CHM 673, Quiz 7   Name_____

1. Formulate the Hohenberg-Kohn theorem.

2. Which term is the least accurate in the Thomas-Fermi-Dirac (orbital-free) density functional theory?

3. What is the main idea of the Kohn-Sham DFT?

4. What is the origin of the self-interaction error in the Kohn-Sham DFT?

?

5. Why are KS-DFT calculations typically much faster than correlated wave function methods?

a. They yield accurate densities
b. They only solve a non-interacting problem   ← ?
c. They use perturbation theory
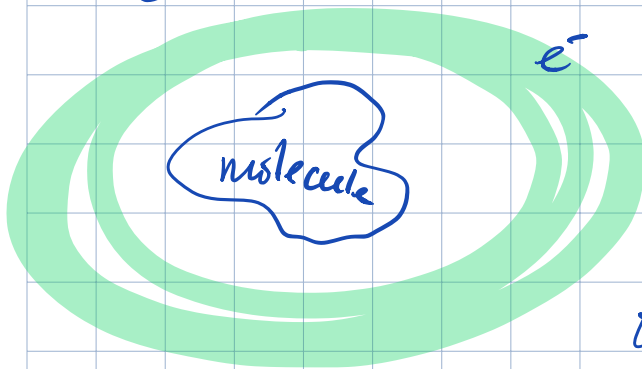d. They use empirical information about the system
                                                 world

functional is
empirical
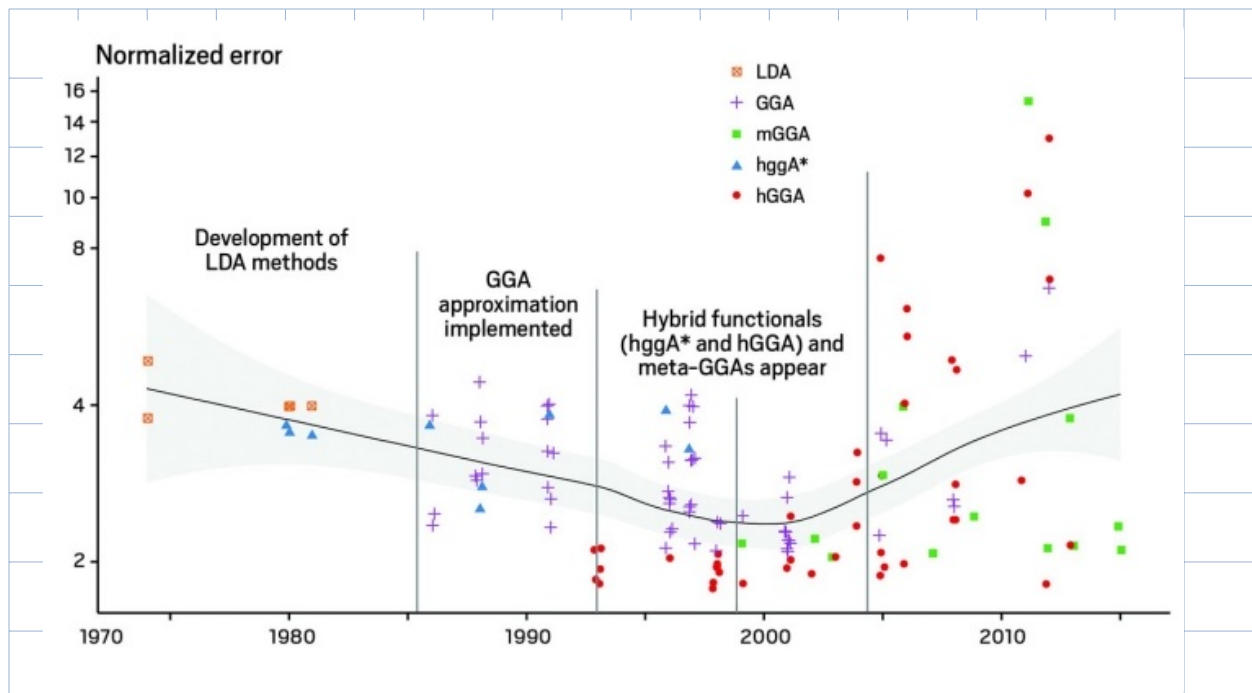
Self-interaction error tends to over-
delocalize electrons
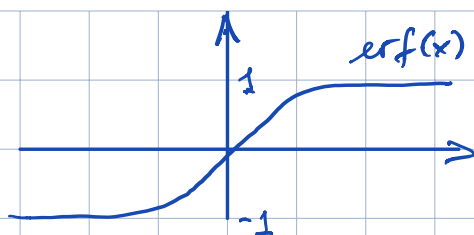
anion

e^- becomes unbound

$\varepsilon_{HOMO} > 0$

especially in diffuse
basis sets electron is
not found

Fix: use range-separated functionals

$$\frac{1}{r_{12}} = \frac{1 - erf(\omega r_{12})}{r_{12}} + \frac{erf(\omega r_{12})}{r_{12}} \leftarrow exact$$

$$erf(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$$

HF
exchange

$\sim 1/r$

erf(x)

1

-1

long-ranged
$\dfrac{erf(r)}{r}$

$\dfrac{1-erf(r)}{r}$
short-ranged
DFT exchange

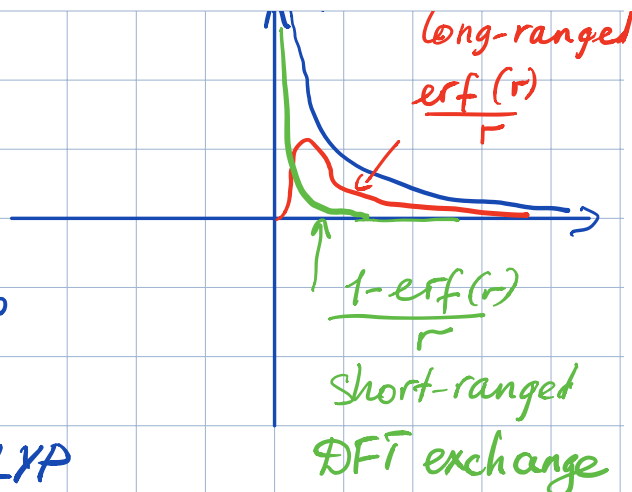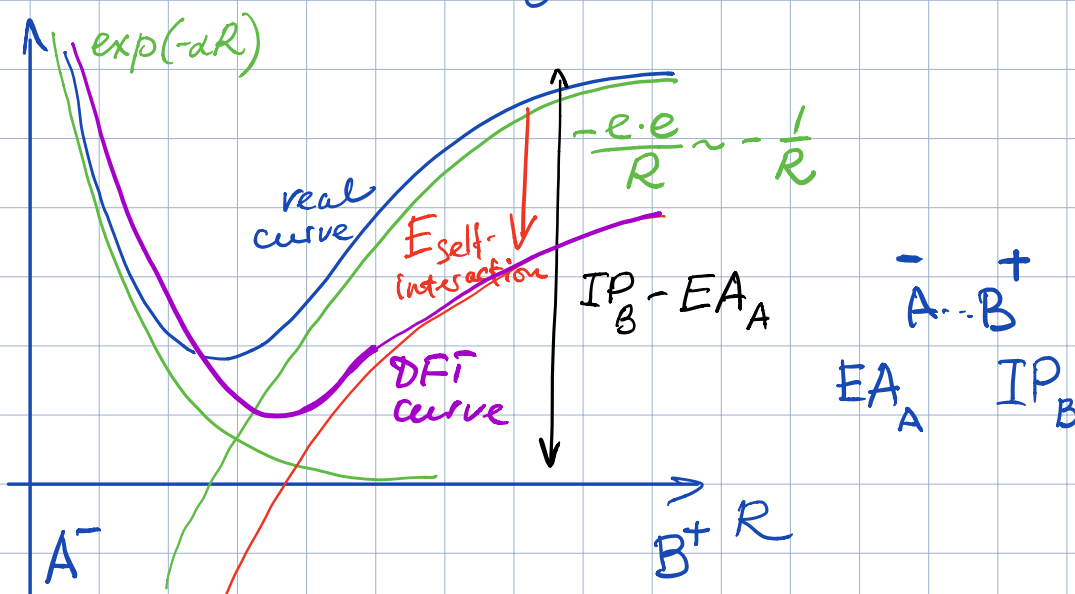| LC	LC-BLYP
ω	ωB97x
CAM	CAM-B3LYP

ω – parameter, ω ~ 0.3–0.5 Bohr
ω is both system dependent & functional dependent

Problem of charge-transfer states

$exp(-\alpha R)$

real curve

$E_{self-intersection}$

$\dfrac{-e \cdot e}{R} \sim -\dfrac{1}{R}$

$IP_B - EA_A$

DFT curve

A^-

$B^+ \ R$

$\bar{A} \cdots B^+$

$EA_A \quad IP_B$

How to describe long-range correlation?
DFT correlation functional accounts for dynamic correlation aka dispersion

$$E_{disp} \sim \frac{C_6}{R^6} + \frac{C_8}{R^8} + \frac{C_{10}}{R^{10}} + \ldots \quad \text{vdW}$$

$$C_6 \sim f(\alpha_A \alpha_B) \qquad \alpha - \text{polarizability}$$

vdW in force fields is described

as $\quad \dfrac{C_{12}}{R^{12}} - \dfrac{C_6}{R_6}$

$$C_6 = \sqrt{C_A C_B} \leftarrow \text{coefs. on atoms}$$
$$\text{A and B}$$

DFT:

$$\Delta E_{disp} = - \sum_{n=6(8,10)} S_n \sum_{AB}^{\text{atoms}} \frac{C_n^{AB}}{R_{AB}^n} f_{damp}(R_{AB})$$

functional-
dependent
coeff.

ensures correct
short-range behavior

D / D2 / D3
→ B3LYP-D2

ΔE_{disp} — additive correction to
KS energy → does not
depend on density

Concerns: - no dependence on density
(not sensitive to multiplets,
excited states)

- f_{damp} not accurate for
strained complexes → large errors

at short separations
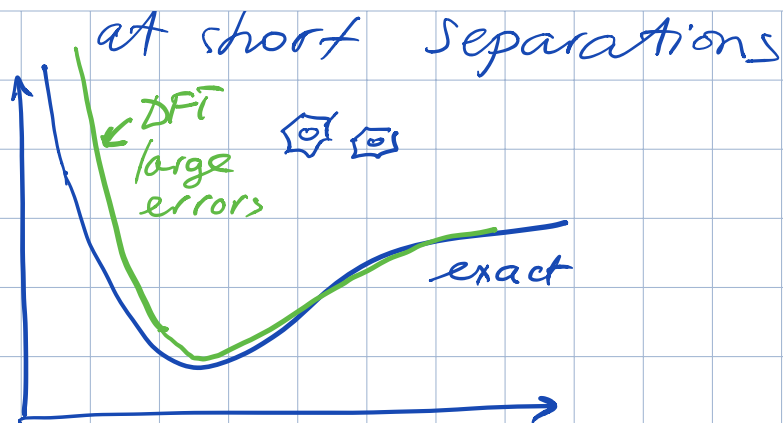
Other ideas: use polarizabilities
either expt. (Tkatchenko f-ns)
or computed (Johnson),
add dependence on density

Computational cost
KS eq-ns are similar to Fock eq-ns
diagonalization of KS matrix $\sim K^3$
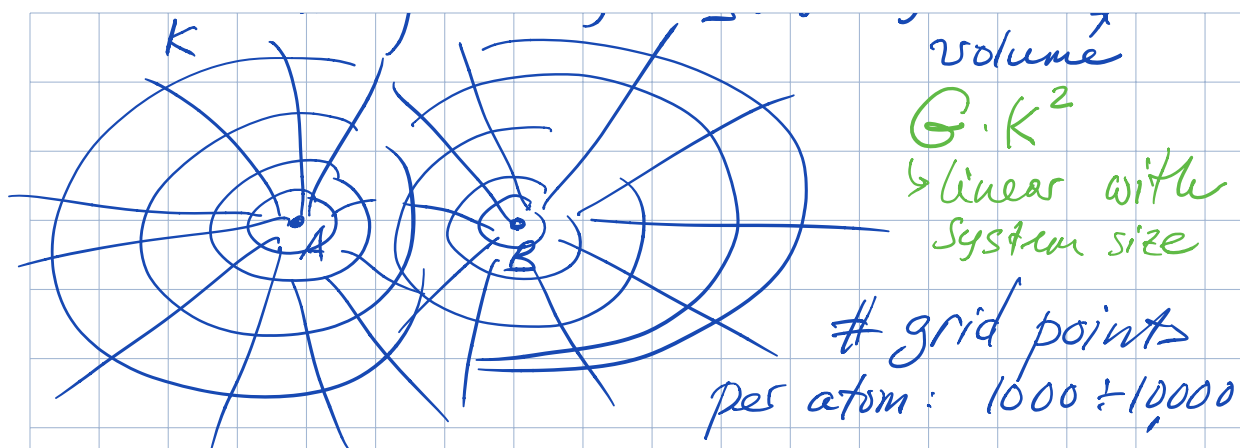   K # of basis functions
Coulomb integrals K^4

$$V_{xc}\left[\rho^{(r)}, \nabla\rho^{(r)}, \nabla^2\rho^{(r)}\right]$$

$$\int \chi_i(r)\, V_{xc}\left[\rho(r), \nabla\rho(r)\right] \chi_i(r)\, dr \qquad \text{no analitical formulas}$$

Numerical integration:
Grid

$$\sum V_{xc}\left[\rho(r_k), \nabla\rho(r_k)\right] \chi_i(r_k)\, \chi_i(r_k)\, \Delta V_k \sim$$

K

volume

$G \cdot K^2$
↳ linear with
system size

grid points
per atom: $1000 \div 10{,}000$

Grid is similar to basis set, tighter
grid gives more accurate energy
Geom. opt. and frequencies need
tighter grid than energy calcs.

Differences between DFT & HF
DFT can be exact theory,
HF - no way
Orbital energies
HF: occ. orbitals - IP
       vir. orbitals - EA   (N electron)

DFT: occ. orbitals $\varepsilon_i = \dfrac{\partial E}{\partial n_i}$

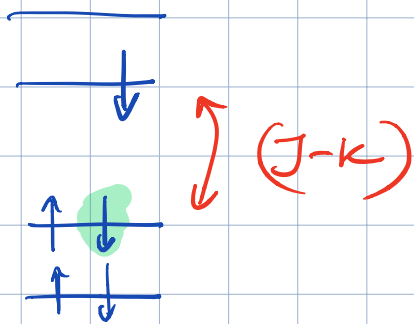(Janak theorem)   ↗ occupation number
                      of the orbital

        still ~ related to IP

virtual orbitals are computed in
field of N-1 electrons ⇒

correspond to exciting (rather
than attaching) an electron
DFT virtuals are lower than HF
                              virtuals

DFT: occ-vir gap ~ excitation energy
HF :   ~ $f(J-K)$

$(J-K)$