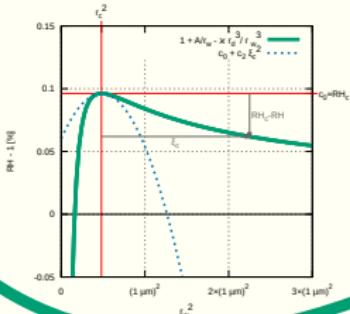
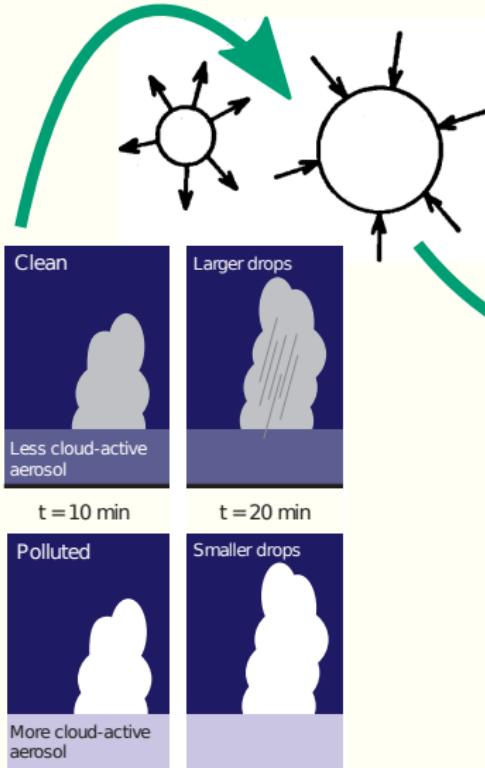
$$\dot{x} = r + x^2$$



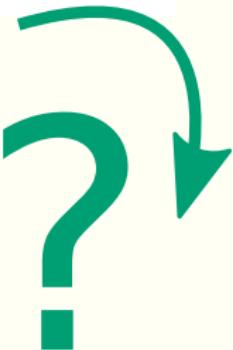
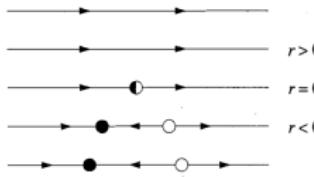
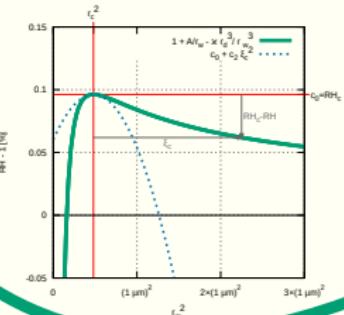
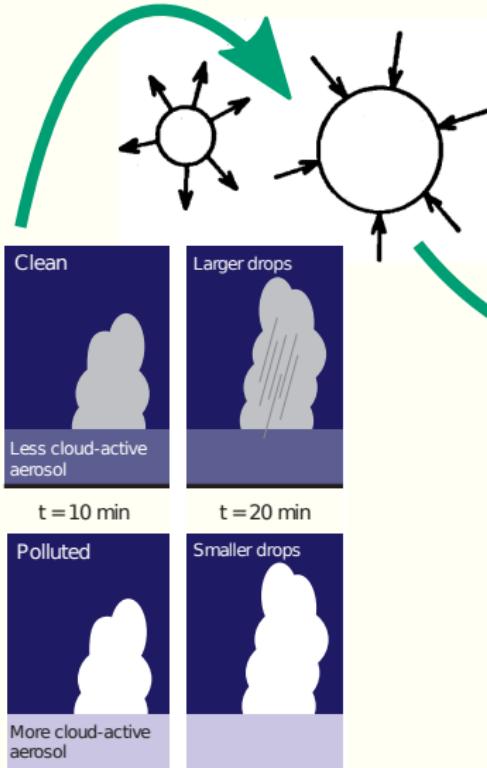
connecting the dots ...



connecting the dots ...



connecting the dots ...



coalescence in the saddle-node bottleneck (sic!)

coalescence in the saddle-node bottleneck (sic!)

Strogatz 2014 (sect. 4.3): *coalescence* of the fixed points
is associated with a passage through a *bottleneck*,

key observation: time of passage through the parabolic *bottleneck*
dominates all other timescales

coalescence in the saddle-node bottleneck (sic!)

Strogatz 2014 (sect. 4.3): *coalescence* of the fixed points
is associated with a passage through a *bottleneck*,

key observation: time of passage through the parabolic *bottleneck*
dominates all other timescales

coalescence in the saddle-node bottleneck (sic!)

Strogatz 2014 (sect. 4.3): *coalescence* of the fixed points
is associated with a passage through a *bottleneck*,

key observation: time of passage through the parabolic *bottleneck*
dominates all other timescales

$$\tau_{act} \approx \int_{-\infty}^{+\infty} \frac{d\xi_c}{\dot{\xi}_c}$$

coalescence in the saddle-node bottleneck (sic!)

Strogatz 2014 (sect. 4.3): *coalescence* of the fixed points
is associated with a passage through a *bottleneck*,

key observation: time of passage through the parabolic *bottleneck*
dominates all other timescales

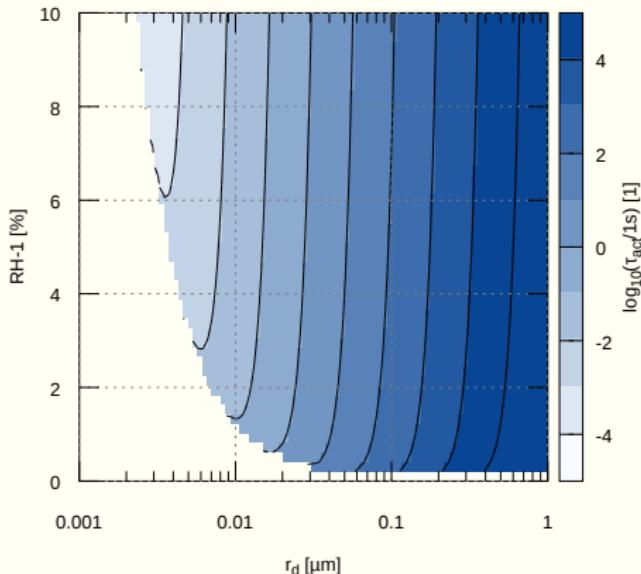
$$\begin{aligned}\tau_{act} &\approx \int_{-\infty}^{+\infty} \frac{d\xi_c}{\dot{\xi}_c} \\ &= \frac{r_c^{5/2}}{\sqrt{A}} \frac{\rho_w/\rho_{vs}}{D_{\text{eff}}} \frac{\pi}{\sqrt{RH - RH_c}}\end{aligned}$$

coalescence in the saddle-node bottleneck (sic!)

Strogatz 2014 (sect. 4.3): *coalescence* of the fixed points
is associated with a passage through a *bottleneck*,

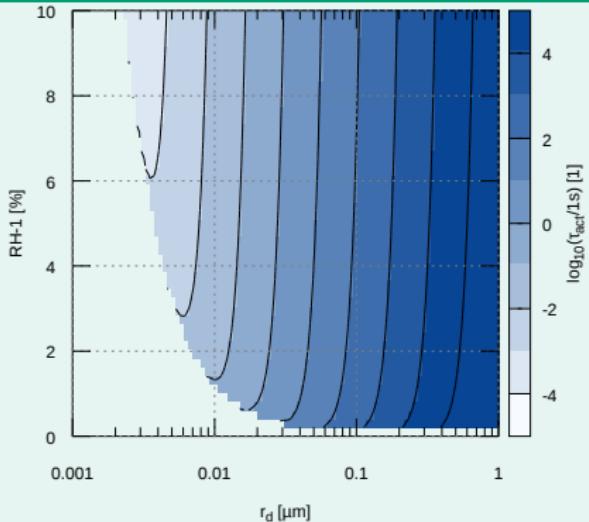
key observation: time of passage through the parabolic *bottleneck*
dominates all other timescales

$$\tau_{act} \approx \int_{-\infty}^{+\infty} \frac{d\xi_c}{\dot{\xi}_c}$$
$$= \frac{r_c^{5/2}}{\sqrt{A}} \frac{\rho_w/\rho_{vs}}{D_{\text{eff}}} \frac{\pi}{\sqrt{RH - RH_c}}$$



activation timescale: analytic vs. numerical

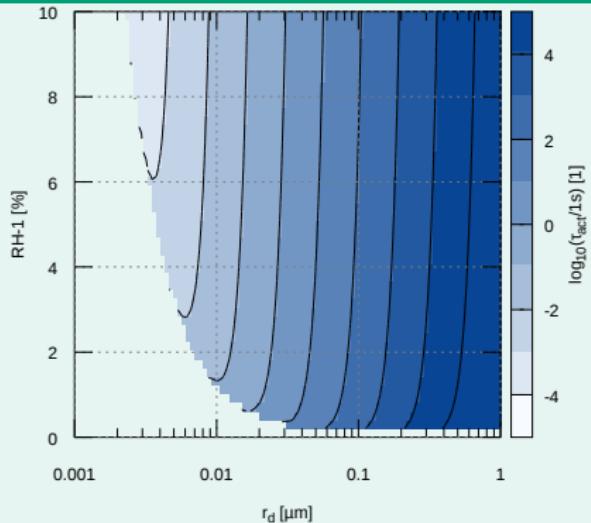
Arabas & Shima 2017



note: axes ranges vs. close-to-equilibrium assumption

activation timescale: analytic vs. numerical

Arabas & Shima 2017



note: axes ranges vs. close-to-equilibrium assumption

Hoffmann, 2016 (MWR)

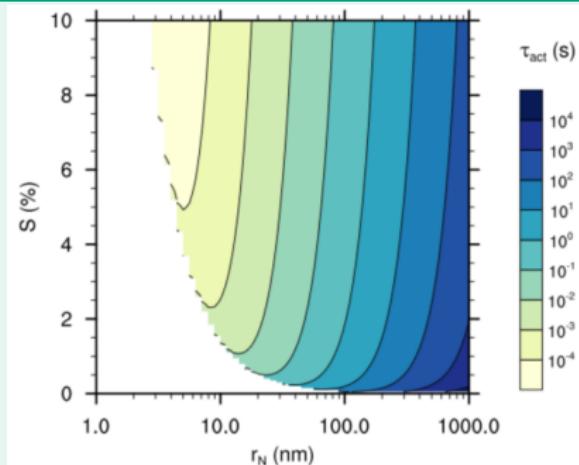
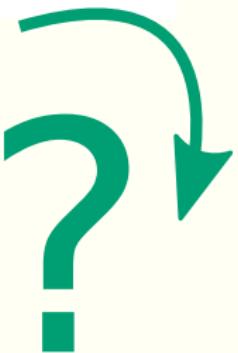
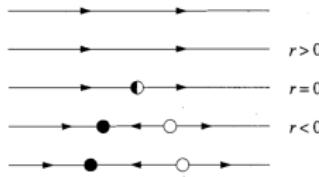
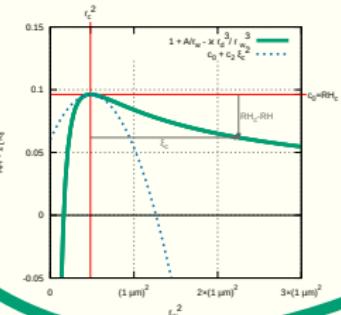
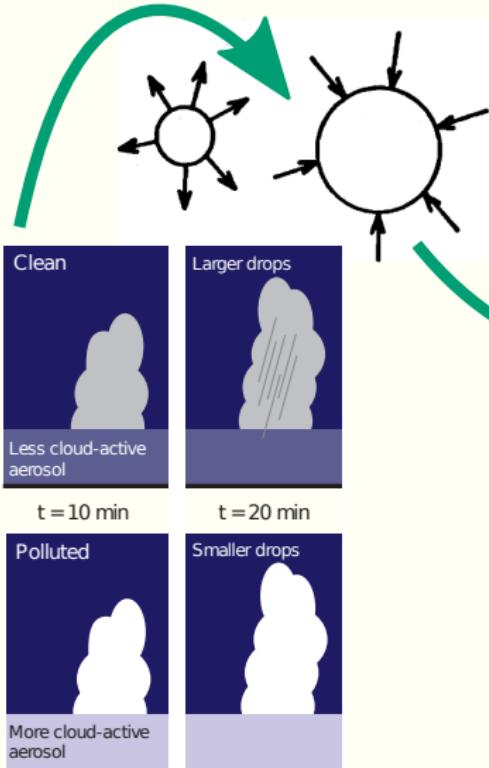


FIG. 2. The activation time scale τ_{act} as a function of dry aerosol radius r_N and supersaturation S . For values of $S < S_{\text{crit}}$ (white areas), τ_{act} does not exist.

$$r \frac{dr}{dt} = \left(S - \frac{A}{r} + \frac{Br_N^3}{r^3} \right) / (F_k + F_D), \quad (10)$$

The second time scale is associated with the activation of particles, for which Köhler theory is essential. This makes an analytic solution for (10) impossible. Numerically calculated values of τ_{act} measuring the time needed for a wetted aerosol to grow beyond its critical radius $r_{\text{crit}} = \sqrt{3Br_N^3/A}$ are given in Fig. 2 as a function of

connecting the dots ...



RH-coupled system & particle concentration as parameter

RH-coupled system & particle concentration as parameter

simple moisture budget (const T,p):

$$\dot{RH} \approx \frac{\dot{\rho}_v}{\rho_{vs}} = -N \underbrace{\frac{4\pi\rho_w}{3\rho_{vs}}}_{\alpha} 3r_w^2 \dot{r}_w$$

RH-coupled system & particle concentration as parameter

simple moisture budget (const T,p):

$$\dot{RH} \approx \frac{\dot{\rho}_v}{\rho_{vs}} = -N \underbrace{\frac{4\pi\rho_w}{3\rho_{vs}}}_{\alpha} 3r_w^2 \dot{r}_w$$

integrating in time:

$$RH = RH_0 - \alpha N r_w^3$$

RH-coupled system & particle concentration as parameter

simple moisture budget (const T,p):

$$\dot{RH} \approx \frac{\dot{\rho_v}}{\rho_{vs}} = -N \underbrace{\frac{4\pi\rho_w}{3\rho_{vs}}}_{\alpha} 3r_w^2 \dot{r}_w$$

integrating in time:

$$RH = RH_0 - \alpha N r_w^3$$

new phase portrait:

$$\dot{\xi} \sim (RH_0 - 1) - \underbrace{\left(\frac{A}{\xi^{\frac{1}{2}}} - \frac{\kappa r_d^3}{\xi^{\frac{3}{2}}} + \alpha N \xi^{\frac{3}{2}} \right)}_f$$

RH-coupled system & particle concentration as parameter

simple moisture budget (const T,p):

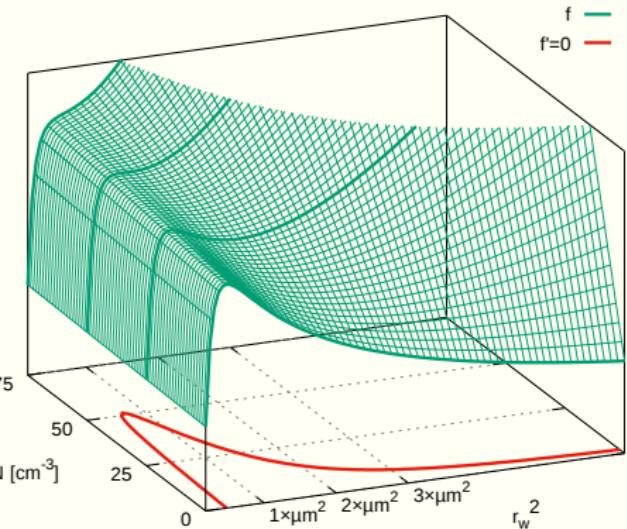
$$\dot{RH} \approx \frac{\dot{\rho_v}}{\rho_{vs}} = -N \underbrace{\frac{4\pi\rho_w}{3\rho_{vs}}}_{\alpha} 3r_w^2 \dot{r}_w$$

integrating in time:

$$RH = RH_0 - \alpha N r_w^3$$

new phase portrait:

$$\dot{\xi} \sim (RH_0 - 1) - \underbrace{\left(\frac{A}{\xi^{\frac{1}{2}}} - \frac{\kappa r_d^3}{\xi^{\frac{3}{2}}} + \alpha N \xi^{\frac{3}{2}} \right)}_f$$



RH-coupled system & particle concentration as parameter

simple moisture budget (const T,p):

$$\dot{RH} \approx \frac{\dot{\rho_v}}{\rho_{vs}} = -N \underbrace{\frac{4\pi\rho_w}{3\rho_{vs}}}_{\alpha} 3r_w^2 \dot{r}_w$$

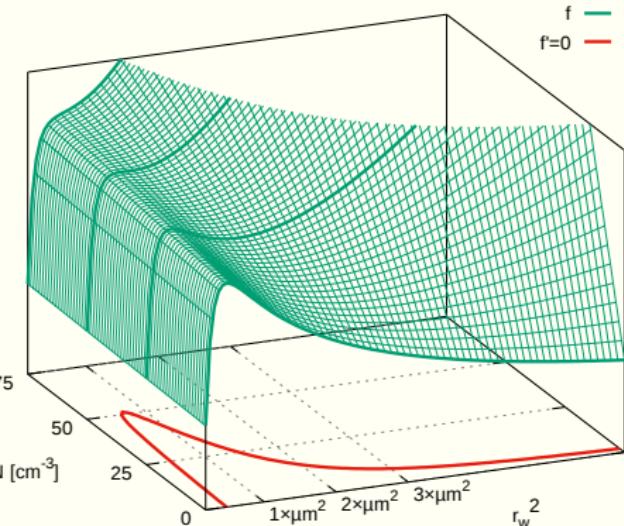
integrating in time:

$$RH = RH_0 - \alpha N r_w^3$$

new phase portrait:

$$\dot{\xi} \sim (RH_0 - 1) - \underbrace{\left(\frac{A}{\xi^{\frac{1}{2}}} - \frac{\kappa r_d^3}{\xi^{\frac{3}{2}}} + \alpha N \xi^{\frac{3}{2}} \right)}_f$$

regime-controlling params: RH, N



RH-coupled system & particle concentration as parameter

simple moisture budget (const T,p):

$$\dot{RH} \approx \frac{\dot{\rho_v}}{\rho_{vs}} = -N \underbrace{\frac{4\pi\rho_w}{3\rho_{vs}}}_{\alpha} 3r_w^2 \dot{r}_w$$

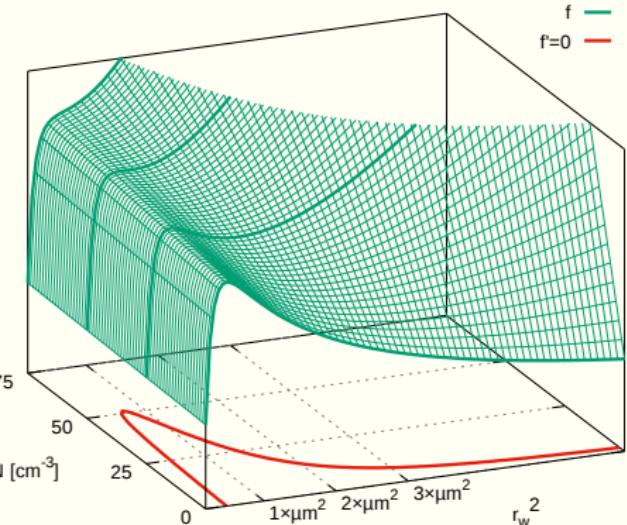
integrating in time:

$$RH = RH_0 - \alpha N r_w^3$$

new phase portrait:

$$\dot{\xi} \sim (RH_0 - 1) - \underbrace{\left(\frac{A}{\xi^{\frac{1}{2}}} - \frac{\kappa r_d^3}{\xi^{\frac{3}{2}}} + \alpha N \xi^{\frac{3}{2}} \right)}_f$$

regime-controlling params: RH, N



$$\text{sgn}(f') = \text{sgn}\left(\kappa r_d^3 - \frac{A}{3} r_w + \alpha N r_w^3\right)$$

bifurcations (and catastrophe) in the RH-coupled system

Prigogine & Stengers 1984

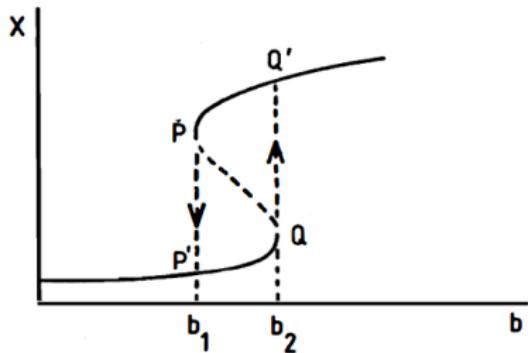


Figure 15. This figure shows how a "hysteresis" phenomenon occurs if we have the value of the bifurcation parameter b first growing and then diminishing. If the system is initially in a stationary state belonging to the lower branch, it will stay there while b grows. But at $b = b_2$, there will be a discontinuity: The system jumps from Q to Q' , on the higher branch. Inversely, starting from a state on the higher branch, the system will remain there till $b = b_1$, when it will jump down to P . Such types of bistable behavior are observed in many fields, such as lasers, chemical reactions or biological membranes.

bifurcations (and catastrophe) in the RH-coupled system

Prigogine & Stengers 1984

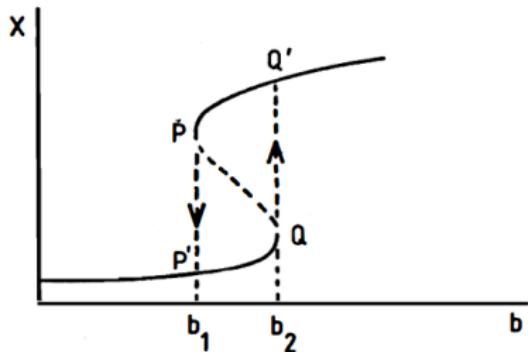
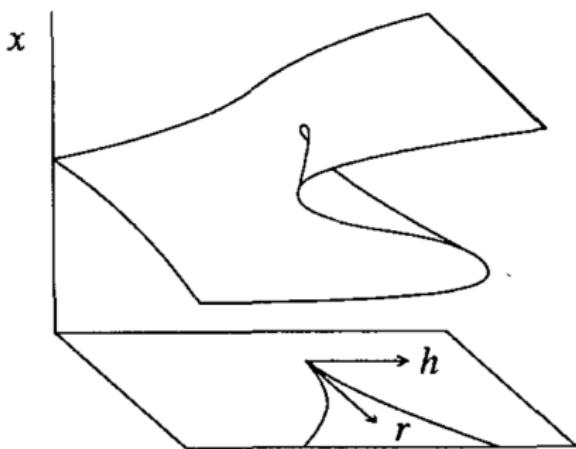


Figure 15. This figure shows how a "hysteresis" phenomenon occurs if we have the value of the bifurcation parameter b first growing and then diminishing. If the system is initially in a stationary state belonging to the lower branch, it will stay there while b grows. But at $b = b_2$, there will be a discontinuity: The system jumps from Q to Q' , on the higher branch. Inversely, starting from a state on the higher branch, the system will remain there till $b = b_1$, when it will jump down to P . Such types of bistable behavior are observed in many fields, such as lasers, chemical reactions or biological membranes.

Strogatz 2014



"cusp catastrophe"

bifurcations (and catastrophe) in the RH-coupled system

Prigogine & Stengers 1984

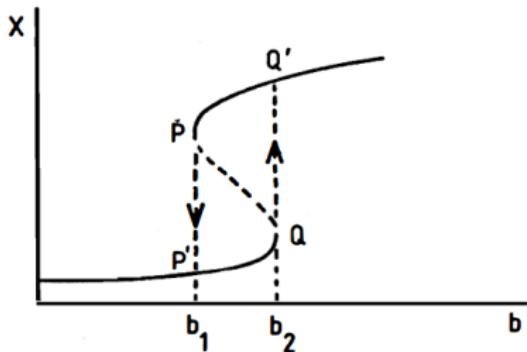
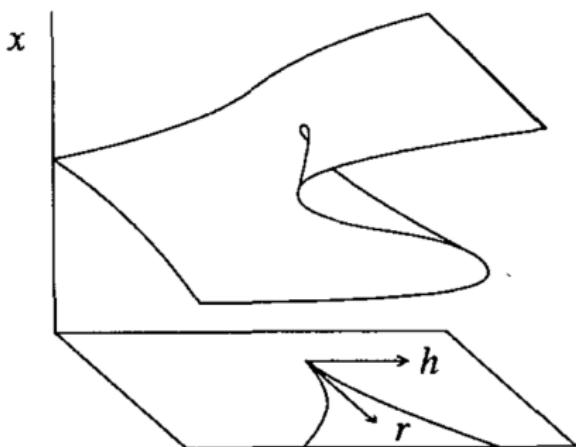


Figure 15. This figure shows how a "hysteresis" phenomenon occurs if we have the value of the bifurcation parameter b first growing and then diminishing. If the system is initially in a stationary state belonging to the lower branch, it will stay there while b grows. But at $b = b_2$, there will be a discontinuity: The system jumps from Q to Q' , on the higher branch. Inversely, starting from a state on the higher branch, the system will remain there till $b = b_1$, when it will jump down to P . Such types of bistable behavior are observed in many fields, such as lasers, chemical reactions or biological membranes.

Strogatz 2014



"cusp catastrophe"

↔ "jumps", hysteretic behaviour (r_w , RH) for small enough N , close to equilibrium (slow process)

hysteresis: activation/deactivation cycle



hysteresis: activation/deactivation cycle



- nomenclature:

hysteresis: activation/deactivation cycle



- nomenclature:
 - CCN activation
 - (heterogeneous) nucleation

hysteresis: activation/deactivation cycle



‣ nomenclature:

- CCN activation
- (heterogeneous) nucleation
- CCN deactivation
- aerosol regeneration / resuspension / recycling
- drop-to-particle conversion
- droplet evaporation

hysteresis: activation/deactivation cycle



➢ nomenclature:

- CCN activation
- (heterogeneous) nucleation
- CCN deactivation
- aerosol regeneration / resuspension / recycling
- drop-to-particle conversion
- droplet evaporation

➢ significance:

hysteresis: activation/deactivation cycle



- nomenclature:
 - CCN activation
 - (heterogeneous) nucleation
 - CCN deactivation
 - aerosol regeneration / resuspension / recycling
 - drop-to-particle conversion
 - droplet evaporation
- significance:
 - aerosol processing by clouds (aqueous chemistry, coalescence)

hysteresis: activation/deactivation cycle



- nomenclature:
 - CCN activation
 - (heterogeneous) nucleation
 - CCN deactivation
 - aerosol regeneration / resuspension / recycling
 - drop-to-particle conversion
 - droplet evaporation
- significance:
 - aerosol processing by clouds (aqueous chemistry, coalescence)
 - spectral broadening (mixing, parcel history, ...)

lifting the constant T-p assumptions: parcel model

lifting the constant T-p assumptions: parcel model

vertically displaced (velocity w , hydrostatic background) adiabatic parcel:
(q : mixing ratio, p_d : bgnd pressure, ρ_d bgnd density, g, l_v, c_{pd} : constants)

$$\begin{bmatrix} \dot{p}_d \\ \dot{T} \\ \dot{r}_w \end{bmatrix} = \begin{bmatrix} -\rho_d g w \\ (\dot{p}_d/\rho_d - \dot{q} l_v)/c_{pd} \\ (D_{\text{eff}}/\rho_w)(\rho_v - \rho_o)/r_w \end{bmatrix}$$

lifting the constant T-p assumptions: parcel model

vertically displaced (velocity w , hydrostatic background) adiabatic parcel:
(q : mixing ratio, p_d : bgnd pressure, ρ_d bgnd density, g, l_v, c_{pd} : constants)

$$\begin{bmatrix} \dot{p}_d \\ \dot{T} \\ \dot{r}_w \end{bmatrix} = \begin{bmatrix} -\rho_d g w \\ (\dot{p}_d/\rho_d - \dot{q} l_v)/c_{pd} \\ (D_{\text{eff}}/\rho_w)(\rho_v - \rho_o)/r_w \end{bmatrix}$$

- $w \rightarrow 0$ (and hence $\dot{p}_d \approx 0$) i.e., slow, close-to-equilibrium evolution of the system relevant to fixed-point analysis (by some means pertinent to formation of non-convective clouds such as fog)

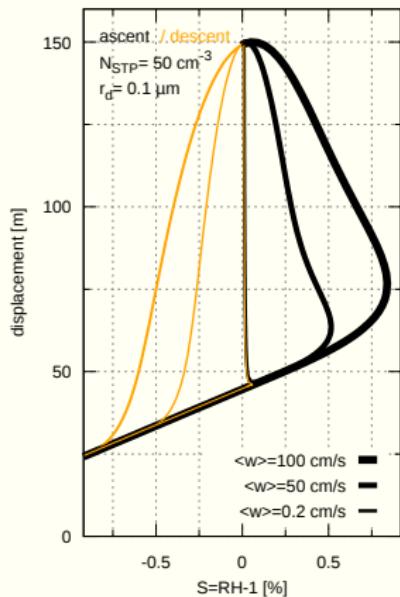
lifting the constant T-p assumptions: parcel model

vertically displaced (velocity w , hydrostatic background) adiabatic parcel:
(q : mixing ratio, p_d : bgnd pressure, ρ_d bgnd density, g, l_v, c_{pd} : constants)

$$\begin{bmatrix} \dot{p}_d \\ \dot{T} \\ \dot{r}_w \end{bmatrix} = \begin{bmatrix} -\rho_d g w \\ (\dot{p}_d/\rho_d - \dot{q} l_v)/c_{pd} \\ (D_{\text{eff}}/\rho_w)(\rho_v - \rho_o)/r_w \end{bmatrix}$$

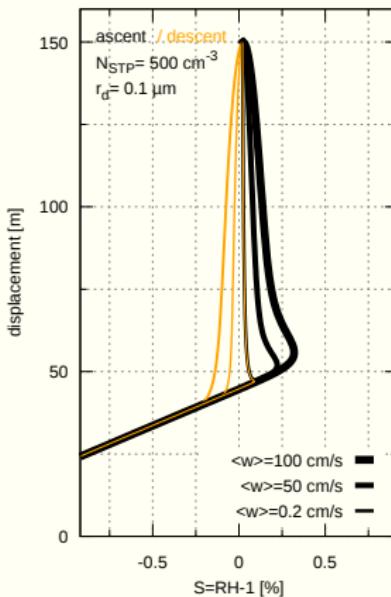
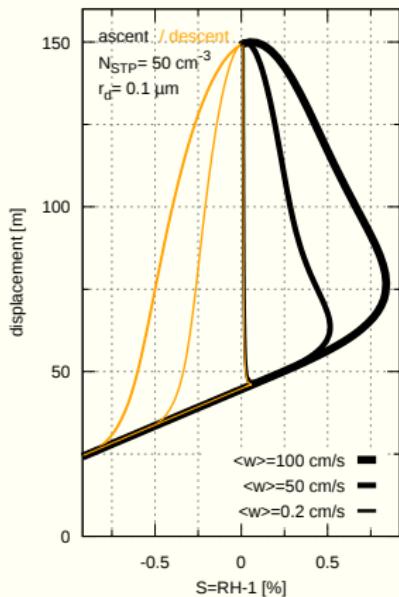
- $w \rightarrow 0$ (and hence $\dot{p}_d \approx 0$) i.e., slow, close-to-equilibrium evolution of the system relevant to fixed-point analysis (by some means pertinent to formation of non-convective clouds such as fog)
- $N \rightarrow 0$ (and hence $\dot{q} \approx 0$) i.e., weak coupling between particle size evolution and ambient thermodynamics (pertinent to the case of low particle concentration).

parcel model: numerical integration (sinusoidal w)



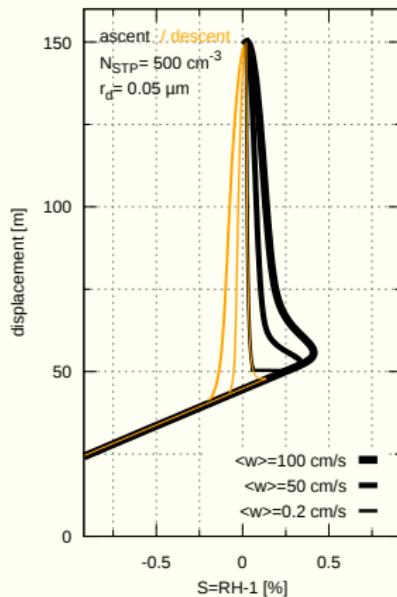
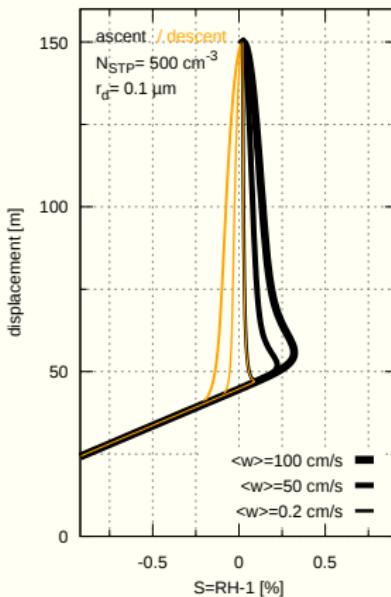
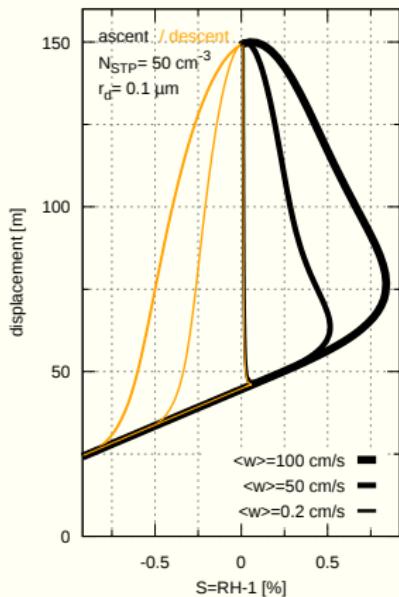
integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

parcel model: numerical integration (sinusoidal w)



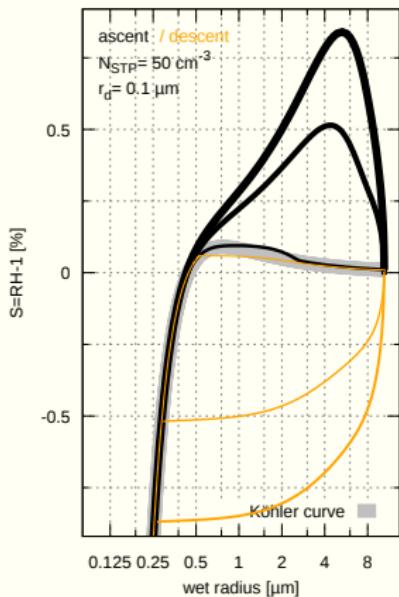
integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

parcel model: numerical integration (sinusoidal w)



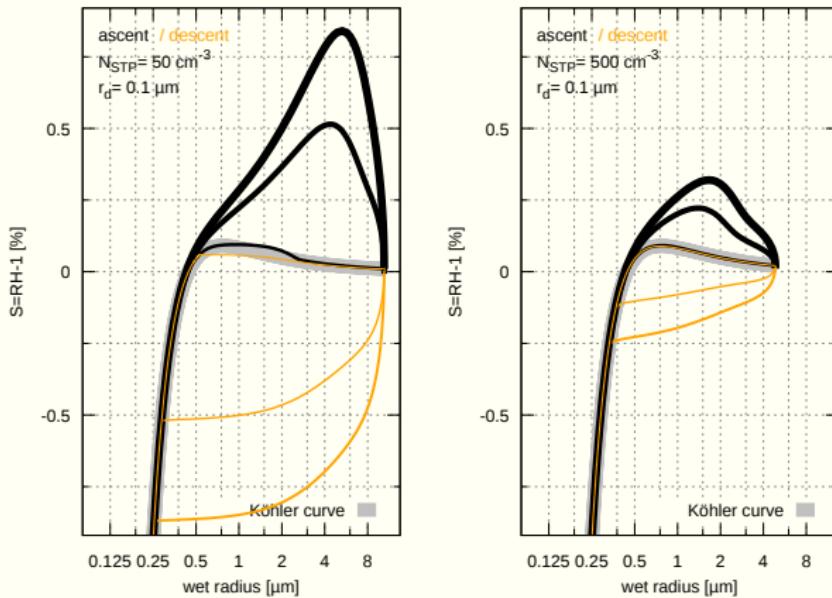
integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

parcel model: numerical integration (sinusoidal w)



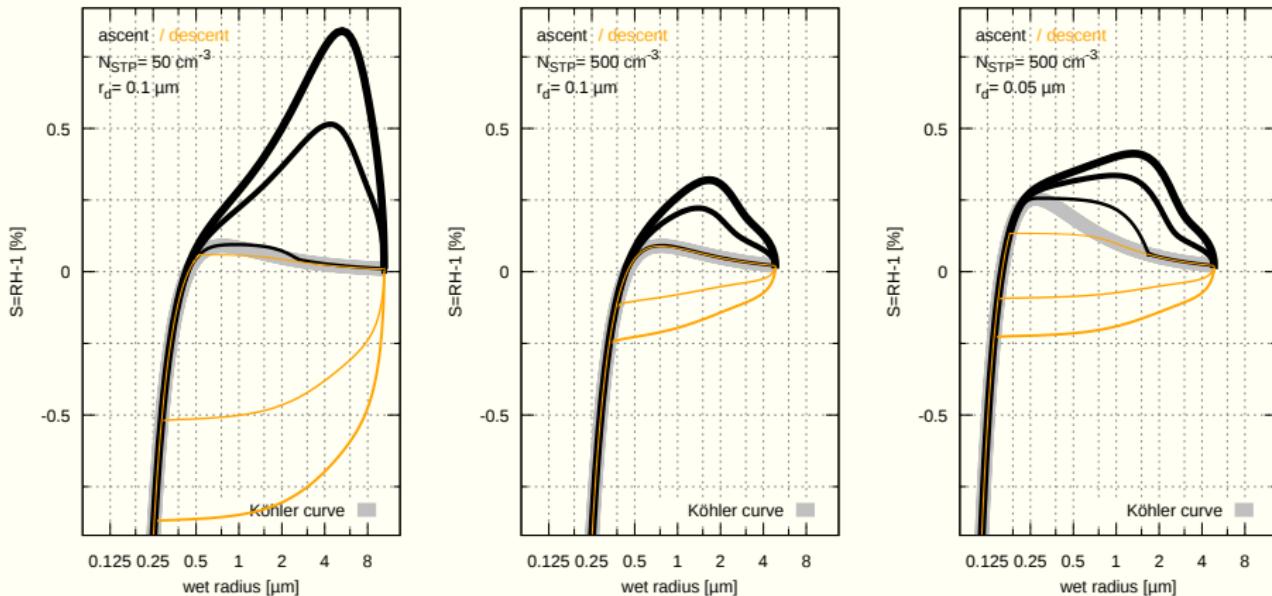
integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

parcel model: numerical integration (sinusoidal w)



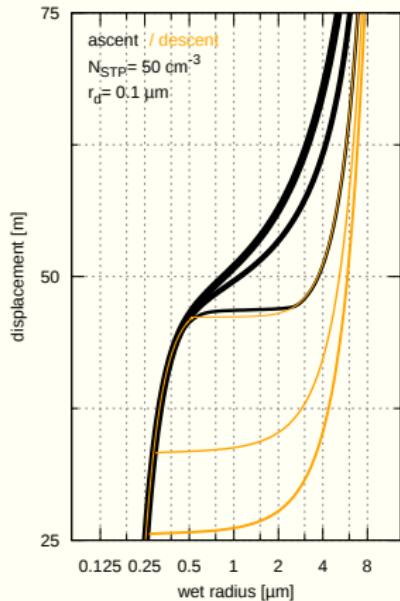
integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

parcel model: numerical integration (sinusoidal w)



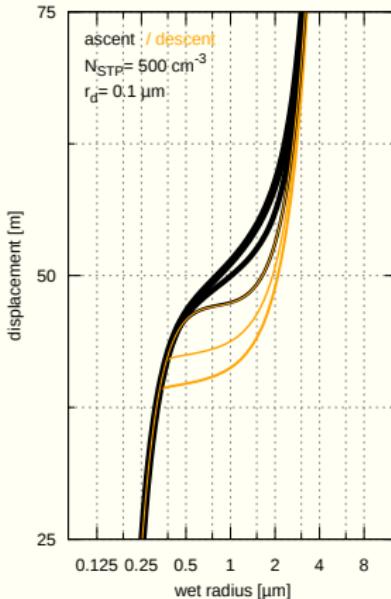
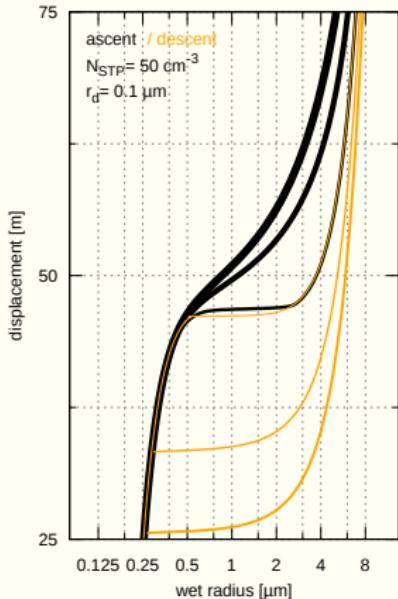
integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

parcel model: numerical integration (sinusoidal w)



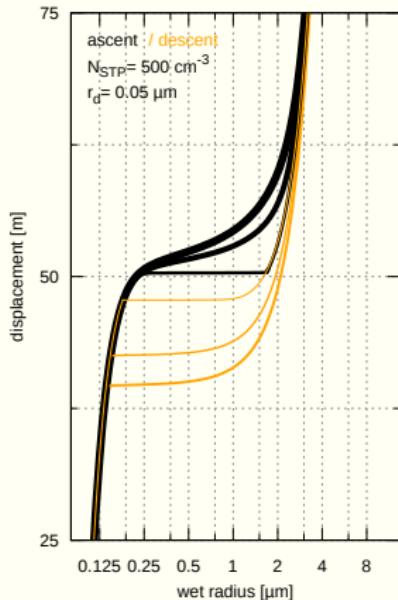
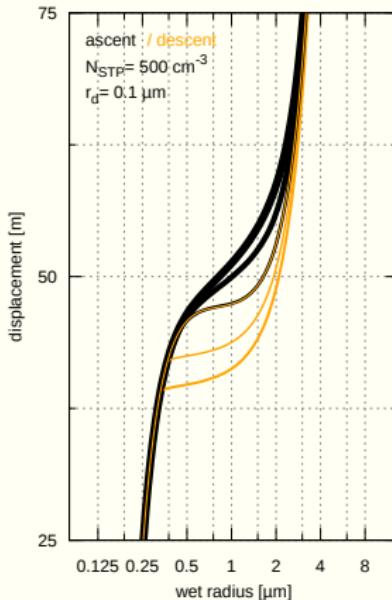
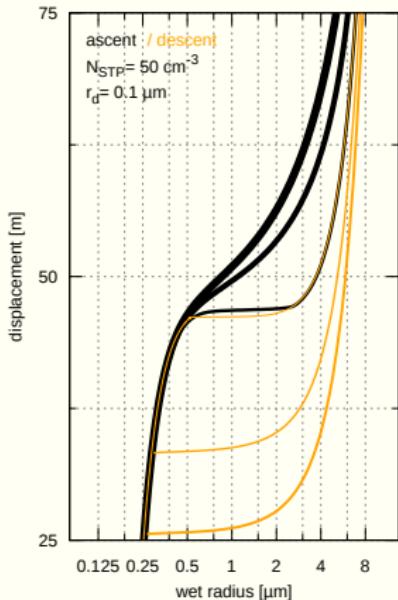
integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

parcel model: numerical integration (sinusoidal w)



integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

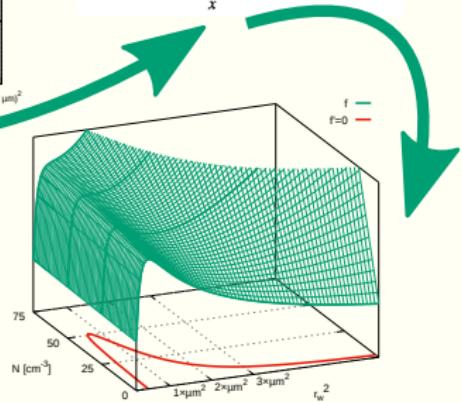
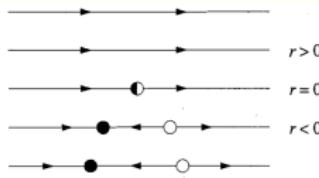
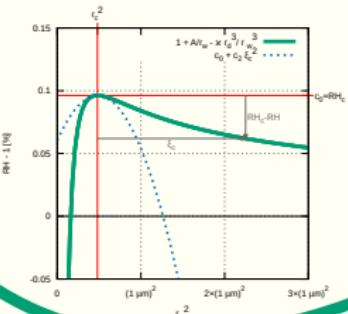
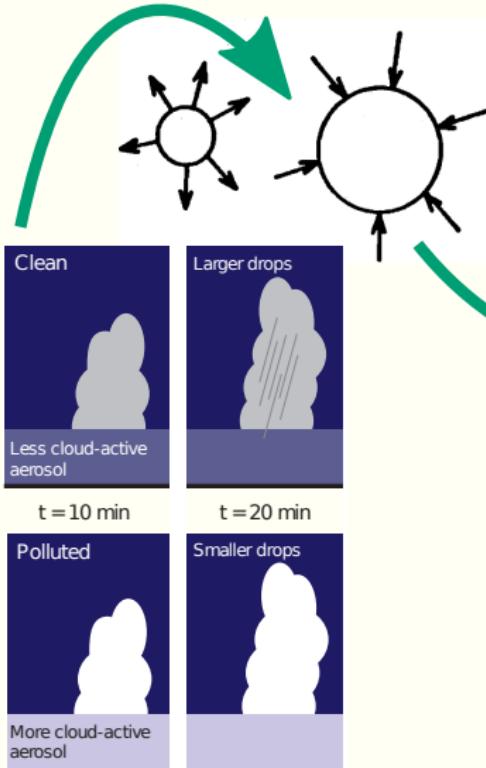
parcel model: numerical integration (sinusoidal w)



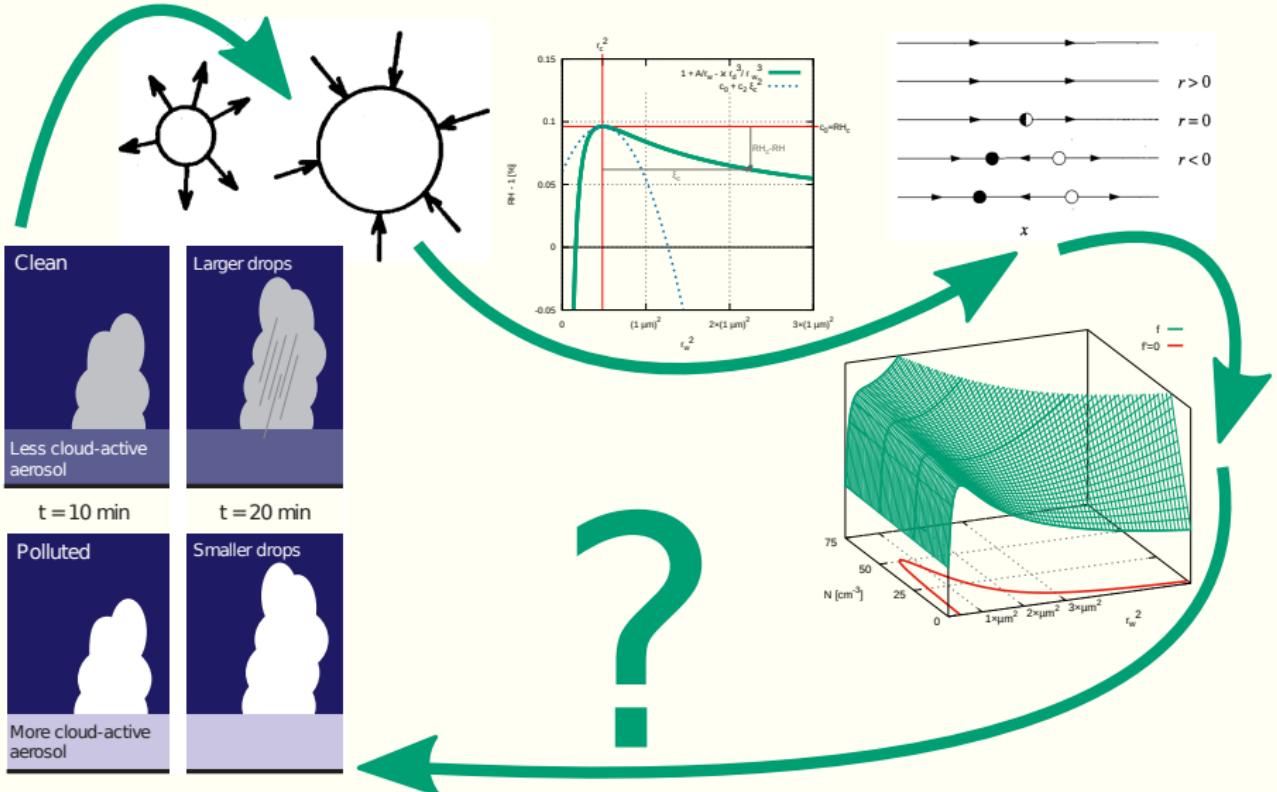
integration using CVODE adaptive solver
open source code (based on libcloudph++) as electronic paper supplement

connecting the dots ...

connecting the dots ...



connecting the dots ...



monodisperse system: limitations and applicability

limitations:

limitations:

- ❖ no spectral width representation
(key for modelling precipitation onset)

limitations:

- ☒ no spectral width representation
(key for modelling precipitation onset)
- ☒ no activated/unactivated partitioning
(documented excitable behaviour)

limitations:

- ☒ no spectral width representation
(key for modelling precipitation onset)
- ☒ no activated/unactivated partitioning
(documented excitable behaviour)
- ☒ no droplet-to-droplet vapour transfer
(Ostwald ripening = Ouzo effect)

limitations:

- ☒ no spectral width representation
(key for modelling precipitation onset)
- ☒ no activated/unactivated partitioning
(documented excitable behaviour)
- ☒ no droplet-to-droplet vapour transfer
(Ostwald ripening = Ouzo effect)

... applicability?

limitations:

- ☒ no spectral width representation
(key for modelling precipitation onset)
- ☒ no activated/unactivated partitioning
(documented excitable behaviour)
- ☒ no droplet-to-droplet vapour transfer
(Ostwald ripening = Ouzo effect)

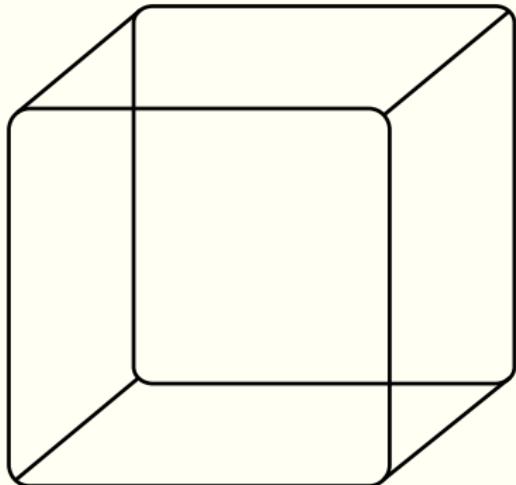
... applicability?

particle-based μ -physics schemes for LES!
(Lagrangian Cloud Models / Super-Droplet Models)

particle-based μ -physics for LES

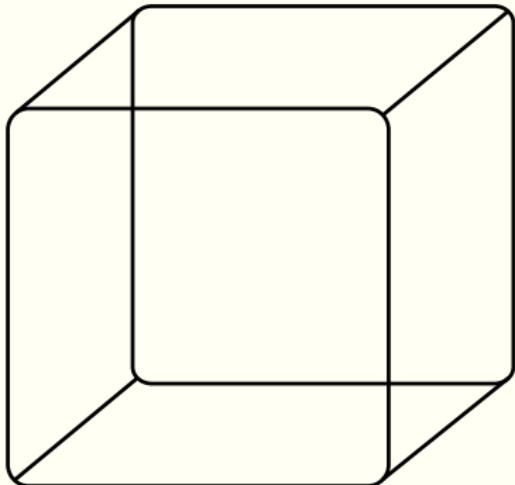
- “information carriers” in LES domain

particle-based μ -physics for LES



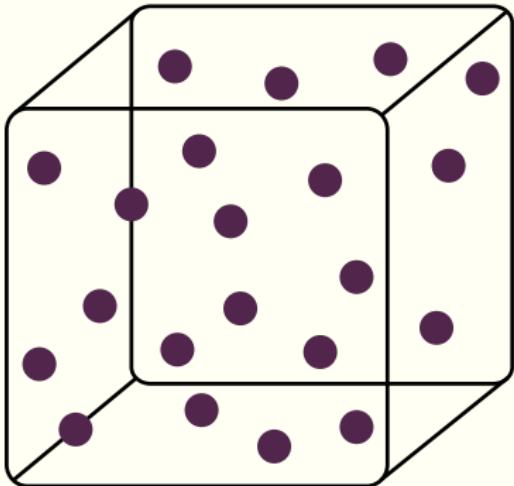
- “information carriers” in LES domain
- ab-initio approach:
particle=aerosol/cloud/rain

particle-based μ -physics for LES



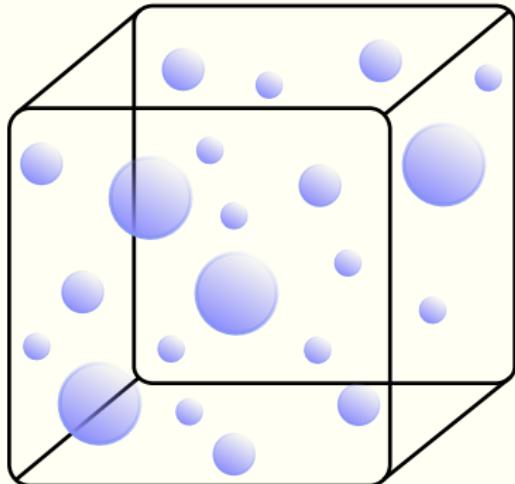
- “information carriers” in LES domain
- ab-initio approach:
particle=aerosol/cloud/rain
- attributes:

particle-based μ -physics for LES



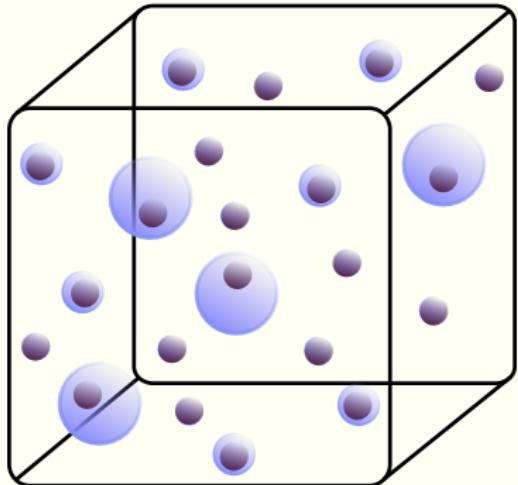
- “information carriers” in LES domain
- ab-initio approach:
particle=aerosol/cloud/rain
- attributes:
 - spatial coordinates

particle-based μ -physics for LES



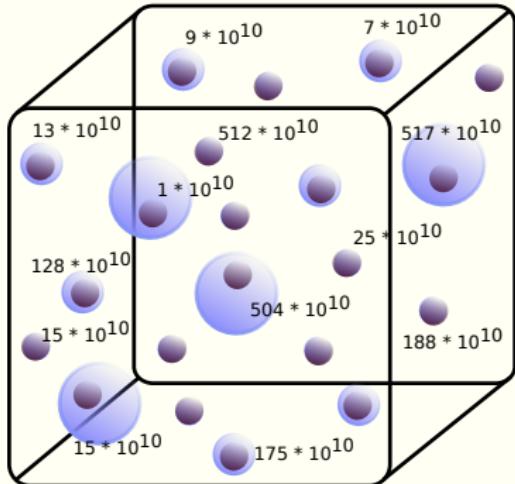
- “information carriers” in LES domain
- ab-initio approach:
particle=aerosol/cloud/rain
- attributes:
 - spatial coordinates
 - wet radius

particle-based μ -physics for LES



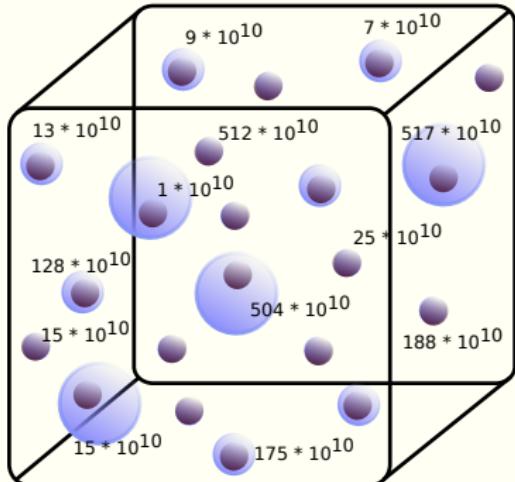
- “information carriers” in LES domain
- ab-initio approach:
particle=aerosol/cloud/rain
- attributes:
 - spatial coordinates
 - wet radius
 - dry radius

particle-based μ -physics for LES



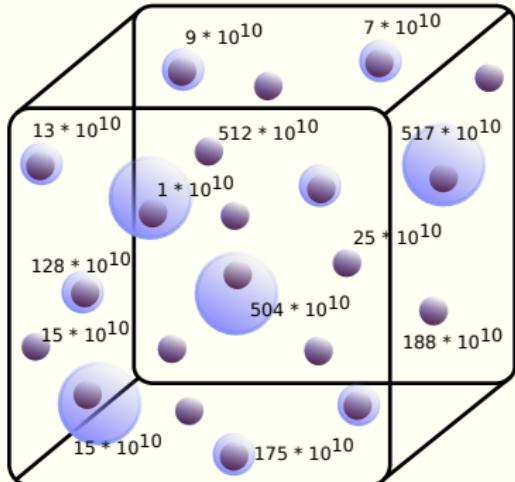
- “information carriers” in LES domain
- ab-initio approach:
particle=aerosol/cloud/rain
- attributes:
 - spatial coordinates
 - wet radius
 - dry radius
 - multiplicity

particle-based μ -physics for LES



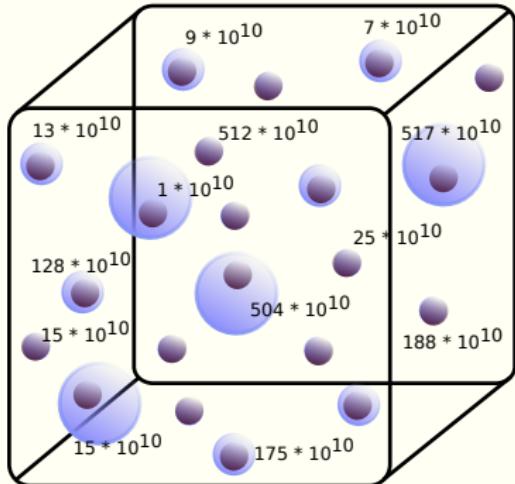
- “information carriers” in LES domain
- ab-initio approach:
particle=aerosol/cloud/rain
- attributes:
 - spatial coordinates
 - wet radius
 - dry radius
 - multiplicity
 - . . .

particle-based μ -physics for LES



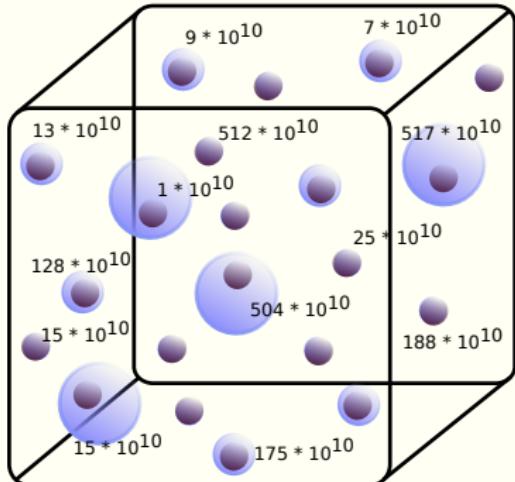
- “information carriers” in LES domain
 - ab-initio approach:
particle=aerosol/cloud/rain
 - attributes:
 - spatial coordinates
 - wet radius
 - dry radius
 - multiplicity
 - ...
 - chemistry
- ~ Jaruga & Pawlowska 2018

particle-based μ -physics for LES



- ☒ “information carriers” in LES domain
- ☒ ab-initio approach:
particle=aerosol/cloud/rain
- ☒ attributes:
 - ☒ spatial coordinates
 - ☒ wet radius
 - ☒ dry radius
 - ☒ multiplicity
 - ☒ ...
 - ☒ chemistry
- ~~ Jaruga & Pawlowska 2018
- ☒ each particle: **monodisperse!**

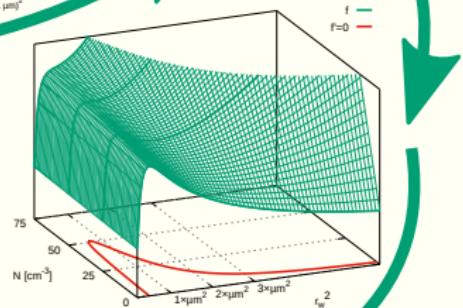
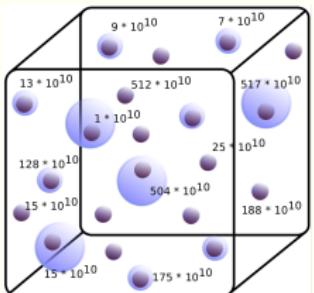
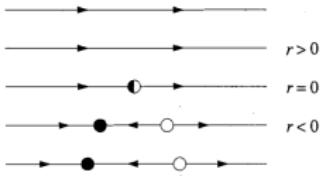
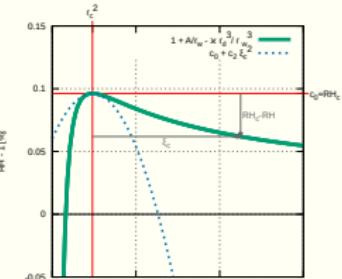
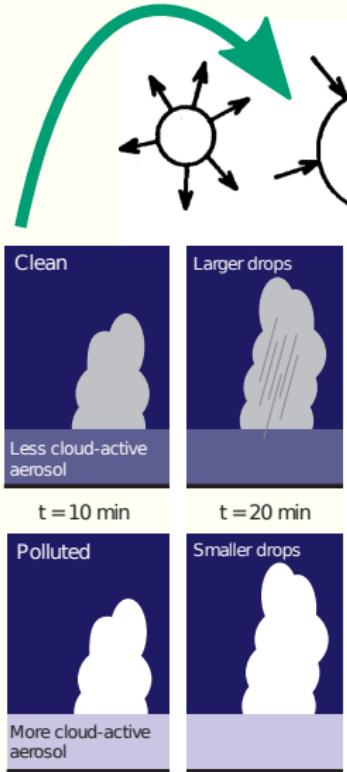
particle-based μ -physics for LES



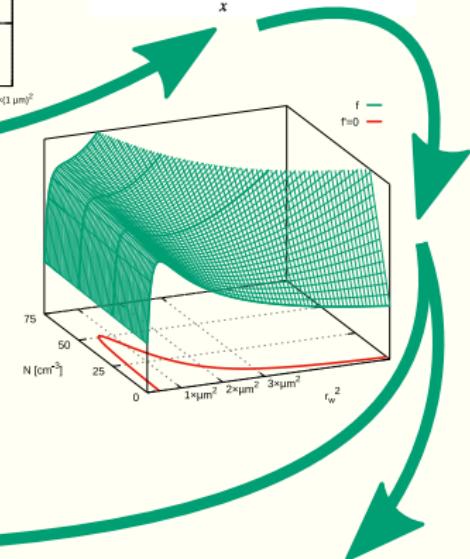
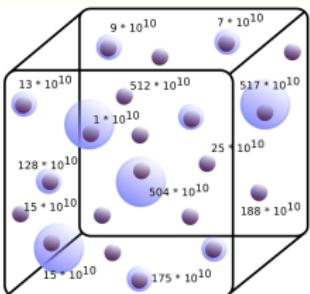
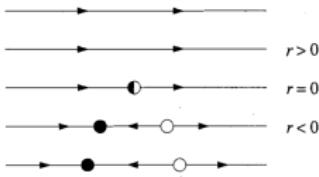
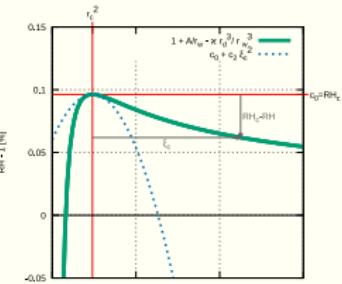
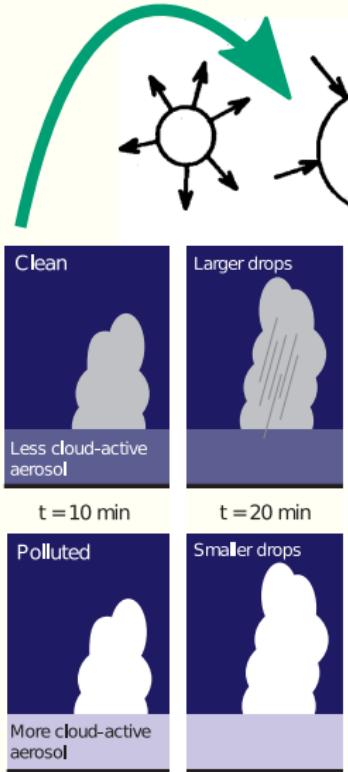
- “information carriers” in LES domain
- ab-initio approach:
particle=aerosol/cloud/rain
- attributes:
 - spatial coordinates
 - wet radius
 - dry radius
 - multiplicity
 - ...
 - **chemistry**
- ~~ Jaruga & Pawlowska 2018
- each particle: **monodisperse!**
- each timestep: **constant RH!**

connecting the dots ...

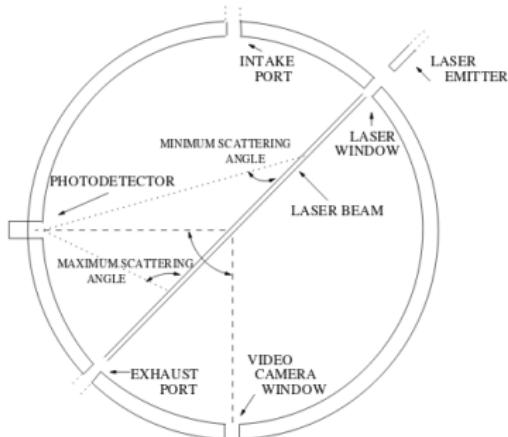
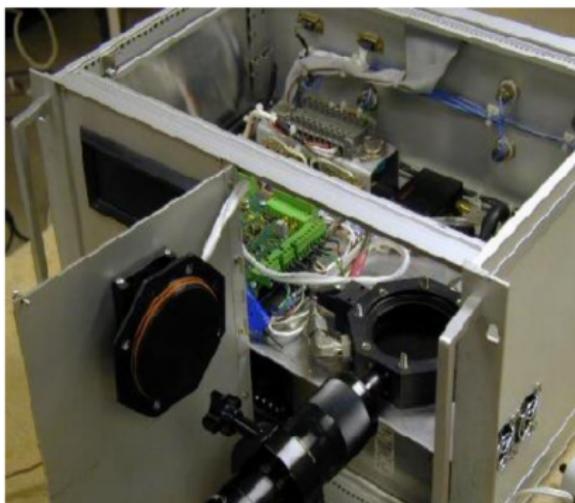
connecting the dots ...



connecting the dots ...



model applicability: CCN instruments? (hypothesis...)

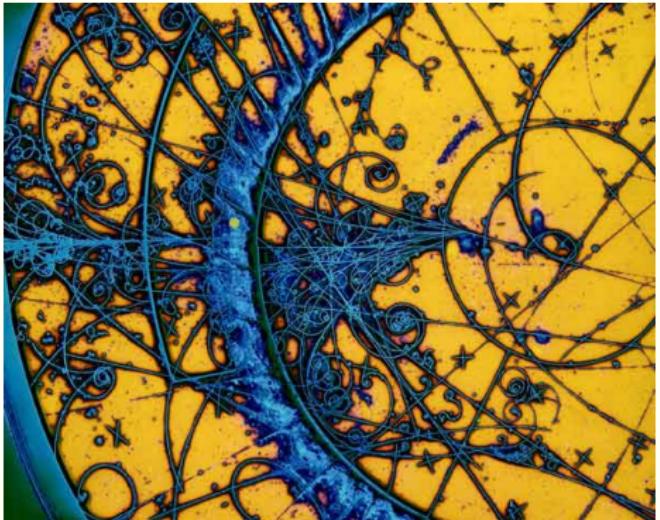
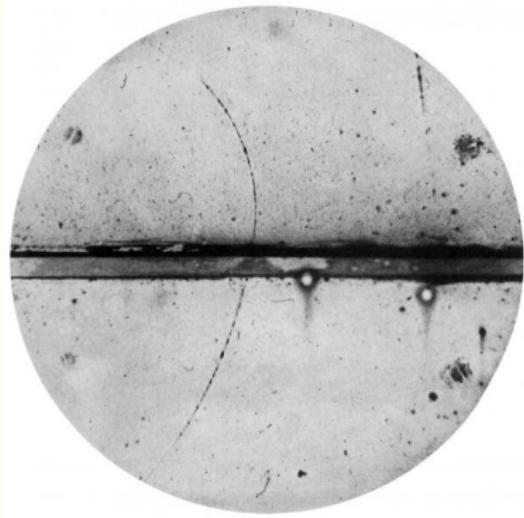


pictured: UWyoming WyoCCN instrument

(photo from DYCOMS-II CCN data report by Jeff Snider et al.)

https://www.eol.ucar.edu/projects/dycoms/dm/archive/docs/snider_ccnreadme.pdf

applicability beyond cloud physics (hypothesis...)



Wilson & bubble chambers

<https://home.cern/about/updates/2015/06/seeing-invisible-event-displays-particle-physics>

conclusions, takeaways, prospects

conclusions, takeaways, prospects

- ❖ CCN (de)activation as a bifurcating dynamical system

conclusions, takeaways, prospects

- CCN (de)activation as a bifurcating dynamical system
- flipped Köhler curve as the phase portrait of the system

conclusions, takeaways, prospects

- CCN (de)activation as a bifurcating dynamical system
- flipped Köhler curve as the phase portrait of the system
- N-constrained vapour budget coupling \rightsquigarrow “non-Köhler” dynamics (two stable equilibria, hysteresis, cusp, no activation for big N)

conclusions, takeaways, prospects

- CCN (de)activation as a bifurcating dynamical system
- flipped Köhler curve as the phase portrait of the system
- N-constrained vapour budget coupling \rightsquigarrow “non-Köhler” dynamics (two stable equilibria, hysteresis, cusp, no activation for big N)
- analytical results: conditions for hysteretic behaviour, timescales

conclusions, takeaways, prospects

- CCN (de)activation as a bifurcating dynamical system
- flipped Köhler curve as the phase portrait of the system
- N-constrained vapour budget coupling \rightsquigarrow “non-Köhler” dynamics (two stable equilibria, hysteresis, cusp, no activation for big N)
- analytical results: conditions for hysteretic behaviour, timescales
- guidance for numerical scheme design (particle-based μ -physics)

conclusions, takeaways, prospects

- ☒ CCN (de)activation as a bifurcating dynamical system
- ☒ flipped Köhler curve as the phase portrait of the system
- ☒ N-constrained vapour budget coupling \rightsquigarrow “non-Köhler” dynamics (two stable equilibria, hysteresis, cusp, no activation for big N)
- ☒ analytical results: conditions for hysteretic behaviour, timescales
- ☒ guidance for numerical scheme design (particle-based μ -physics)
- ☒ extensions: response to fluctuations, bi-/poly-disperse spectra, ...

conclusions, takeaways, prospects

- ☒ CCN (de)activation as a bifurcating dynamical system
- ☒ flipped Köhler curve as the phase portrait of the system
- ☒ N-constrained vapour budget coupling \rightsquigarrow “non-Köhler” dynamics (two stable equilibria, hysteresis, cusp, no activation for big N)
- ☒ analytical results: conditions for hysteretic behaviour, timescales
- ☒ guidance for numerical scheme design (particle-based μ -physics)
- ☒ extensions: response to fluctuations, bi-/poly-disperse spectra, ...
- ☒ applications: CCN instrumentation modelling, non-cloud appl...

last slide

last slide

Big thanks for the Foundation for Polish Science!

Big thanks for the Foundation for Polish Science!

Thank you for your attention!

<https://doi.org/10.5194/npg-24-535-2017>