VV285 RC Part VII

Vector Fields<u>Flux. Circulation and Fundamental Theorem of Calculus</u>

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Outline



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- Green's Identity

Overview



Flux, circulation, divergence, rotation, gradient,... The concepts we learned last week have practical significance in subject such as *fluid dynamics* and *electromagnetics*. It is important for us to be familiar with these concepts.



Figure: Electromagnetics

Flux and Divergence I



Let $F: \mathbb{R}^n \to \mathbb{R}^n$ be a vector field defined in a neighborhood of a hypersurface S with parametrization $\varphi: \Omega \to \mathbb{R}^n, \Omega \subset \mathbb{R}^{n-1}$. Then we define the *flux of F through* S by

$$\int_{\mathcal{S}} F \, d\vec{A} := \int_{\mathcal{S}} \langle F, N \rangle \, dA \qquad (= \int_{\mathcal{S}} \langle F, d\vec{A} \rangle)$$
$$= \int_{\Omega} \langle F \circ \varphi(x), N \circ \varphi(x) \rangle \sqrt{g(x)} \, dx_{1} ... dx_{n-1}$$

And

$$\operatorname{div} F := \frac{\partial F_1}{\partial x_1} + \dots + \frac{\partial F_n}{\partial x_n}$$

is called the *divergence* of *F*.

Remark: The total flux density at a point x is given by the divergence of the field at x.

Flux and Divergence II



In triangle calculus, the divergence of a vector field F can be expressed as

$$\operatorname{div} F = \langle \nabla, F \rangle = \left\langle \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \vdots \\ \frac{\partial}{\partial x_n} \end{pmatrix} \begin{pmatrix} F_1 \\ \vdots \\ F_n \end{pmatrix} \right\rangle = \frac{\partial F_1}{\partial x_1} + \dots + \frac{\partial F_n}{\partial x_n}.$$

In some other notations, one also uses $\nabla \cdot F$ to indicates the divergence of F. This type of notation is more common in a physical textbooks.

Exercise 1: Flux



Find the flux of the vector field $F: \mathbb{R}^3 \to \mathbb{R}^3$, F(x,y,z) = (x,1,z), along the surface $S = \{(x,y,z) \in \mathbb{R}^3 : x^2 + z^2 \le 9, y = x^2 + z^2\}$ Assume y-axis to be negative orientation.

Circulation and Rotation/Curl I



Let $\Omega \subset \mathbb{R}^n$ be an open set, $F : \Omega \to \mathbb{R}^n$ a continuously differentiable vector field and \mathcal{C}^* an oriented closed curve in \mathbb{R}^n . Then

$$\int_{\mathcal{C}^*} \langle F, T \rangle \, ds$$

is called the (total) *circulation* of F along C. Then the anti-symmetric, bilinear form

$$rotF|_{x}: \mathbb{R}^{n} \times \mathbb{R}^{n} \to \mathbb{R}, \qquad rotF|_{x}(u,v) := \langle \, DF|_{x}u \,, \, v \, \rangle - \langle \, DF|_{x}v \,, \, u \, \rangle$$

is called the *rotation* (in mainland Europe) or *curl* (in anglo-saxon countries) of the vector field F at $x \in \mathbb{R}^n$.

Circulation and Rotation/Curl II



Let $\Omega \subset \mathbb{R}^2$ be open and $F:\Omega \to \mathbb{R}^2$ a continuously differentiable vector field. Then here exists a uniquely defined continuous **potential** function rot $F:\Omega \to \mathbb{R}$ such that

$$rot F|_{x}(u, v) = rot F(x) \cdot det(u, v).$$

Moreover,

$$\operatorname{rot} F(x) = \frac{\partial F_2}{\partial x_1} - \frac{\partial F_1}{\partial x_2}.$$

Remark: rot F represents the circulation density in \mathbb{R}^2 .

Circulation and Rotation/Curl III



Let $\Omega \subset \mathbb{R}^3$ be open and $F:\Omega \to \mathbb{R}^3$ a continuously differentiable vector field. Then here exists a uniquely defined continuous **vector** field rot $F:\Omega \to \mathbb{R}^3$ such that

$$rot F|_{X}(u, v) = \langle rot F(x), u \times v \rangle = det(rot F(x), u, v).$$

Moreover,

$$\operatorname{rot} F(x) = \begin{pmatrix} \frac{\partial F_3}{\partial x_2} - \frac{\partial F_2}{\partial x_3} \\ \frac{\partial F_1}{\partial x_3} - \frac{\partial F_3}{\partial x_1} \\ \frac{\partial F_2}{\partial x_1} - \frac{\partial F_1}{\partial x_2} \end{pmatrix}.$$

Remark: rot F represents the circulation density in \mathbb{R}^3 .

A simply connected innovational vector field is a potential field.

Circulation and Rotation/Curl IV



In triangle calculus, the rotation of a vector field F can be formally written as

$$\operatorname{rot} F = \nabla \times F(x) = \operatorname{det} \begin{pmatrix} e_1 & e_2 & e_3 \\ \frac{\partial}{\partial x_1} & \frac{\partial}{\partial x_2} & \frac{\partial}{\partial x_3} \\ F_1 & F_2 & F_3 \end{pmatrix},$$

where e_1, e_2, e_3 are the standard basis vectors in \mathbb{R}^3 .

Remark: The triangle notation is a good mnemonics to help us memorize these formulas. (Why?)

Exercise 2 & 3: Circulation & Potential



Let
$$G(x,y) = \left(2xye^{x^2} + \frac{y}{x}, e^{x^2} + \ln x + 6y\right)$$
 be a vector field in \mathbb{R}^2 , $G: \{(x,y): x > 0, y > 0\} \to \mathbb{R}^2$.

- ▶ Is *G* a potential field?
- ▶ If so, find the form of its potential function. Otherwise, prove it's not conservative.

Let $F(x, y, z) = (e^z, 1, xe^z)$ be a vector field in \mathbb{R}^3 .

- ► Is *F* a potential field?
- ▶ If so, find the form of its potential function. Otherwise, prove it's not conservative.

Green's Theorem



Let $R \subset \mathbb{R}^2$ be a bounded, simple region and $\Omega \supset R$ an open set containing R. Let $F:\Omega \to \mathbb{R}^2$ be a continuously differentiable vector field. Then

$$\int_{\partial R^*} F \, d\vec{s} = \int_R \left(\frac{\partial F_2}{\partial x_1} - \frac{\partial F_1}{\partial x_2} \right) \, dx$$

where ∂R^* denotes the boundary curve of R with positive (counter-clockwise) orientation.

Remark: Green's theorem is basically a particular case of Stokes's theorem. A clever use of it is to calculate the area of a simple region. For $F(x_1, x_2) = (-x_2, x_1)$ we obtain

$$|R| = \int_R 1 dx = \frac{1}{2} \int_{\partial R} {-x_2 \choose x_1} d\vec{s}.$$

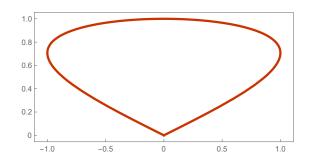
For a bounded simple region, the value along the boundary contains the information of the area.

Exercise 4: Green's Theorem



Find the area of region R that is bounded by the curve C, whose parametrization is

$$\gamma(t) = \begin{pmatrix} \sin 2t \\ \sin t \end{pmatrix}, \quad t \in [0, \pi].$$



Gauß's Theorem



Let $R \subset \mathbb{R}^n$ be an admissible region and $F : \overline{R} \to \mathbb{R}^n$ a continuously differentiable vector field. Then

$$\int_{R} \operatorname{div} F \, dx = \int_{\partial R^*} F \, d\vec{A}.$$

Exercise 5: Gauß's Theorem



Calculate the integral

$$\int_{S^2} (xy + yz + zx) d\sigma$$

where $S^2=\left\{(x,y,z)\in\mathbb{R}^3: x^2+y^2+z^2=1\right\}$ is the unit sphere.

Exercise 6: Gauß's Theorem



Let
$$F: \mathbb{R}^3 \to \mathbb{R}^3$$
, $F(x, y, z) = (x, y, z)$ and

$$U = \{(x, y, z) \in \mathbb{R}^3 : x^2 + y^2 \le 1, x + z \le 2, z \ge 0, x \ge 0\}$$

Denote by $S = \partial U$ the surface of U and by N the outward-pointing unit normal vector on S. Let S^+ denote the surface S with positive orientation.

- Sketch U.
- ► Calculate the flux $\int_{S^+} \langle F, N \rangle d\sigma$

Stokes's Theorem



Let $\Omega \subset \mathbb{R}^3$ be an open set, $\mathcal{S} \subset \Omega$ a parametrized, admissible surface in \mathbb{R}^3 with boundary $\partial \mathcal{S}$ and let $F:\Omega \to \mathbb{R}^3$ be a continuously differentiable vector field. Then

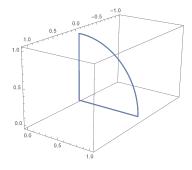
$$\int_{\partial \mathcal{S}^*} F \, d\vec{s} = \int_{\mathcal{S}^*} \operatorname{rot} F \, d\vec{A}$$

where the orientations of the boundary curve $\partial \mathcal{S}^*$ and the surface \mathcal{S}^* are chosen according to *right hand law*.

Exercise 7: Stokes's Theorem



Let \mathcal{C} be the curve illustrated below.



For F(x, y, z) = (y, z, x), compute the circulation along C using Stokes's theorem.

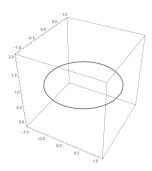
Exercise 8: Stokes's Theorem



Let

$$F(x,y,z) = \left(\sin x - \frac{y^3}{3}, \cos y + \frac{x^3}{3}, xyz\right).$$

Calculate the circulation of F along \mathcal{C} , where \mathcal{C} is the oriented unit circle in the plane z=1. (The orientation is counter-clockwise if viewing from positive z-axis).



Application 1: Electromagnetics



Integral Equations

$$\oint_{\partial\Omega} E \cdot dS = \frac{1}{\varepsilon_0} \int_{\Omega} \rho dV \qquad \nabla \cdot E = \frac{\rho}{\varepsilon_0}$$

$$\oint_{\partial\Omega} B \cdot dS = 0 \qquad \nabla \cdot B = 0$$

$$\oint_{\partial\Sigma} E \cdot d\ell = -\frac{d}{dt} \int_{\Sigma} B \cdot dS \qquad \nabla \times E = -\frac{\partial B}{\partial t}$$

$$\oint_{\partial\Sigma} B \cdot d\ell = \mu_0 \left(\int_{\Sigma} J \cdot dS + \varepsilon_0 \frac{d}{dt} \int_{\Sigma} E \cdot dS \right) \qquad \nabla \times B = \mu_0 \left(J + \varepsilon_0 \frac{\partial E}{\partial t} \right)$$

Differential Equations

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\varepsilon_0}$$

$$\nabla \cdot B = 0$$

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

$$\nabla \times B = \mu_0 \left(J + \varepsilon_0 \frac{\partial E}{\partial t} \right)$$

The famous equations above are called the *Maxwell equations*. They together describe fundamental behaviors of electromagnetic fields. Can you find the relations between the integral equations and differential equations?

Exercise 9: *Vector Calculus Identity



As we have discussed last week, physicists really love using triangle notations because this good notation makes life easier. See those common equalities in vector calculus, and think about what they indicate/why they hold/why the triangle notations are user-friendly:

- $\nabla \times (\nabla f) = 0$
- $\nabla \cdot (\nabla \times F) = 0$

Remark: In electromagnetics, we define *magnetic vector potential A* to be a vector field such that

$$B = \nabla \times A$$

and *electric potential* ϕ to be scalar field such that

$$E = -\nabla \phi - \frac{\partial A}{\partial t}.$$

How does such a definition correspond with Maxwell equations?

Green's Identity



Let $R \subset \mathbb{R}^n$ be an admissible region and $u, v : \overline{R} \to \mathbb{R}$ be twice continuously differentiable potential functions. Then

$$\int_{R} \langle \nabla u, \nabla v \rangle dx = -\int_{R} u \cdot \Delta v dx + \int_{\partial R^{*}} u \frac{\partial v}{\partial n} dA$$
 (1)

and

$$\int_{R} (u \cdot \Delta v - v \cdot \Delta u) \, dx = \int_{\partial R^*} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \, dA. \tag{2}$$

(1) is commonly called *Green's first identity* and (2) *Green's second identity*.

Application 2: Fluid Dynamics



The properties of the Laplace operator Δ acting on smooth functions defined on some set $\Omega \subset \mathbb{R}^n$ are important in the theory of differential equations. Green's identities are your friend here:

Let $\Omega \subset \mathbb{R}^2$ be an open, bounded, connected set and $u:\Omega \to \mathbb{R}$ a twice differentiable function such that

$$\frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2} = 0 \quad and \quad \frac{\partial u}{\partial n} \Big|_{\partial \Omega} = 0$$

- 1. Prove that u is constant, i.e., there exists some $c \in \mathbb{R}$ such that u(x) = c for all $x \in \Omega$.
- 2. Interpret the statement of this exercise physically in terms of fluid flow (potential flow) in a container.

Extra



Enjoy the following!