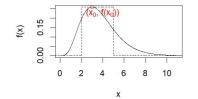
## **Algorithms and Data Structures**

# Quadrature Laplace's method





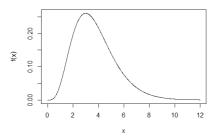
#### Learning goals

Laplace's method

**Target:** Approximate integral of function *f* with the following properties:

- The mass concentrates on a small area around a center and the function has very rapidly decreasing tails ("similarity" to the density of a normal distribution)
- The function we want to integrate is the density of a random variable that is approximately normally distributed





In particular, we assume that f

- Can only be positive
- Is two times continuously differentiable
- Has a global maximum at  $x_0$

We could approximate the area underneath the graph of the function with a staircase function and represent the integral with a very simple formula that depends on  $f(x_0)$ :

$$\int f(x) dx \approx f(x_0) \cdot c$$

$$(x_0, f(x_0))$$

$$0 \quad 2 \quad 4 \quad 6 \quad 8 \quad 10$$

$$x$$



But instead of the step function we would like to choose a function that approximates *f* **better** and which has well-known properties.

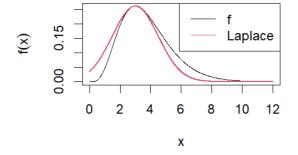
**Idea:** Approximate the integral using the density function of the normal distribution!

**How?** We center and scale the density function of the normal distribution such that it approximates *f* "best possible".



In other words: We determine **expectation** and **standard deviation** of a normal distribution such that the corresponding density function fits best possible to the function f we are interested in.





#### Mathematical derivation:

Let there be a function f with a global maximum at  $x_0$ .

We define  $h(x) := \log f(x)$  as the logarithmized function and rewrite the integral

$$\int_{a}^{b} f(x) dx = \int_{a}^{b} \exp(\underbrace{\log f(x)}_{:=h(x)}) dx$$

Using Taylor's theorem around  $x_0$  we obtain

$$\int_{a}^{b} \exp(h(x)) \approx \int_{a}^{b} \exp\left(h(x_{0}) + h'(x_{0})(x - x_{0}) + \frac{1}{2}h''(x_{0})(x - x_{0})^{2}\right) dx$$



 $x_0$  is also the maximum of  $h(x) = \log(f(x))$ . Hence,  $h'(x_0) = 0$  and the second summand disappears:

$$\int_{a}^{b} \exp(h(x)) \ dx \approx \int_{a}^{b} \exp\left(h(x_{0}) + \frac{1}{2}h''(x_{0})(x - x_{0})^{2}\right) dx$$

We take advantage of the fact that exp(x + y) = exp(x) exp(y)

$$\int_a^b \exp(h(x_0)) \cdot \exp\left(\frac{1}{2}h''(x_0)(x-x_0)^2\right) dx$$

and pull the constant  $\exp(h(x_0))$  out of the integral

$$\exp(h(x_0)) \cdot \int_a^b \exp\left(\frac{1}{2}h''(x_0)(x-x_0)^2\right) dx$$



Within the integral there is now an expression which "almost" corresponds to the density of a normal distribution with expectation  $\mu := x_0$  and variance  $\sigma^2 := -h''(x_0)^{-1}$ :

$$\int_{a}^{b} f(x)dx \approx \exp(h(x_0)) \cdot \int_{a}^{b} \exp\left(\frac{1}{2}h''(x_0)(x - x_0)^2\right) dx$$

$$= \exp(h(x_0)) \cdot \int_{a}^{b} \exp\left(-\frac{1}{2}\frac{(x - x_0)^2}{-h''(x_0)^{-1}}\right) dx$$

$$= \exp(h(x_0)) \cdot \int_{a}^{b} \exp\left(-\frac{1}{2}\frac{(x - \mu)^2}{\sigma^2}\right) dx$$

 $-h''(x_0)^{-1}$  must be truly positive to correspond to the variance of a normal distribution. Since h(x) has a global maximum in  $x_0$ , the second derivative at this point is negative and therefore  $-h''(x_0)^{-1} > 0$ .



If we add (and cancel) the multiplicative constant  $c=\frac{1}{\sqrt{2\pi\sigma^2}}$ , we obtain

$$\int_{a}^{b} f(x)dx \approx \frac{1}{c} \cdot \exp(h(x_{0})) \cdot \int_{a}^{b} \underbrace{c \cdot \exp\left(-\frac{1}{2} \frac{(x-\mu)^{2}}{\sigma^{2}}\right)}_{\text{Density ND}} dx$$

$$= \frac{1}{c} \underbrace{\exp(h(x_{0}))}_{f(x_{0})} \cdot \int_{a}^{b} \phi_{\mu,\sigma^{2}}(x) dx$$

$$= \frac{1}{c} f(x_{0}) \cdot (\Phi_{\mu,\sigma^{2}}(b) - \Phi_{\mu,\sigma^{2}}(a))$$

where  $\phi_{\mu,\sigma^2}(x)$  denotes the density and  $\Phi_{\mu,\sigma^2}(x)$  the distribution function of a normal distribution with expectation  $\mu$  and variance  $\sigma^2$ .



For integration limits  $b = \infty$  and  $a = -\infty$  Laplace's method of f is then

$$\int_{-\infty}^{\infty} f(x) dx \approx \frac{1}{c} \cdot f(x_0) \cdot \left( \Phi_{\mu,\sigma^2}(+\infty) - \Phi_{\mu,\sigma^2}(-\infty) \right)$$
$$= \sqrt{-\frac{2\pi}{h''(x_0)}} \cdot f(x_0)$$

with 
$$h(x) = \log f(x)$$
.

Laplace's method thus corresponds to a value that only depends on the maximum of the function  $f(x_0)$  and the curvature of the logarithmic function  $h''(x_0)$ .

Laplace's method also works well in higher dimensions. For  $f:\mathbb{R}^m\to\mathbb{R}$  with global maximum in  $\mathbf{x}_0$  the generalized form is given by

$$I(f) \approx (2\pi)^{m/2} \det(-H_f(x_0)^{-1})^{1/2} \exp(f(x_0))$$

where  $H_f(x_0)$  denotes the Hessian matrix of f at  $x_0$ . Since  $x_0$  is a global maximum,  $H_f(x_0)$  is negative definite.

The problem of integration is reduced to

- Solving an optimization problem → find x<sub>0</sub>
- Determining the second derivative h''(x) (or generally the Hessian matrix  $H_f(\mathbf{x})$ ) at the optimal position  $x_0$ .

Instead of integration, an optimization problem must now be solved, which is often much easier and faster.



**Application example:** Bayesian computation

Given:

$$x|\lambda \sim \mathsf{Poisson}(\lambda)$$
 (Likelihood)  
 $\lambda \sim \mathsf{Gamma}(\alpha,\beta)$  (Prior)

**Wanted**: Posterior density of the parameter  $\lambda$  given n observations  $\mathbf{x} = (x^{(1)}, x^{(2)}, ..., x^{(n)})$ 

$$p(\lambda|oldsymbol{x}) = rac{egin{array}{c} ext{Likelihood} & ext{Prior} \ p(oldsymbol{x}|\lambda) \cdot \pi(\lambda) \ \end{pmatrix}}{\int p(oldsymbol{x}|\lambda) \cdot \pi(\lambda) \ d\lambda}$$

The density of the gamma distribution is given by  $\pi_{\alpha,\beta}(\lambda) = \frac{1}{\beta^{\alpha}\Gamma(\alpha)}\lambda^{\alpha-1}\exp(-\lambda\beta)$ 



To keep the calculations simple, we calculate the posterior density for only **one** observation x.

The posterior density of  $\lambda$  given the observation x is (except for one constant)

$$p(\lambda|x) \propto \lambda^{x+\alpha-1} \exp\left(-\frac{\lambda}{1/\beta+1}\right) =: f(\lambda).$$

So to determine the posterior density  $p(\lambda|x)$  exactly, we search for the normalization constant c, which ensures that  $\int c \cdot f(\lambda) d\lambda = 1$ , hence

$$c \cdot \int f(\lambda) d\lambda = 1$$

$$c = \frac{1}{\int f(\lambda) d\lambda}$$



**Goal:** Approximation of 
$$\int f(\lambda) \ d\lambda$$
 with  $f(\lambda) = \lambda^{x+\alpha-1} \exp\left(-\frac{\lambda}{1/\beta+1}\right)$ 

We calculate  $h(\lambda) = \log f(\lambda)$ 

$$h(\lambda) = \log\left(\lambda^{x+\alpha-1} \cdot \exp(-\frac{\lambda}{1/\beta+1})\right)$$
$$= (x+\alpha-1)\log\lambda - \frac{\lambda}{1/\beta+1}$$
$$h'(\lambda) = \frac{x+\alpha-1}{\lambda} - \frac{1}{1/\beta+1}$$
$$h''(\lambda) = -\frac{x+\alpha-1}{\lambda^2}$$



To approximate the integral using Laplace's method, we need  $\lambda_0 := \arg \max f(\lambda)$  and  $h''(\lambda_0)$ , where  $h(\lambda) := \log(f(\lambda))$ .

The maximum of  $f(\lambda)$  is the same as the maximum of  $h(\lambda)$  (easier to calculate)

$$\frac{x+\alpha-1}{\lambda} - \frac{1}{1/\beta+1} = 0$$

$$\lambda_0 = \frac{x+\alpha-1}{1/\beta+1}$$

and thus

$$h''(\lambda_0) = -\frac{(1/\beta + 1)^2}{x + \alpha - 1}$$



We insert  $\lambda_0 = \frac{x+\alpha-1}{1/\beta+1}$  and  $h''(\lambda_0)$  into the formula for Laplace's method and obtain

$$\int f(\lambda)d\lambda \approx \sqrt{-\frac{2\pi}{h''(\lambda_0)} \cdot f(\lambda_0)}$$
$$= \sqrt{2\pi} \cdot \frac{\sqrt{x+\alpha-1}}{1/\beta+1} \cdot f(\lambda_0)$$

Hence, the normalization constant c can be approximated by

$$c = \frac{1}{\int f(\lambda) d\lambda} pprox \frac{1}{\sqrt{2\pi}} \frac{1/\beta + 1}{\sqrt{x + \alpha - 1}} \cdot \frac{1}{f(\lambda_0)}$$



When calculating posterior distributions, Laplace's method provides a good approximation if

- The number *n* of observations is large
- The posterior distributions are roughly symmetric

