# **The Weekly Progress**

#### Since Nov. 10, 2018

In the autumn semester, hopefully I'll record my progress every Thursday morning

### **SO FAR**

- 1. Understand why X-ray cause profound effects to ISM
- 2. Aquire the X-ray spectra emitted the SMBH in the center of Milky Way Galaxy several mys ago
- 3. Decide to focus on a cold dense cloud 8 kpc from the SMBH, thus fixing the flux
- 4. Figure out how KROME works, especially its X-ray module (details on <a href="https://github.com/slowdivePTG/X">https://github.com/slowdivePTG/X</a>
  <a href="https://github.com/slowdivePTG/X">-ray-chemistry/tree/master</a>)
- 5. Be able to calculate primary X-ray ionizations according to a X-ray spectra and the cross sections of certain species.

$$\zeta_k^p = \int_{E_{min}}^{E_{max}} rac{F_X(E)}{E} \sigma_k(E) \mathrm{d}E$$

$$F_X(E): \ \mathrm{erg}\cdot \ \mathrm{cm}^{-2}\cdot \ \mathrm{s}^{-1}$$

6. Learn about the cosmic ray ionization

## Nov. 10, 2018

1. Trace the most important reactions using the Jacobian Matrix pd(i,j)

$$pd(i,j) = rac{\partial}{\partial n(j)} igg(rac{\mathrm{d}n(i)}{\mathrm{d}t}igg)$$

$$\left(\frac{\mathrm{d}n(i)}{\mathrm{d}t}\right)_i = n(j) \cdot pd(i,j)$$

2. Test a simple model but consider a mixture of H/He rather than  ${
m H_2/He}$ .  $J_{X21}=0.1$  for 1 mys then turn off the X-ray radiation. Plot the abundance of  ${
m CH_4}$  with time.

For this molecular cloud test we choose the  $osu_01_2007$  network and the initial conditions proposed by  $\it Wakelam~\&~Herbst~(2008)$ : a constant temperature of  $\rm T=10~K, H_2$  density of  $\rm 10^4~cm^{-3}$ , cosmic rays ionization rate of  $\rm 1.3\times10^{-17}~s^{-1}$ , and a visual extinction of 10. The initial conditions of the species are listed in Table 7 and correspond to the EA2 model of  $\it Wakelam~\&~Herbst~(2008)$ , an high-metal environment observed in the diffuse cloud  $\it \zeta~Ophiuchi$ 

```
Tgas = 1d1 !gas temperature (K)
xH = 2d4 !Hydrogen density
!user commons for opacity and CR rate
call krome_set_user_av(1d1) !opacity Av (#)E11.3,
```

```
call krome set user crate(1.3d-17) !CR rate (1/s)
  call krome set user gas dust ratio(7.57d11) !gas/dust
 call krome init()
 j21xs = 0d0
 call krome_set_J21xray(j21xs)
 print *,"J21 Xray = ",j21xs
 x(:) = 1.d-20
 !initial densities (model EA2 Wakelam+Herbst 2008)
 x(KROME idx H2) = 0.5d0 * xH !In this case we actually use <math>x(KROME idx H) =
0.5d0 * xH, which is not true
 x(KROME idx He) = 9d-2 * xH
 x (KROME idx N) = 7.6d-5 * xH
 x(KROME idx O) = 2.56d-4 * xH
 x(KROME idx Cj) = 1.2d-4 * xH
 x(KROME idx Sj) = 1.5d-5 * xH
 x(KROME_idx_Sij) = 1.7d-6 * xH
 x(KROME idx Fej) = 2d-7 * xH
 x(KROME_idx_Naj) = 2d-7 * xH
 x(KROME_idx_Mgj) = 2.4d-6 * xH
 x(KROME idx Clj) = 1.8d-7 * xH
 x(KROME idx Pj) = 1.17d-7 * xH
 x(KROME idx Fj) = 1.8d-8 * xH
 !calculate elctrons (neutral cloud)
 x(KROME_idx_e) = krome_get_electrons(x(:))
```

# Nov. 15, 2018

1. (From Prof. Du) The secondary ionization rate of molecular hydrogen can be approximately considered as 2 hydrogen atoms.

Related readings:

X-ray chemistry in the envelopes around young stellar objects, Stäuber 2005

X-Ray--irradiated Molecular Gas. I. Physical Processes and General Results, Maloney 1996

According to Maloney 1996, the main process in the X-ray (secondary) ionization of  $H_2$  is:

$$\mathrm{H_2} + \mathrm{e^-} \longrightarrow \mathrm{H_2^+} + 2\,\mathrm{e^-}$$

We can assume the rate coefficient is twice the ionization rates of H, and thus include it in the network.

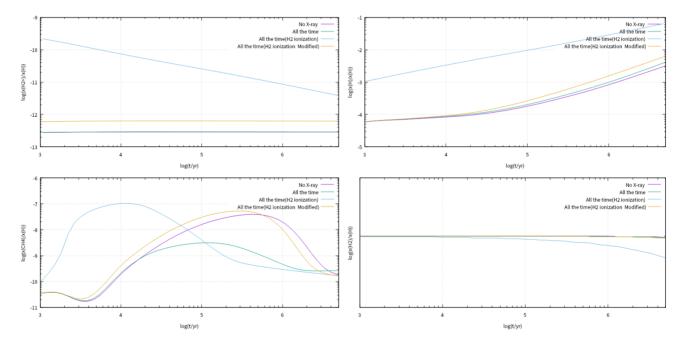
**INCLUDE** means when calculating the secondary ionization rates, we take the abundance of H:

$$n_t(\mathrm{H}) = n(\mathrm{H}) + 2n(\mathrm{H}_2)$$

The modification of Fortran code in krome subs.f90 is shown here

```
!H2 -> H2+ + E
```

## Some results are:



Obviously the ionization of  ${\bf H}_2$  becomes more significant, and the profile of  ${\bf CH}_4$  changes a lot. If we don't take  $n_t({\bf H})$  the effect is even more profound.

• (NOT AT PRESENT) For heavier atoms or other molecules, we can apply the approximation with the energy *E* fixed, neglecting the dependence of the cross sections on energy:

$$\zeta_i = \zeta_{
m H_2} rac{\sigma_i(E)}{\sigma_{
m H_2}(E)}$$

2. Species observants most interested in:

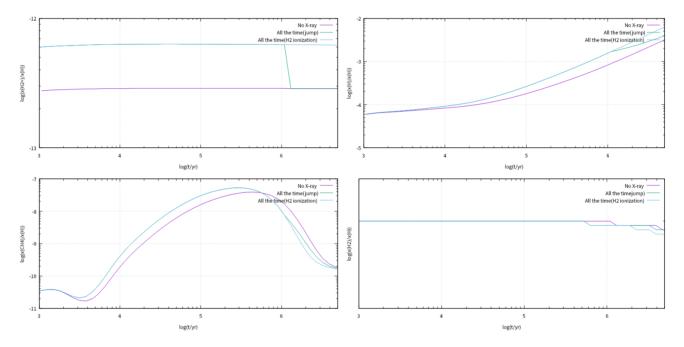
经常观测的分子包括CO、HCO<sup>+</sup>、 $N_2H^+$ 、HCN、各种碳氢化合物 $C_2H$ 、 $C_3H_3$ 等等。可以关注一下代表了电离度的电子丰度 $x_e$ ,以及 $H_3^+$ 、C<sup>+</sup>、OH<sup>+</sup>、 $H_2O^+$ 等离子的丰度;OH、CH、NH $_3$ 、 $H_2O$ 等也可以看看。还可以看更复杂一些的CH $_3$ CN、CH $_3$ OH等

 $x_e$  is important because it decides the chemical composition and whether the heating process is sufficient

- 3. The simple H/He model used before:
  - Abundance of CH<sub>4</sub> too low
  - $\circ~$  Not enough time to form  ${\rm H_2}$
  - Does not include H<sub>2</sub> formation on **grains**

Therefore it is not reasonable enough, and I shall focus on the new  $\rm H_2/He$  model to see what will happen if I turn off the X-ray after 1 mys.

Without  ${
m H_2}$  formation on **grains**, there seems no jump in the abundance of  ${
m CH_4}$ , which is shown in the plot.



# Nov. 22, 2018

1. More details on secondary ionization

The secondary ionization rate for H ( Stäuber 2005, Maloney 1996, ):

$$egin{align} \zeta_H^{(2)} &= \int_{E_{min}}^{E_{max}} rac{F(E)}{E} N_{sec}(E,x_e) \sigma(E) \mathrm{d}E \ N_{sec}(E,x_e) &= rac{E-E_{th}}{W(E)} pprox rac{E}{W(E)} \end{split}$$

 $N_{sec}(E)$  is the number of secondary **ionizations** per unit energy produced by primary photoelectrons W(E) is the mean energy expended to produce an ion pair through rapid electron process

#### A BIG MISTAKE HERE LAST WEEK!

 $N_{sec}(E)$  is not the number of secondary electrons but the number of ion pairs produced by a secondary electron. **WE DON'T COUNT ELECTRONS HERE!** 

#### More comments on the formula

- $\circ \ W(E)$  is nearly independent of E if  $E>200~{
  m eV}$  , which is easy to satisfied in our case
- o  $N_{sec}(E,x_e)$  is dependent on  $x_e$ . When  $x_e>1\%$ , the energy in photoelectrons will be mainly lost to Coulomb Interaction with ambient thermal electrons. In our case,  $n_e/n_{\rm H}<10^{-4}$ . For weakly ionized gases,  $\approx 40\%$  of the primary electron energy will be used for ionization (Krolik 1983, Shull 1985), then for a secondary electron with energy  $1~{\rm keV}$ , the number of  ${\rm H}$  it can ionize is:

$$N = rac{0.40 imes 1 ext{ keV}}{13.6 ext{ eV}} pprox 30$$

More precise calculation ( Glassgold & Langer 1973, Dalgarno 1999, Glassgold 2012 ) shows that in  ${
m H_2}+{
m He}$  neutral gas, 47% of energy goes to ionization, 47% in  ${
m H_2}$  and  ${
m He}$  excitation, the rest goes for heating.  $N_{sec}(E)\approx 27$  and thus we can take  $W(E)\approx 37~{
m eV}$ .

## The formula I actually use

$$\zeta_{ ext{tot}}^i = \zeta_{ ext{p}}^i + \sum_{j= ext{H.He}} rac{n_j}{n_i} \zeta_{ ext{p}}^j \left\langle \phi^i 
ight
angle$$

The numbers of secondary ionization of H and He per primary ionization are taken from Shull
 1985

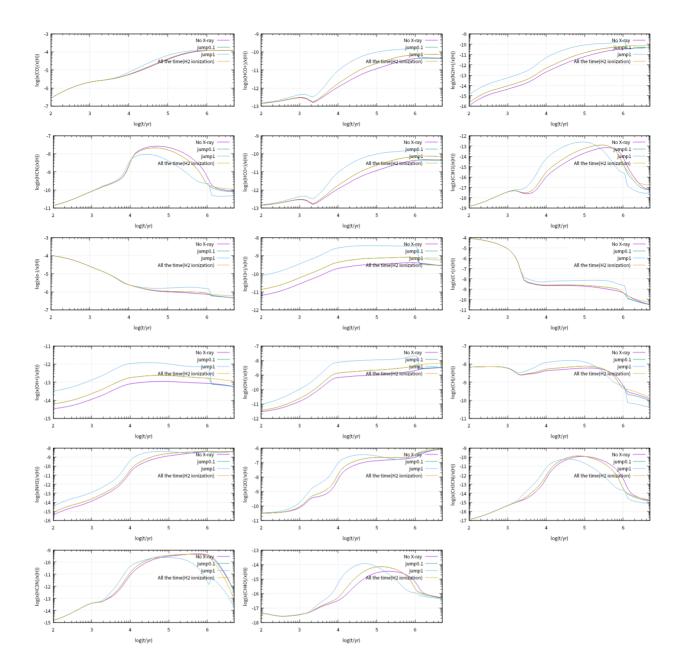
$$\phi^{ ext{H}}\left(E,x_{ ext{e}}
ight) = \left(rac{E}{13.6 ext{eV}}-1
ight) 0.3908 ig(1-x_{ ext{e}}^{0.4092}ig)^{1.7592} \ \phi^{ ext{He}}\left(E,x_{ ext{e}}
ight) = \left(rac{E}{24.6 ext{eV}}-1
ight) 0.0554 ig(1-x_{ ext{e}}^{0.4014}ig)^{1.660}$$

o Can be rewritten like this:

$$egin{aligned} \zeta_{tot}^{
m H} &= \zeta_p^{
m H} (1 + \langle \phi^{
m H} 
angle) + \zeta_p^{
m He} \, rac{n_{
m He}}{n_{
m H}} \langle \phi^{
m H} 
angle \ \zeta_{tot}^{
m He} &= \zeta_p^{
m He} (1 + \langle \phi^{
m He} 
angle) + \zeta_p^{
m H} \, rac{n_{
m H}}{n_{
m He}} \langle \phi^{
m He} 
angle \end{aligned}$$

which average  $\phi$  over the X-ray spectrum F(E)

2. Test all the species listed last week



# Nov. 29, 2018

1. Recalculate the ionization rate using the primal formula

$$\zeta_{tot}^{\mathrm{H}} = \zeta_{p}^{\mathrm{H}} + \frac{n_{\mathrm{He}}}{n_{\mathrm{H}}} \int_{E_{min}}^{E_{max}} \frac{F(E)}{E} e^{-\tau(E)} \phi^{\mathrm{H}} \sigma_{\mathrm{He}}(E) \mathrm{d}E + \int_{E_{min}}^{E_{max}} \frac{F(E)}{E} e^{-\tau(E)} \phi^{\mathrm{H}} \sigma_{\mathrm{H}}(E) \mathrm{d}E$$

rather than the simplified one

$$\zeta_{tot}^{
m H} = \zeta_p^{
m H} (1 + \langle \phi^{
m H} 
angle) + \zeta_p^{
m He} rac{n_{
m He}}{n_{
m H}} \langle \phi^{
m H} 
angle$$

and compare the results.

When we set  $n_{
m H}=2 imes 10^4~{
m cm^{-3}}, n_{
m He}=2 imes 10^3~{
m cm^{-3}}, J_{X21}=0.1, T=160~{
m K}$  , we get

$$\begin{split} \zeta_{\rm H} &= 4.09 \times 10^{-18} \ {\rm s}^{-1}, \quad \zeta_{\rm He} = 3.61 \times 10^{-18} \ {\rm s}^{-1} \quad \text{(Primal)} \\ \zeta_{\rm H}' &= 7.23 \times 10^{-18} \ {\rm s}^{-1}, \quad \zeta_{\rm He}' = 6.07 \times 10^{-18} \ {\rm s}^{-1} \quad \text{(Simplified)} \\ \frac{\zeta_{\rm H}}{\zeta_{\rm H}'} &\approx \frac{\zeta_{\rm He}}{\zeta_{\rm He}'} \approx 0.56 \end{split}$$

Since we take  $E_{min}=2~{\rm keV}, E_{max}=10~{\rm keV}$  , and in our case,  $F(E)\propto E^{-1.5}$  ,  $\sigma(E)\propto E^{-3}$  ,  $\phi(E,x_e)\propto E$  ,  $e^{-\tau(E)}\sim 1$ , we have

$$\zeta^{(2)} = \int_{E}^{E_{max}} \frac{F(E)}{E} \phi(E) \sigma(E) \mathrm{d}E = F_0 \phi_0 \sigma_0 \cdot \int_{2}^{10} E^{-4.5} \mathrm{d}E = \frac{2^{-3.5} - 10^{-3.5}}{3.5} F_0 \phi_0 \sigma_0 \approx 0.025 F_0 \phi_0 \sigma_0$$

and

$$egin{aligned} \int_{E_{min}}^{E_{max}} rac{F(E)}{E} \sigma(E) \mathrm{d}E &= F_0 \sigma_0 \int_2^{10} E^{-5.5} \mathrm{d}E pprox 0.0098 F_0 \sigma_0 \ \int_{E_{min}}^{E_{max}} F(E) \phi(E) \mathrm{d}E &= F_0 \phi_0 \int_2^{10} E^{-0.5} \mathrm{d}E pprox 3.50 F_0 \phi_0 \ \int_{E_{min}}^{E_{max}} F(E) \mathrm{d}E &= F_0 \int_2^{10} E^{-1.5} \mathrm{d}E pprox 0.078 F_0 \ &\Rightarrow \zeta^{(2)'} pprox rac{0.0098 F_0 \sigma_0 \cdot 3.50 F_0 \phi_0}{0.078 F_0} pprox 0.044 F_0 \phi_0 \sigma_0 pprox 0.57 \zeta^{(2)} \end{aligned}$$

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origin/master

- 2. Why does KROME use a simplified formula to calculate the total ionization rate?
  - $\circ \zeta_p^{\mathrm{H}}, \zeta_p^{\mathrm{He}}$  are stored in data-files rateH.dat and rateHe.dat (column density *v.s.* primary ionization rate)
  - To calculate  $\zeta^{(2)}$  the program needs to read  $x_e$ ,  $n_{\rm H}$ ,  $n_{\rm He}$  for each calculation
  - o By using the simplified formula, we may avoid integration

Here is part of the codes ( krome subs.f90 ):

```
user xray H anytabz(:,:), &
        user_xray_H_anytabxmul, &
        user xray H anytabymul, &
       logH,logHe-logH)
        !linear interpolation to find the proper rate for the column density
   phiH = .3908d0*(1e0-xe**.4092)**1.7592* 327.832286034056d0! avoid
integration
   ratexH = 1d1**user xray H
   user xray He = fit anytab2D(user xray He anytabx(:), &
       user xray He anytaby(:), &
       user xray He anytabz(:,:), &
       user xray He anytabxmul, &
       user xray He anytabymul, &
       logH, logHe-logH)
   phiHe = .0554d0*(1d0-xe**.4614)**1.666*180.793458763612d0
   ratexHe = 1d1**user xray He
   k(:) = small !inizialize coefficients
   !H -> H+ + E
   k(1) = small + ((ratexH &
       * (1d0+phiH) + n(idx He)&
       /(n(idx_H)+1d-40) * ratexHe * phiH)&
        * J21xray)
   !HE -> HE+ + E
   k(2) = small + ((ratexHe &
        * (1d0+phiHe) + n(idx H)&
        /(n(idx He)+1d-40) * ratexH * phiHe)&
        * J21xray)
   print '(a10,E15.8)', "XrateH:", log10(k(1))
   print '(a10,E15.8)', "XrateHe:", log10(k(2))
   coe(:) = k(:) !set coefficients to return variable
```

# Dec. 5, 2018

- 1. Try to repeat some results in former papers...but pretty difficult to do
  - $\circ$  Krolik 1983 did not consider the X-ray ionization of  $H_2$ 
    - Dec. 12 review: Actually,  ${
      m H_2}$  ionization is considered with an approximate model (Brown and Gould's 1970)
  - $\circ$  Those who have considered the x-ray ionization of  ${
    m H_2}$  usually have too complex physical problems and models, which means they are not easy to fully understand and repeat within hours
  - The model we consider is simple and does not include heating/cooling process and totally ignore
     UV-photon process
- 2. Add molecular hydrogen formation process to the network (from react\_primordial, primordial rates from several authors mostly based on Abel et al. 1997 + Grassi et al. 2014)

# Dec. 12, 2018

1. Reread and try to repeat the results in Krolik 1983.

#### A Review on Krolik 1983

- Modify the codes
  - Initial abundances
  - Longer time period (steady-state solutions)  $\Delta t=10^8~{
    m yrs}$  (Strange tendency appears when  $\Delta t>10^8~{
    m yrs}$ , abundances of some species changes rapidly)
  - Set two branches for  $H_2$  ionization & dissociation
  - $\blacksquare \ \ \mathsf{Add} \ \mathsf{H}_2$  formation process on grains
- Step 1. Still use the power law spectrum,
  - It is the total ionization rate rather than the detail of the spectrum that is the most important to the chemistry
- 2. A new paper on the enhanced molecular ionization in the central molecular zone of our Galaxy Molecular Ionization Rates and Ultracompact Dark Matter Minihalos, Silk, J. 2018, PRL

# Dec. 20, 2018

- 1. Grain Process
  - $\circ~$  It's important to find out how  $\rm H_2$  molecules form on the grain surface as well as how KROME deal with the process

$$H + H + dust \longrightarrow H_2 + dust$$

o Theoretically, KROME uses

$$rac{\mathrm{d}n_{\mathrm{H}_{2}}}{\mathrm{d}t}=rac{\pi}{2}n_{\mathrm{H}}v_{g}\sum_{j\in\left[\mathrm{C,Si}
ight]}\sum_{i}n_{ij}a_{ij}^{2}\epsilon_{j}\left(T,T_{i}
ight)lpha\left(T,T_{i}
ight)$$

where  $v_g$  is the gas thermal velocity,  $a_{ij}$  is the size of the grain,  $T_i$  is the dust temperature,  $\epsilon_j$  depends on the type of the dust and  $\alpha$  is the sticking coefficient

• The rate is defined in the file krome dust.f90

```
function krome_H2_dust(nndust, Tgas, Tdust, nH, H2_eps_f, myvgas)
!H2 formed on dust (1/cm3/s)
use krome_constants
use krome_commons
real*8::H2_dust, krome_H2_dust, Tgas, Tdust(:)
real*8::myvgas, H2_eps, nndust(:), nH, H2_eps_f
integer::i

H2_dust = 0.d0
do i = 1, size(Tdust)
```

as well as related functions. Again in  $krome_ode.f90$  this rate is added to the total rate of  $H_2$  formation. The process is shown in kromeobj.py where

```
#add H2 formation on dust
if(self.useDustH2const):
   if("H"==specs[idnw].name): x += " - 2d0*nH2dust"
   if("H2"==specs[idnw].name): x += " + nH2dust"
```

#### 2. Slightly modify the X-ray ionization

- In Krolik 1983, there are two branches for molecular Hydrogen ionization (See X-ray module in the <u>review</u>)
- Modify the krome subs.f90

```
!H2 -> H+ + H + E
k(36) = rateEvaluateOnce(36)
!H2 -> H2+ + E
k(38) = rateEvaluateOnce(38)
!H -> H+ + E
   k(4452) = small + ((ratexH &
        * (1d0+phiH) + n(idx He)&
        /(n(idx_H)+2*n(idx_H2)+1d-40) * ratexHe * phiH) &
        * J21xray)
    k(36) = k(36) + 2 * k(4452) * 0.2
    k(38) = k(38) + 2 * k(4452) * 0.8
!HE -> HE+ + E
    k(4453) = small + ((ratexHe &
        * (1d0+phiHe) + (n(idx H)+2*n(idx H2))&
        /(n(idx He)+1d-40) * ratexH * phiHe)&
        * J21xray)
```

#### 3. Repeat Krolik's results

- Cosmic ray
  - well matched, except for several species such as  $\mathrm{HNO}^+$  which has a larger lpha
  - lacktriangle The difference in lpha between stronger and weaker ionization rates is larger