

# **Super-droplet $\mu$ -physics: a probabilistic look at modelling clouds and the immersion freezing process**

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**Sylwester Arabas**

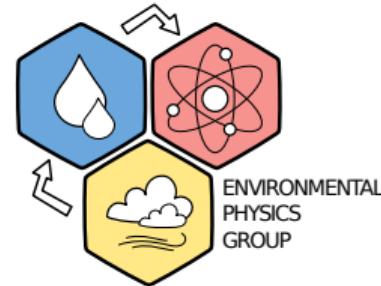
AGH University of Krakow, Poland

July 3 2025 (Institute for Physics of the Atmosphere @uni-mainz.de)



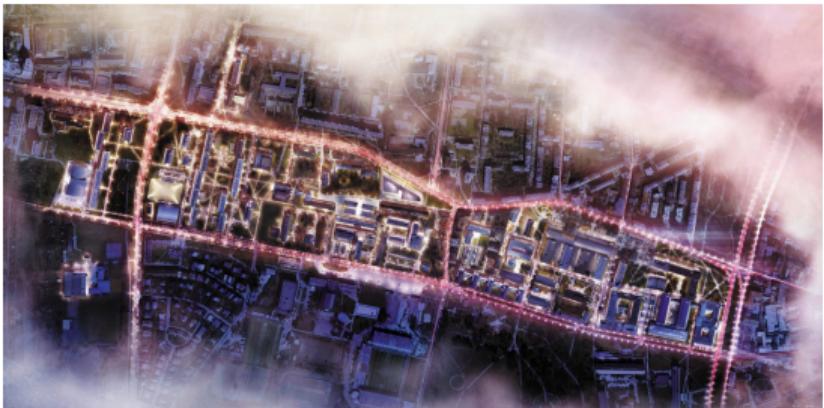
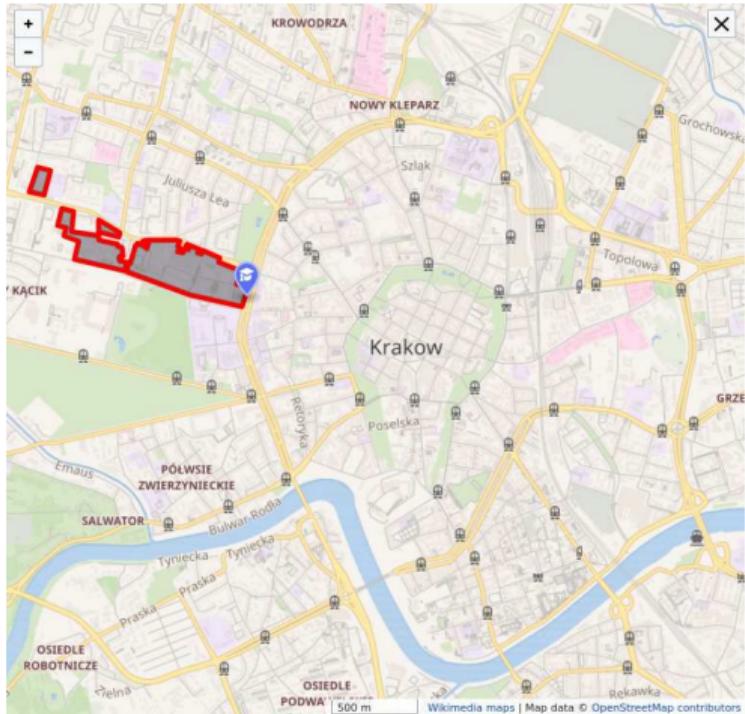


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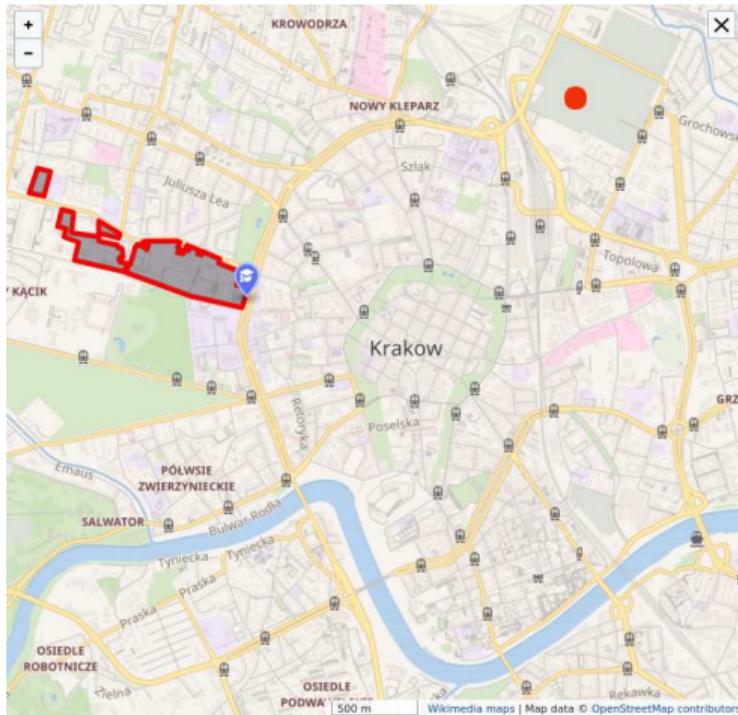


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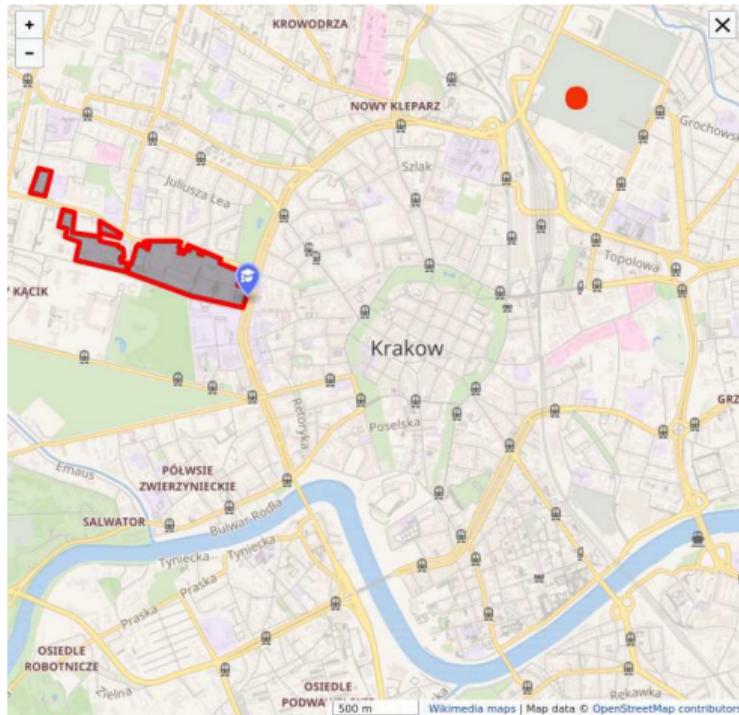
# AGH campus in Kraków and the "neighbourhood" ..



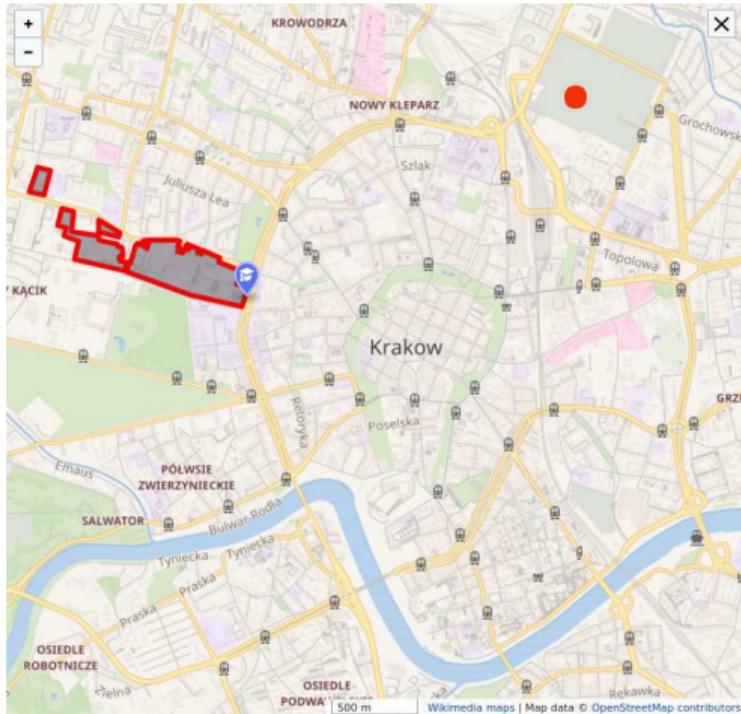
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**cloud  $\mu$ -physics models in 1910s?**

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# Smoluchowski & Rudzki

## Versuch einer mathematischen Theorie der Koagulationskinetik kolloider Lösungen.

Von

M. v. Smoluchowski.

(Mit 3 Figuren im Text.)

(Eingegangen am 8. 9. 16.)

### I. Einleitung.

So sehr auch bis heute die Literatur über Koagulation kolloider Lösungen angewachsen ist, sind doch unsere Kenntnisse betreffs des quantitativen Verlaufs, sowie betreffs des Mechanismus des Koagulationsprozesses äußerst mangelhaft. Die meisten Forscher begnügen sich mit qualitativen Beobachtungen oder stellen ihre Messungsreihen in Tabellen oder Kurvenform<sup>1)</sup> dar, da die mathematische Wiedergabe derselben auf aussergewöhnliche Schwierigkeiten stösst.

In den interessanten Arbeiten<sup>2)</sup> von S. Miyazawa, N. Ishizaka, H. Freundlich, J. A. Gann wird allerdings eine formelmässige Zusammenfassung des empirischen Versuchsmaterials, sowie eine Aufklärung desselben nach Analogie mit den Gesetzen der chemischen Kinetik angestrebt. Aber klare Gesetzmässigkeiten haben sich bisher auf diese Weise nicht ergeben, und wurden sogar gewisse, anfangs aufgestellte Gesetzformeln (Paine, Freundlich und Ishizaka) bei exakterer Prüfung (Freundlich und Gann) als unhaltbar zurückgenommen<sup>3)</sup>.

Die Erfolglosigkeit der bisherigen Versuche, auf dem empirisch-induktiven Wege zu einem Verständnis der hier geltenden Gesetze zu gelangen, kann man nur als einen Grund auffassen, einmal den um-

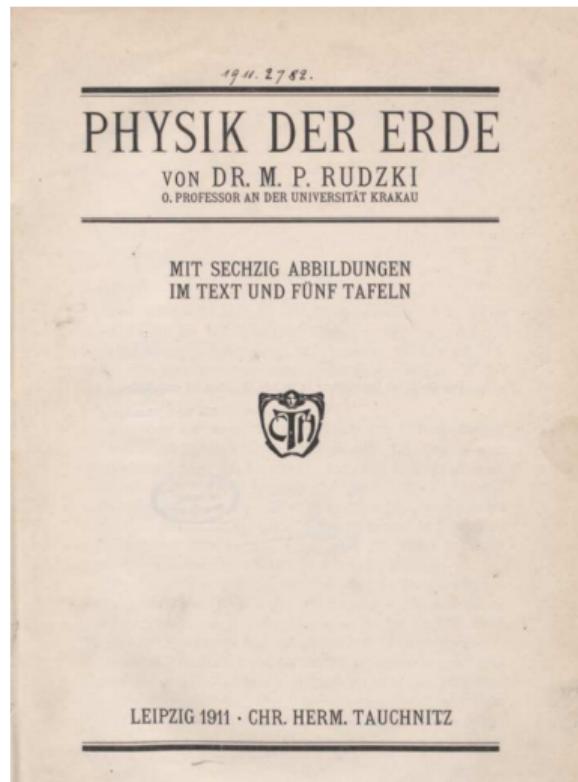
<sup>1)</sup> Vgl. z. B.: A. Galecki, Zeitschr. f. anorg. Chemie **74**, 174 (1912); Kolloid-Zeitschr. **10**, 169 (1912); A. Lottermoser, Kolloid-Zeitschr. **15**, 145 (1914); H. H. Paine, Kolloidchem. Beihefte **4**, 24 (1912); Kolloid-Zeitschr. **11**, 115 (1912).

<sup>2)</sup> S. Miyazawa, Journ. Chem. Soc. Tokio **83**, 1179, 1210 (1912); N. Ishizaka, Zeitschr. f. physik. Chemie **83**, 97 (1913); H. Freundlich u. N. Ishizaka, ebendort **85**, 393 (1913); Kolloid-Zeitschr. **12**, 230 (1913); J. Gann, Kolloidchem. Beihefte **8**, 64 (1916).

<sup>3)</sup> Siehe Abschnitt VI.  
Zeitschr. f. physik. Chemie. XCII.

(Smoluchowski 1918,

DOI:10.1515/zpch-1918-9209)



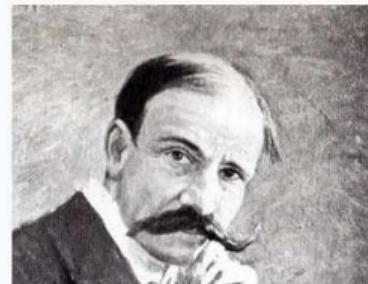
(available at U. Mainz Zentralbibliothek)

# Maurycy Pius Rudzki

From Wikipedia, the free encyclopedia

**Maurycy Pius Rudzki** (b. 1862, d. 1916) was the first person to call himself a professor of geophysics. He held the Chair of Geophysics at the Jagiellonian University in Kraków, and established the Institute of Geophysics there in 1895. His research specialty was elastic anisotropy, as applied to wave propagation in the earth, and he established many of the fundamental results in that arena. <sup>[1]</sup>

**Maurycy Pius Rudzki**

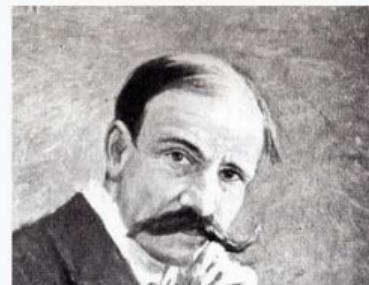


# Maurycy Pius Rudzki

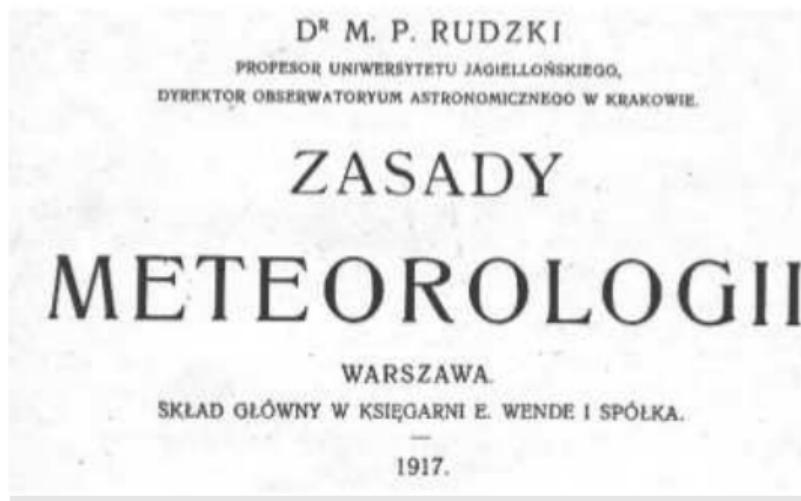
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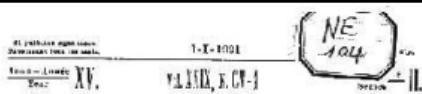
**Maurycy Pius Rudzki**



## “Principles of Meteorology” book (1917)



# Rudzki 1917: Principles of Meteorology



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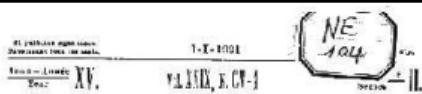
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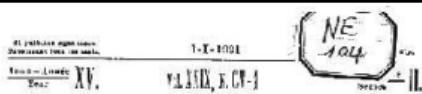
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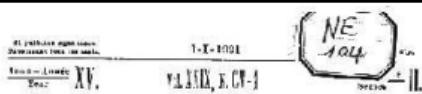
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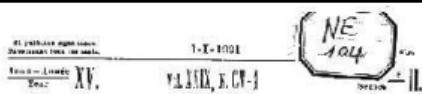
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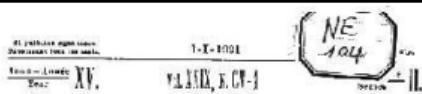
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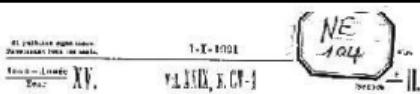
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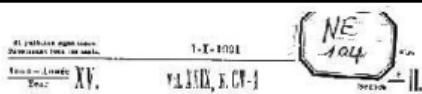
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## **modelling coagulation: SCE & SDM**

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## Smoluchowski's coagulation equation (SCE)

droplet concentration:  $\textcolor{red}{c}(x, t) : \mathbb{R}^+ \times \mathbb{R}^+ \rightarrow \mathbb{R}^+$

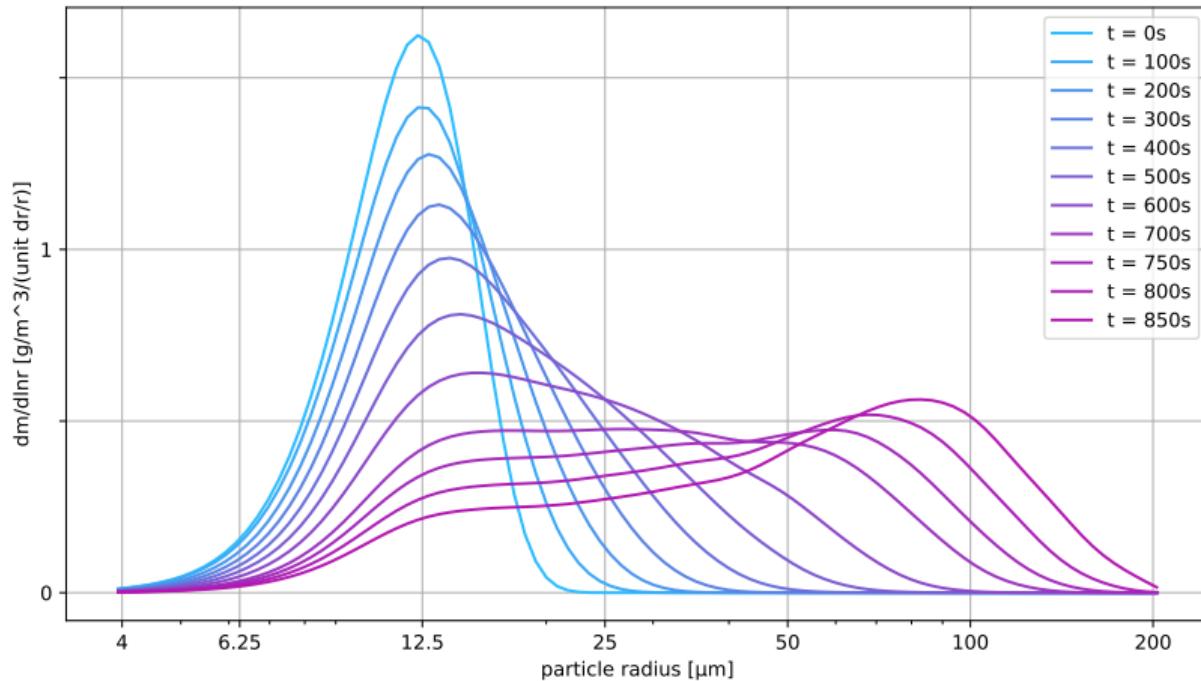
collision kernel:  $\textcolor{blue}{a}(x_1, x_2) : \mathbb{R}^+ \times \mathbb{R}^+ \rightarrow \mathbb{R}^+$

$$\dot{\textcolor{red}{c}}(x) = \frac{1}{2} \int_0^x \textcolor{blue}{a}(y, x-y) \textcolor{red}{c}(y) \textcolor{red}{c}(x-y) dy - \int_0^\infty \textcolor{blue}{a}(y, x) \textcolor{red}{c}(y) \textcolor{red}{c}(x) dy$$

droplet concentration:  $\textcolor{red}{c}_i = \textcolor{red}{c}(x_i)$

$$\dot{\textcolor{red}{c}}_i = \frac{1}{2} \sum_{k=1}^{i-1} \textcolor{blue}{a}(x_k, x_{i-k}) \textcolor{red}{c}_k \textcolor{red}{c}_{i-k} - \sum_{k=1}^\infty \textcolor{blue}{a}(x_k, x_i) \textcolor{red}{c}_k \textcolor{red}{c}_i$$

# cloud droplet collisional growth





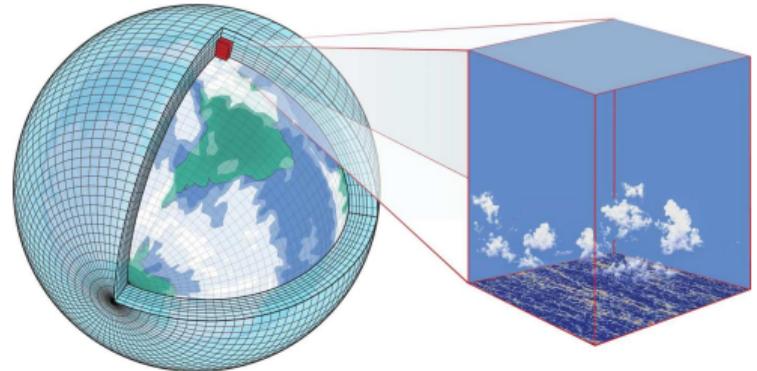
"Cloud and ship. Ukraine, Crimea, Black sea, view  
from Ai-Petri mountain"

(photo: Yevgen Timashov / National Geographic)



"Cloud and ship. Ukraine, Crimea, Black sea, view from Ai-Petri mountain"

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"Grid cells in a global climate model and a large-eddy simulation of shallow cumulus clouds at 5 m resolution"

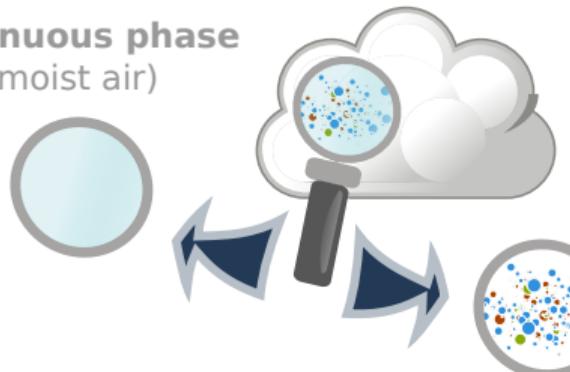
(fig. from Schneider et al. 2017)

## Eulerian vs. Lagrangian microphysics



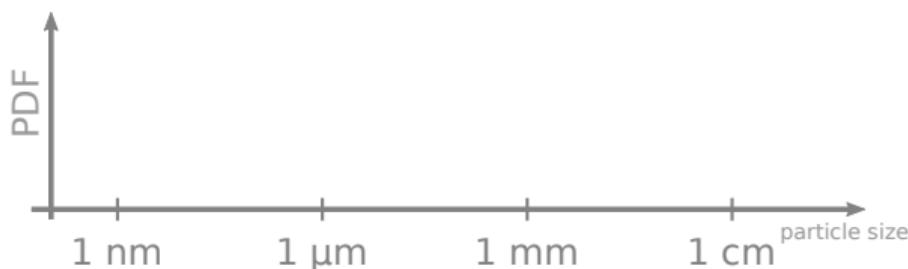
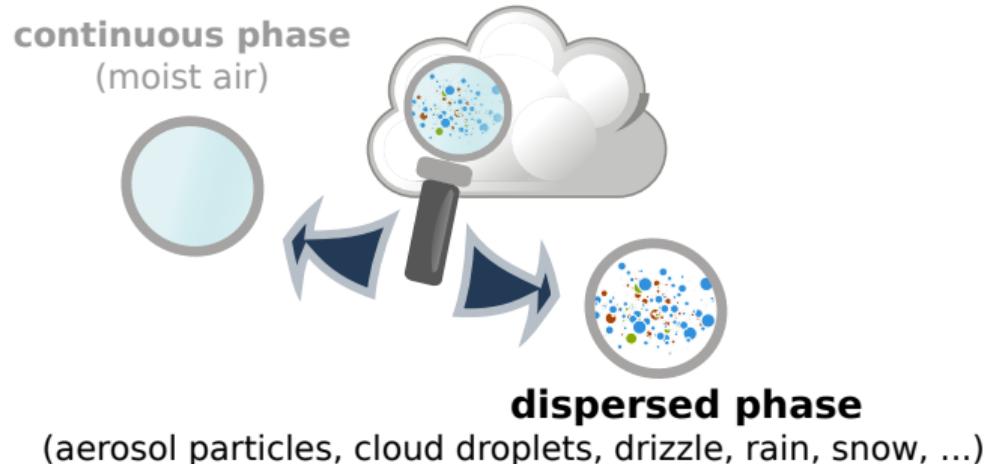
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**continuous phase**  
(moist air)



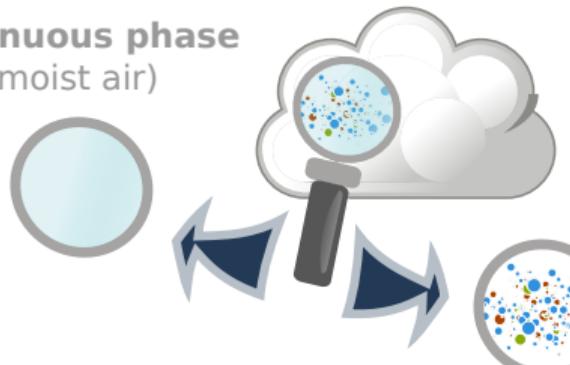
**dispersed phase**  
(aerosol particles, cloud droplets, drizzle, rain, snow, ...)

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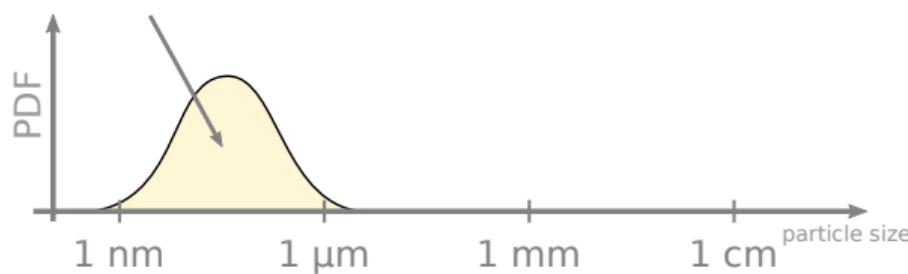


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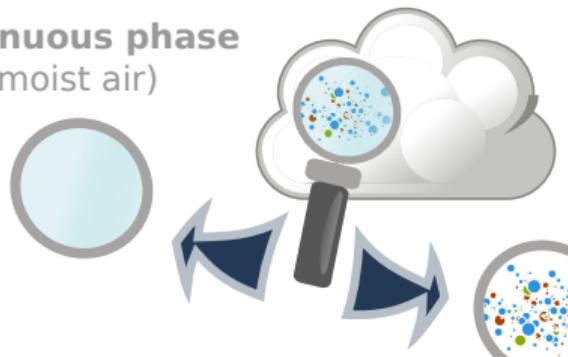


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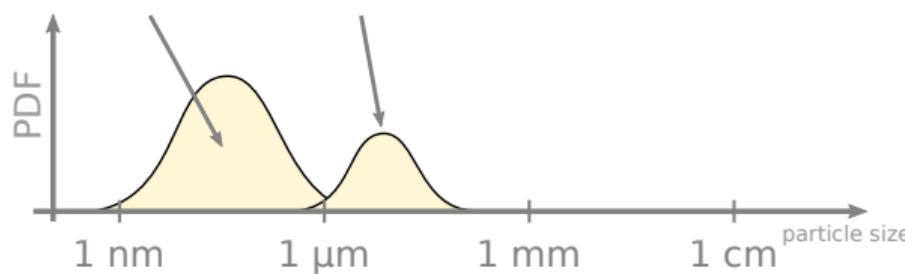


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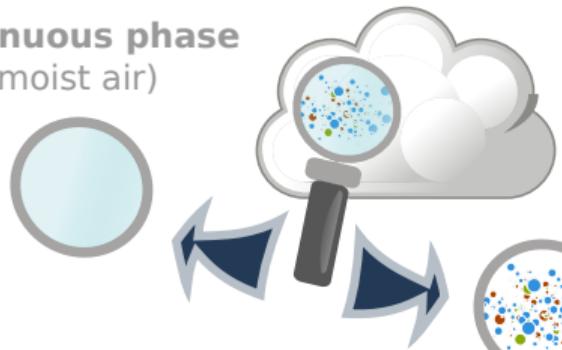


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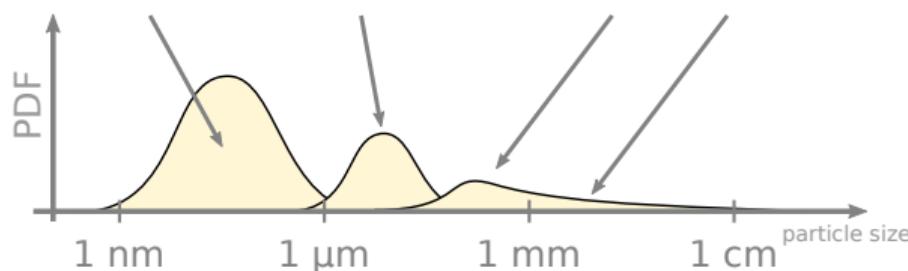


# Eulerian vs. Lagrangian microphysics

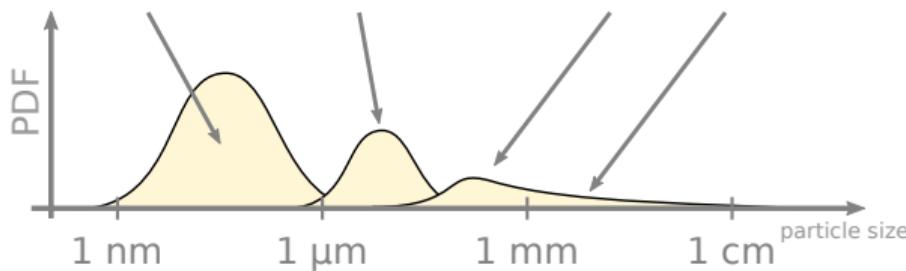
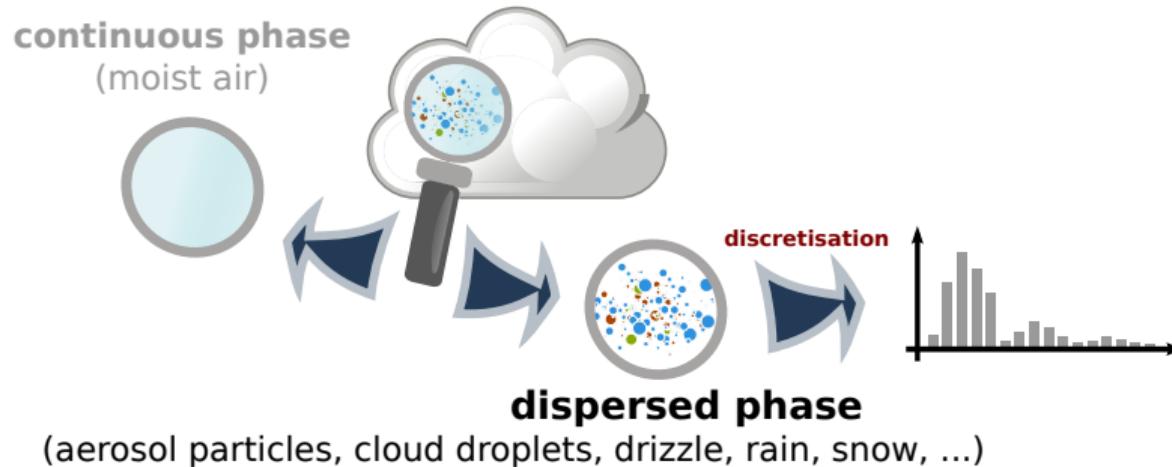
**continuous phase**  
(moist air)



**dispersed phase**  
(aerosol particles, cloud droplets, drizzle, rain, snow, ...)



# Eulerian vs. Lagrangian microphysics

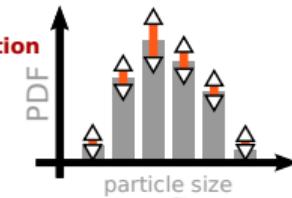


# Eulerian vs. Lagrangian microphysics

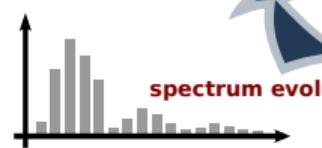
**continuous phase**  
(moist air)



Eulerian representation

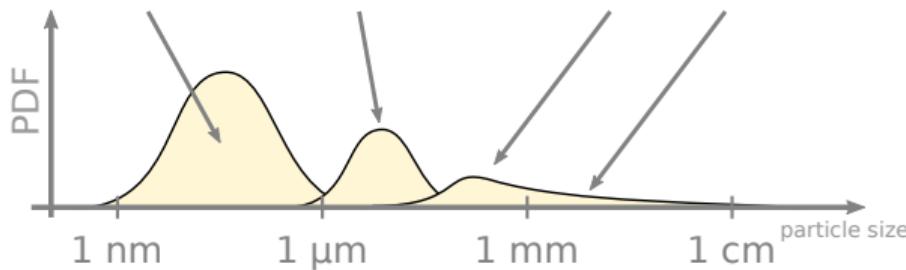


spectrum evolution



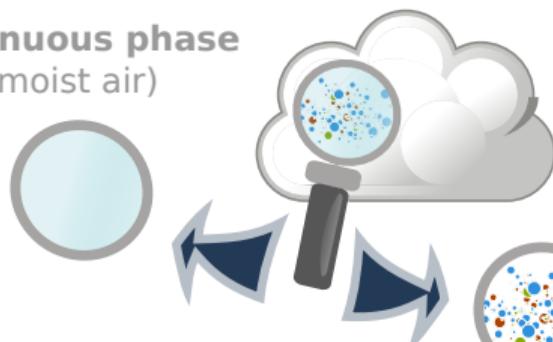
## dispersed phase

(aerosol particles, cloud droplets, drizzle, rain, snow, ...)

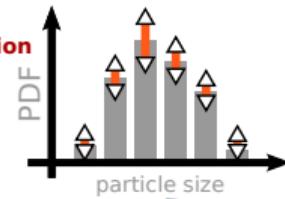


# Eulerian vs. Lagrangian microphysics

continuous phase  
(moist air)

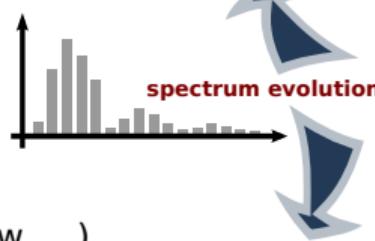


Eulerian representation



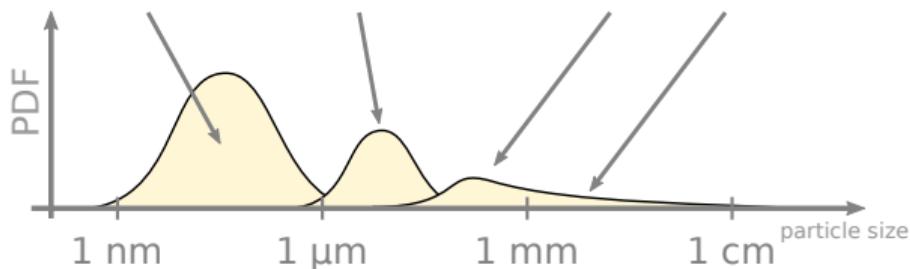
PDEs

discretisation

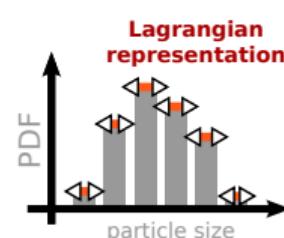


dispersed phase

(aerosol particles, cloud droplets, drizzle, rain, snow, ...)



Lagrangian representation



ODEs

# Lagrangian microphysics: early works (0D)

JOURNAL OF METEOROLOGY

## THE GROWTH OF CLOUD DROPS IN UNIFORMLY COOLED AIR

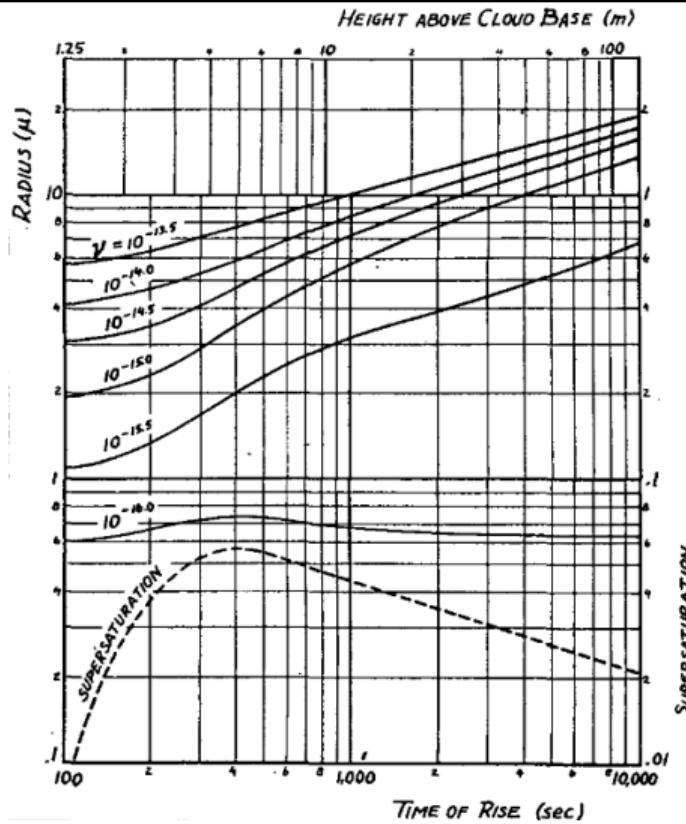
By Wallace E. Howell<sup>1</sup>

Blue Hill Meteorological Observatory, Harvard University<sup>2</sup>

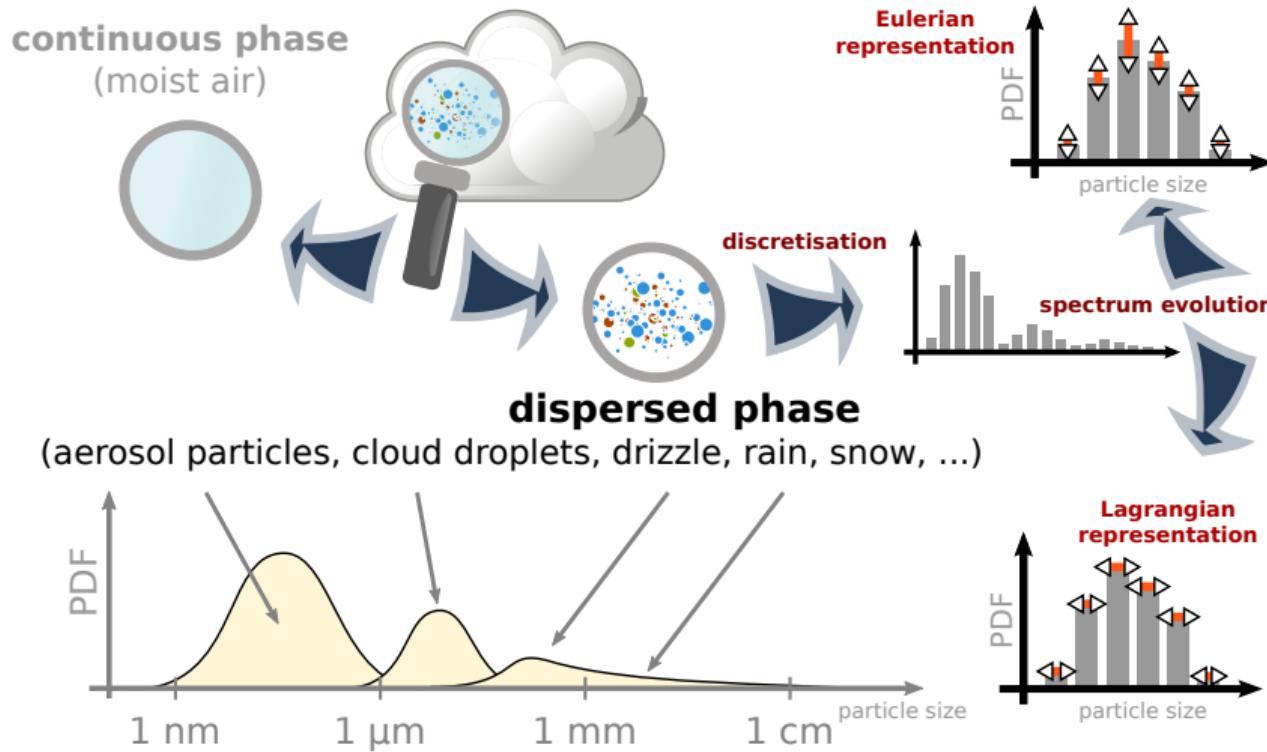
(Manuscript received 10 June 1948)

### ABSTRACT

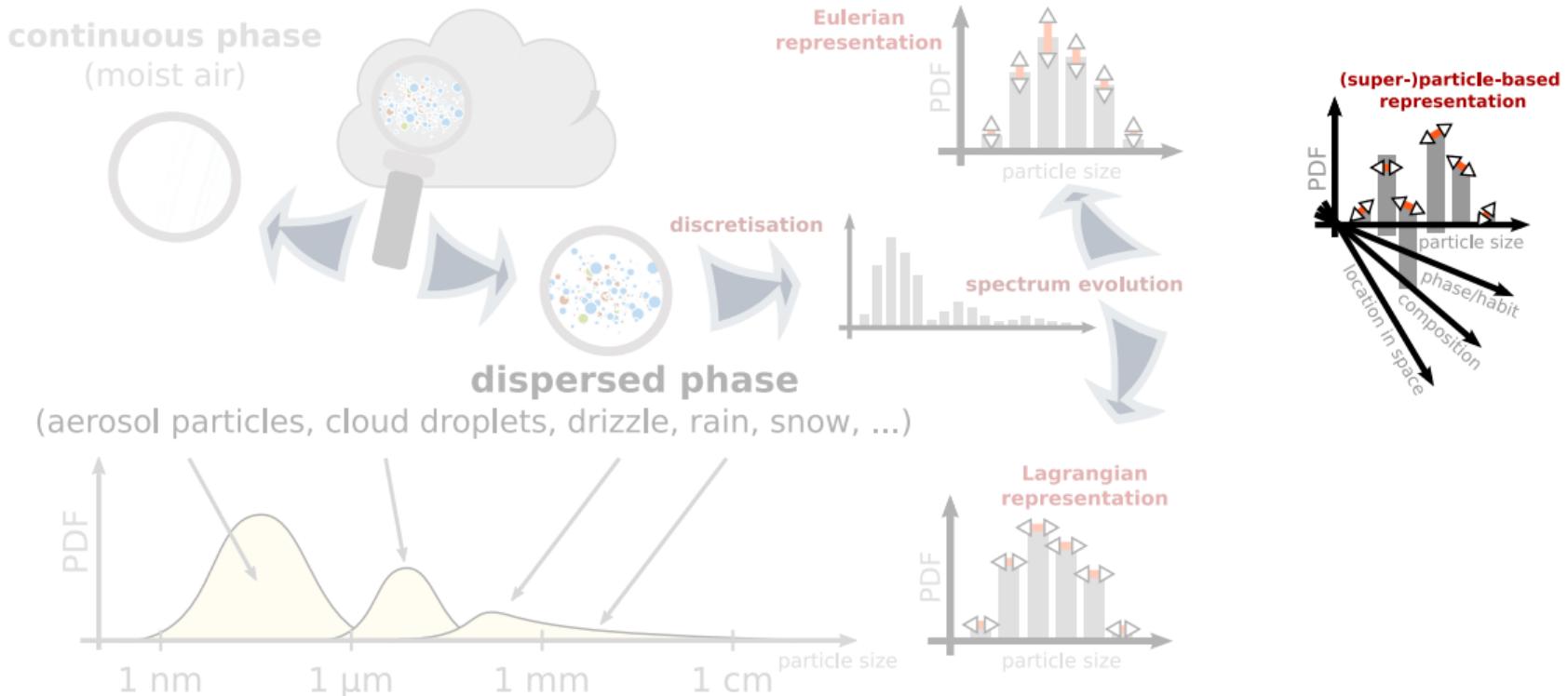
Recent studies of precipitation, aircraft icing, and visibility through fog have focussed attention on the physical constitution of clouds, a subject to which knowledge of the drop-size spectrum and its origin would be an important contribution. The drop-size spectrum resulting when air containing condensation nuclei is uniformly cooled may be computed, leading to a differential equation for the growth of a cloud drop which cannot be integrated analytically. A numerical method of integration is therefore employed.



# Eulerian vs. Lagrangian microphysics



# Eulerian vs. Lagrangian microphysics

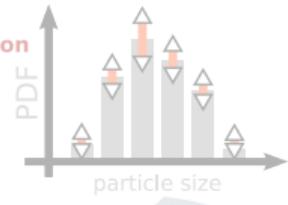


# Eulerian vs. Lagrangian microphysics

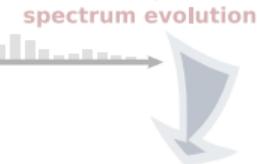
continuous phase  
(moist air)



Eulerian representation

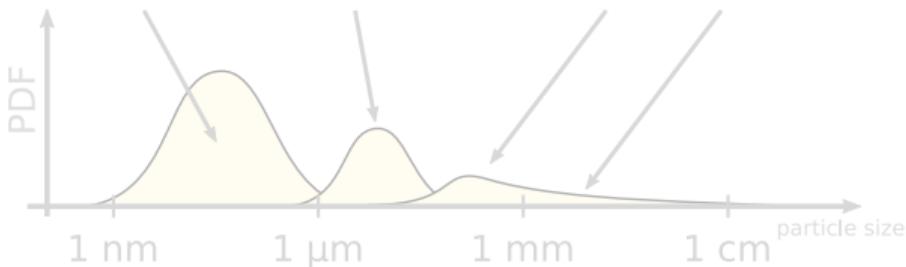


discretisation

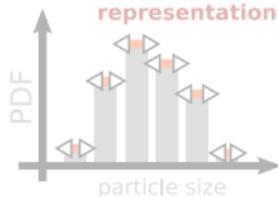


dispersed phase

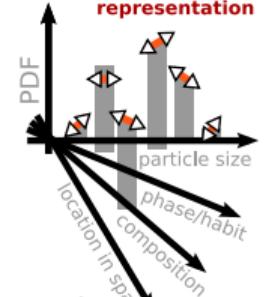
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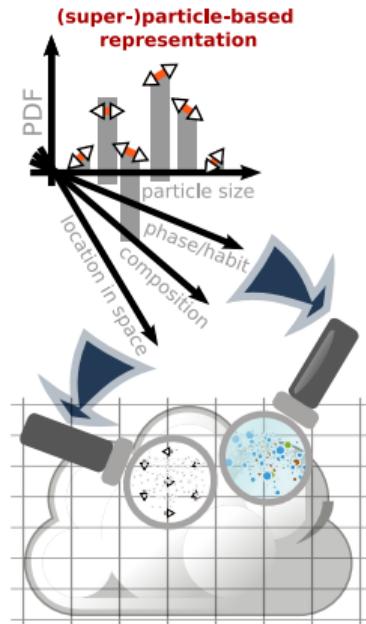
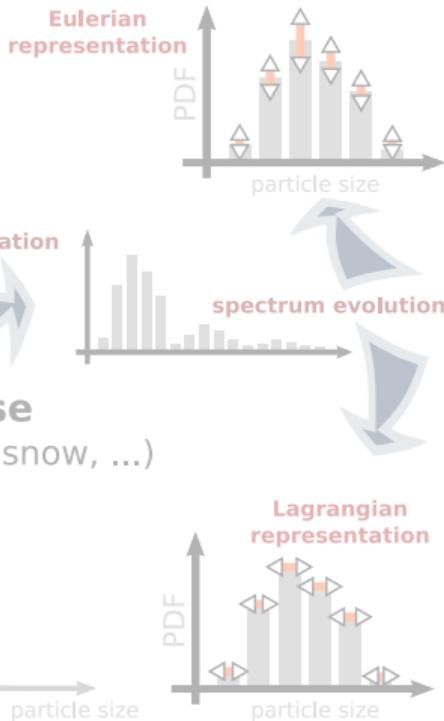
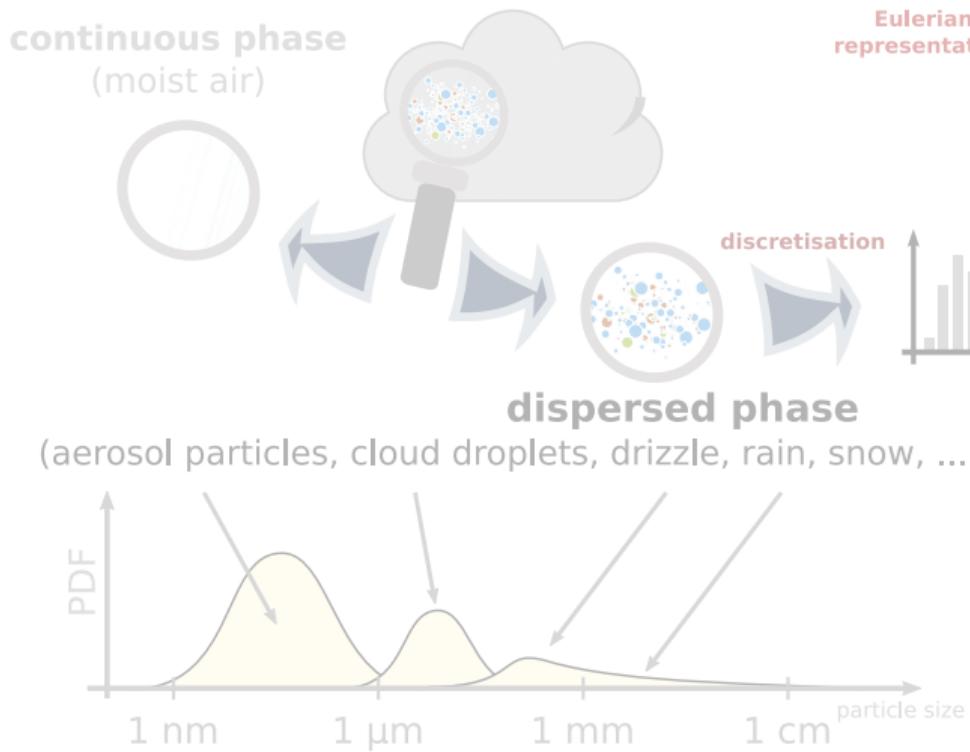
Lagrangian representation



(super-)particle-based representation



# Eulerian vs. Lagrangian microphysics



## A Numerical Experiment on Stochastic Condensation Theory

TERRY L. CLARK AND W. D. HALL

*National Center for Atmospheric Research,<sup>1</sup> Boulder, CO 80307*

(Manuscript received 30 August 1978, in final form 20 November 1978)

### ABSTRACT

A three-dimensional numerical model is used to study the effect of small-scale supersaturation fluctuations on the evolving droplet distribution in the first 150 m above cloud base. The primary purpose of this research is to determine whether the irreversible coupling between the thermodynamics and dynamics due to finite phase relaxation time scales  $\tau_S$  is sufficient to produce significant small-scale horizontal variations in supersaturation. Thus, the paper is concerned only with this internal source for thermodynamic variability. All other source terms, such as the downgradient flux of the variance of thermodynamic fields, have purposely been neglected.

Lagrangian particle experiments were run in parallel with the basic Eulerian model. The purpose of these experiments is to relax some of the microphysical parameterization assumptions with respect to assumed distribution shape and as a result add credibility to the results of distribution broadening.

# Eulerian vs. Lagrangian microphysics: a (probabilistic) breakthrough

## pre-2009:

„advantage of the full-moving size structure is that core particle material is preserved during growth ... second advantage ... it eliminates numerical diffusion ... [but] nucleation, coagulation ... cause problems ... the full-moving structure is **not used in three-dimensional models**“<sup>a</sup>

„the use of a fixed grid allows for an easy implementation of collision processes, which is not possible for a moving grid (Lagrangian) approach“<sup>b</sup>

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<sup>a</sup> Jacobson 2005: Fundamentals of Atmospheric Modeling

<sup>b</sup> Simmel & Wurzler 2006: Condensation and activation in sectional cloud microphysical models

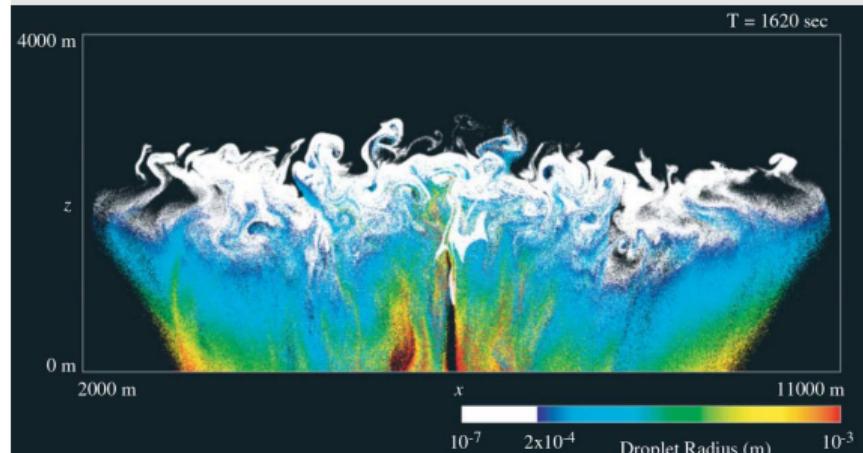
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## Shima 2009: Monte-Carlo particle-based collision algorithm for cloud simulations



Super-droplet simulation of a shallow convective cloud  
(figure: Shima et al. 2009, QJRMS)

<sup>a</sup> Jacobson 2005: Fundamentals of Atmospheric Modeling

<sup>b</sup> Simmel & Wurzler 2006: Condensation and activation in sectional cloud microphysical models

## SCE (naïve impl)

## SDM

### method type

mean-field, deterministic

Monte-Carlo, stochastic

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all (i,j) pairs

random set of  $n_{sd}/2$  non-overlapping pairs,  
probability up-scaled by  $(n_{sd}^2 - n_{sd})/2$  to  $n_{sd}/2$  ratio

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collide all of  $\min\{\xi_{[i]}, \xi_{[j]}\}$  ("all or nothing")

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collisions triggered	every time step	by comparing probability with a random number
collisions	colliding a fraction of $\xi_{[i]}, \xi_{[j]}$	collide all of $\min\{\xi_{[i]}, \xi_{[j]}\}$ ("all or nothing")
interpretation	concentration " $c_i$ " in size bin " $i$ "	besides $c_i$ , each "particle" $i$ carries other physicochemical attributes, e.g. position $(x_i, y_i, z_i)$

# super-particles as an alternative to bulk or bin $\mu$ -physics

## Confronting the Challenge of Modeling Cloud and Precipitation Microphysics

Hugh Morrison✉, Marcus van Lier-Walqui, Ann M. Fridlind, Wojciech W. Grabowski, Jerry Y. Harrington,

Corinna Hoose, Alexei Korolev, Matthew R. Kumjian, Jason A. Milbrandt, Hanna Pawlowska, Derek J.

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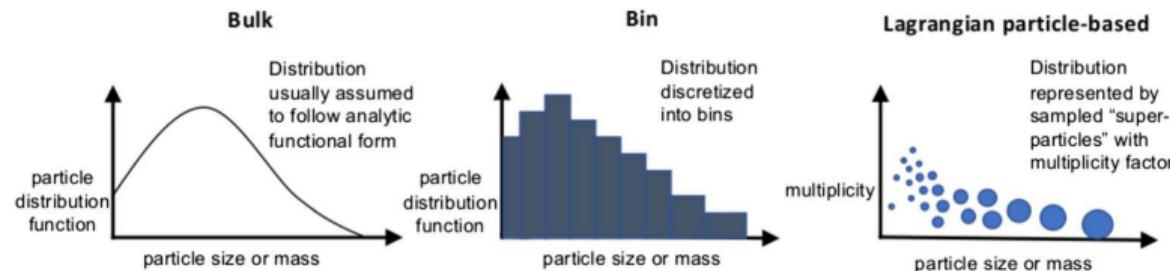
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Journal of Advances in Modeling Earth Systems [10.1029/2019MS001689](https://doi.org/10.1029/2019MS001689)



**Figure 3.** Representation of cloud and precipitation particle distributions in the three main types of microphysics

# **PySDM: SDM implementation for reproducible research & active learning**

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PySDM

## PySDM goals:

---

- implementation of SDM + particle-based/Monte-Carlo models of other cloud processes

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- applicable in research on aerosol-cloud-interactions (and beyond)

KPI: reproduction of results from classic literature, use in new research

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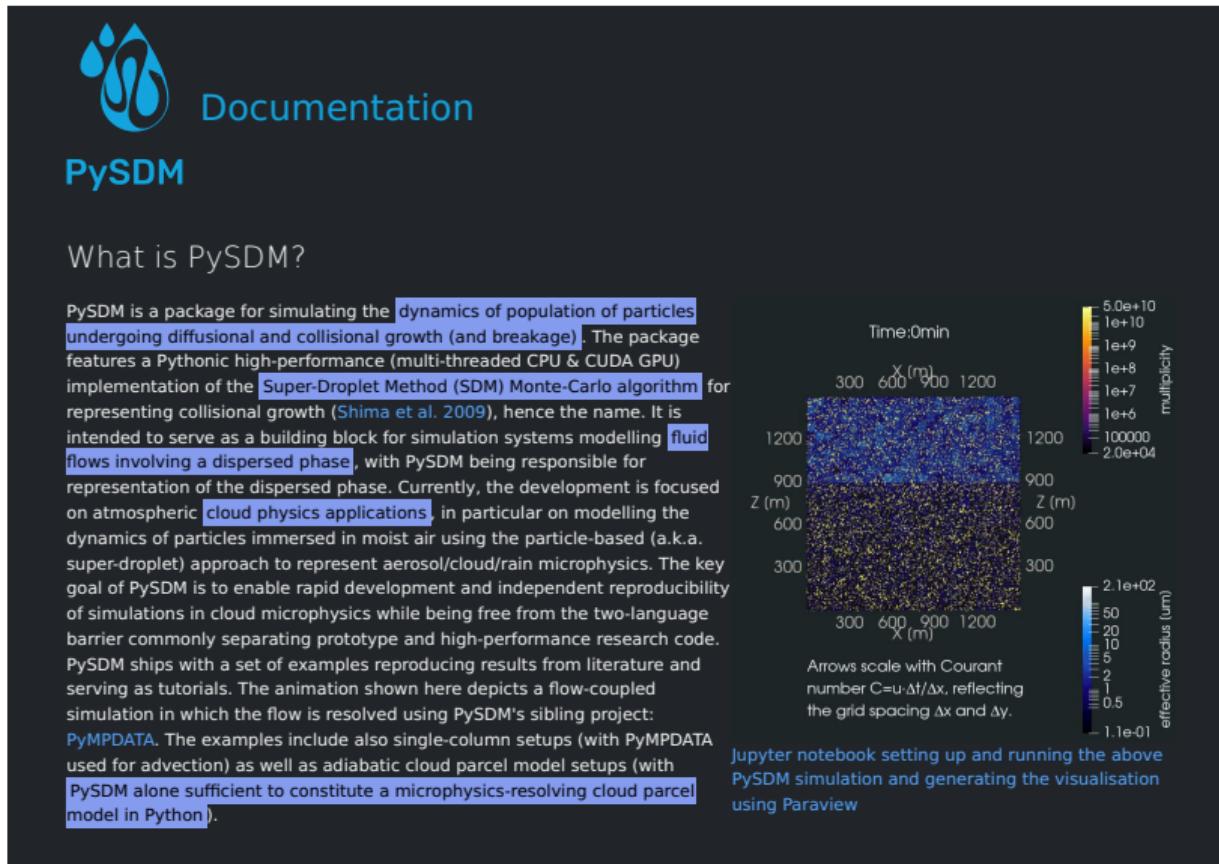
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KPI: reproduction of results from classic literature, use in new research
- **easy to reuse**: code (100% Python), documentation, examples/tutorials (Jupyter), extensibility (modular, high test coverage), interoperability (other languages, i/o), leveraging modern hardware (GPUs, multi-core CPUs)  
KPI: user feedback & multi-institutional contributions

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KPI: continuous integration on all targeted platforms
- **curation**: open licensing (GPL), public versioned development (Github), archival (Zenodo)  
KPI: instant and anonymous execution of arbitrary version in commodity environments





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**Emily de Jong @caltech.edu (now @llnl.gov)** collisional breakup, ...

# Monte-Carlo collisional breakup (constant super-droplet number formulation)

Geosci. Model Dev., 16, 4193–4211, 2023  
<https://doi.org/10.5194/gmd-16-4193-2023>  
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the Creative Commons Attribution 4.0 License.



Geoscientific  
Model Development  
Open Access

Development and technical paper

## Breakups are complicated: an efficient representation of collisional breakup in the superdroplet method

Emily de Jong<sup>1</sup>, John Ben Mackay<sup>2,a</sup>, Oleksii Bulenok<sup>3</sup>, Anna Jaruga<sup>4</sup>, and Sylwester Arabas<sup>5,b,c</sup>

<sup>1</sup>Department of Mechanical and Civil Engineering, California Institute of Technology, Pasadena, CA, USA

<sup>2</sup>Scripps Institution of Oceanography, San Diego, CA, USA

<sup>3</sup>Faculty of Mathematics and Computer Science, Jagiellonian University, Kraków, Poland

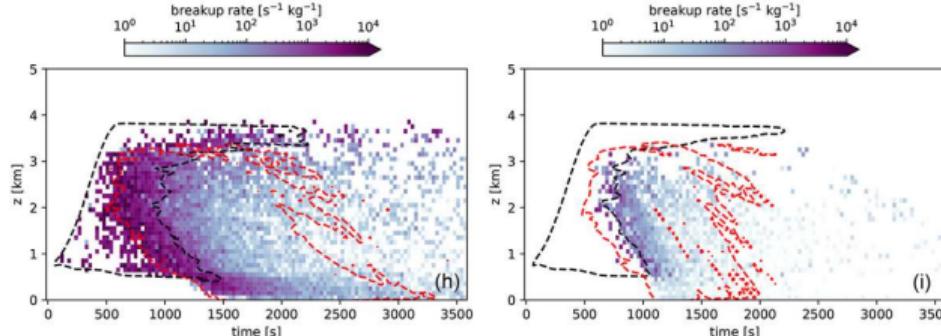
<sup>4</sup>Department of Environmental Science and Engineering, California Institute of Technology, Pasadena, CA, USA

<sup>5</sup>Faculty of Physics and Applied Computer Science, AGH University of Krakow, Kraków, Poland

<sup>a</sup>formerly at: Department of Environmental Science and Engineering, California Institute of Technology, Pasadena, CA, USA

<sup>b</sup>formerly at: Department of Atmospheric Sciences, University of Illinois Urbana-Champaign, Urbana, IL, USA

<sup>c</sup>formerly at: Faculty of Mathematics and Computer Science, Jagiellonian University, Kraków, Poland





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**Jatan Buch @columbia.edu** cloud seeding, ...



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**Jason Barr @washington.edu** coagulation kernels, ...



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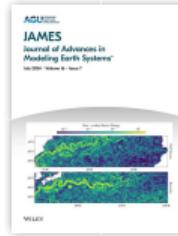
**Jason Barr @washington.edu** coagulation kernels, ...

**Sajjad Azimi & Anna Jaruga @caltech.edu** calibration of bulk model using PySDM, ...

## Training Warm-Rain Bulk Microphysics Schemes Using Super-Droplet Simulations

Sajjad Azimi , Anna Jaruga, Emily de Jong, Sylvester Arabas, Tapio Schneider

First published: 26 July 2024 | <https://doi.org/10.1029/2023MS004028>



Volume 16, Issue 7

July 2024

e2023MS004028

This article also appears in:  
The CliMA Earth System Model

## Abstract

Cloud microphysics is a critical aspect of the Earth's climate system, which involves processes at the nano- and micrometer scales of droplets and ice particles. In climate modeling, cloud microphysics is commonly represented by bulk models, which contain simplified process rates that require calibration. This study presents a framework for calibrating warm-rain bulk schemes using high-fidelity super-droplet simulations that provide a more accurate and physically based representation of cloud and precipitation processes. The calibration framework employs ensemble Kalman methods including Ensemble Kalman Inversion and Unscented Kalman Inversion to calibrate bulk microphysics schemes with probabilistic super-droplet simulations. We demonstrate the framework's effectiveness by calibrating a single-moment bulk scheme, resulting in a reduction of data-model mismatch by more than 75% compared to the model with initial parameters. Thus, this study demonstrates a powerful tool for enhancing the accuracy of bulk microphysics schemes in atmospheric models and improving climate modeling.



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**Emily de Jong @caltech.edu (now @llnl.gov)** collisional breakup, ...

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**SA @illinois.edu (now @agh.edu.pl)** immersion freezing, ...

# **particle-based modelling of immersion freezing**

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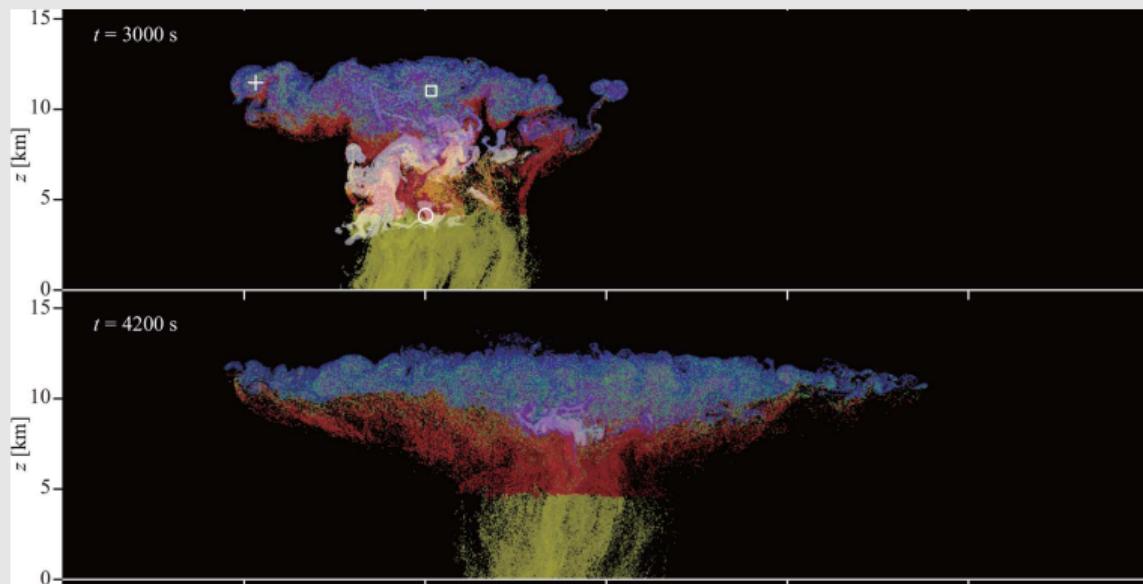
## Monte-Carlo SCE alternatives: e.g., SDM by Shima et al.

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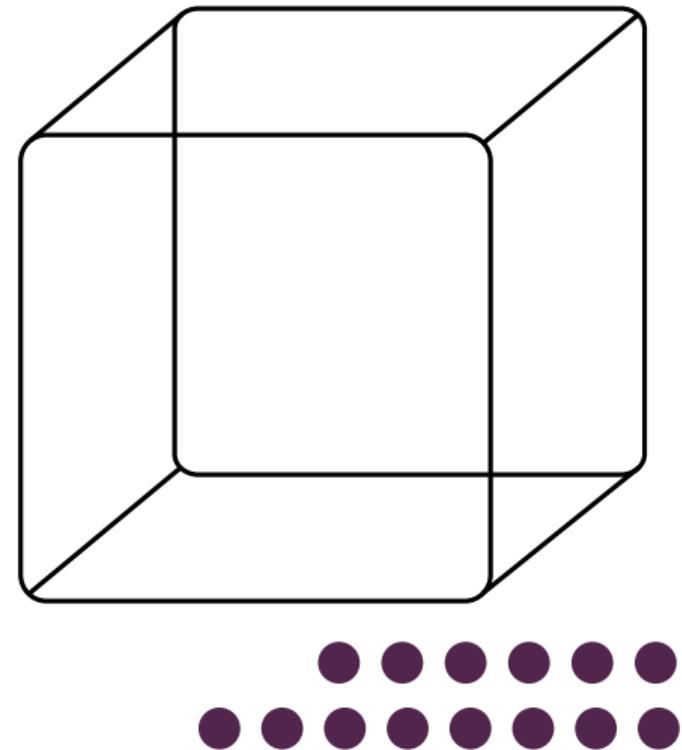
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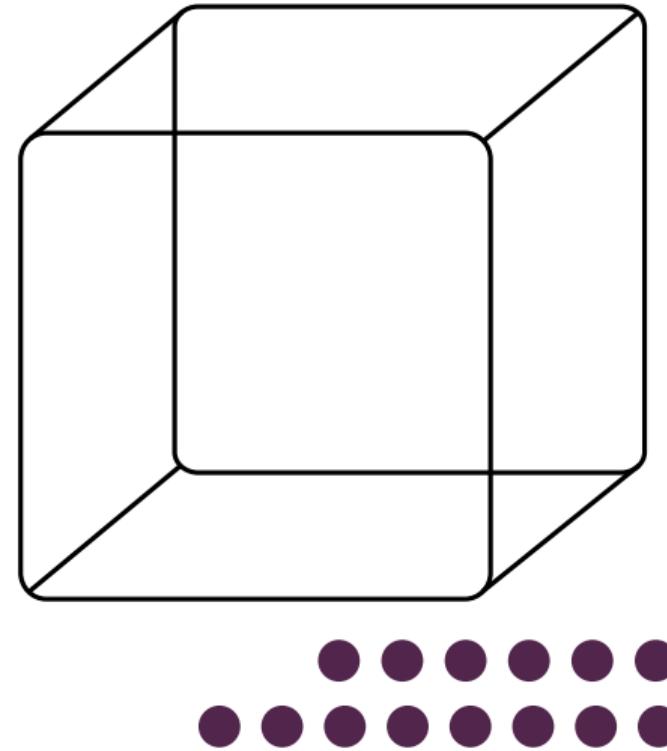
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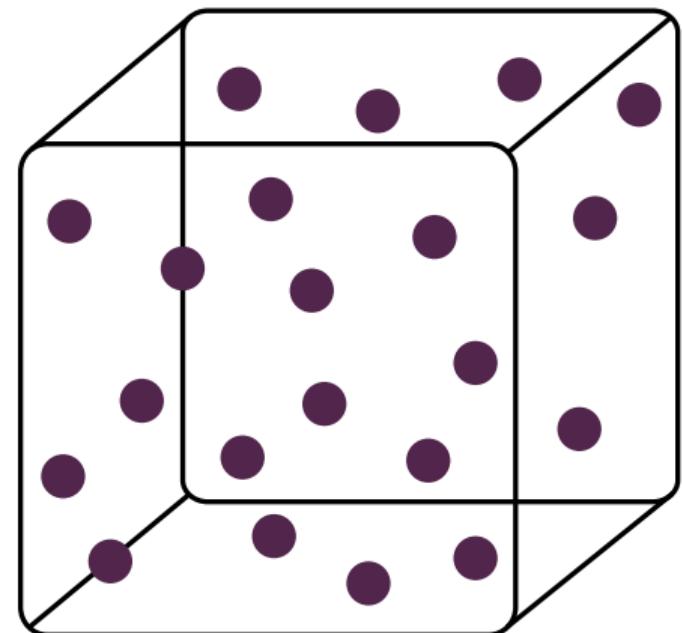
**Figure 1.** Typical realization of CTRL cloud spatial structures at  $t = 2040, 2460, 3000, 4200$ , and  $5400$  s. The mixing ratio of cloud water, rainwater, cloud ice, graupel, and snow aggregates are plotted in fading white, yellow, blue, red, and green, respectively. The symbols indicate examples of unrealistic predicted ice particles (Sects. 7.3 and 9.1). See also Movie 1 in the video supplement.

Domain randomly populated with " $\mu$ -physics information carriers"  
(super particles / super droplets)





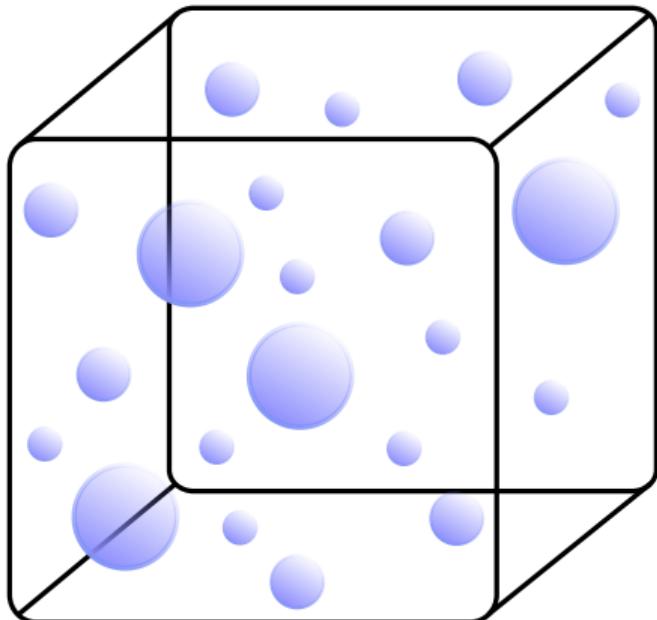
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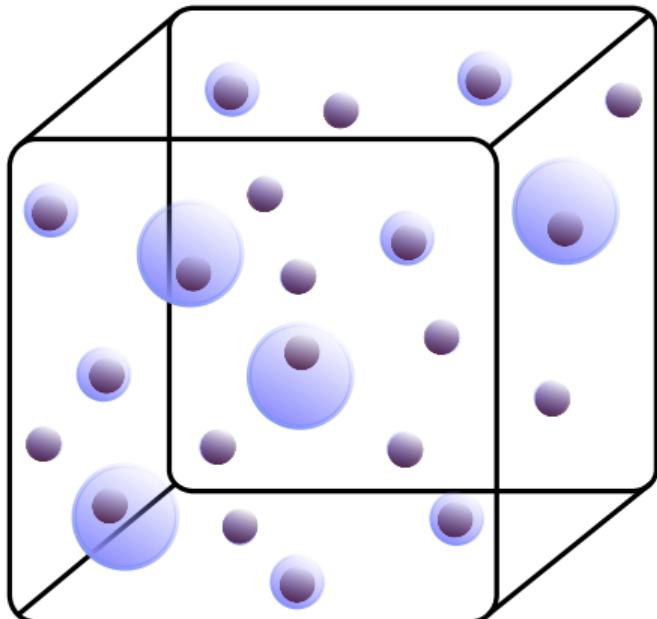
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carrier attributes:

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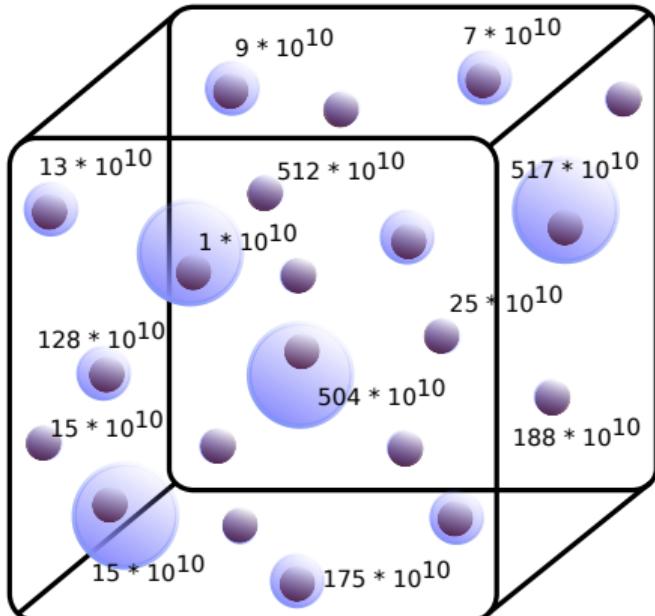
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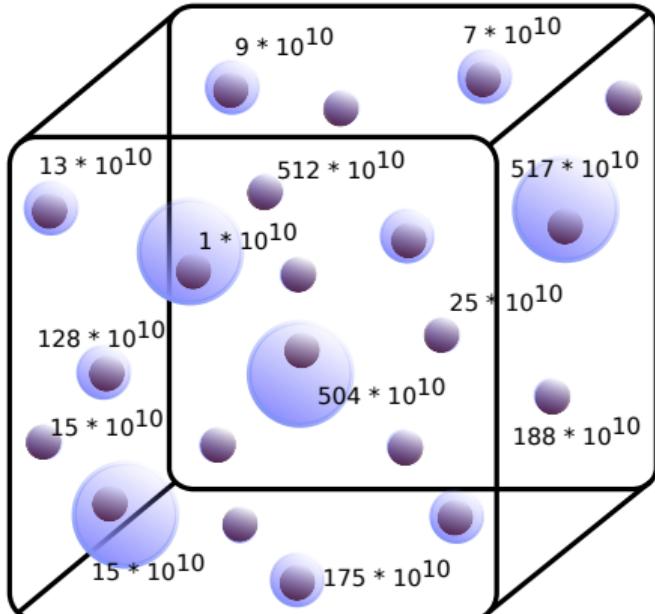
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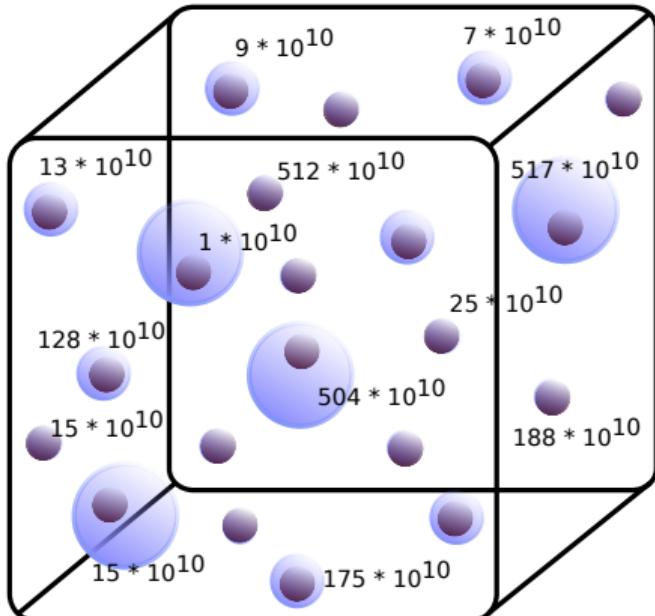
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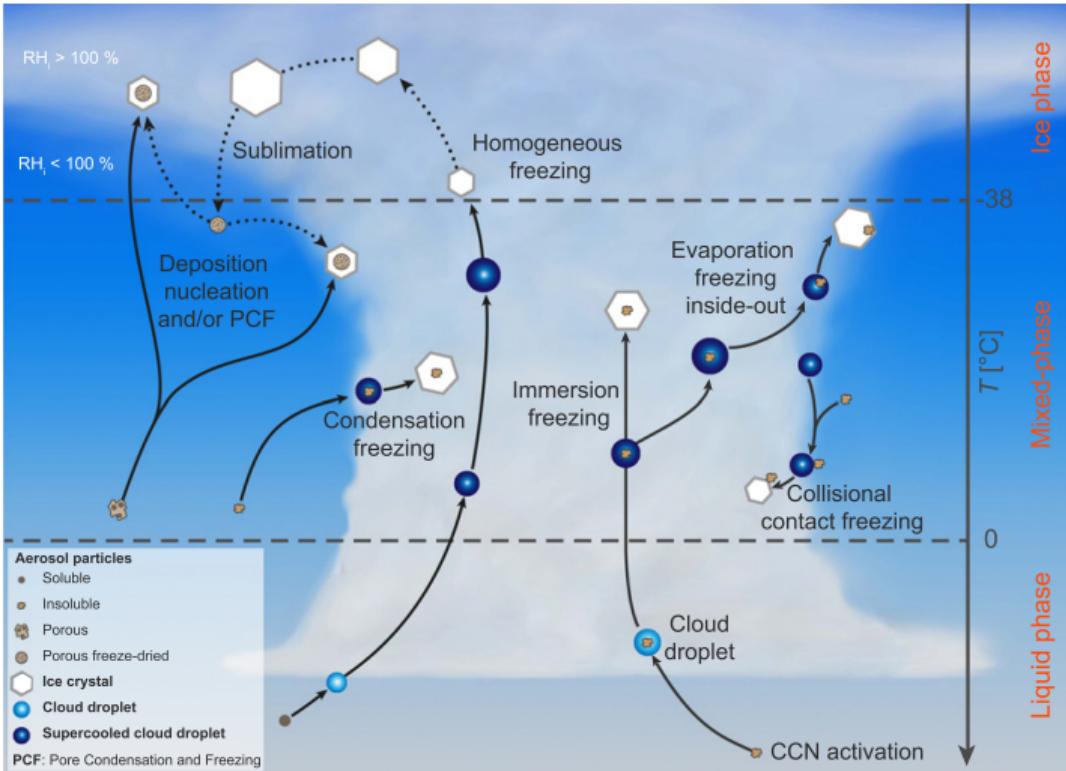


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carrier attributes:

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- ...
- freezing temperature?
- freezing rate (temperature)?





Kanji et al. 2017, graphics F. Mahrt, <https://doi.org/10.1175/AMSMONOGRAPHS-D-16-0006.1>

## RESEARCH ARTICLE

10.1002/2016JD025251

### Key Points:

- Very ice active Snomax protein aggregates are fragile and their ice nucleation ability decreases over months of freezer storage
- Partitioning of ice active protein aggregates into the immersion oil reduces the droplet's measured freezing temperature

- Caution is warranted in the use of

## The unstable ice nucleation properties of Snomax® bacterial particles

Michael Polen<sup>1</sup>, Emily Lawlis<sup>1</sup>, and Ryan C. Sullivan<sup>1</sup>

<sup>1</sup>Center for Atmospheric Particle Studies, Carnegie Mellon University, Pittsburgh, Pennsylvania, USA

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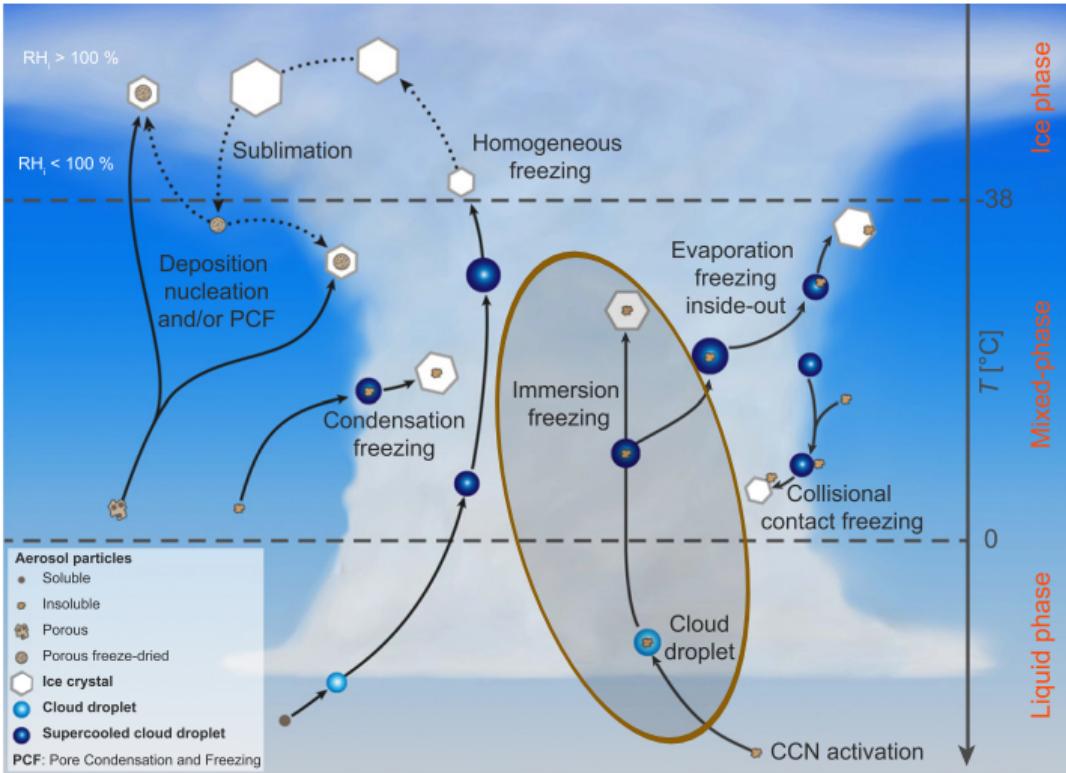
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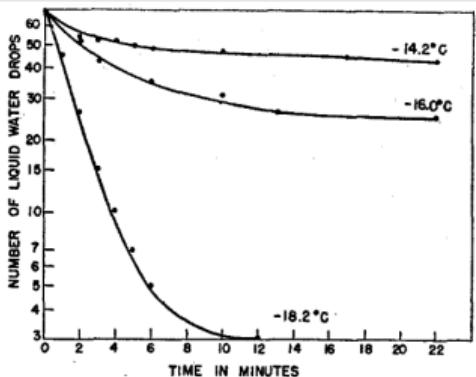


[https://www.reuters.com/markets/commodities/  
making-snow-stick-wind-challenges-winter-games-slope-makers-2021-11-29/](https://www.reuters.com/markets/commodities/making-snow-stick-wind-challenges-winter-games-slope-makers-2021-11-29/)

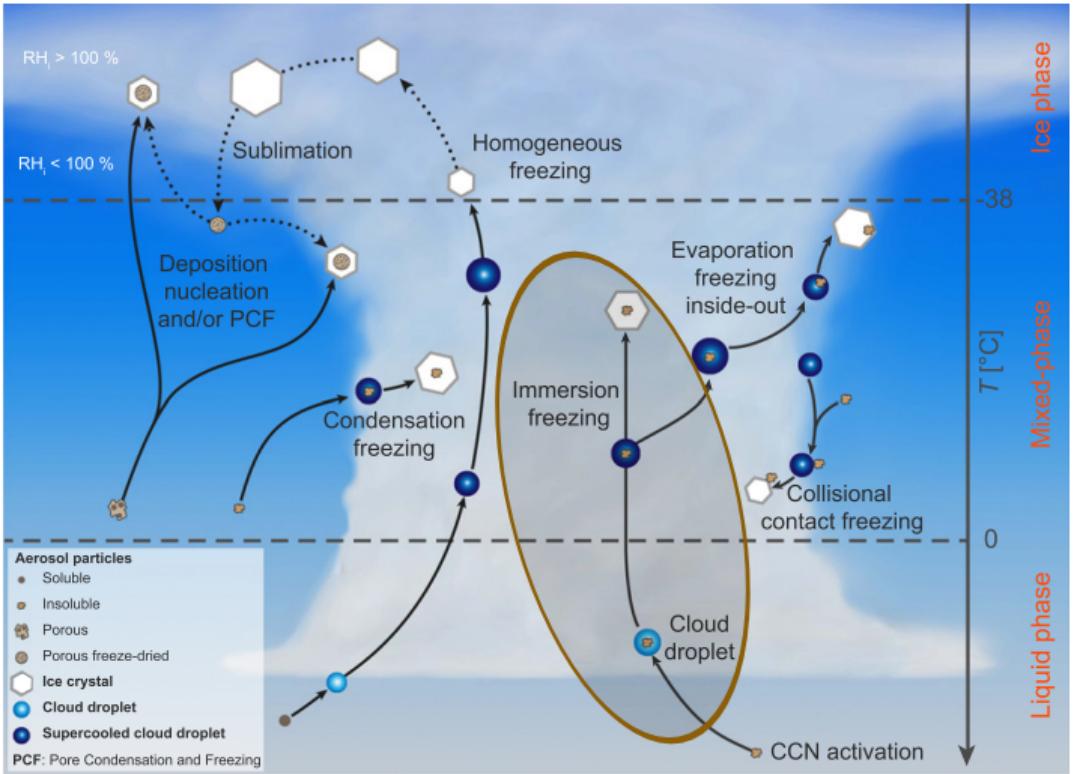


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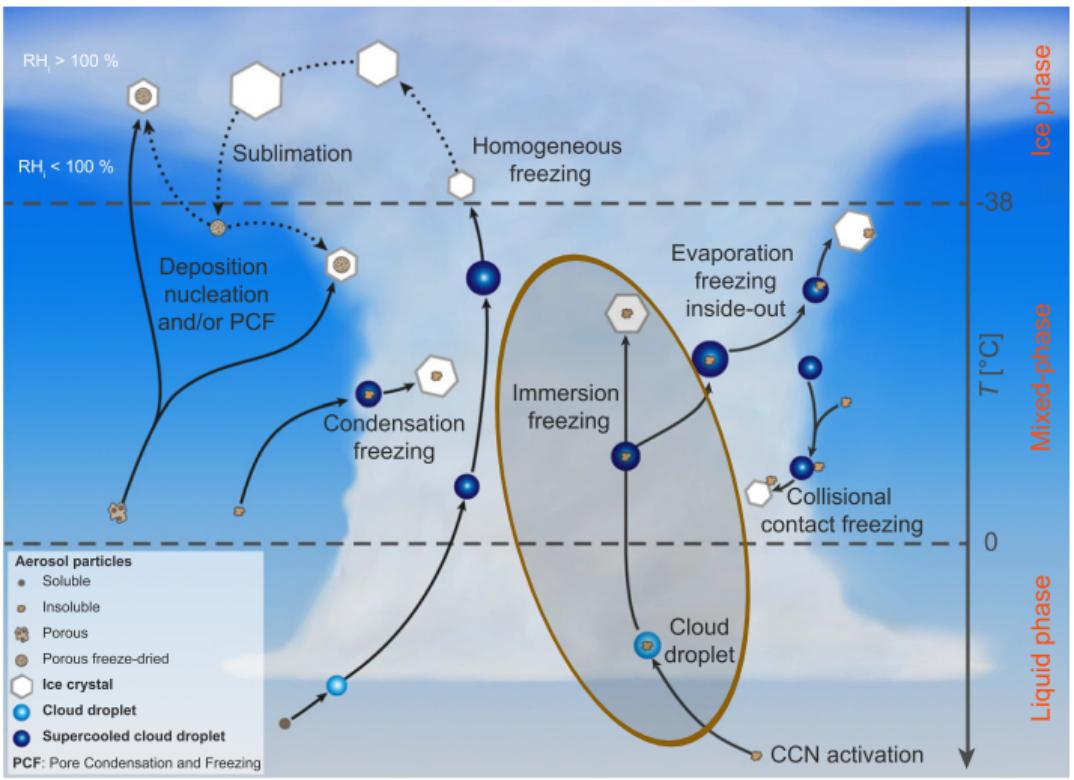
## Vonnegut 1948 (J. Colloid Sci.)



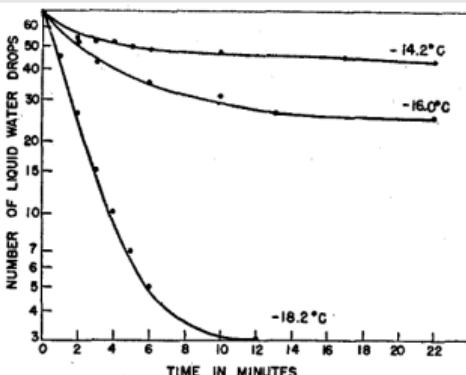
Fraction of water drops remaining unfrozen as a function of time.



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## Vali 2014 (ACP)

*"Interpretations of the experimental results face considerable difficulties ... two separate ways of interpreting the same observations; one assigned primacy to time the other emphasized the temperature-dependent impacts of the impurities ... dichotomy – the stochastic and singular models"*

# Heterogeneous Nucleations is a Stochastic Process

by

J. S. MARSHALL

McGill University, Montreal, Canada.

*Presented at the International Congress on the Physics of Clouds (Hailstorms)  
at Verona 9-13 August 1960.*

[http://cma.entecra.it/Astro2\\_sito/doc/Nubila\\_1\\_1961.pdf](http://cma.entecra.it/Astro2_sito/doc/Nubila_1_1961.pdf)

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Poisson counting process with rate  $r$ :

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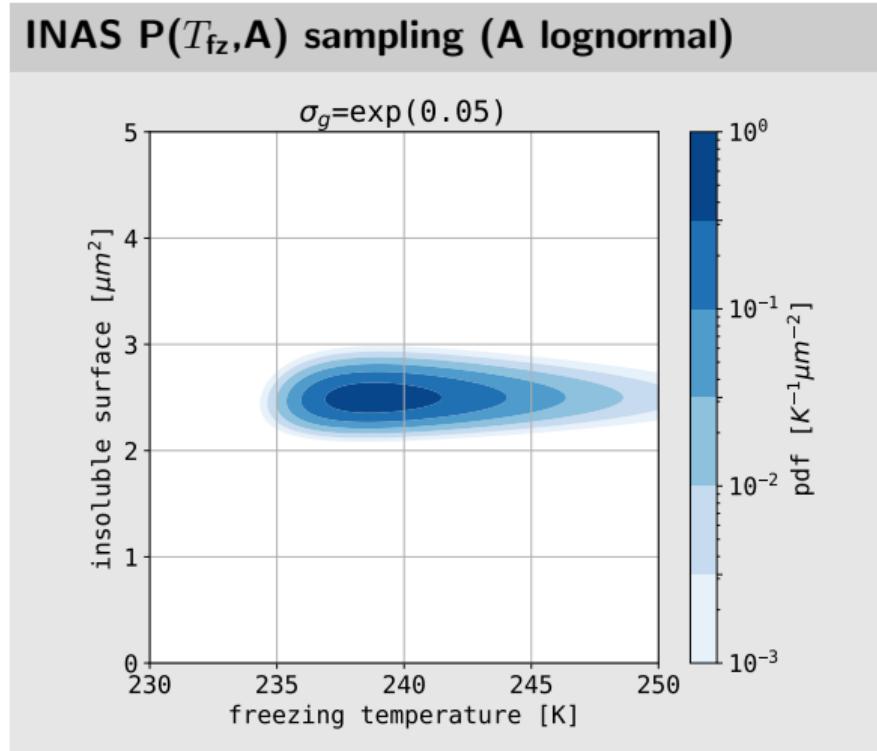
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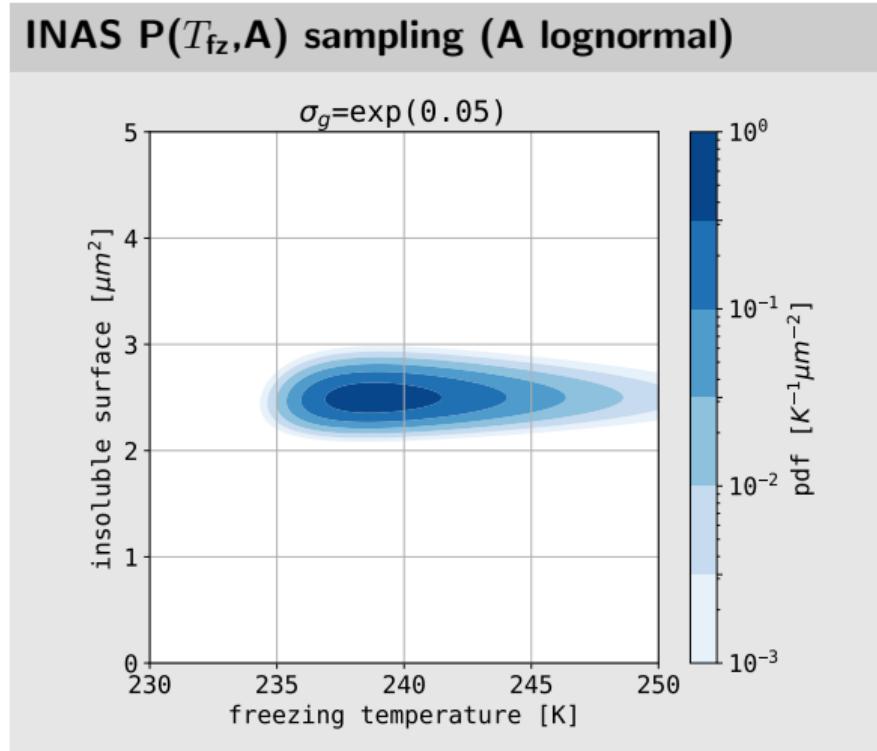
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experimental  $n_s(T)$  fits: e.g., Niemand et al. 2012

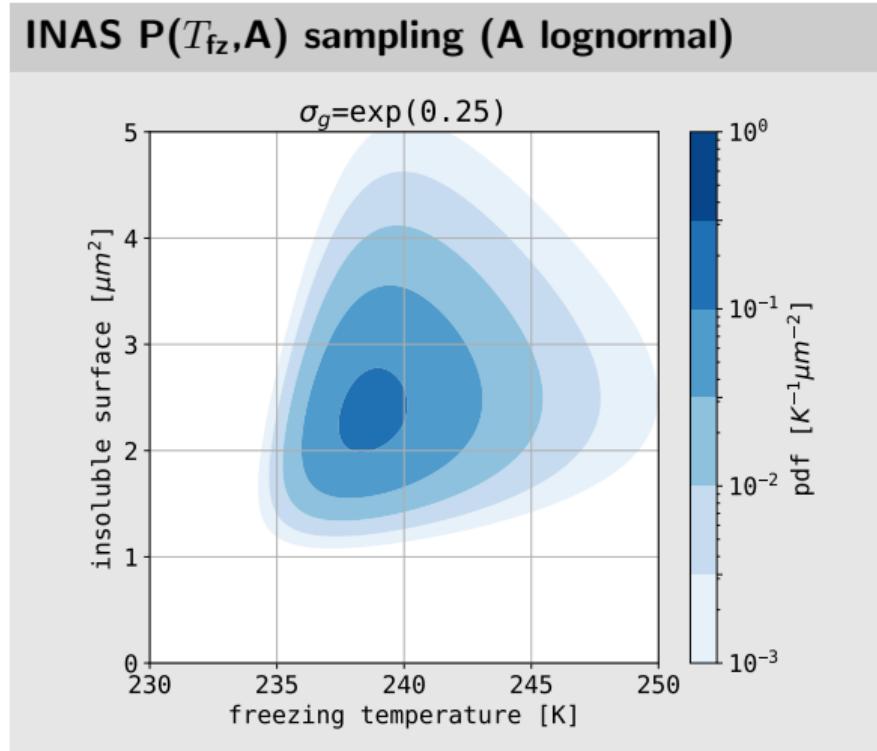
## freezing temperature $T_{\text{fz}}$ as a super-particle attribute: initialisation



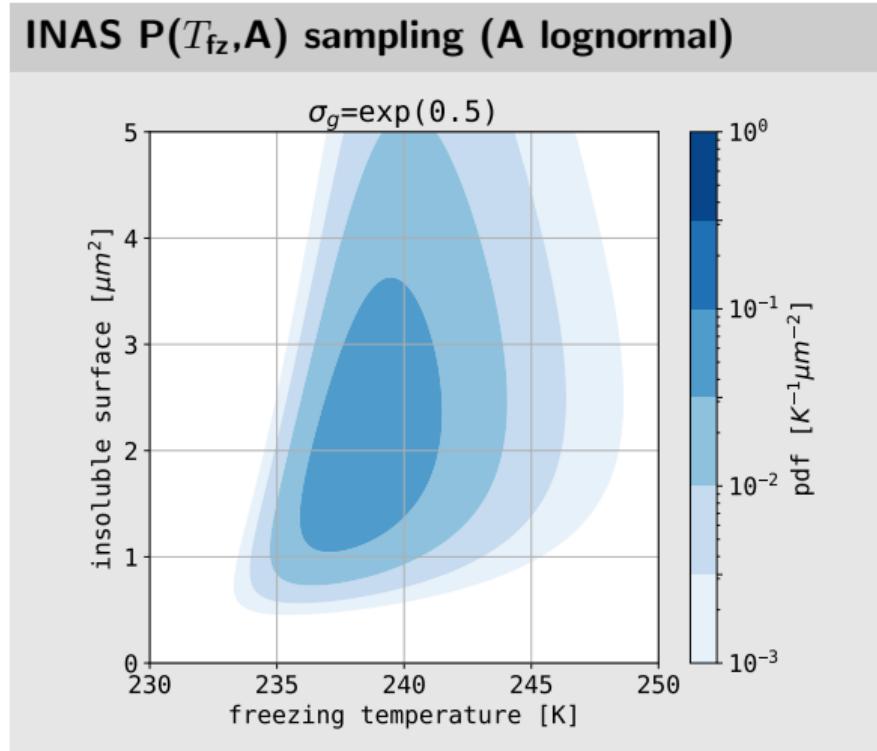
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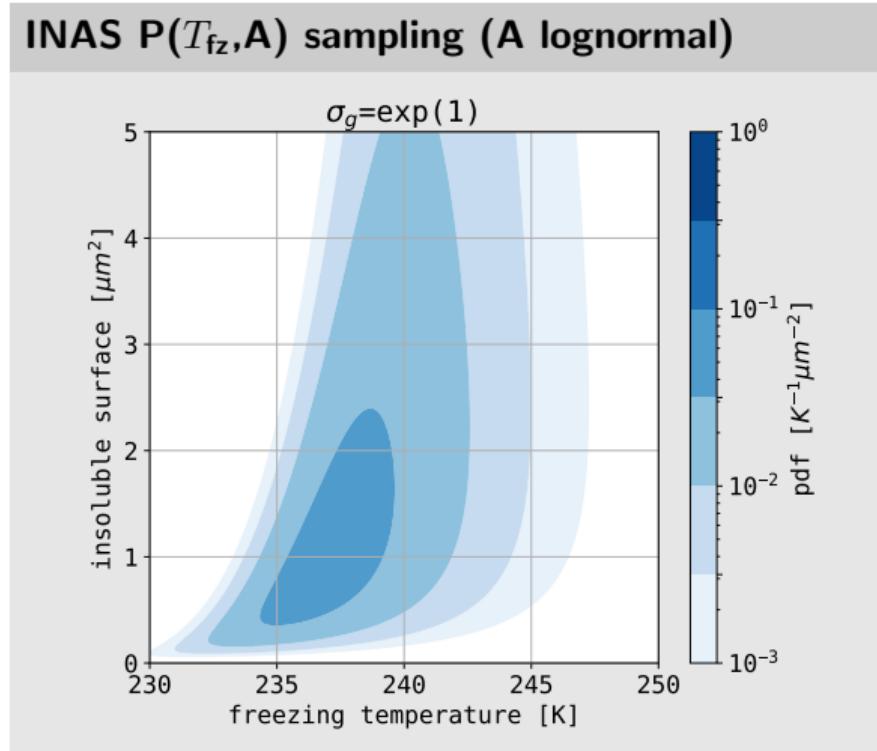
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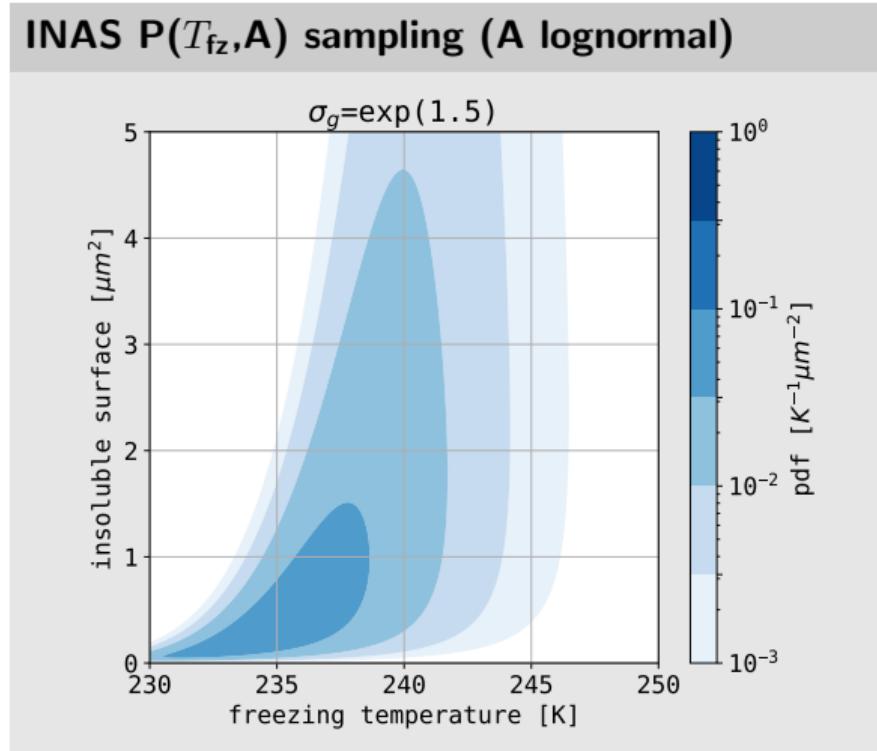
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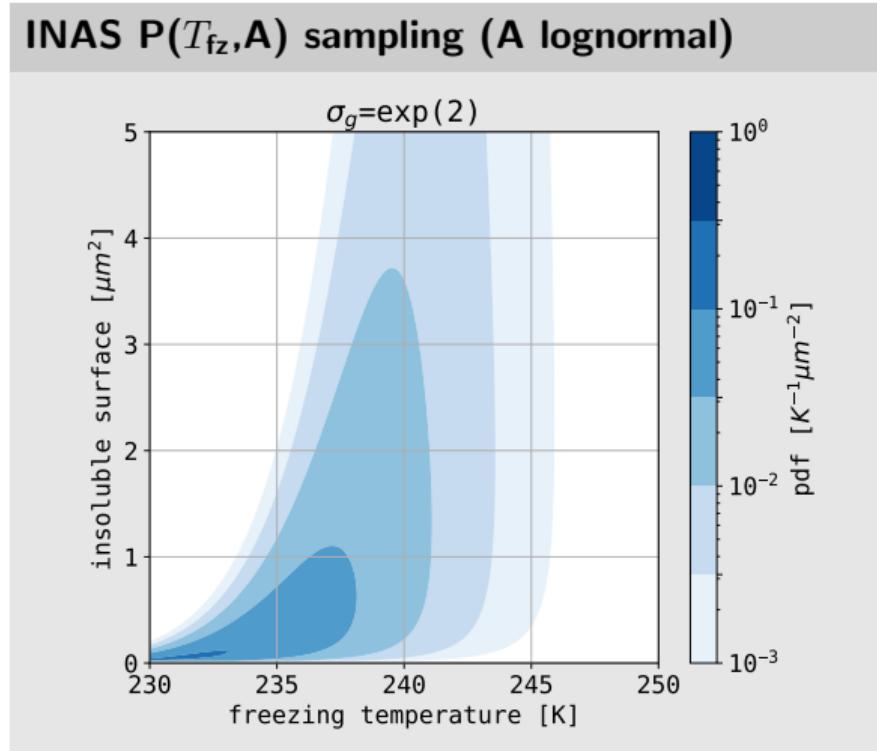
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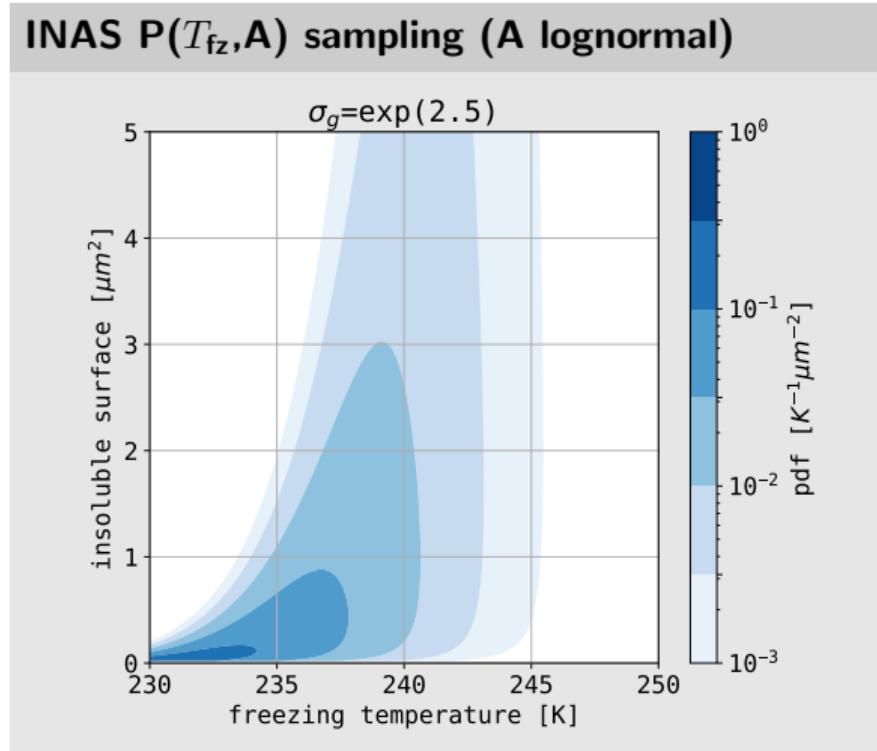
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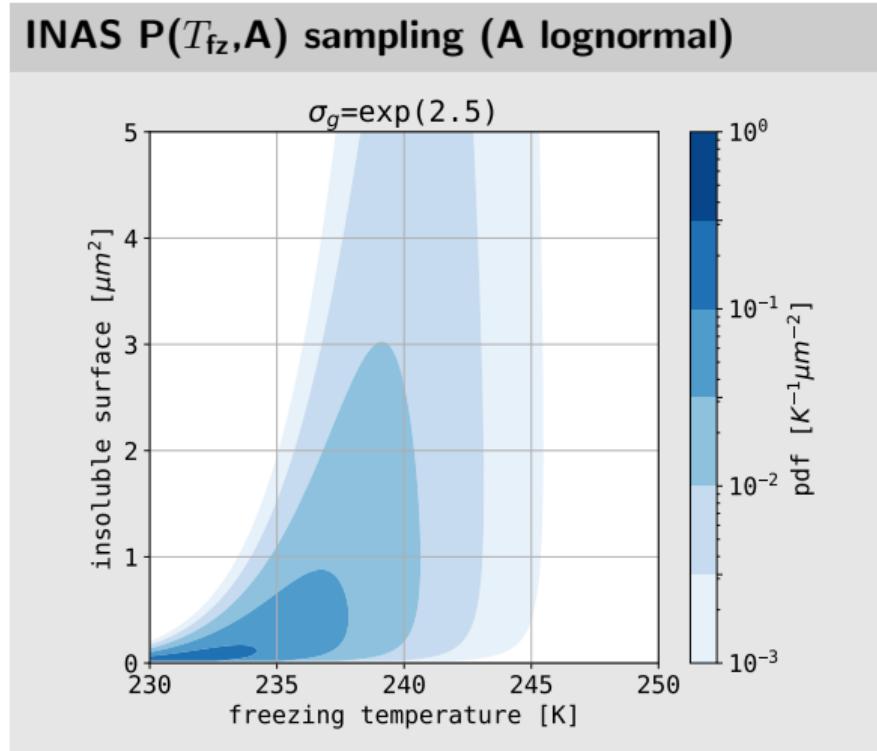
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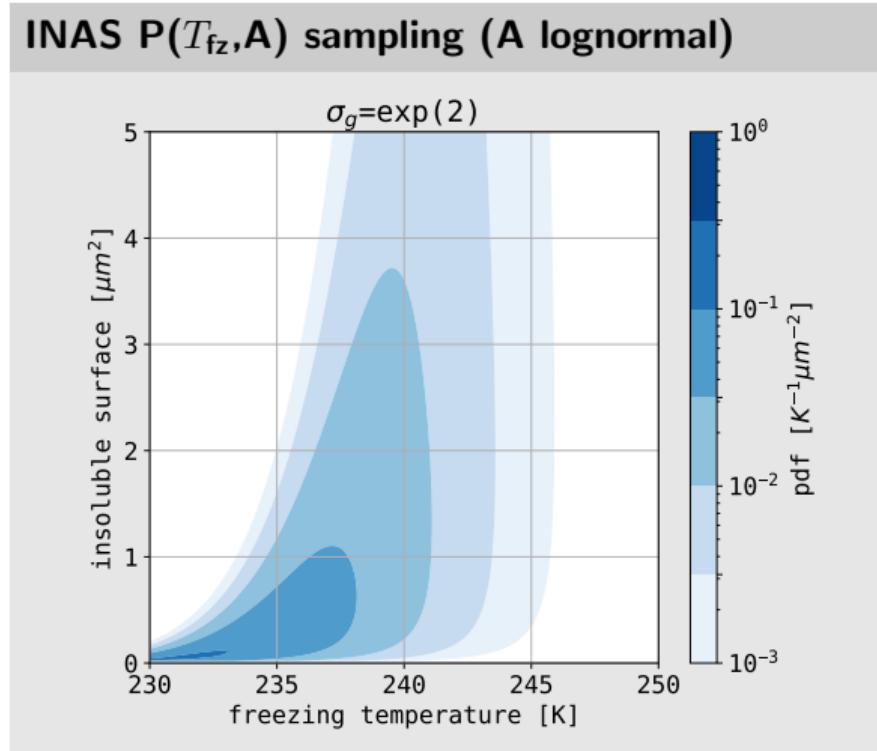
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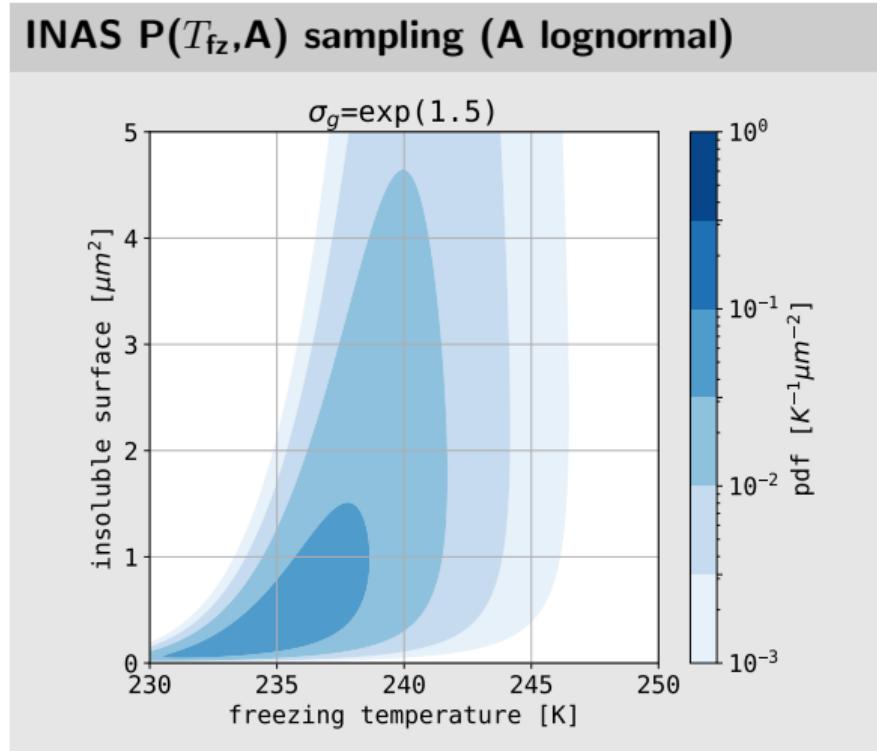
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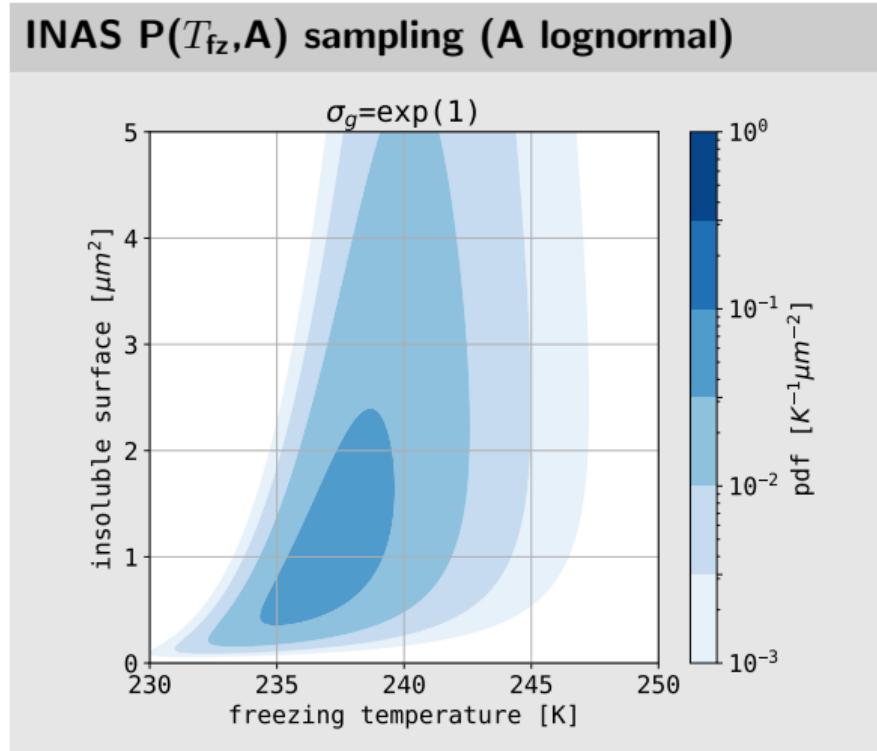
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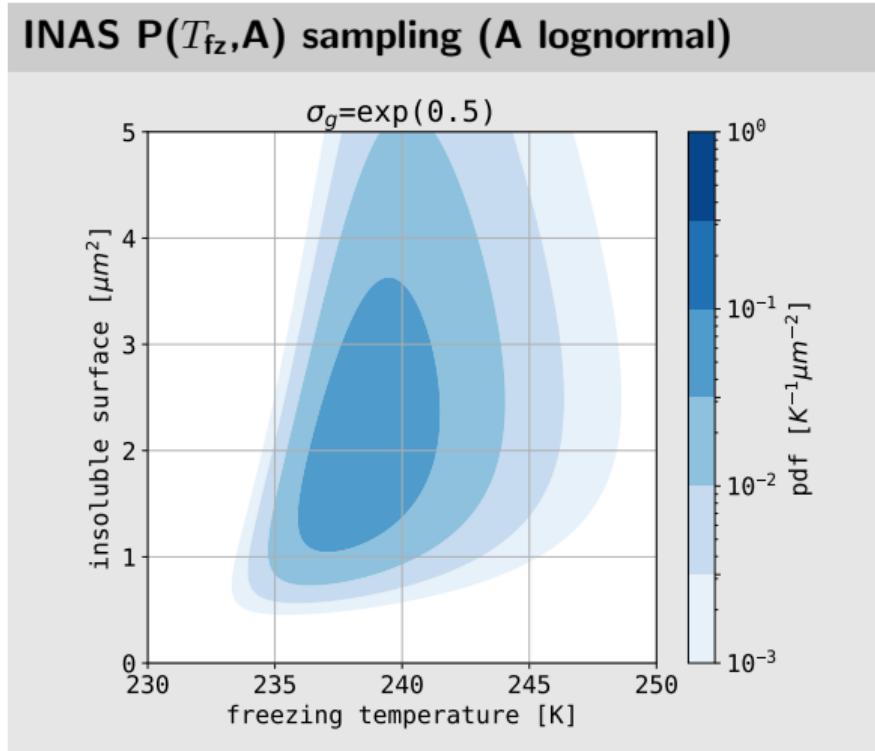
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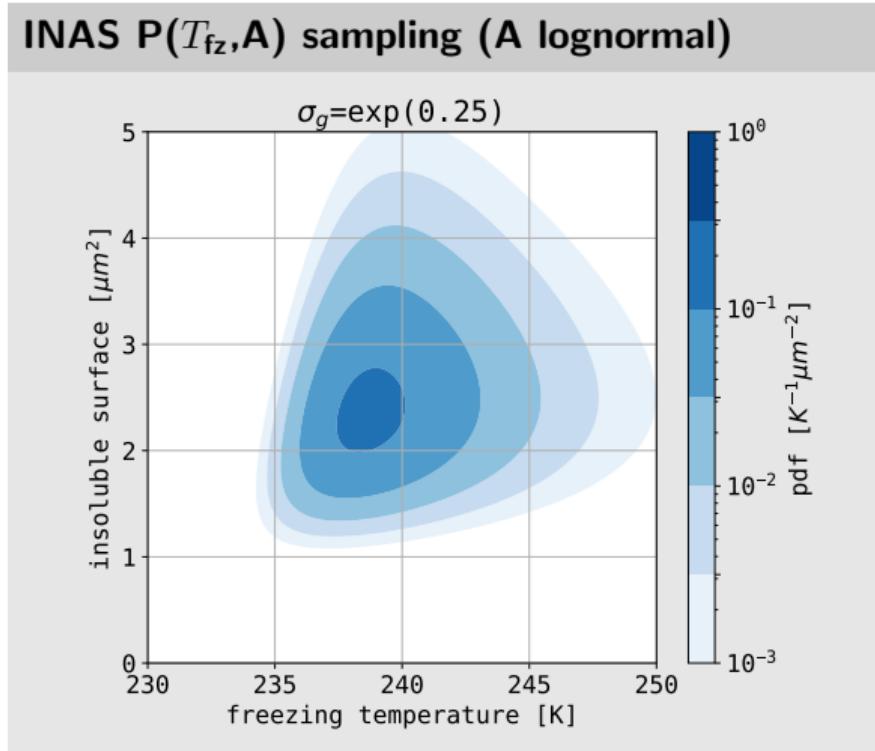
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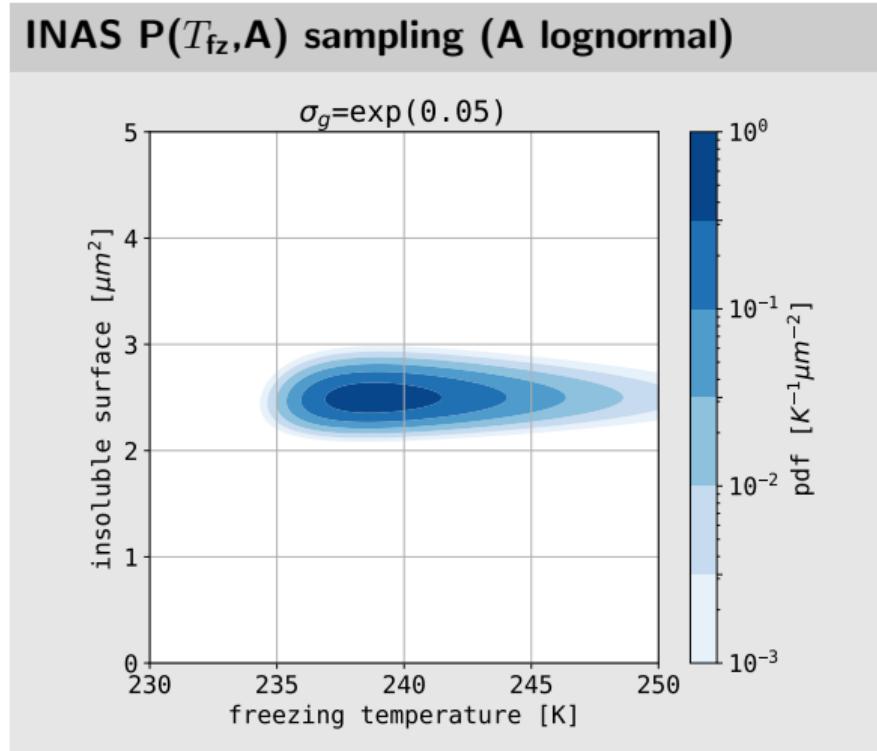
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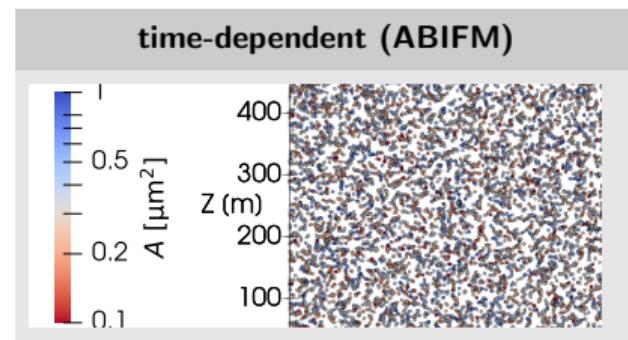
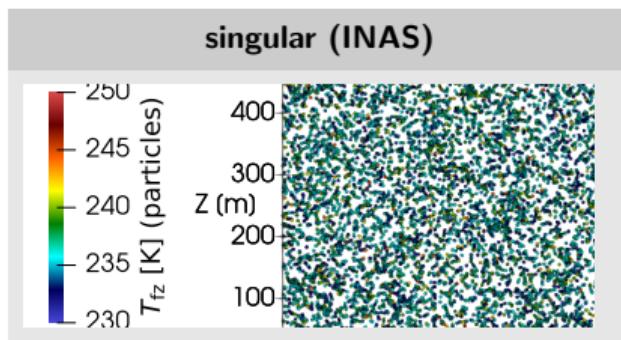
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**time-dependent:**  $A$  as attribute (randomly sampled from the same lognormal)

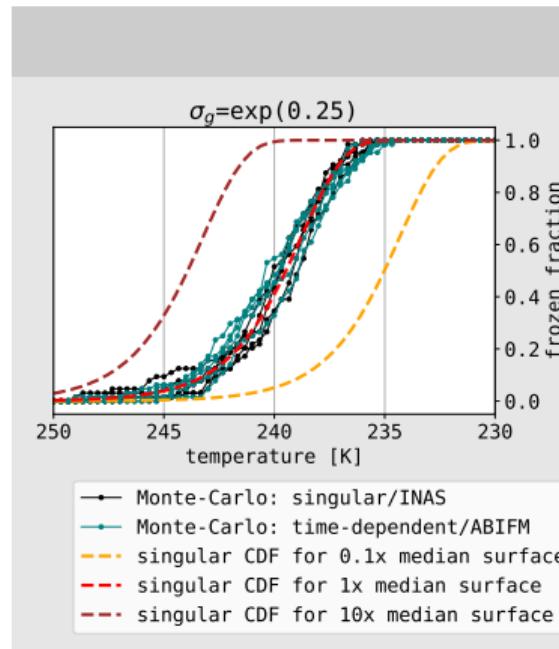
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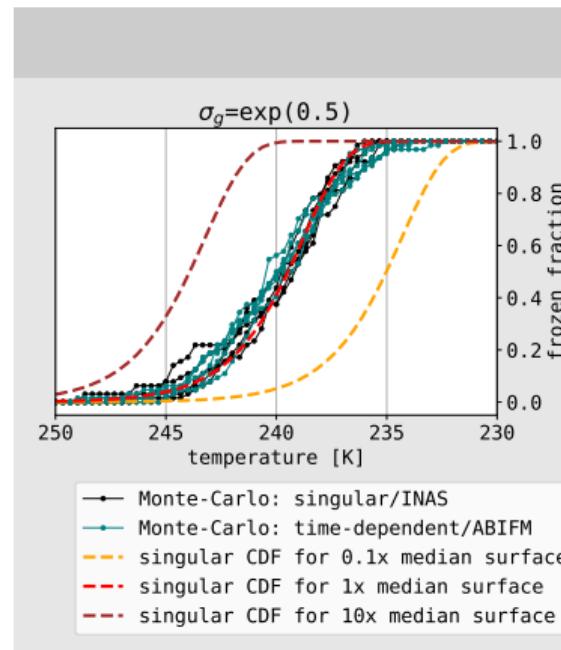
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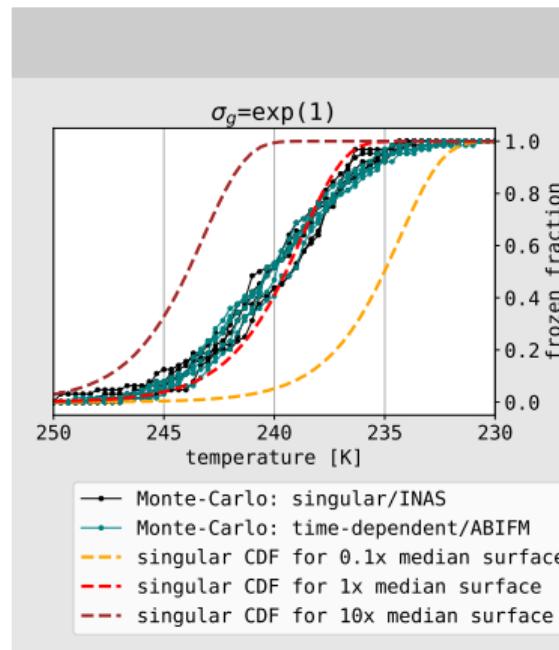
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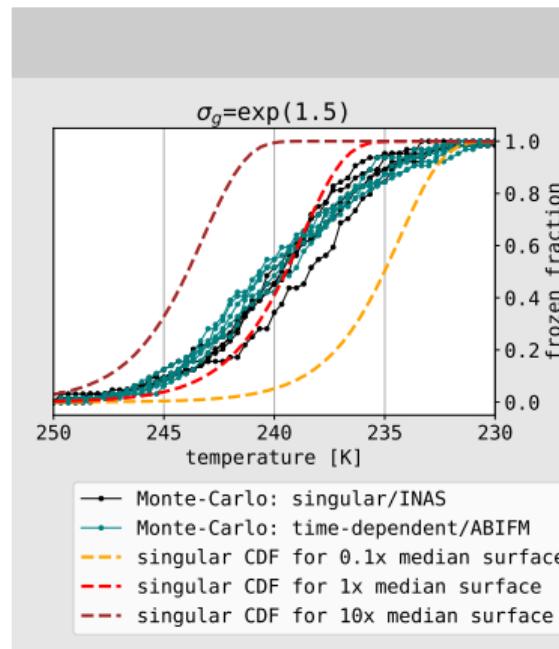
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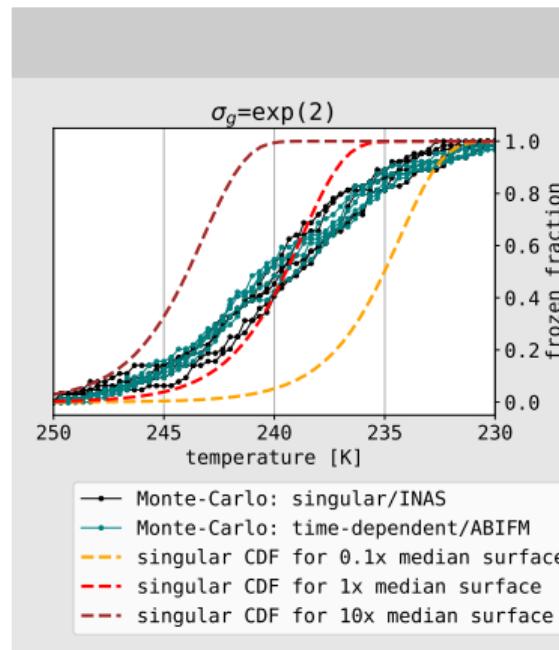
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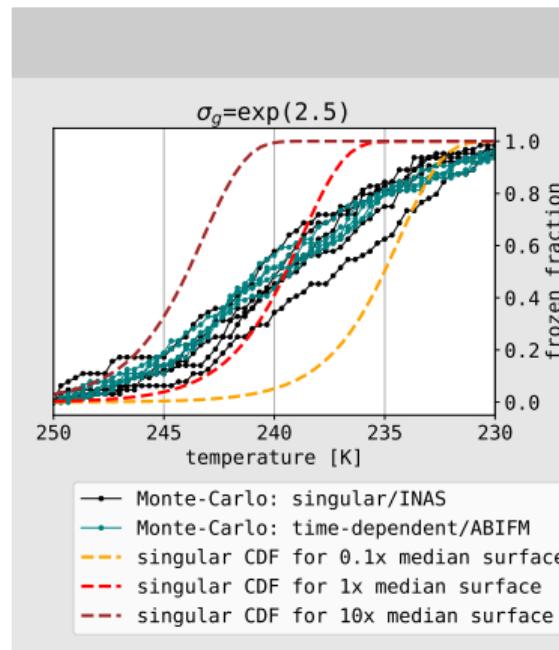
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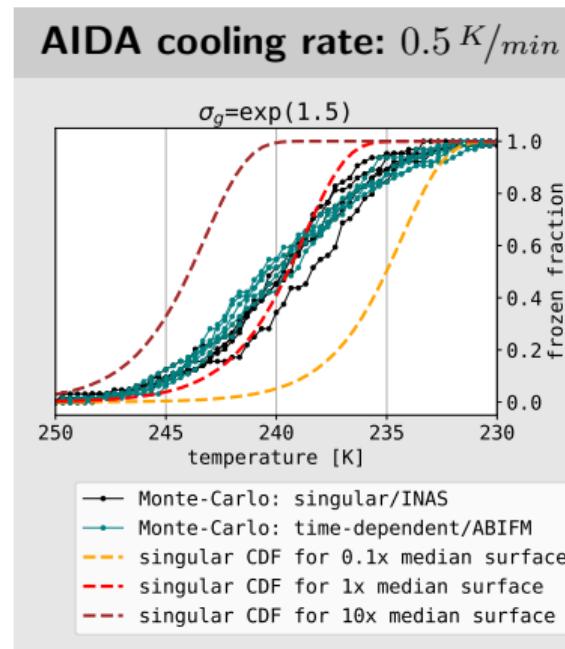
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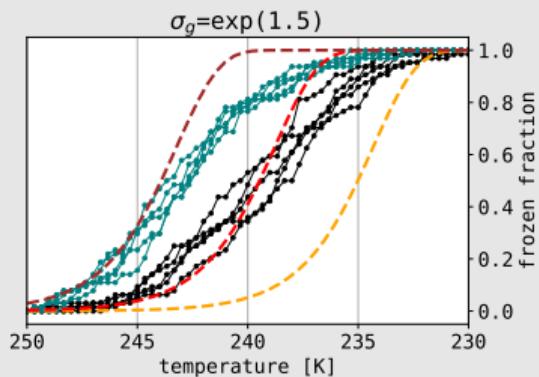


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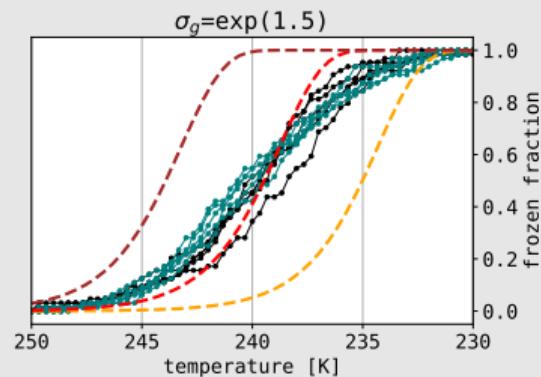
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**cooling rate:**  $0.1 \text{ K/min}$



**AIDA cooling rate:  $0.5 \text{ K/min}$**



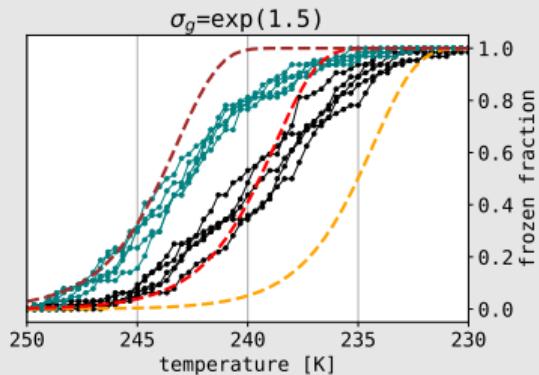
- Monte-Carlo: singular/INAS
- Monte-Carlo: time-dependent/ABIFM
- singular CDF for  $0.1x$  median surface
- - - singular CDF for  $1x$  median surface
- - - singular CDF for  $10x$  median surface

**singular:** INAS  $T_{fz}$  as attribute; initialisation by random sampling from  $P(T_{fz}, A)$  with lognormal  $A$  ( $A$  is not an attribute, initialisation only); freezing if  $T(t) < T_{fz}(t = 0)$

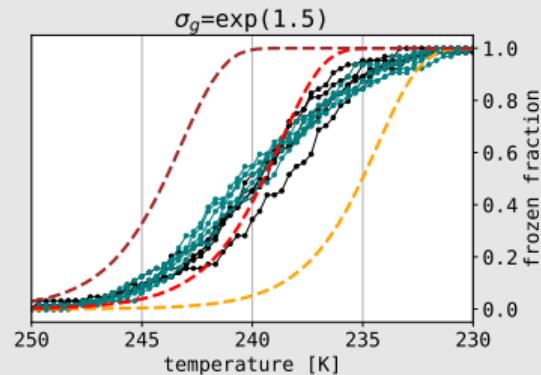
**time-dependent:**  $A$  as attribute (randomly sampled from the same lognormal)

Monte-Carlo freezing trigger using  $P(J_{\text{het}}(T(t)))$

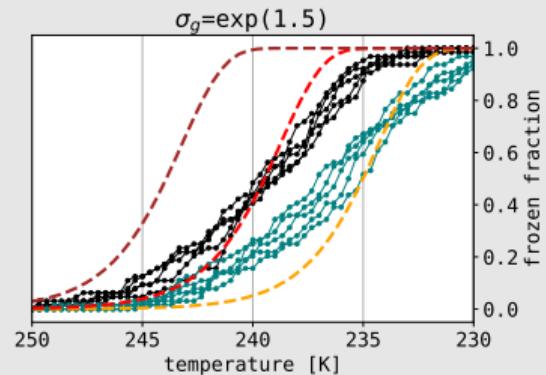
**cooling rate:  $0.1 \text{ K/min}$**



**AIDA cooling rate:  $0.5 \text{ K/min}$**



**cooling rate:  $2.5 \text{ K/min}$**



- Monte-Carlo: singular/INAS
- Monte-Carlo: time-dependent/ABIFM
- singular CDF for  $0.1x$  median surface
- singular CDF for  $1x$  median surface
- singular CDF for  $10x$  median surface

## theory (in modern notation)

(Bigg '53, Langham & Mason '58, Carte '59, Marshall '61)

Poisson counting process with rate  $r$ :

$$P^*(k \text{ events in time } t) = \frac{(rt)^k \exp(-rt)}{k!}$$

$$P(\text{one or more events in time } t) = 1 - P^*(k = 0, t)$$

$$\ln(1 - P) = -rt$$

introducing  $J_{\text{het}}(T)$ ,  $T(t)$  and INP surface  $A$ :

$$\ln(1 - P(A, t)) = -A \underbrace{\int_0^t J_{\text{het}}(T(t')) dt'}_{I(T)}$$

**INAS:**  $I(T) = n_s(T) = \exp(a \cdot (T - T_0{}^\circ C) + b)$

experimental  $n_s(T)$  fits: e.g., Niemand et al. 2012

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$$\ln(1 - P(A, t)) = -\frac{A}{c} \int_{T_0}^{T_0 + ct} J_{\text{het}}(T') dT' = -A \cdot I(T)$$

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$\underbrace{\hspace{10em}}$   
 $I(T)$

INAS:  $I(T) = n_s(T) = \exp(a \cdot (T - T_0 {}^\circ C) + b)$

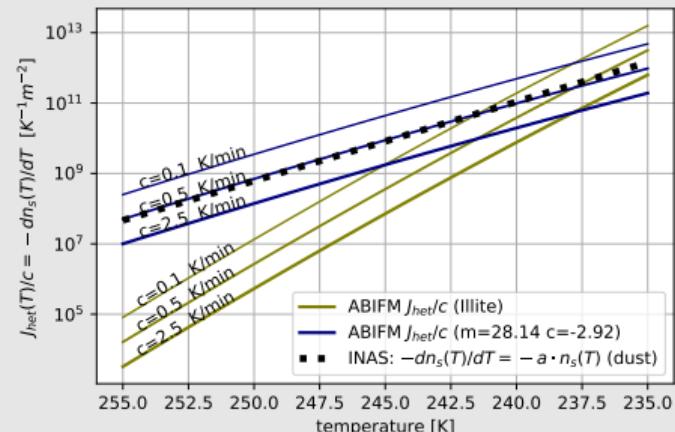
experimental  $n_s(T)$  fits: e.g., Niemand et al. 2012

for a constant cooling rate  $c = dT/dt$ :

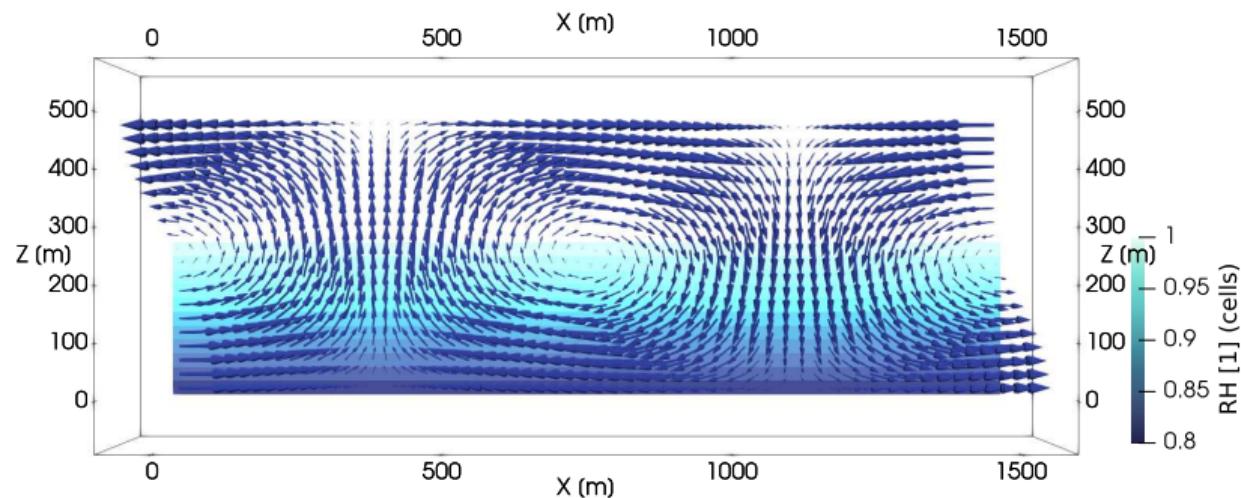
$$\ln(1 - P(A, t)) = -\frac{A}{c} \int_{T_0}^{T_0+ct} J_{\text{het}}(T') dT' = -A \cdot I(T)$$

$$\frac{dn_s(T)}{dT} = a \cdot n_s(T) = -\frac{1}{c} J_{\text{het}}(T)$$

experimental fits: INAS  $n_s$  (Niemand et al. '12)  
ABIFM  $J_{\text{het}}$  (Knopf & Alpert '13)

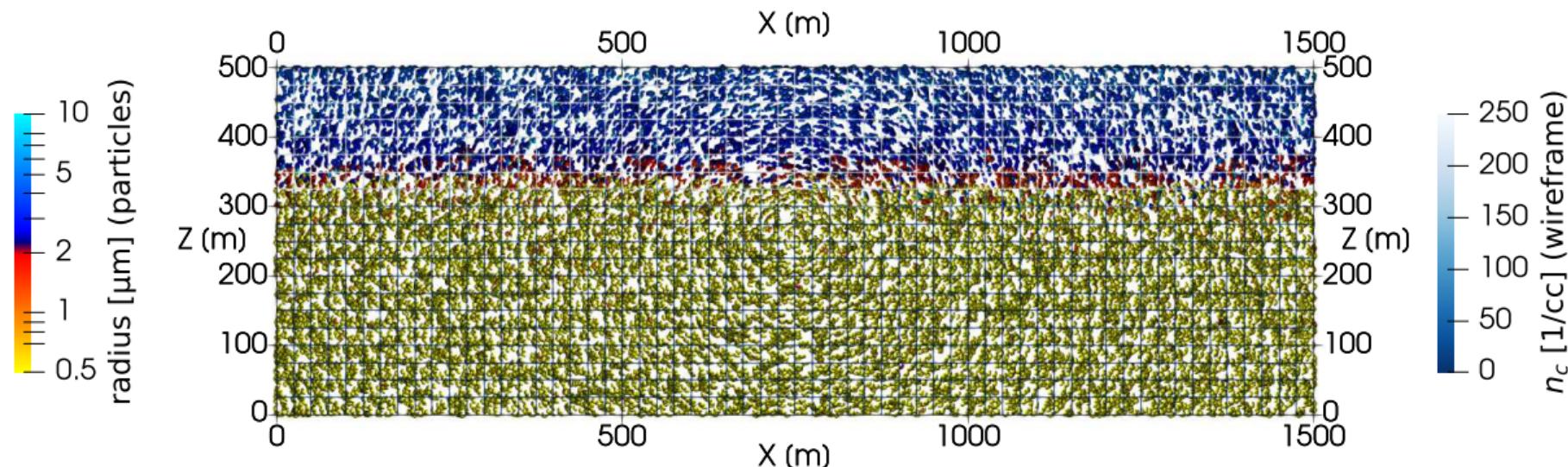


## immersion freezing: singular vs. time-dependent in flow-coupled simulation



# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 60 s (spin-up till 600.0 s)



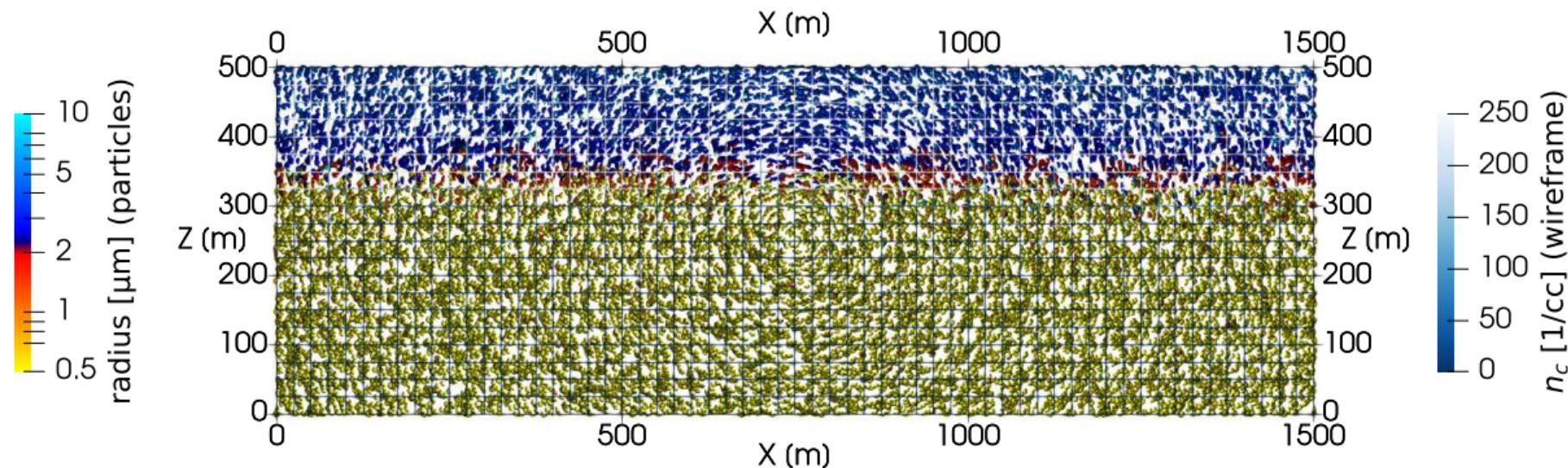
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 90 s (spin-up till 600.0 s)



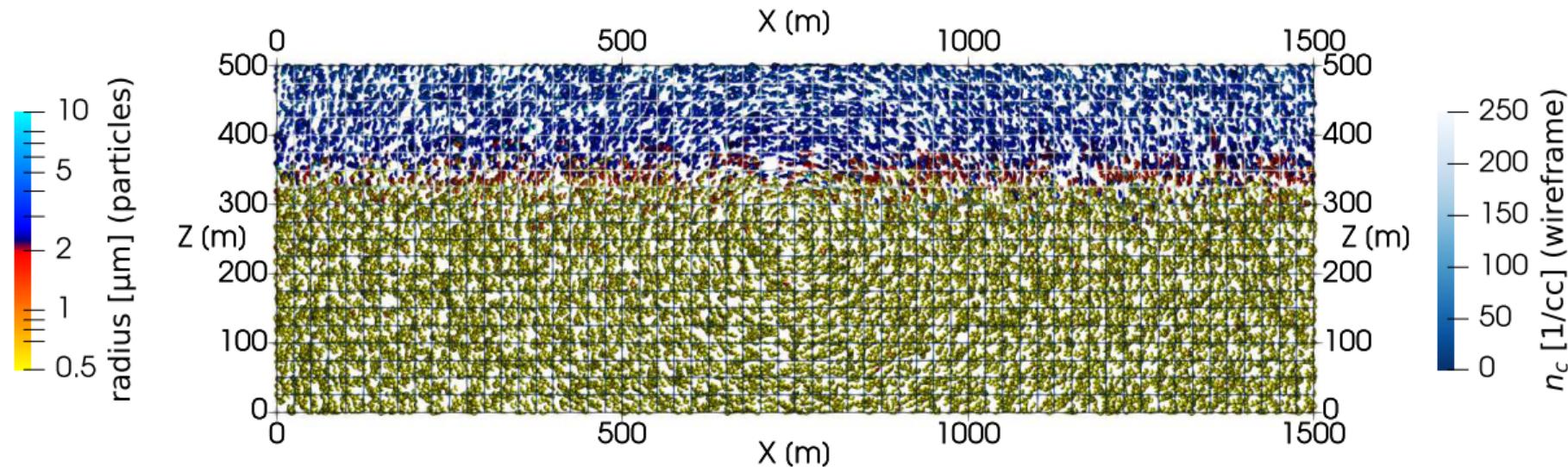
16+16 super-particles/cell for INP-rich + INP-free particles

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spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 120 s (spin-up till 600.0 s)



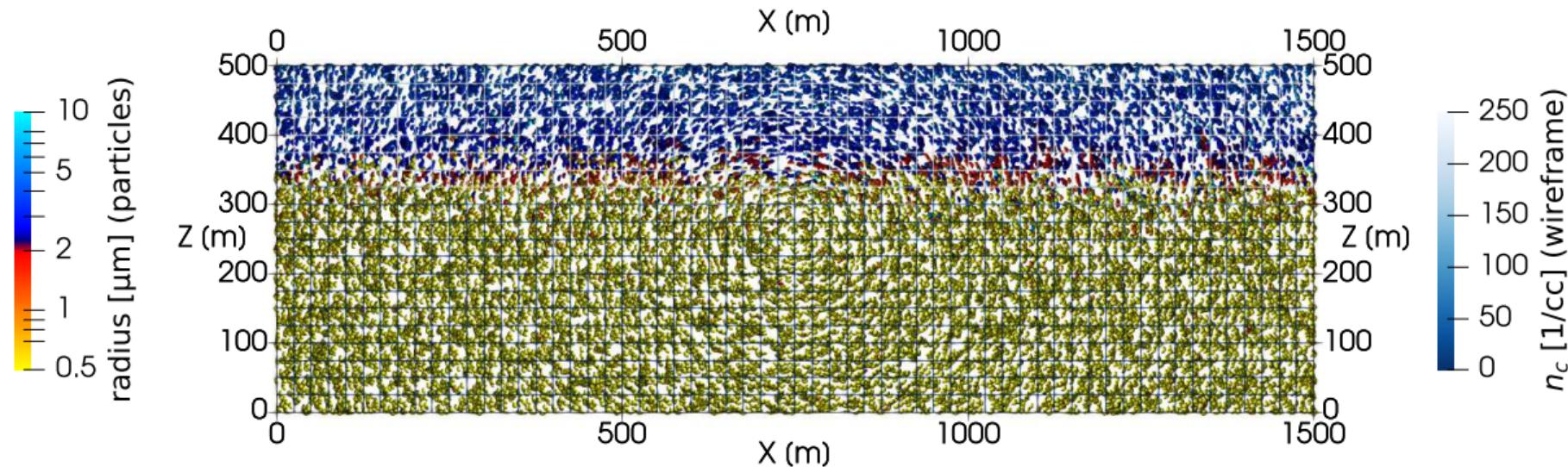
16+16 super-particles/cell for INP-rich + INP-free particles

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spin-up = freezing off; subsequently frozen particles act as tracers

## immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 150 s (spin-up till 600.0 s)



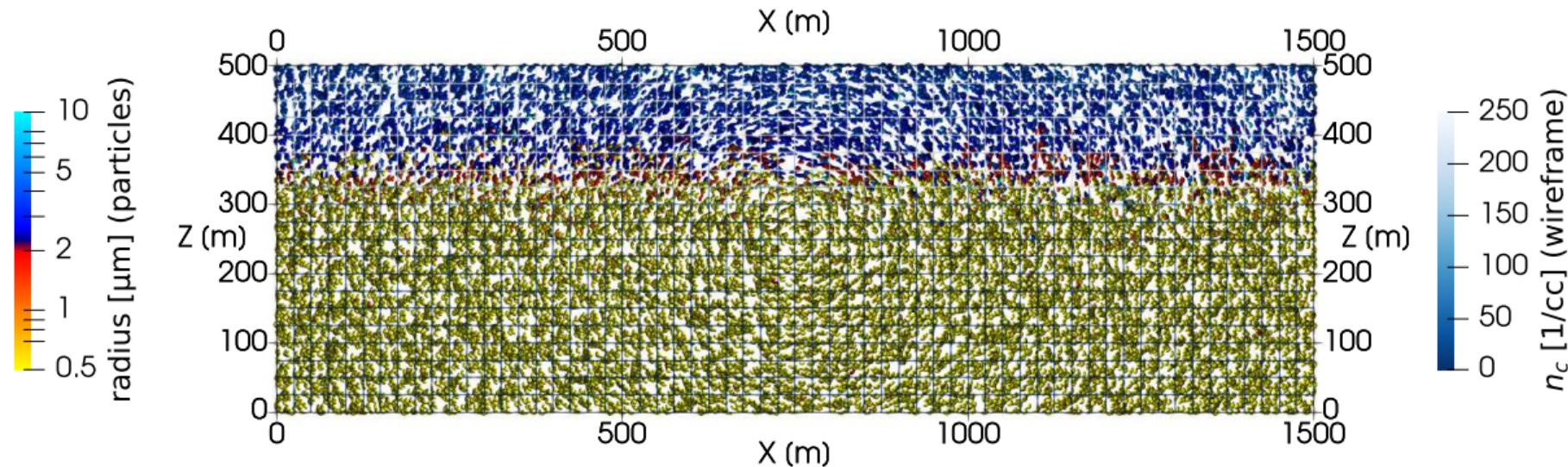
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

## immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 180 s (spin-up till 600.0 s)



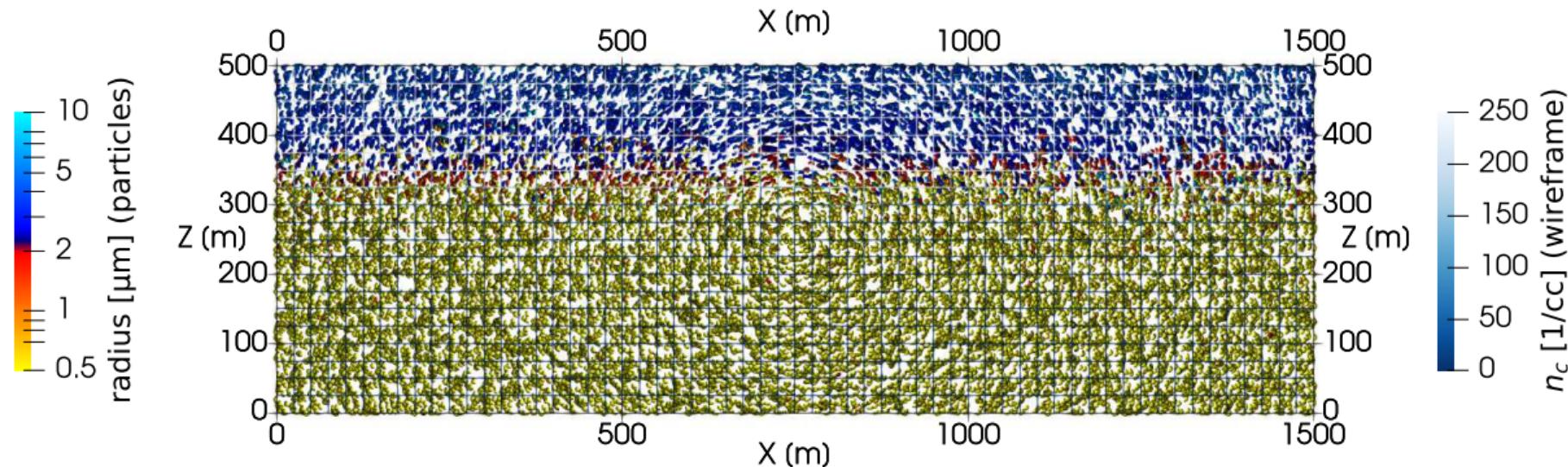
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 210 s (spin-up till 600.0 s)



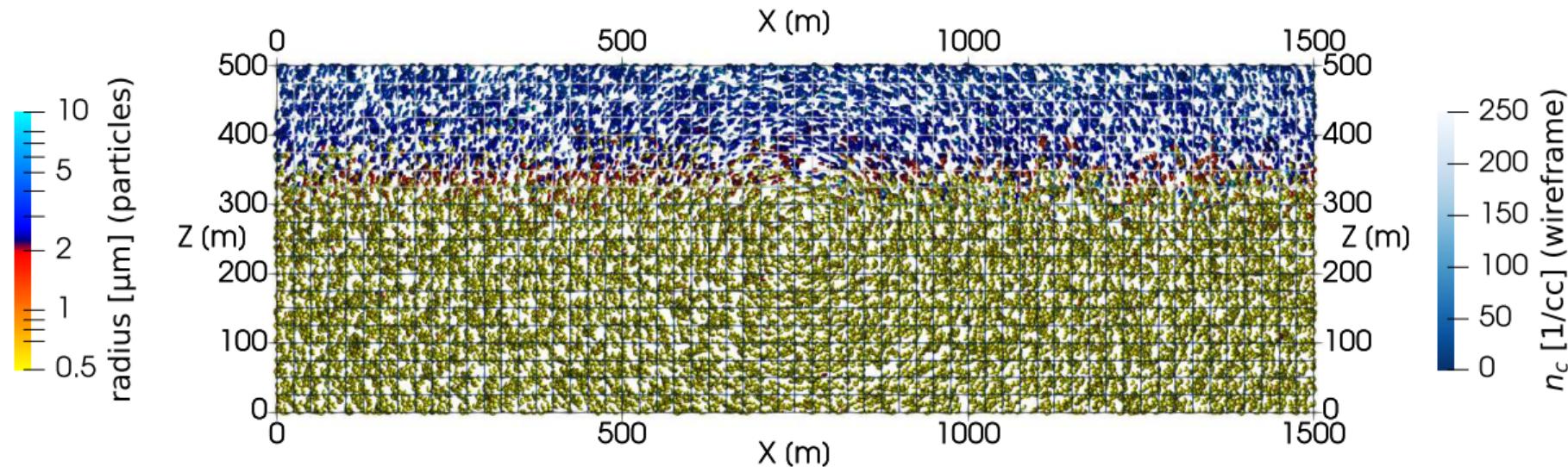
16+16 super-particles/cell for INP-rich + INP-free particles

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spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 240 s (spin-up till 600.0 s)



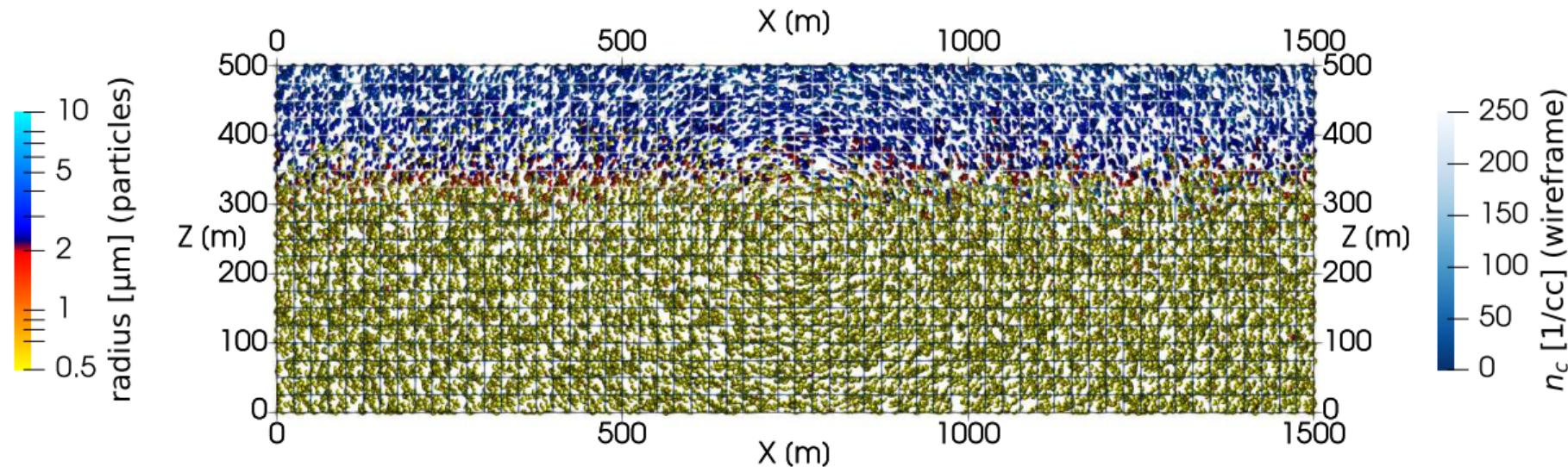
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)    $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 270 s (spin-up till 600.0 s)



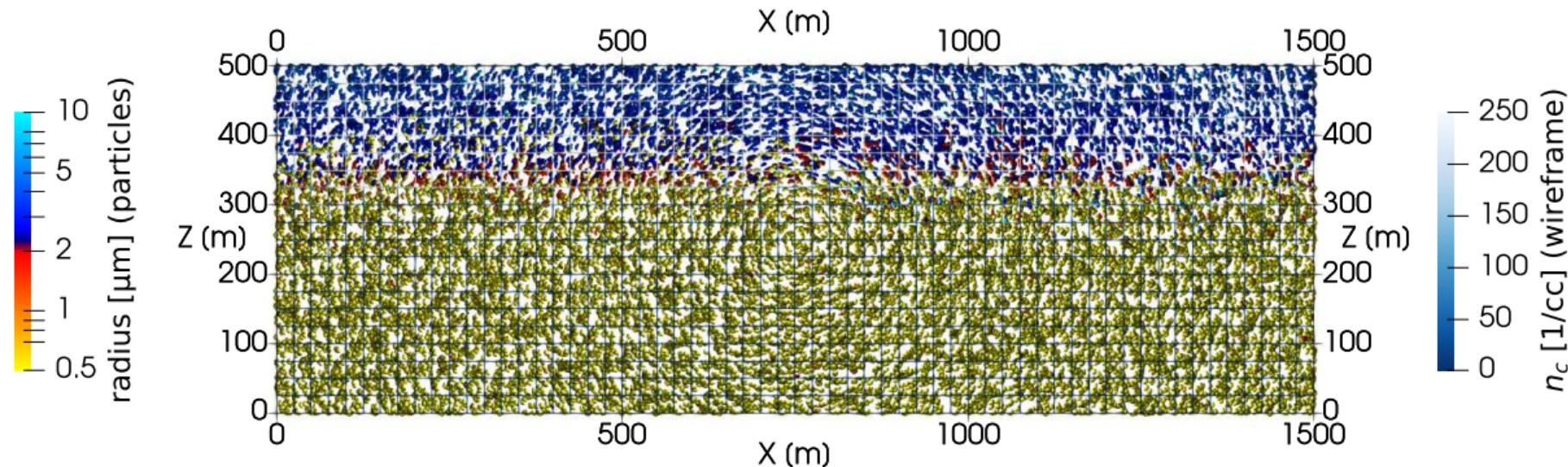
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 300 s (spin-up till 600.0 s)



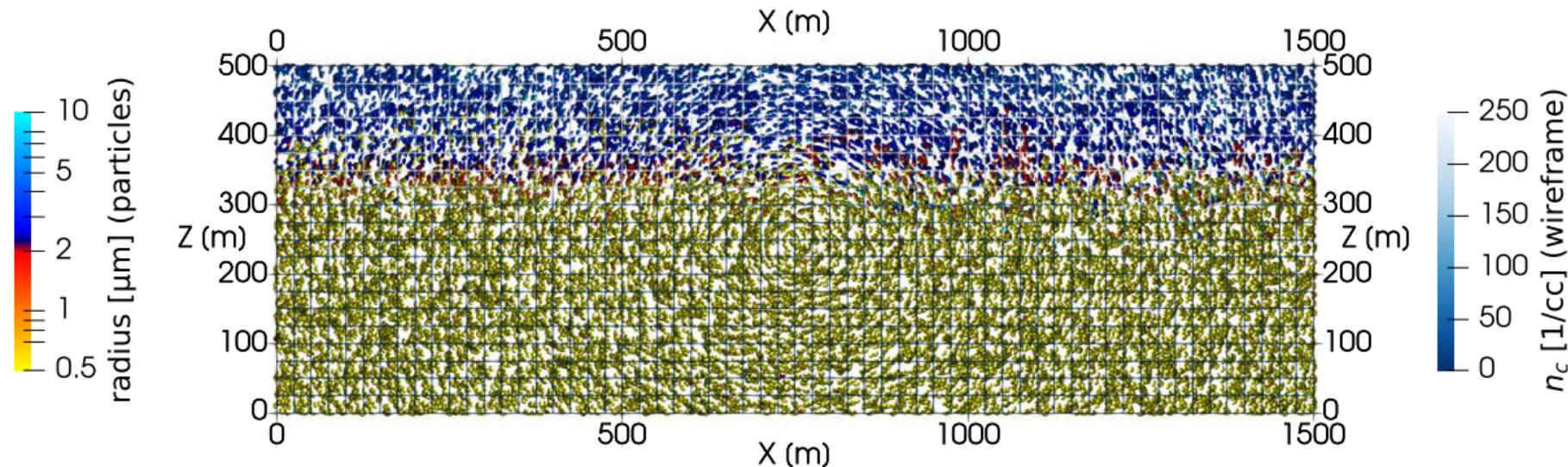
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 330 s (spin-up till 600.0 s)



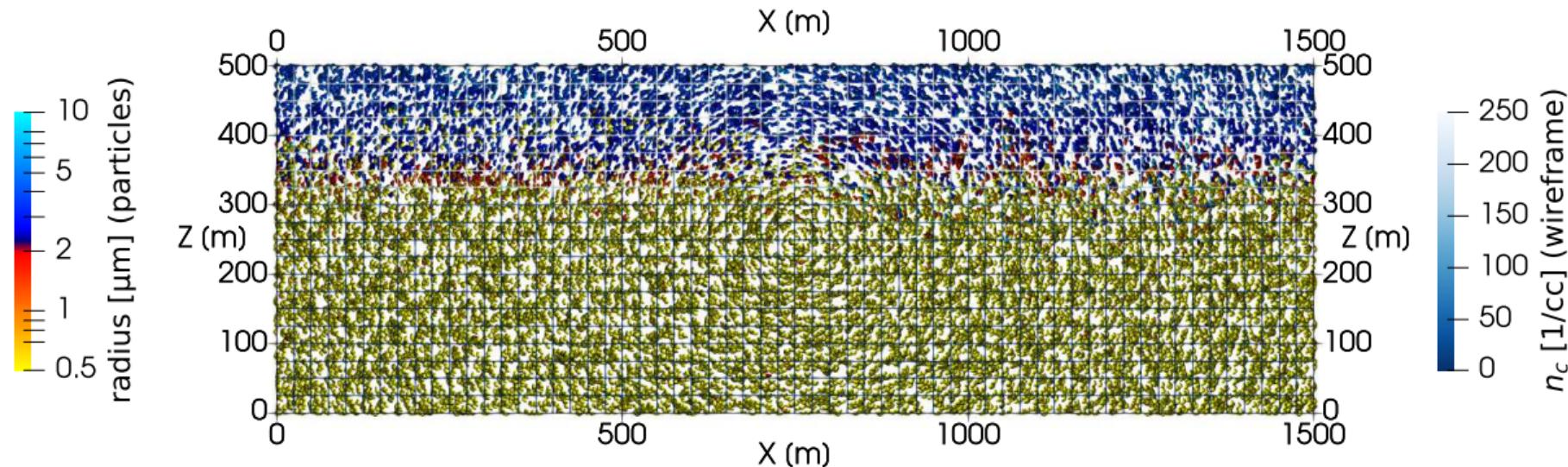
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 360 s (spin-up till 600.0 s)



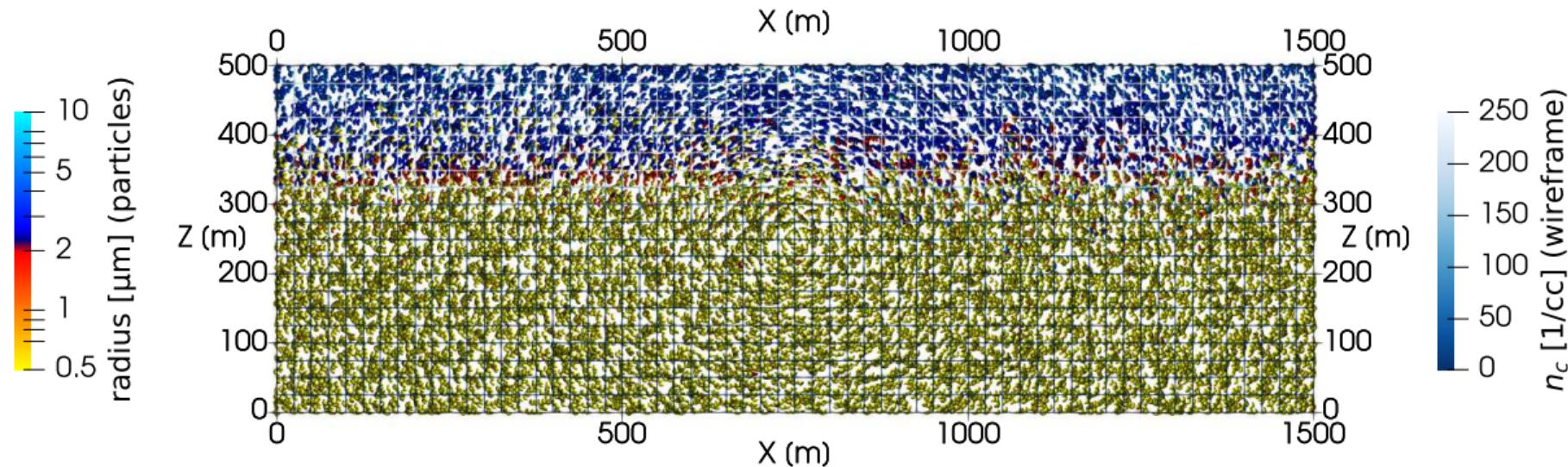
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spin-up = freezing off; subsequently frozen particles act as tracers

## immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 390 s (spin-up till 600.0 s)



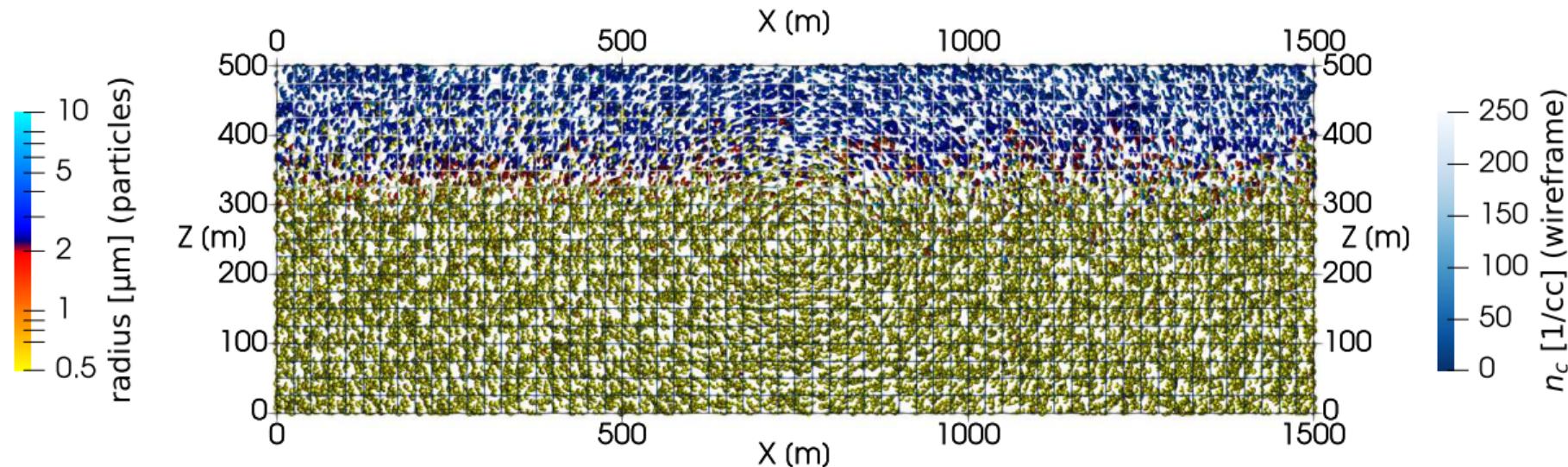
16+16 super-particles/cell for INP-rich + INP-free particles

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spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 420 s (spin-up till 600.0 s)



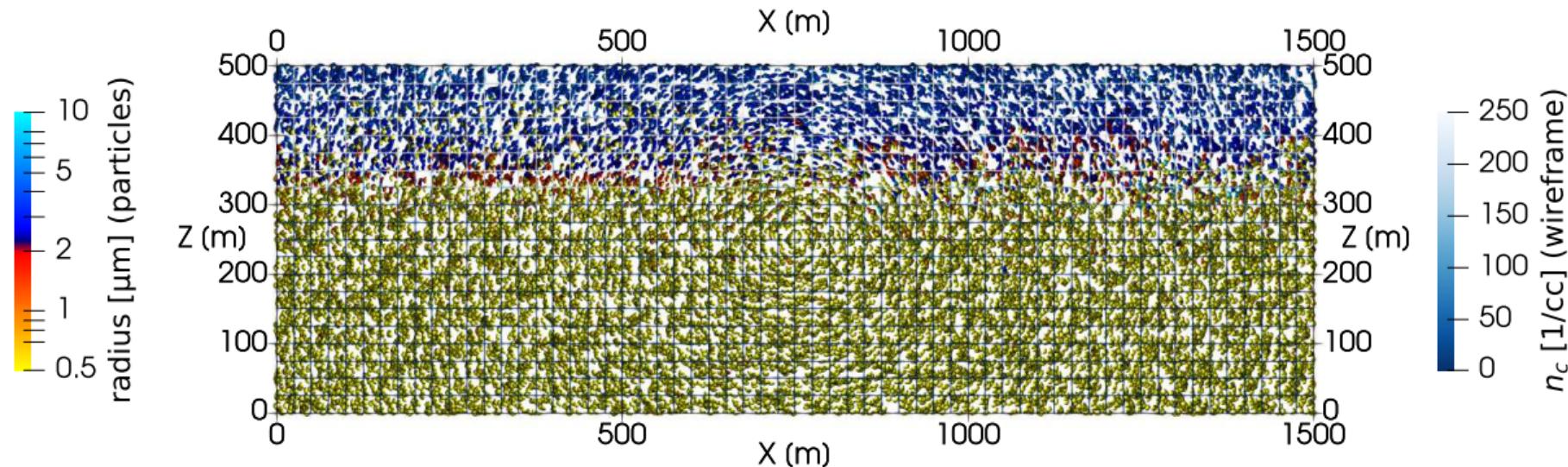
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 450 s (spin-up till 600.0 s)



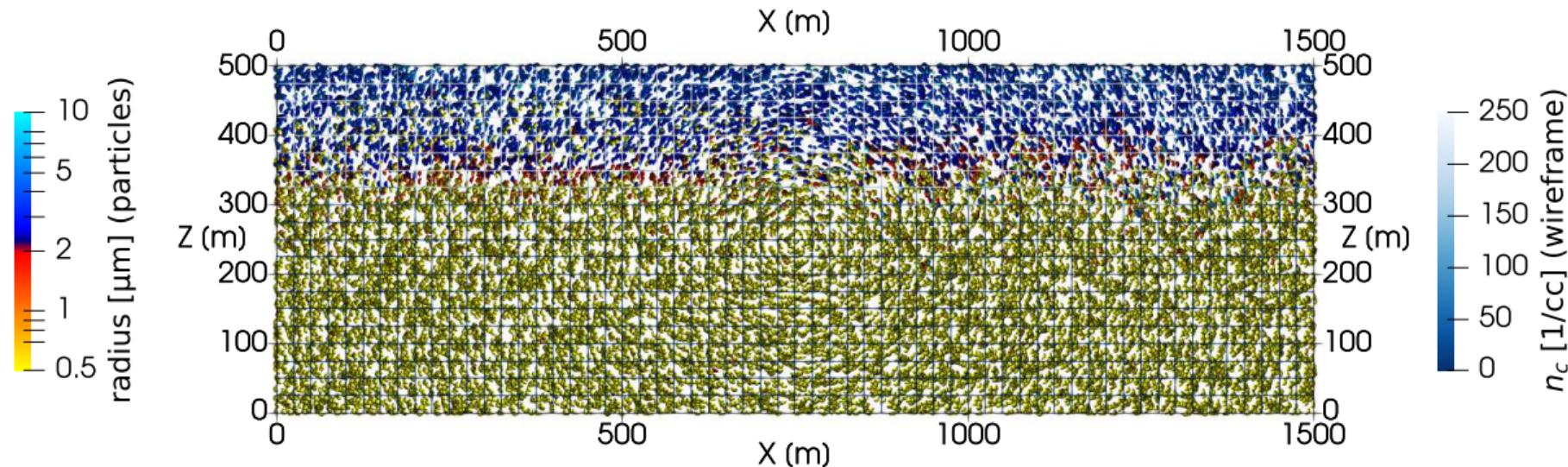
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 480 s (spin-up till 600.0 s)



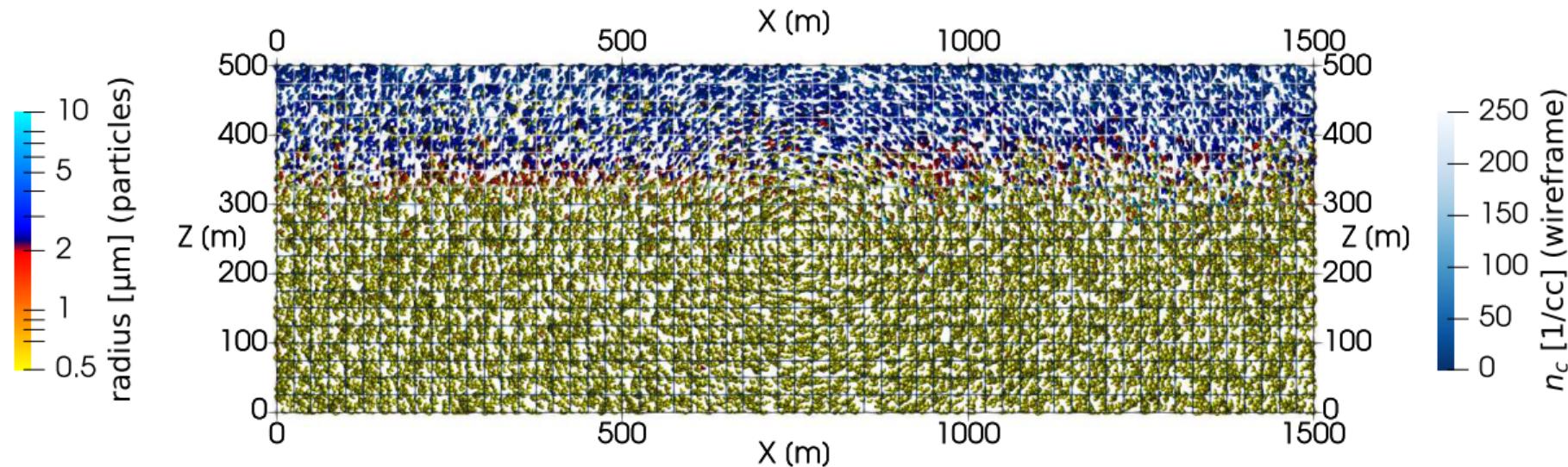
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 510 s (spin-up till 600.0 s)



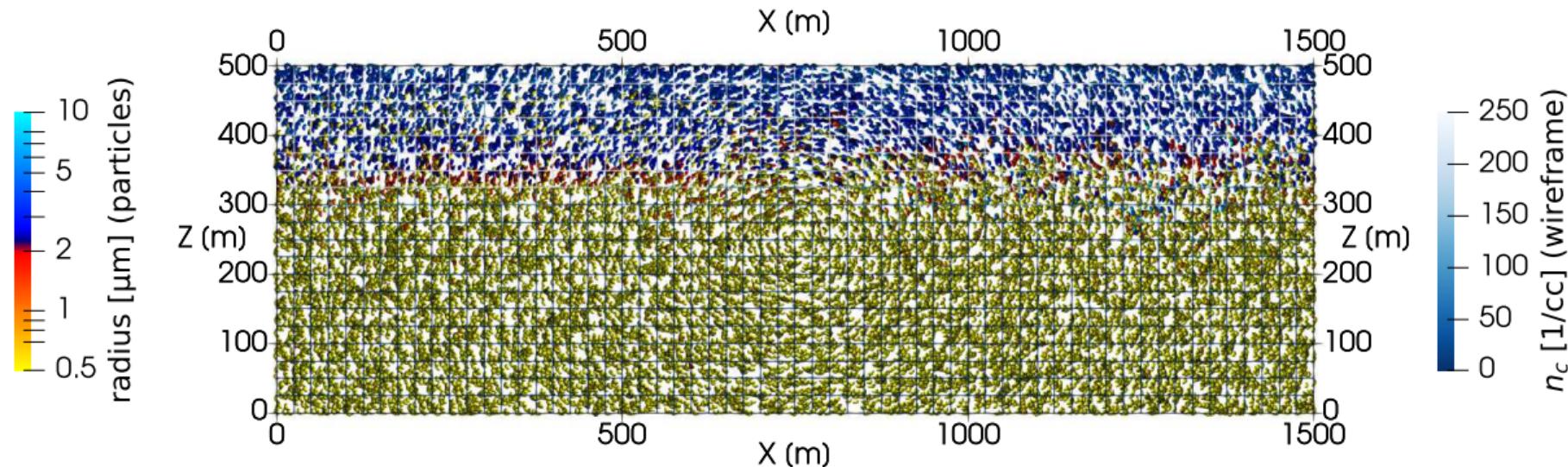
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$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)    $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 540 s (spin-up till 600.0 s)



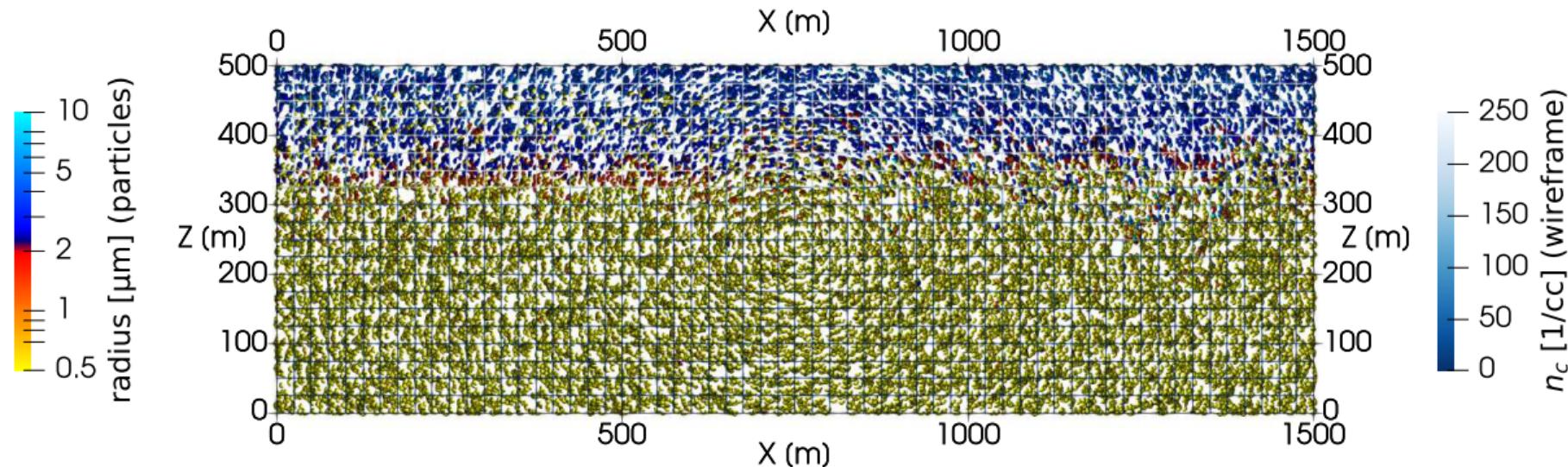
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spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 570 s (spin-up till 600.0 s)



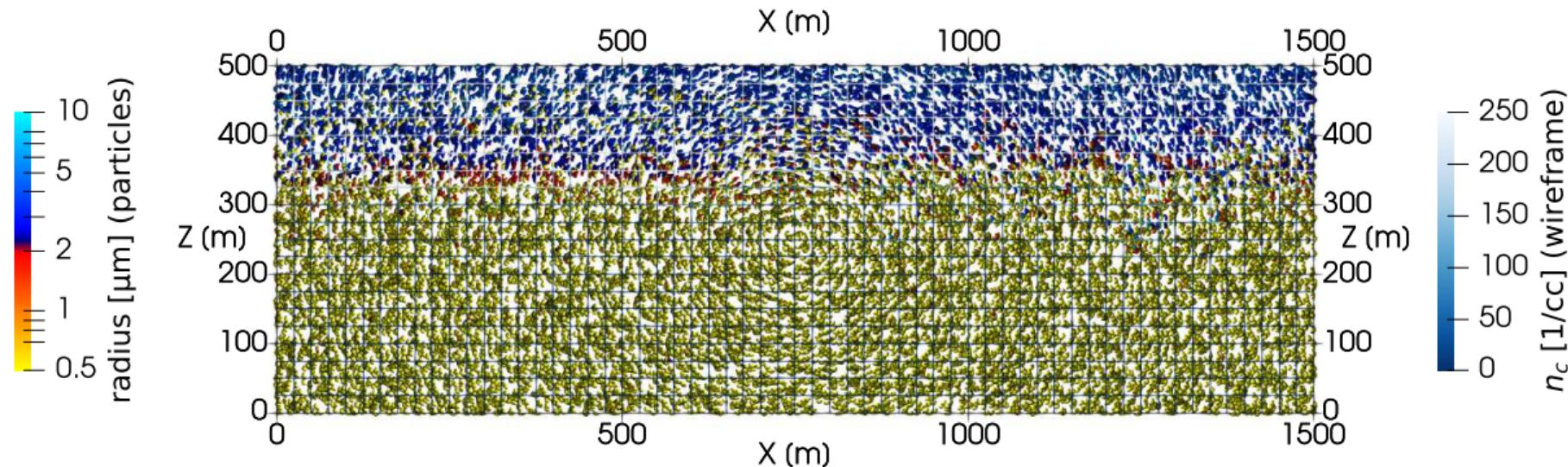
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spin-up = freezing off; subsequently frozen particles act as tracers

## immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 600 s (spin-up till 600.0 s)



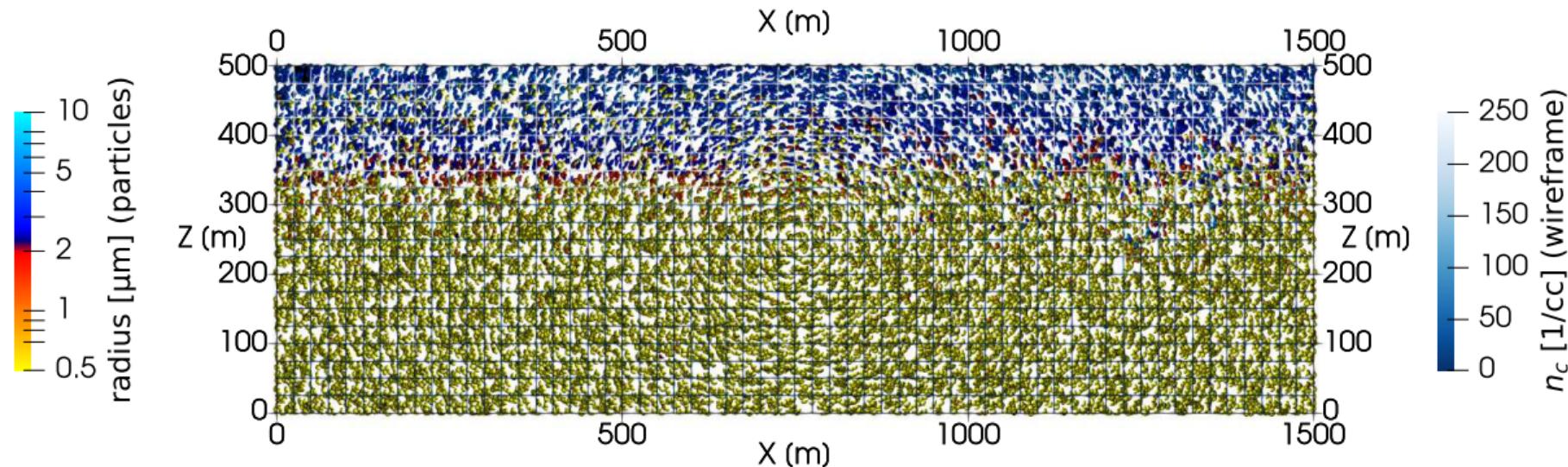
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 630 s (spin-up till 600.0 s)



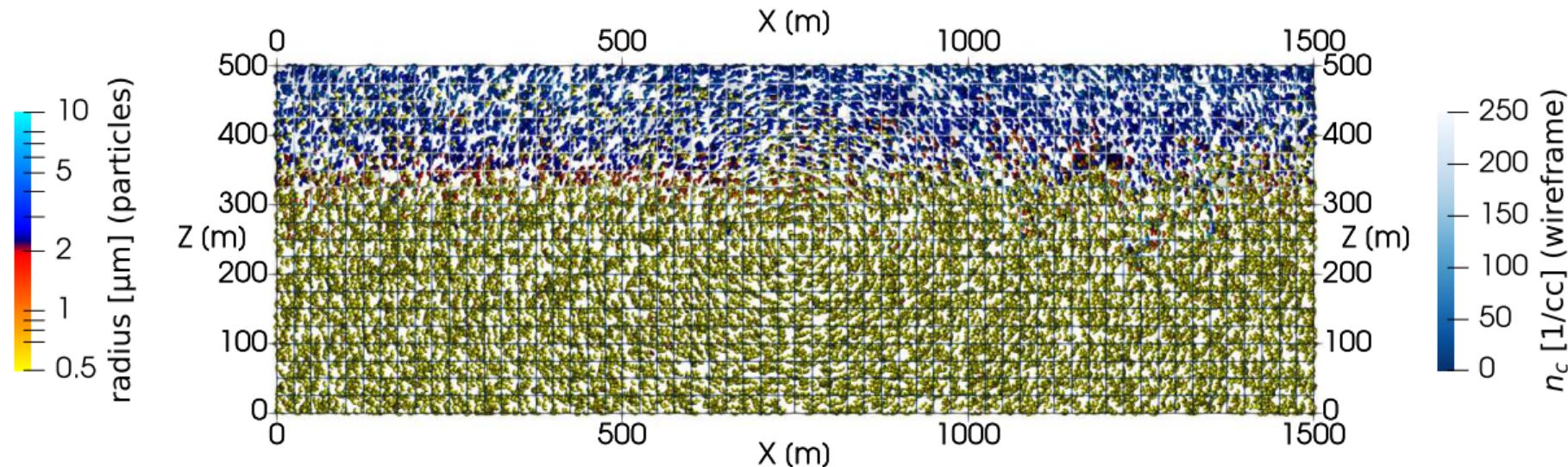
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 660 s (spin-up till 600.0 s)



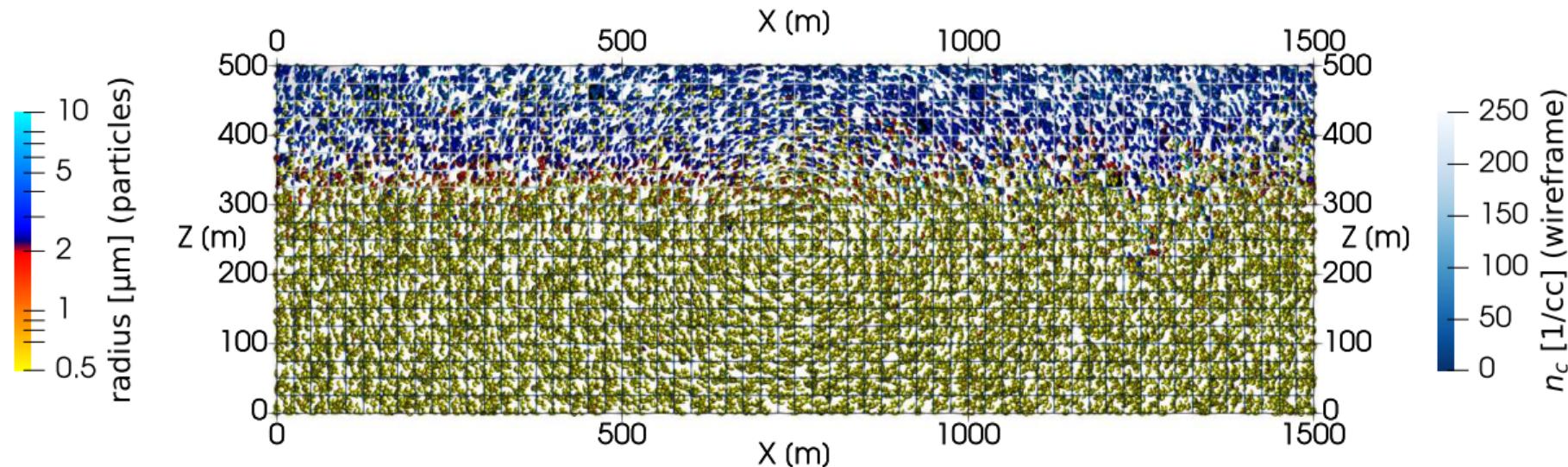
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 690 s (spin-up till 600.0 s)



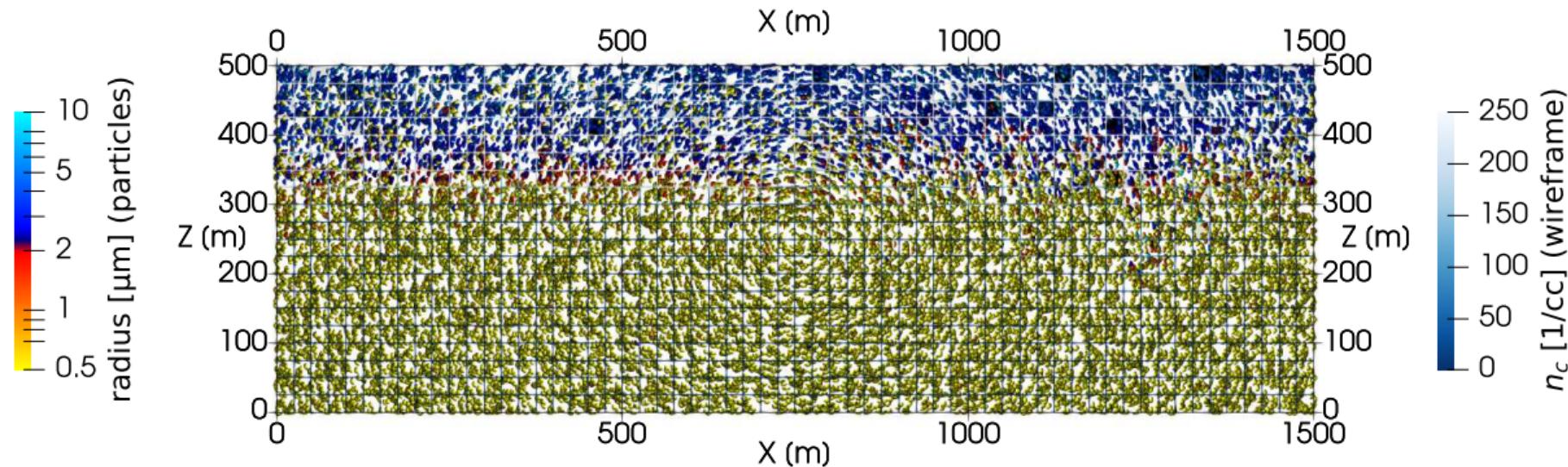
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 720 s (spin-up till 600.0 s)



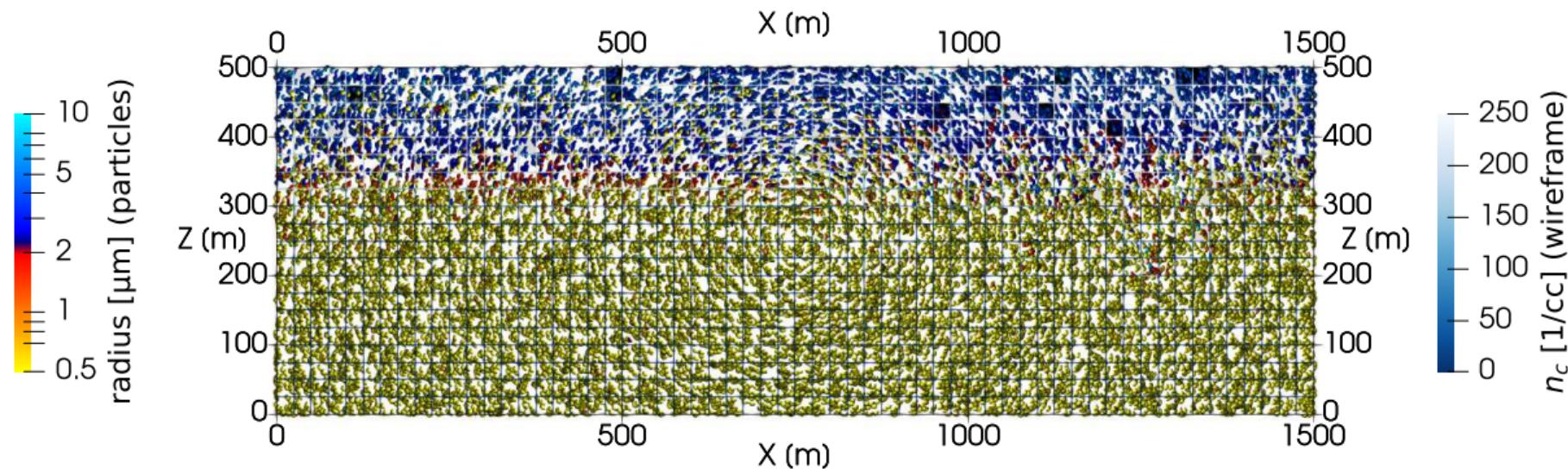
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 750 s (spin-up till 600.0 s)



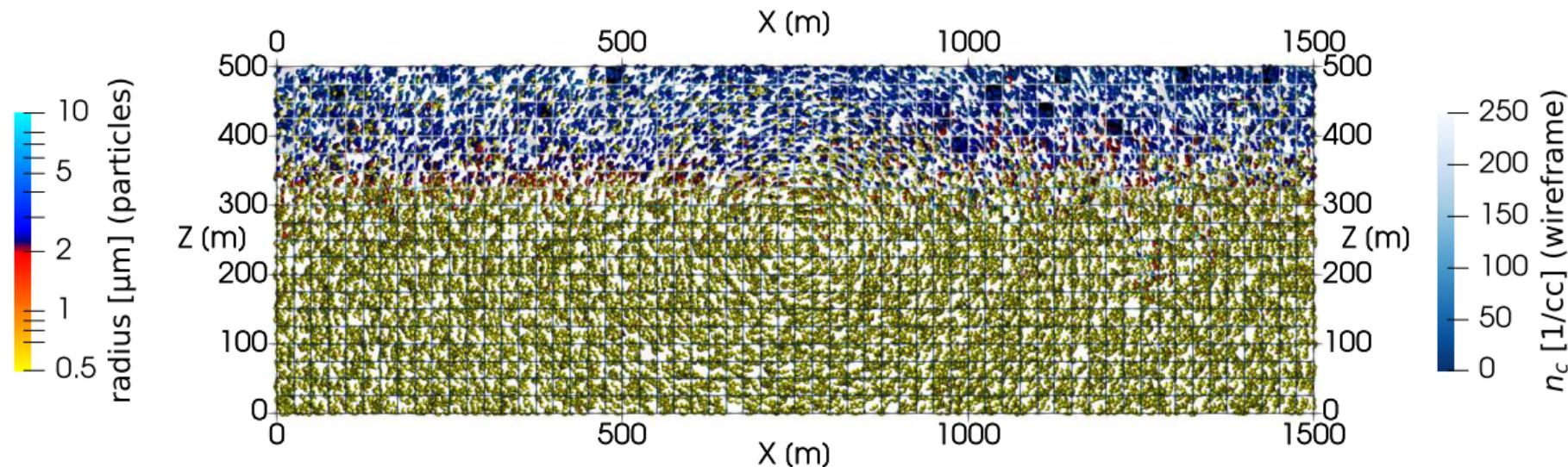
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 780 s (spin-up till 600.0 s)



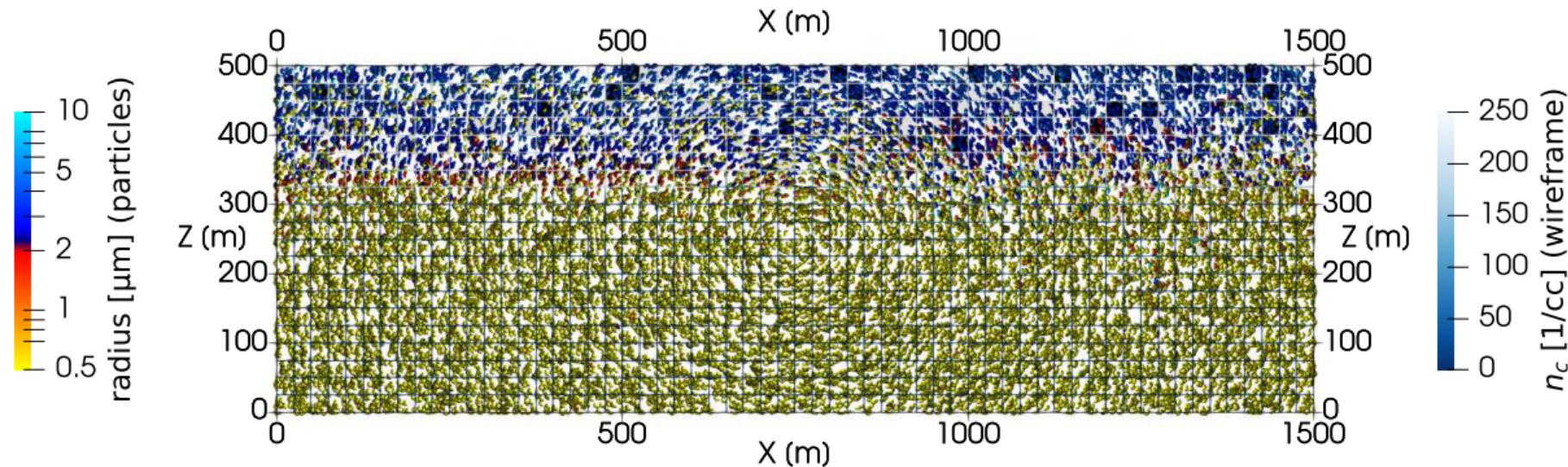
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 810 s (spin-up till 600.0 s)



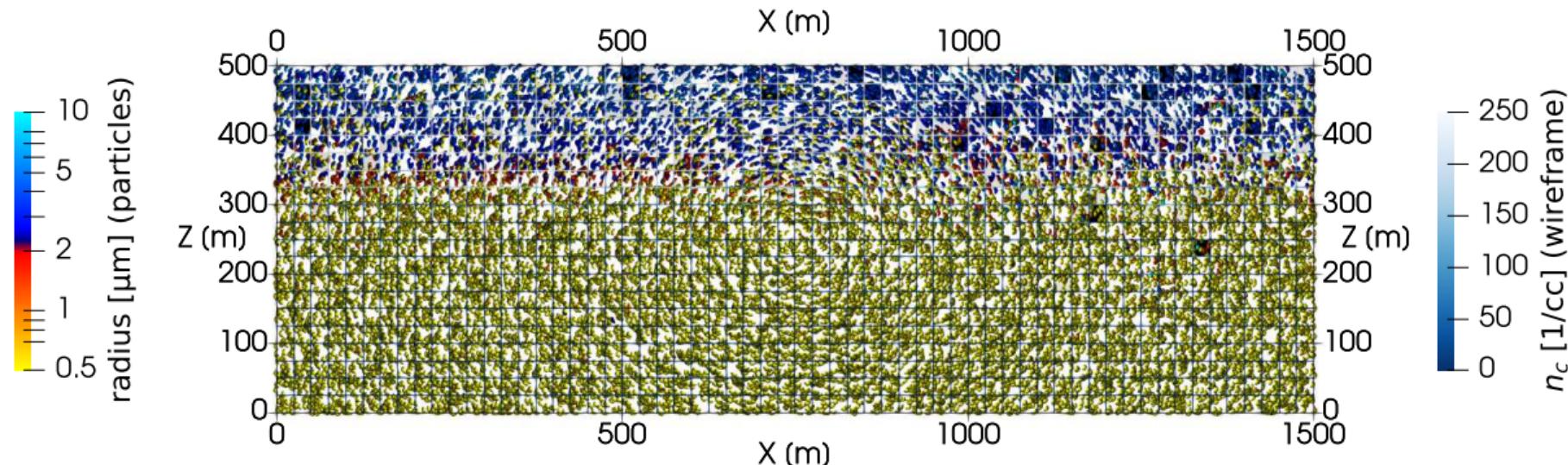
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 840 s (spin-up till 600.0 s)



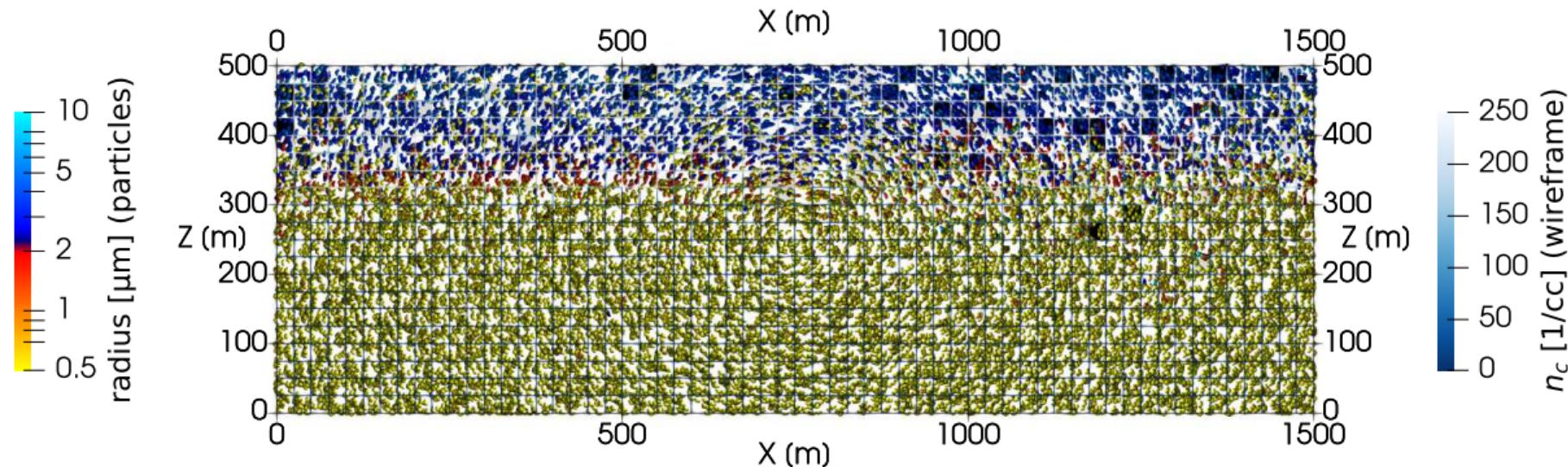
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 870 s (spin-up till 600.0 s)



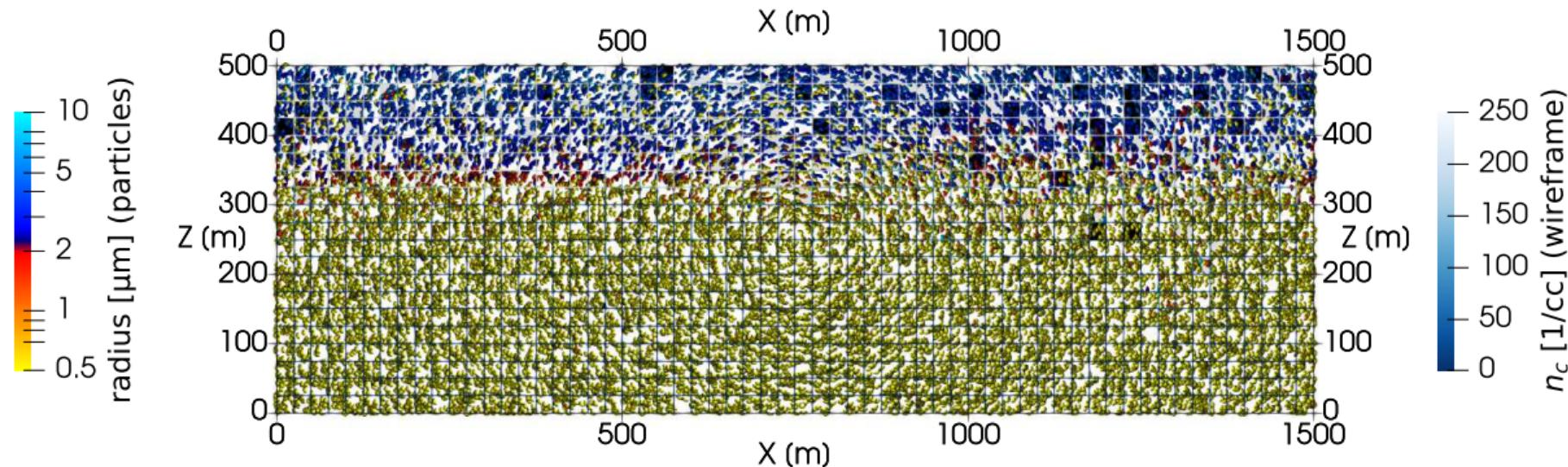
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$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 900 s (spin-up till 600.0 s)



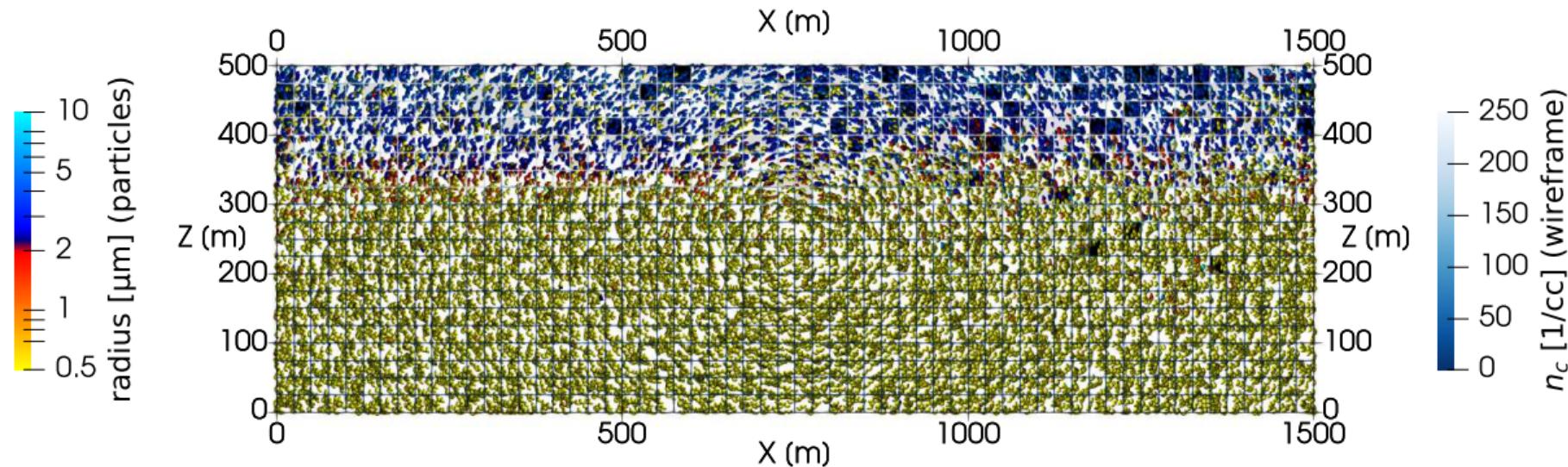
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 930 s (spin-up till 600.0 s)



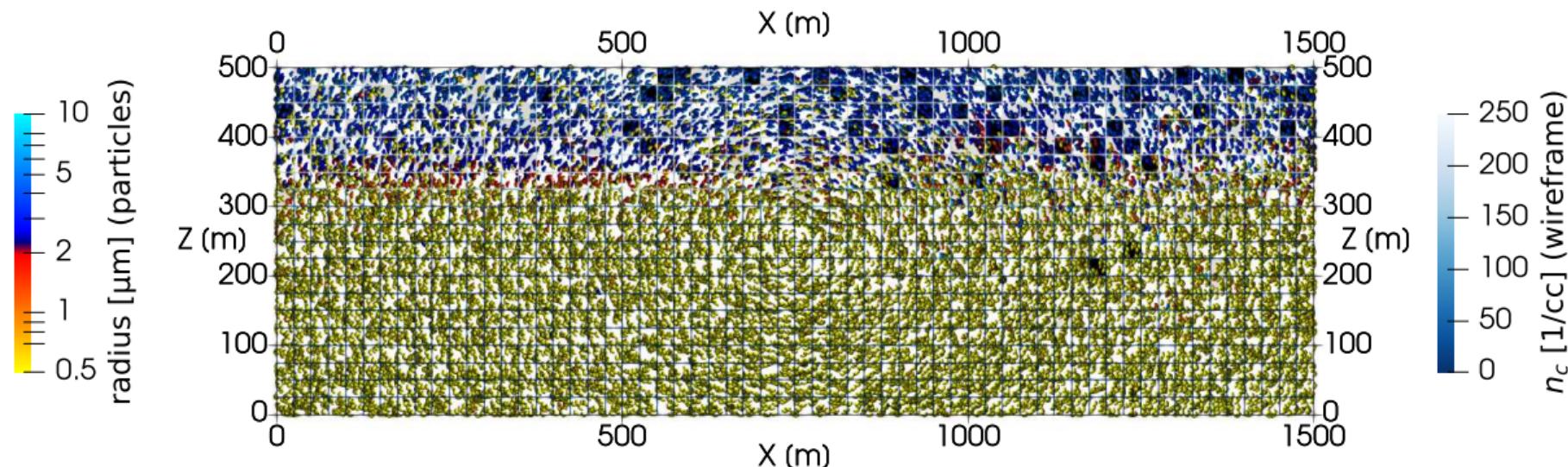
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 960 s (spin-up till 600.0 s)



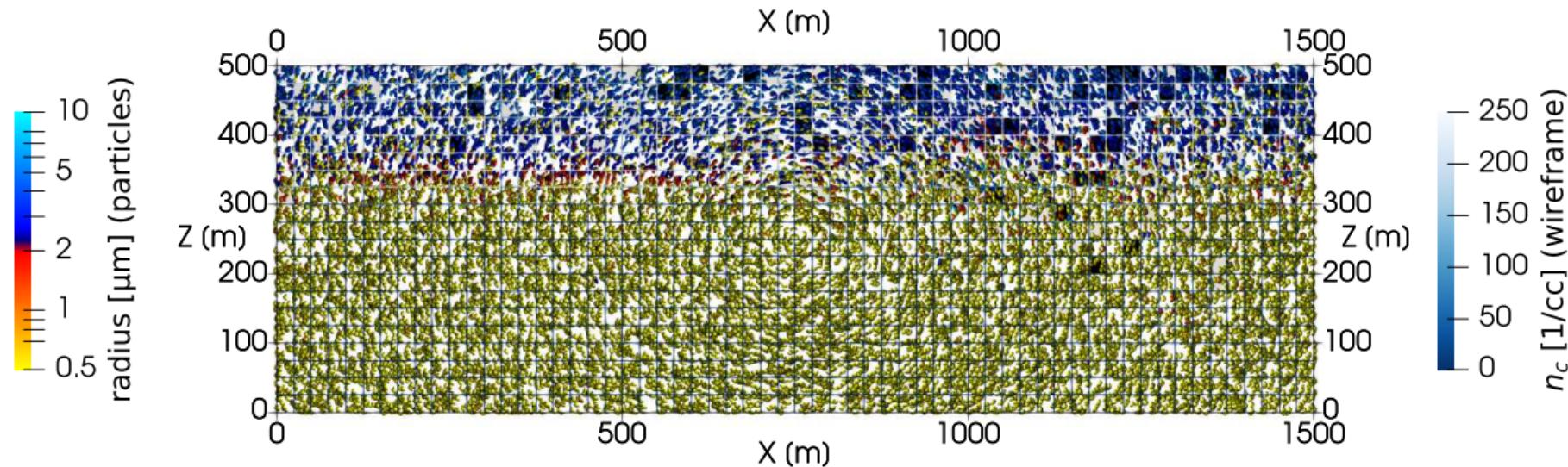
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 990 s (spin-up till 600.0 s)



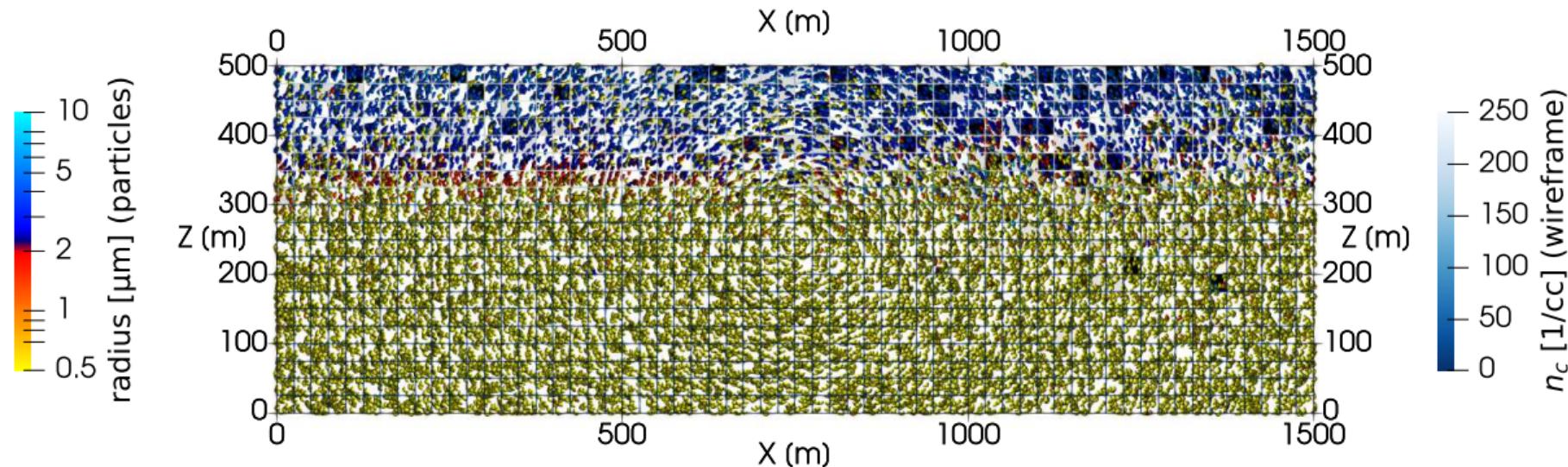
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 1020 s (spin-up till 600.0 s)



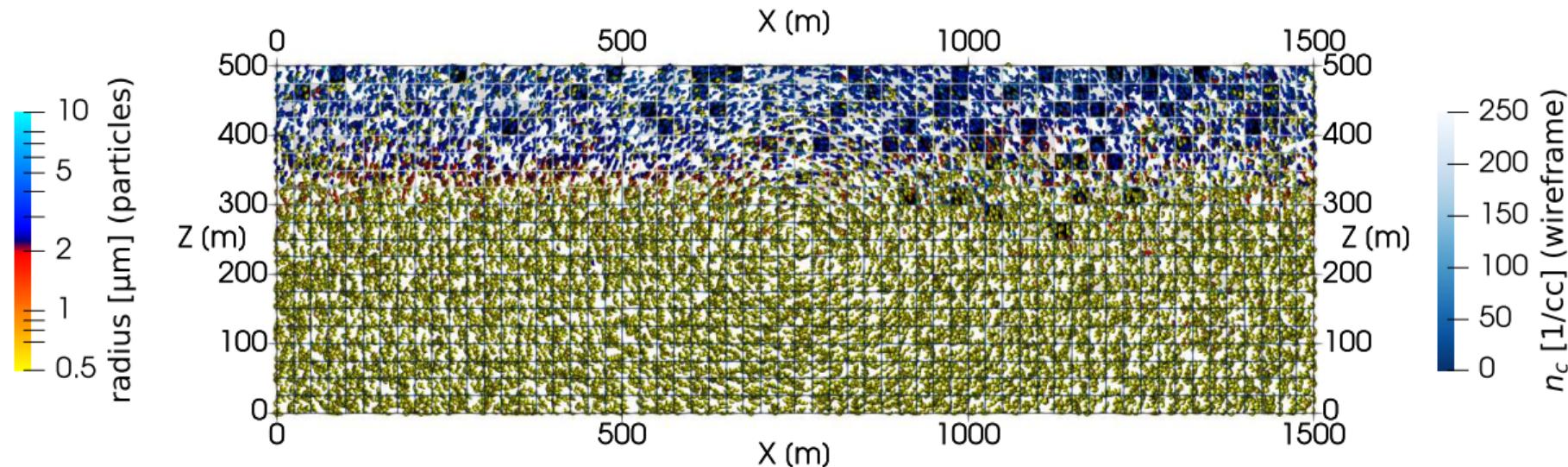
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 1050 s (spin-up till 600.0 s)



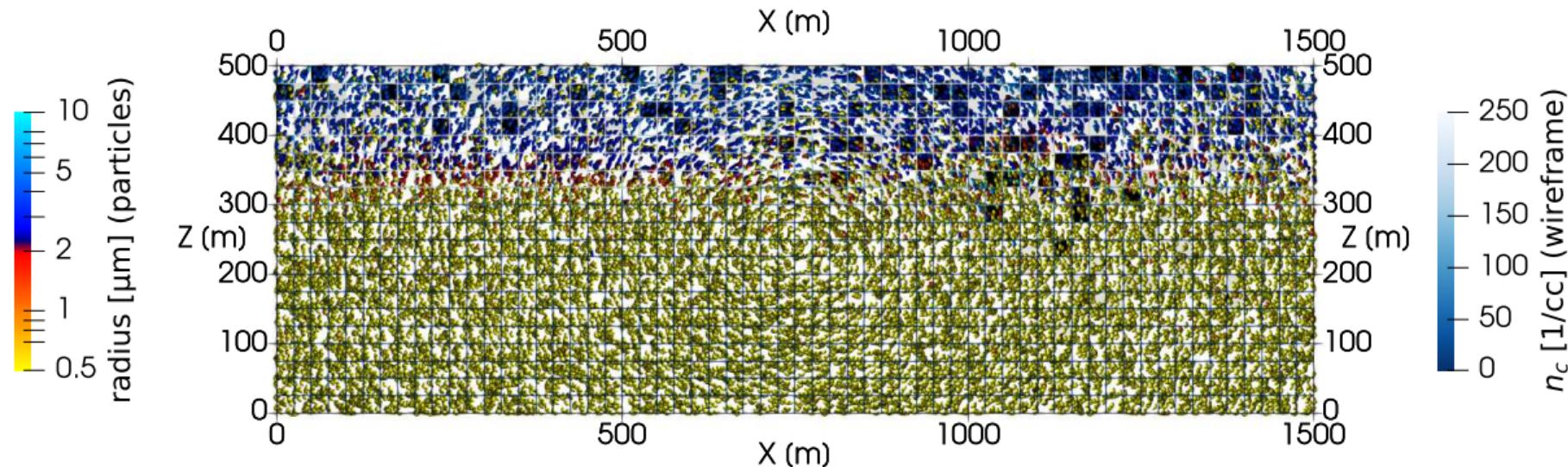
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 1080 s (spin-up till 600.0 s)



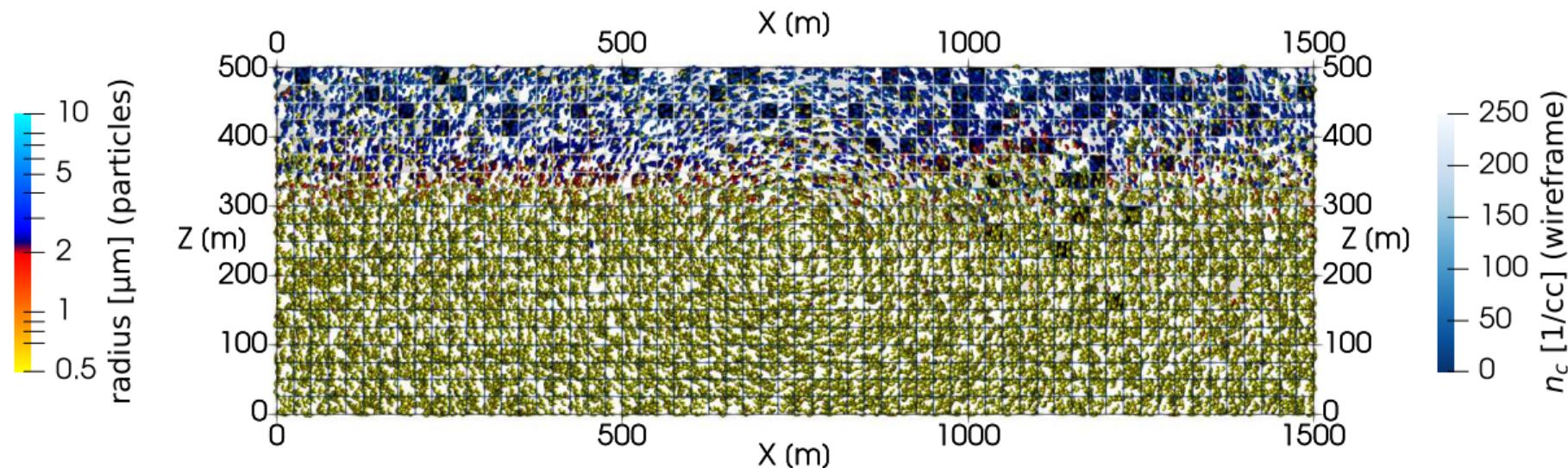
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 1110 s (spin-up till 600.0 s)



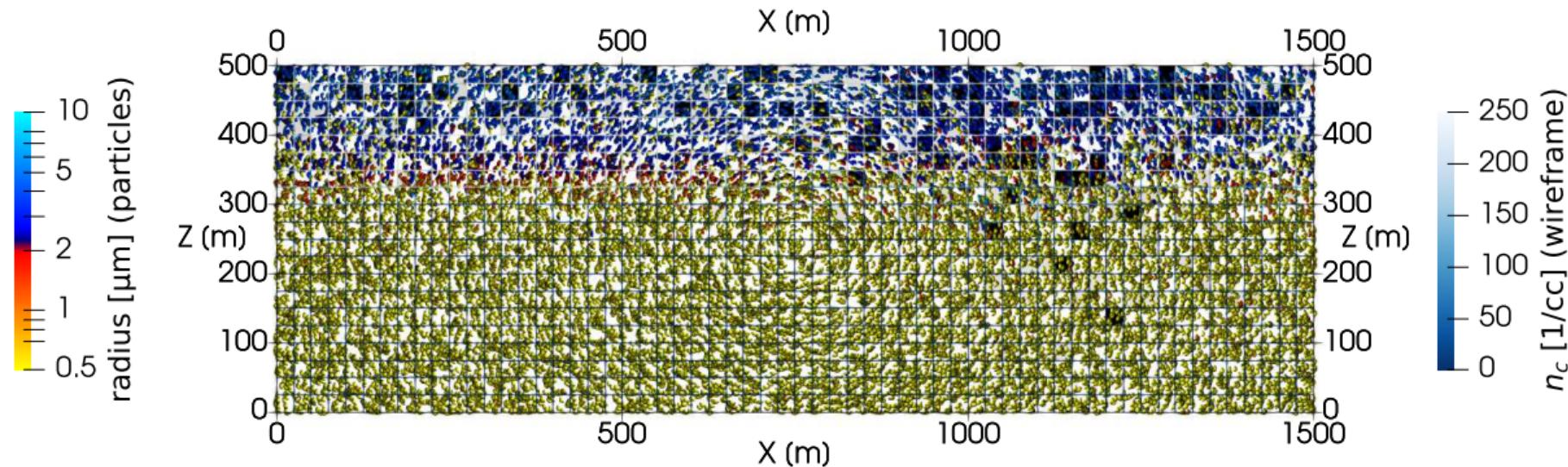
16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 1140 s (spin-up till 600.0 s)



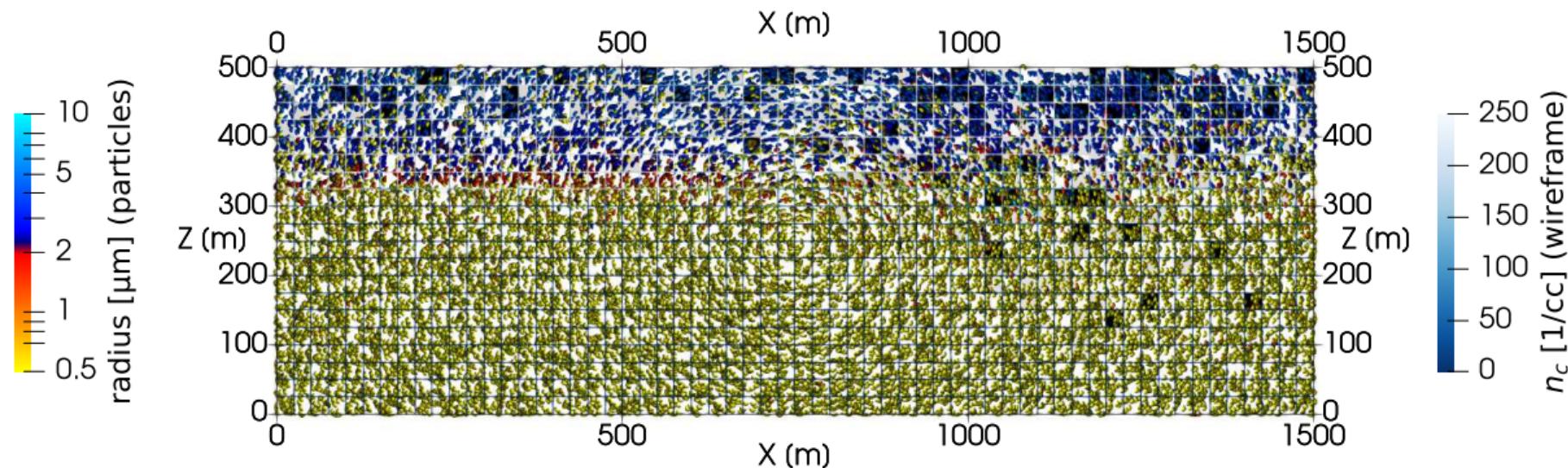
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$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 1170 s (spin-up till 600.0 s)



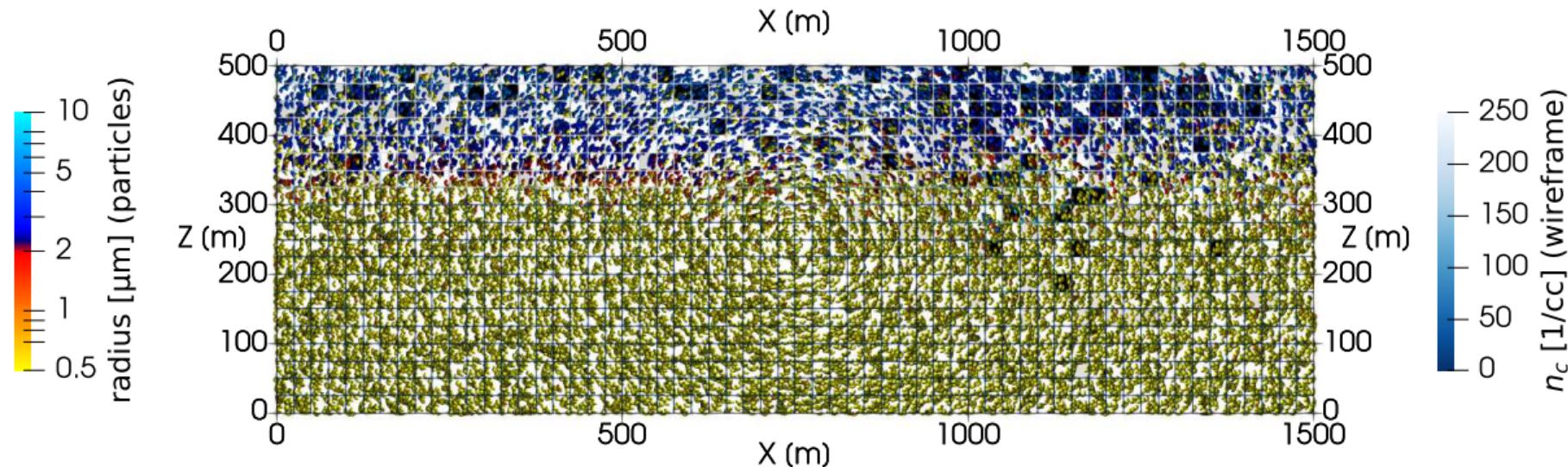
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spin-up = freezing off; subsequently frozen particles act as tracers

# immersion freezing: singular vs. time-dependent in flow-coupled simulation

Time: 1200 s (spin-up till 600.0 s)



16+16 super-particles/cell for INP-rich + INP-free particles

$N_{\text{aer}} = 300/\text{cc}$  (two-mode lognormal)     $N_{\text{INP}} = 150/L$  (lognormal,  $D_g = 0.74 \mu\text{m}$ ,  $\sigma_g = 2.55$ )

spin-up = freezing off; subsequently frozen particles act as tracers

# **new paper in JAMES using PySDM's singular & time-dependent immersion freezing**

[Journal of Advances in Modeling Earth Systems / Volume 17, Issue 4 / e2024MS004770](#)

## **Immersion Freezing in Particle-Based Aerosol-Cloud Microphysics: A Probabilistic Perspective on Singular and Time-Dependent Models**

Sylwester Arabas , Jeffrey H. Curtis, Israel Silber, Ann M. Fridlind, Daniel A. Knopf,  
Matthew West, Nicole Riemer 

First published: 12 April 2025

<https://doi.org/10.1029/2024MS004770>

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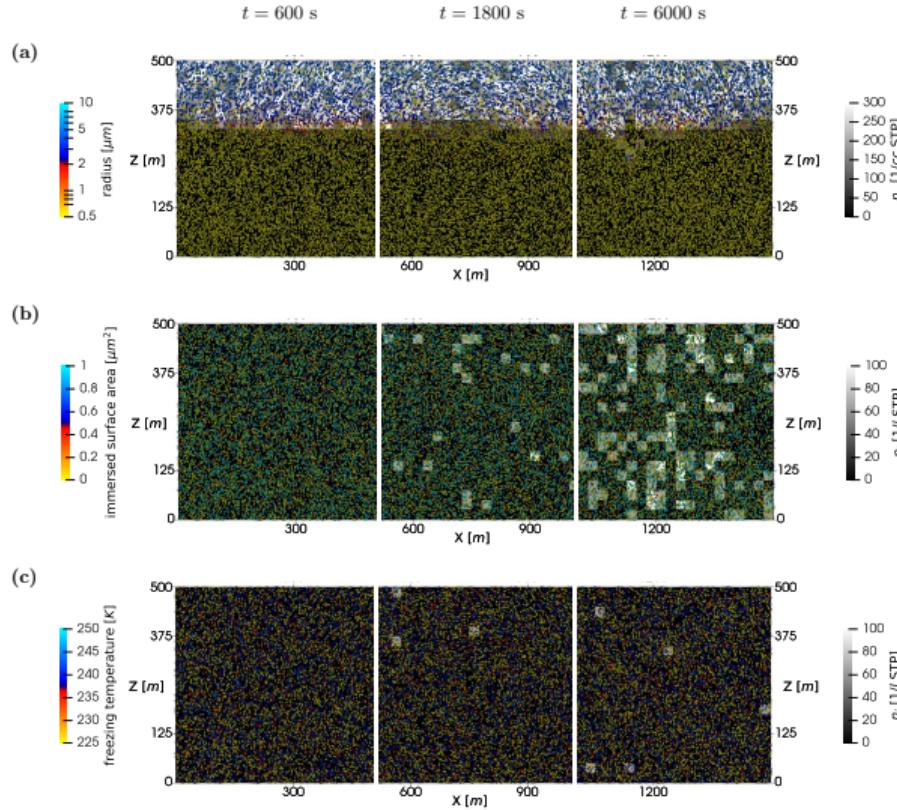
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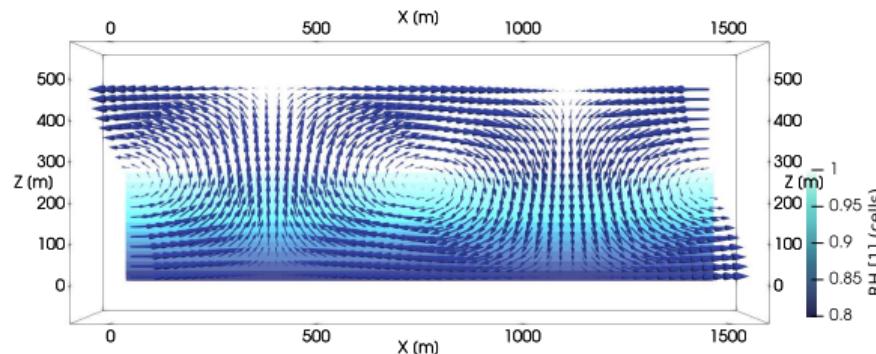
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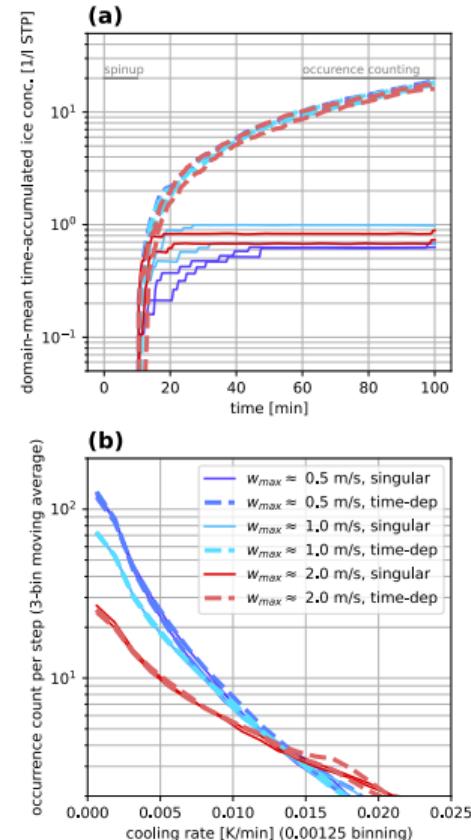
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- ~~ singular model not robust to flow regimes
- ~~ working with Tim on extending to ice growth



PySDM@Kraków funding:



Mainz visit funding:



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6-month postdoc position at AGH available

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6-month postdoc position at AGH available  
AMS Annual Meeting @Houston (25-29 Jan 2026) session



## 18th Symposium on Aerosol-Cloud-Climate Interactions

## Third Symposium on Cloud Physics

Abstracts are due by **14 August 2025 at 5:00 PM ET**

[SUBMIT ABSTRACT](#)

### Joint Sessions

- Advances in numerical modeling of aerosol-cloud interactions: moment-, bin- and particle-resolved methods and beyond

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**Thank you for your attention!**

[sylwester.arabas@agh.edu.pl](mailto:sylwester.arabas@agh.edu.pl)