

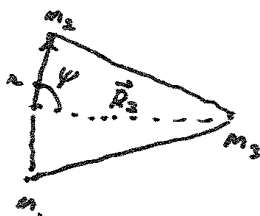
Sec. Dyn. Reading Group

13 July 2010 ①

Jean + Sara, Murray & Dermott Ch 7

Secular perturbations: those that never vanish on $\langle \rangle_{\text{orbit}}$
(as opposed to resonant pert., which are only nonzero $\langle \rangle$ when
in resonance).

disturbing fcn: $R = \frac{G m_1 m_2 m_3}{2(m_1 + m_2) R_3} \left(\frac{R_2}{R_3}\right)^2 (3 \cos^2 \psi - 1)$



Want $\langle R \rangle = \frac{1}{T} \int_0^T R dT \stackrel{\text{define angle } \phi}{=} \frac{1}{2\pi} \int_0^{2\pi} R d\phi$

(Effectively average time: once over inner orbit, once over
outer orbit.) See, e.g. 6.164 from M & D, which comes from
6.107.

Recall D'Alembert condition, $\ddot{q}_j = 0$ (j keeps track of the
orders of the various angles in disturbing function's Fourier series.)
If set $\ddot{q}_j = 0$, have coverage over λ ; recall remaining constraint

Let's do it $\frac{2}{3}$

Recall Poincaré variables: eg 2.179 $(\lambda, \ell, z, \Lambda, P, \tilde{z})$.

$\gamma = -\omega \cdot R$; $P = \alpha(1 - \sqrt{1 - e^2}) \approx e^2/2 \Rightarrow e = \sqrt{2P}$

$R = K_1 \frac{A_1}{R_1^3} \frac{A_2}{R_2^3} \frac{A_3}{R_3^3} \cos(\psi) \approx k_1 e^2 + k_2 e'^2 + k_3 e e' \cos(\tilde{\omega} - \tilde{\omega}')$

Equation of motion (Hamiltonian):

$\frac{d\ell}{dt} = \partial R / \partial P$; $\frac{dP}{dt} = -\partial R / \partial \gamma$

13 July 2018 (2)

Transform to a complex canonical variable:

$$z = \sqrt{\rho} e^{i\delta} \Rightarrow \frac{dz}{dt} = +i \partial H / \partial z^* \quad (\text{note } z^* = \sqrt{\rho} e^{-i\delta})$$

Apply to $1/2^n$ system: $z = \sqrt{\rho'} e^{i\delta'}$ or $z = \frac{\rho'}{\sqrt{2}} e^{i\bar{\omega}}$

$$\text{then } \langle R \rangle = 2k_1 |z|^2 + k_3 e' \frac{\sqrt{2}}{2} \left\{ e^{-i(\delta' + \bar{\omega}')} + e^{i(\delta' + \bar{\omega}')} \right\} |z|$$

↑ Note: 2 terms collapse into one because $k_1 = k_2$.

$$\langle R \rangle = \dots + k_3 e' \frac{\sqrt{2}}{2} \left\{ z e^{i\bar{\omega}'} + z^* e^{-i\bar{\omega}'} \right\}$$

$$\Rightarrow \frac{dz}{dt} = i \left[2k_1 z + \frac{\sqrt{2}}{2} k_3 e' e^{-i\bar{\omega}'} \right] \quad (\text{combine (6.170), (6.171) in complex form})$$

Re-write as $dz/dt = ia z + ib z'$, and from there to

$$e^{-i\omega_0 t} \frac{d}{dt} \left\{ z e^{i\omega_0 t} \right\} = ib z' \quad \text{where } i\omega_0 = ia = 2ik_1.$$

Now seek particular solution for this equation:

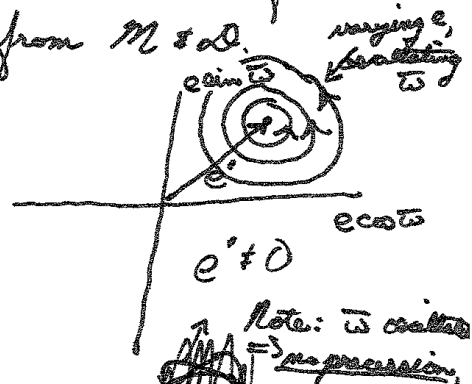
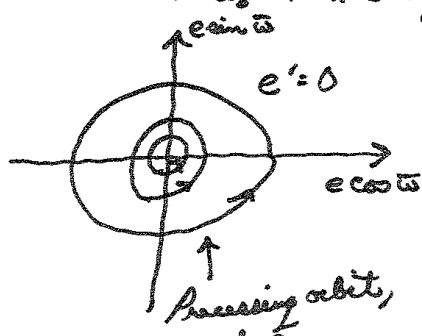
$$z = \frac{b}{\omega_0} z' + C_0 e^{-i\omega_0 t}$$

This gives the motion of the orbit as a specific oscillation (free) plus a forced term $\frac{b}{\omega_0} z'$ from the effects of perturbation e .

\Rightarrow if $e' = 0$, then e rotates in orbit has const magnitude;

if $e' \neq 0$ then e magnitude fluctuates about forced

solution $z = \frac{b}{\omega_0} z'$. See fig. 7.2 from M & A.



13 July 2018 (3)

Now force outer planet to precess:

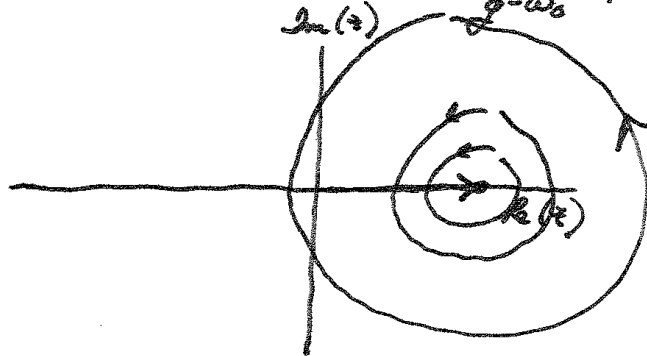
$$\frac{dz}{dt} = i(\omega_0 z + e' e^{igt}) \quad g \text{ is precession rate of pericenter planet.}$$

$$\text{Solution is } z(t) = k e^{i\omega_0 t} + \frac{\omega_0 e'}{g - \omega_0} e^{igt}$$

The denominator $g - \omega_0$ allows for resonance if precession of pericenter (g) occurs at the same rate as free precession of planet (ω_0).

If we go into "rotating frame" and plot $e^{-igt} z$:

$$e^{-igt} z = k e^{i(\omega_0 - g)t} + \frac{\omega_0 e'}{g - \omega_0}$$



\Rightarrow small Δe : librate about peric. } modified freq $\omega_0 - g$
 large Δe : Rotate

$\Delta e = |z| - |z'|$ controls the distance between fixed pt and circle.

[Yoram told a story about nonlinear resonances that I don't have down here.] Yoram is preparing a paper about Mercury, Jupiter in the framework (all plus non-linear terms) and

13 July 2010

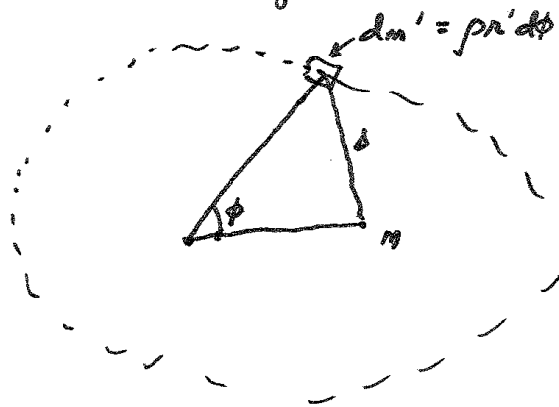
(4)

Gauss' Method

Test bodies as rings distributed over their orbits
proportional to the time they spend @ each location.

Then apply perturbative forces (see Ch 2) due to
ring on orbit. (e.g. 2.165)

Gauss' method calculates the forces: $dF = \bar{R}\hat{r} + \bar{T}\hat{\theta} + \bar{N}\hat{z}$ on
the orbit.



See 7.98. Gauss' method does $V = \int dm \frac{1}{\delta}$, $F = -\nabla V$.

When done in book, find 7.123. Can check prior expansions
in e, i , etc. by expanding Gauss' approach (i.e. Gauss does
not assume that $e \ll 1$ or $i \ll 1$).