

Ex 1.

We first parametrize the surface,

$$\varphi : [0, 3] \times [0, 2\pi] \rightarrow \mathbb{R}^3 \quad \Rightarrow \quad F \circ \varphi(r, \theta) = (r \cos \theta, r^2, r \sin \theta)$$

$$\varphi(r, \theta) = (r \cos \theta, r^2, r \sin \theta)$$

The tangent vectors are

$$\begin{aligned} t_r(r, \theta) &= (\cos \theta, 2r, \sin \theta) \\ t_\theta(r, \theta) &= (-r \sin \theta, 0, r \cos \theta) \end{aligned} \quad \left. \begin{array}{l} \\ \end{array} \right\} n(r, \theta) = t_r \times t_\theta(r, \theta)$$

$$d\bar{A} = (2r^2 \cos \theta, -r, 2r^2 \sin \theta) dr d\theta \leftarrow \text{the vectorial surface element.}$$

$$\int_S \langle F, d\bar{A} \rangle = \int_0^{2\pi} \int_0^3 (2r^3 - r) dr d\theta = 72\pi$$

Ex 2.

i) \mathbf{G} is a potential field, because the region is simple-connected, and

$$\operatorname{rot} \mathbf{G} = \frac{\partial G_y}{\partial x} - \frac{\partial G_x}{\partial y} = 2xe^{x^2+\frac{1}{x}} - (2xe^{x^2+\frac{1}{x}}) = 0$$

ii). $\int \underline{G_x dx} = \int (2xye^{x^2} + \frac{y}{x}) dx = \underline{ye^{x^2} + y \ln x + C_1(y)}$

$$\underline{\int G_y dy} = \int (e^{x^2} + \ln x + by) dy = \underline{ye^{x^2} + y \ln x + \frac{3}{2}y^2 + C_2(x)}$$

$$\Rightarrow g(x,y) = \underline{ye^{x^2} + y \ln x + \frac{3}{2}y^2 + C}, \quad C \in \mathbb{R} \text{ is a constant.}$$

$$\frac{\partial}{\partial x} C = 0$$

Ex 3.

i) \vec{F} is a potential field because \mathbb{R}^3 is simply-connected, and.

$$\text{rot } \vec{F} = \begin{pmatrix} 0 & 0 \\ e^z & -e^z \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad (\Rightarrow \text{irrotational})$$

ii)

$$\int F_x dx = xe^z + C_1(y, z)$$

$$\int F_y dy = y + C_2(x, z)$$

$$\int F_z dz = xe^z + C_3(x, y)$$

$$\Rightarrow f(x, y, z) = xe^z + y + C, \text{ where } C \in \mathbb{R} \text{ is a constant.}$$

Ex 4

We apply Green's theorem,

$$\begin{aligned}|R| &= \int_R 1 dx = \frac{1}{2} \int_{\partial R} \begin{pmatrix} -x_2 \\ x_1 \end{pmatrix} d\vec{s} \\&= \frac{1}{2} \int_0^{\pi} \left\langle \begin{pmatrix} -\sin t \\ \sin 2t \end{pmatrix}, \begin{pmatrix} 2\cos 2t \\ \cos t \end{pmatrix} \right\rangle dt \\&= \frac{1}{2} \int_0^{\pi} (\sin 2t \cos t - 2\sin t \cos 2t) dt \\&= \frac{1}{2} \int_0^{\pi} 2 \sin^3 t dt \\&= \int_0^{\pi} \sin^3 t dt \\&= \frac{4}{3}\end{aligned}$$

Ex 5.

If we define $\underline{F(x,y,z)} = (y, z, x)$

we have

$$\int_{S^2} (xy + yz + zx) d\sigma = \int_{S^2} \langle F, N \rangle d\sigma \quad (= \int_{S^2} \langle F, d\bar{\omega} \rangle)$$

By Gauss's theorem,

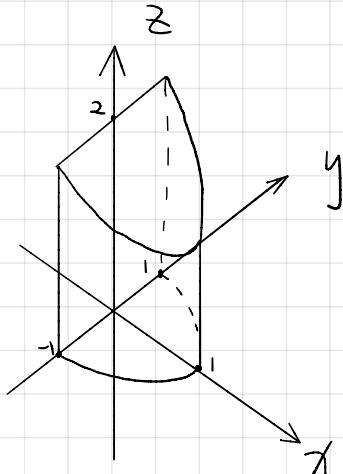
$$\int_{S^2} \langle F, N \rangle d\sigma = \int_{B(0)} \operatorname{div} F dV. = 0$$

where $B(0)$ is the unit ball centered at 0 , and

$$\operatorname{div} F = 0 + 0 + 0 = 0$$

Ex 6.

i) Sketch :



By Gauss's theorem,

$$\text{i)} \int_{S^+} \langle \mathbf{F}, \mathbf{N} \rangle d\sigma = \int_V \operatorname{div} \mathbf{F} dV$$

$$= \int_V 3 dV$$

Using cylindrical coordinates,

$$\begin{aligned} \int_V 3 dV &= 3 \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \int_0^1 \int_0^{2-r\cos\phi} r dz dr d\phi \\ &= 3 \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \int_0^1 (2-r\cos\phi) r dr d\phi \\ &= 3 \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} [r^2 - \frac{r^3}{3} \cos\phi]_0^1 d\phi \\ &= 3 \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} (1 - \frac{1}{3} \cos\phi) d\phi \\ &= 3\pi - (\sin\phi)_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \\ &= 3\pi - 2 \end{aligned}$$

Ex 7. We parametrize the surface, $\varphi: [0,1] \times [0, \frac{\pi}{2}] \rightarrow S$

$$\varphi(r, \theta) = \begin{pmatrix} 0 \\ r\cos\theta \\ r\sin\theta \end{pmatrix},$$

$$\text{rot } F = \begin{vmatrix} \mathbf{i}_1 & \mathbf{i}_2 & \mathbf{i}_3 \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ y & z & x \end{vmatrix} = \begin{pmatrix} -1 \\ -1 \\ -1 \end{pmatrix}$$

The tangent vector: $t_r(r, \theta) = \begin{pmatrix} \cos\theta \\ \sin\theta \\ 0 \end{pmatrix}$, $t_\theta(r, \theta) = \begin{pmatrix} 0 \\ r\sin\theta \\ r\cos\theta \end{pmatrix}$

Normal vector: $t_r \times t_\theta(r, \theta) = \begin{pmatrix} r \\ 0 \\ 0 \end{pmatrix} \xrightarrow[\text{by orientation}]{\substack{\text{change sgn}}} \begin{pmatrix} -r \\ 0 \\ 0 \end{pmatrix}$

Apply Stokes theorem,

$$\int_{\partial S^*} \langle F, d\ell \rangle = \int_S \langle \text{rot } F, d\bar{A} \rangle = \int_0^1 \int_0^{\frac{\pi}{2}} \left\langle \begin{pmatrix} -1 \\ -1 \\ -1 \end{pmatrix}, \begin{pmatrix} -r \\ 0 \\ 0 \end{pmatrix} \right\rangle dr d\theta$$

$$= \frac{\pi}{4}$$

Ex8. We parametrize the surface: $\varphi [0,1] \times [0,2\pi] \rightarrow S$

$$\varphi(r, \theta) = (r \cos \theta, r \sin \theta, 1)$$

$$\text{rot } F = \begin{vmatrix} e_1 & e_2 & e_3 \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \sin x - \frac{y^3}{3} & \cos y + \frac{x^2}{3} & xy^2 \end{vmatrix} = \begin{pmatrix} xz \\ -yz \\ x^2 + y^2 \end{pmatrix} = \begin{pmatrix} r \cos \theta \\ -r \sin \theta \\ r^2 \end{pmatrix}$$

The tangent vector: $t_r(r, \theta) = \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix}$, $t_\theta(r, \theta) = \begin{pmatrix} -r \sin \theta \\ r \cos \theta \\ 0 \end{pmatrix}$

normal vector: $\begin{pmatrix} 0 \\ 0 \\ r \end{pmatrix}$ (verify direction is correct)

Apply Stokes's theorem,

$$\int_{\partial S^*} \langle F, d\bar{l} \rangle = \int_S \langle \text{rot } F, d\bar{A} \rangle = \int_0^1 \int_0^{2\pi} \left\langle \begin{pmatrix} r \cos \theta \\ -r \sin \theta \\ r^2 \end{pmatrix}, \begin{pmatrix} 0 \\ 0 \\ r \end{pmatrix} \right\rangle d\theta dr$$
$$= 2\pi \int_0^1 r^3 dr = \frac{\pi}{2}$$

Ap 1.

Using Gaup's theorem,

$$\oint_{\partial\Omega} \underline{E} \cdot d\underline{S} = \int_{\Omega} \nabla \cdot \underline{E} dV = \int_{\Omega} \frac{\rho}{\epsilon_0} dV = \frac{1}{\epsilon_0} \int \rho dV$$

$$\oint_{\partial\Omega} \underline{B} \cdot d\underline{S} = \int_{\Omega} \nabla \cdot \underline{B} dV = \int_{\Omega} 0 \cdot dV = 0$$

$$\oint_{\partial\Omega} \underline{E} \cdot d\underline{l} = \int_{\Sigma} \nabla \times \underline{E} \cdot d\underline{S} = \int_{\Sigma} -\frac{\partial B}{\partial t} \cdot d\underline{S} = -\frac{d}{dt} \int_{\Sigma} \underline{B} \cdot d\underline{S}$$

12.9. Theorem Slide 329.

$$\begin{aligned} \oint_{\partial\Omega} \underline{B} \cdot d\underline{l} &= \int_{\Sigma} \nabla \times \underline{B} \cdot d\underline{S} = \int_{\Sigma} \mu_0 (J + \epsilon_0 \frac{\partial \underline{E}}{\partial t}) \cdot d\underline{S} \\ &= \mu_0 \left(\int_{\Sigma} J \cdot d\underline{S} + \epsilon_0 \frac{d}{dt} \int_{\Sigma} \underline{E} \cdot d\underline{S} \right) \end{aligned}$$

Ex9.

$$\mathbf{B} = \nabla \times \mathbf{A} \Rightarrow$$

$$\nabla \cdot (\nabla \times \mathbf{A}) = 0 \Leftrightarrow \nabla \cdot \mathbf{B} = 0$$

$$\mathbf{E} = -\nabla \phi - \frac{\partial \mathbf{A}}{\partial t} \Rightarrow$$

$$\nabla \times \mathbf{E} = \nabla \times \left(-\nabla \phi - \frac{\partial \mathbf{A}}{\partial t} \right)$$

$$= -\nabla \times \nabla \phi - \frac{\partial}{\partial t} \frac{\nabla \times \mathbf{A}}{\partial t}$$

$$= 0 - \frac{\partial}{\partial t} \mathbf{B} = -\frac{\partial \mathbf{B}}{\partial t}$$

..... (Learn more in VE230!)

Ap 2.

i) By Green's first identity,

$$\int_{\Omega} \langle \nabla u, \nabla u \rangle dx = - \underbrace{\int_{\Omega} u \Delta u dx}_{=0} + \underbrace{\int_{\partial\Omega^*} u \frac{\partial u}{\partial n} dA}$$

Since $\Delta u = 0$, and $\frac{\partial u}{\partial n} = 0$, we have

$$\int_{\Omega} \langle \nabla u, \nabla u \rangle dx = 0 \Rightarrow \int_{\Omega} \|\nabla u\|^2 dx = 0$$

Therefore, $\nabla u(x) = 0$ a.e. on Ω .

Since $u \in C^2(\Omega) \Rightarrow \nabla u(x) = 0$ for all $x \in \Omega$

In summary, u is constant in Ω .

- ii).
1. The flow is irrotational (Because we model it using a potential field. $\boldsymbol{v} = \nabla u$)
 2. There are no sources/sinks, and the fluid is incompressible, (Because $\Delta u = \operatorname{div} \boldsymbol{v} = 0$)
 3. Here u represents the potential of the velocity $\boldsymbol{v} = \nabla u$ of the fluid flow.
 4. The boundary condition $\frac{\partial u}{\partial n}|_{\partial \Omega} = 0$ indicates $v_n|_{\partial \Omega} = 0$. So that no fluid flows through Ω 's boundary \Rightarrow "solid-wall" boundary
 5. u is constant $\Rightarrow \boldsymbol{v} = 0$. The fluid is stationary in Ω .
- * This means fluid moving in a closed container cannot be both irrotational & incompressible. Otherwise it would be stationary.