

On constraining the spin of the MBH the GC via star orbits: the effects of stellar perturbations



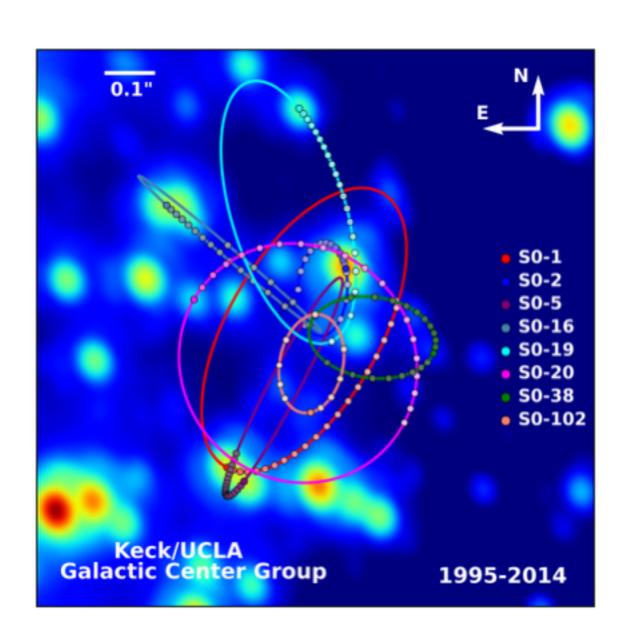
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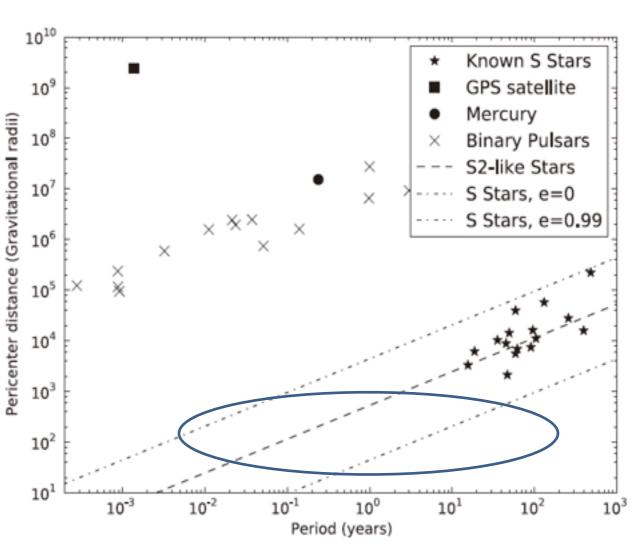
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2016. 2. 11, Aspen Center for Physics

Strong field GR test and the GC S-stars

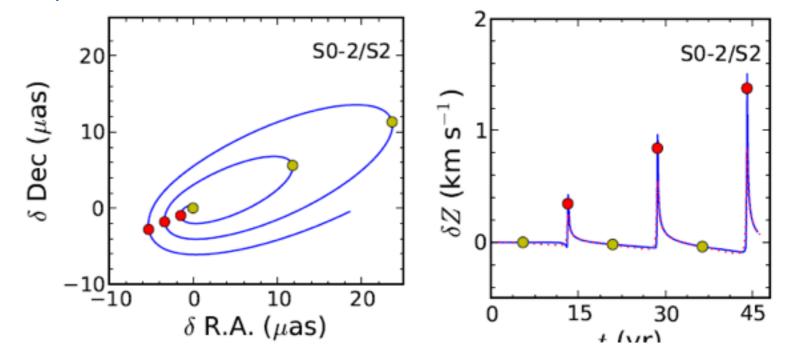
- Clusters of young stars in the GC
- Provide a unique environment of testing GC by stellar orbits



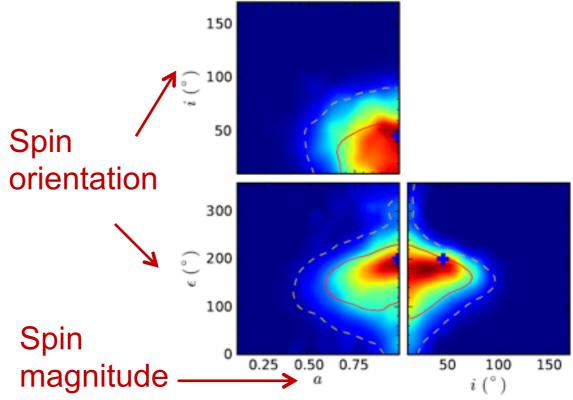


Our previous works

- The constraints of the spin parameters by observing the trajectories and the redshifts of the S-stars by future facilities (Zhang, Lu, & Yu 2015; Yu, Zhang, & Lu, submitted)
 - Full GR treatment
 - MCMC fitting
 - Magnitude and direction of spin, 6 orbital elements, MBH and R_{GC}



- We can constraint the spin by observing the orbits of S2 or other inner S-stars
- But stellar perturbations are not considered



Perturbations

- Stars: Early and late type stars (Bartko, et, al. 2010)
- Stellar remnants
 - Stellar mass black holes: Mass segregation (Freitag, et, al. 2006)
 - Neutron stars, pulsars, white dwarfs: (Morris 1993)
- Intermediate mass black hole(s): 100-1000 solar mass, distance>200
 AU (Yu & Tremaine 2003, Gualandris & Merritt 2009; etc)
- Dark matter
- Distinguish
 - Gravitational perturbations from background sources
 - Spin-induced perturbations
 - Different predictions from other gravity theories (e.g., f(R) theory)

Previous studies

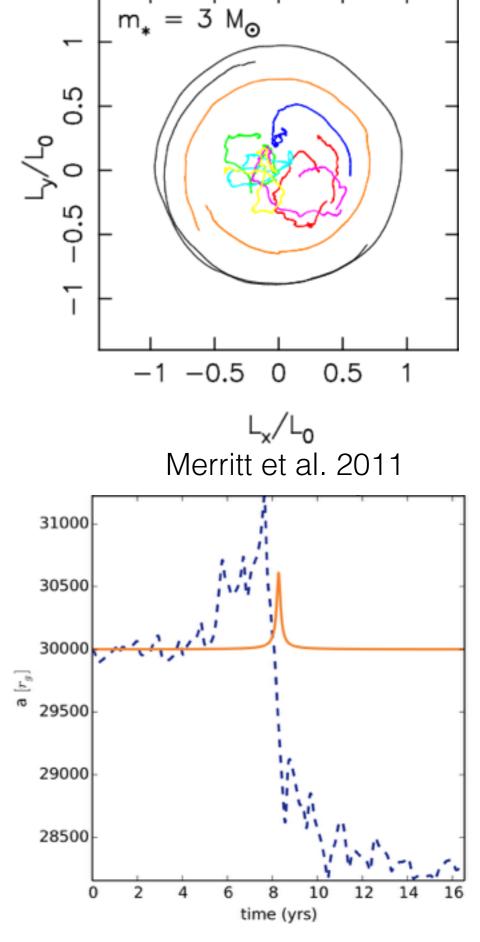
- Post-NW approximation (Merritt et al. 2010)
 - Frame-dragging obscured beyond 0.5mpc

$$\begin{aligned} \mathbf{a}_{J,1} &= -\frac{3G^2M_{\bullet}}{c^3} \sum_{j \neq 1} \frac{m_j}{r_{1j}^3} \bigg\{ \left[\mathbf{v}_{1j} - \left(\mathbf{n}_{1j} \cdot \mathbf{v}_{1j} \right) \mathbf{n}_{1j} \right] \times \boldsymbol{\chi} - 2\mathbf{n}_{1j} \left(\mathbf{n}_{1j} \times \mathbf{v}_{1j} \right) \cdot \boldsymbol{\chi} \bigg\}, \\ \mathbf{a}_{J,j} &= \frac{2G^2M_{\bullet}^2}{c^3r_{1j}^3} \bigg\{ \left[2\mathbf{v}_{1j} - 3\left(\mathbf{n}_{1j} \cdot \mathbf{v}_{1j} \right) \mathbf{n}_{1j} \right] \times \boldsymbol{\chi} - 3\mathbf{n}_{1j} \left(\mathbf{n}_{1j} \times \mathbf{v}_{1j} \right) \cdot \boldsymbol{\chi} \bigg\}, \\ \dot{\boldsymbol{\chi}} &= \frac{G}{2c^2} \sum_{j \neq i} \frac{m_j}{r_{ij}^2} \left[\mathbf{n}_{1j} \times (3\mathbf{v}_1 - 4\mathbf{v}_j) \right] \times \boldsymbol{\chi}, \\ r_{ij} &= |\mathbf{x}_i - \mathbf{x}_j|, \quad \mathbf{x}_{ij} \equiv \mathbf{x}_i - \mathbf{x}_j, \quad \mathbf{n}_{ij} = \mathbf{x}_{ij}/r_{ij}, \quad \mathbf{v}_{ij} \equiv \mathbf{v}_i - \mathbf{v}_j. \end{aligned}$$

- Hamiltonian perturbation (Angelil & Saha 2014)
 - Frame-Wavelet decomposition

$$H_{\text{stellar}} = \sum_{j} \frac{m_{j}}{M} \left(\frac{\boldsymbol{x} \cdot \boldsymbol{x}_{j}}{|\boldsymbol{x}_{j}|^{3}} - \frac{1}{|\boldsymbol{x} - \boldsymbol{x}_{j}|} \right),$$

Orbital perturbation theories
 Sadeghian & Will 2011; Iorio 2011; etc



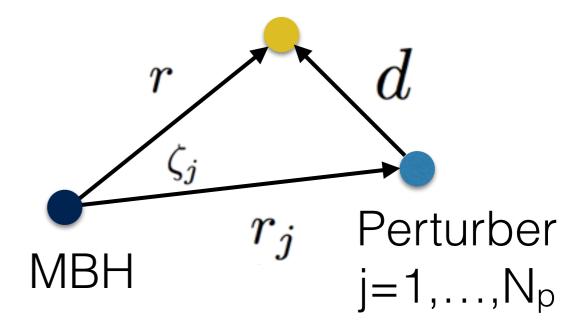
Angelil & Saha 2014

Motion of the perturbed target star

 Hamiltonian contributed by perturbation (Angelil & Saha 2014, Wisdom & Holman 1991)

target star

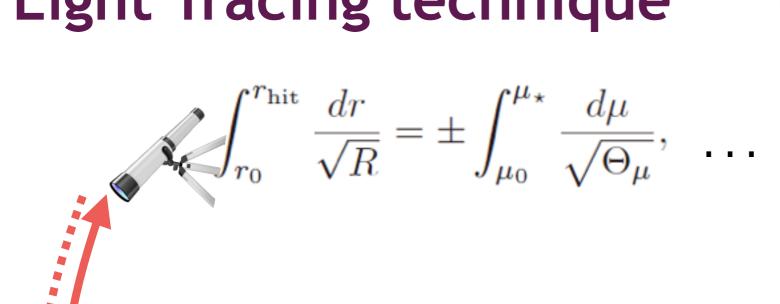
$$H_{\rm p} = \sum_{j}^{N_p} m_{{\rm p},j} \left(\frac{r}{r_j^2} \cos \zeta - \frac{1}{d} \right)$$



- Simplification
 - The multure interactions between perturber are ignored
 - The target star is a test particle (mass=0)
- Motions of the perturbers follows the unperturbed motion equation

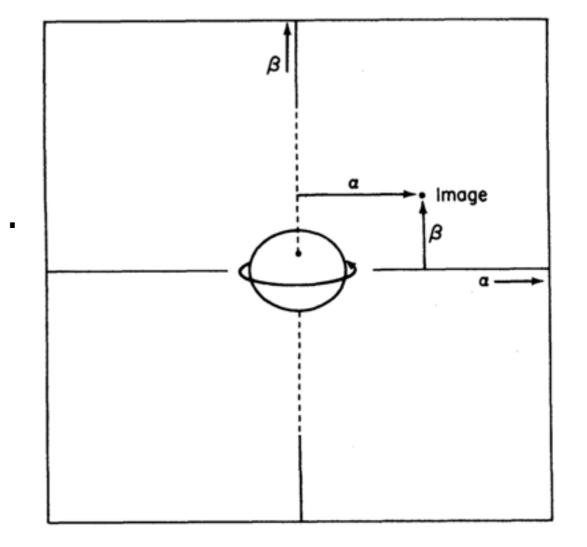
$$H=H_0+H_p$$

Light Tracing technique

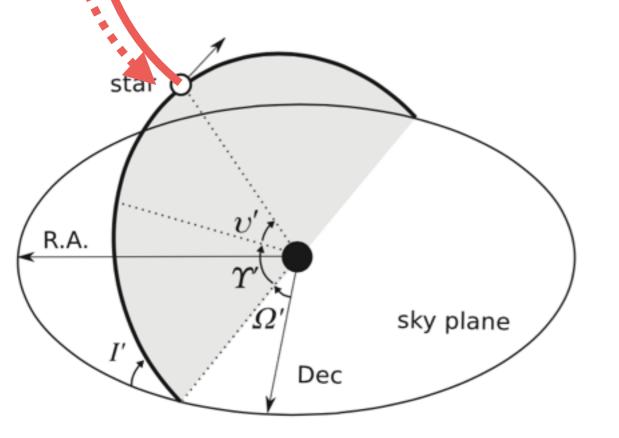


$$\lambda = -\alpha \sin i,$$

$$q^2 = \beta^2 + (\alpha^2 - a^2)\cos^2 i.$$



Cunningham & Bardeen 1972



$$lpha r_{
m g}/R_{
m GC}$$
 R.A. $eta r_{
m g}/R_{
m GC}$. Dec

$$Z = \frac{\boldsymbol{p}_{\text{hit}} \cdot \boldsymbol{u}_{\star}}{\boldsymbol{p}_{\text{o}} \cdot \boldsymbol{u}_{\text{o}}} - 1 = -\frac{\boldsymbol{p}_{\text{hit}} \cdot \boldsymbol{u}_{\star}}{E_{\text{o}}} - 1.$$

Perturbations on the observational quantities

Positions of the star in the sky at time t

$$\delta R.A.$$

$$\delta \mathrm{Dec}$$

$$\delta R.A.$$
 δDec $\delta R(t) = \sqrt{\delta R.A.^2 + \delta Dec^2}$

Redshift at time t

$$\delta Z(t)$$

Root mean square value (in three orbits)

$$\delta R_{\rm rms} = \sqrt{\frac{1}{T_{\rm ob}} \int^{T_{\rm ob}} \delta R(t)^2 dt}$$

$$\delta R_{\rm rms} = \sqrt{\frac{1}{T_{\rm ob}}} \int^{T_{\rm ob}} \delta R(t)^2 dt$$
 $\delta Z_{\rm rms} = \sqrt{\frac{1}{T_{\rm ob}}} \int_0^{T_{\rm ob}} \delta Z(t)^2 dt,$

Spin-induced effects:

unperturbed target star, *a=0.99*



unperturbed target star, **a=0.0**

Stellar perturbations:

perturbed target star, a=0.0



unperturbed target star, a=0.0

Total perturbations:

perturbed target star, a=0.99



unperturbed target star, a=0.0

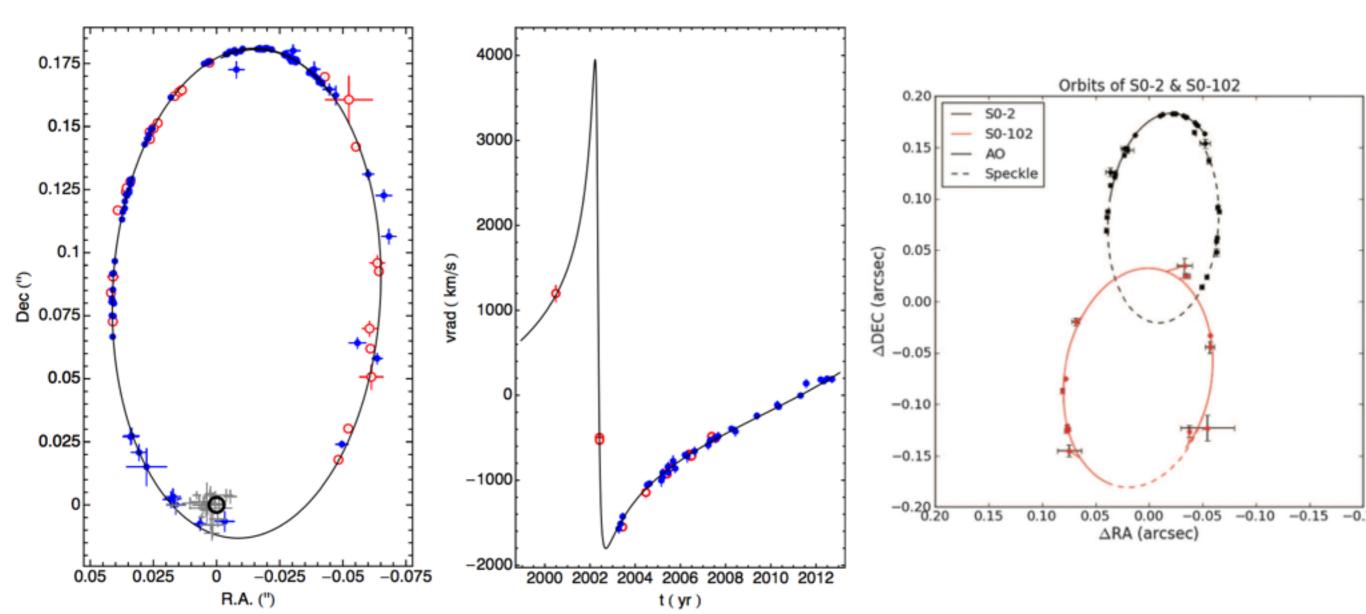
Results: Single Perturber

S2/S0-2 and S0-102

- S2/S0-2
 - Orbital period of 15 years
 - Pericenter distance of 100AU

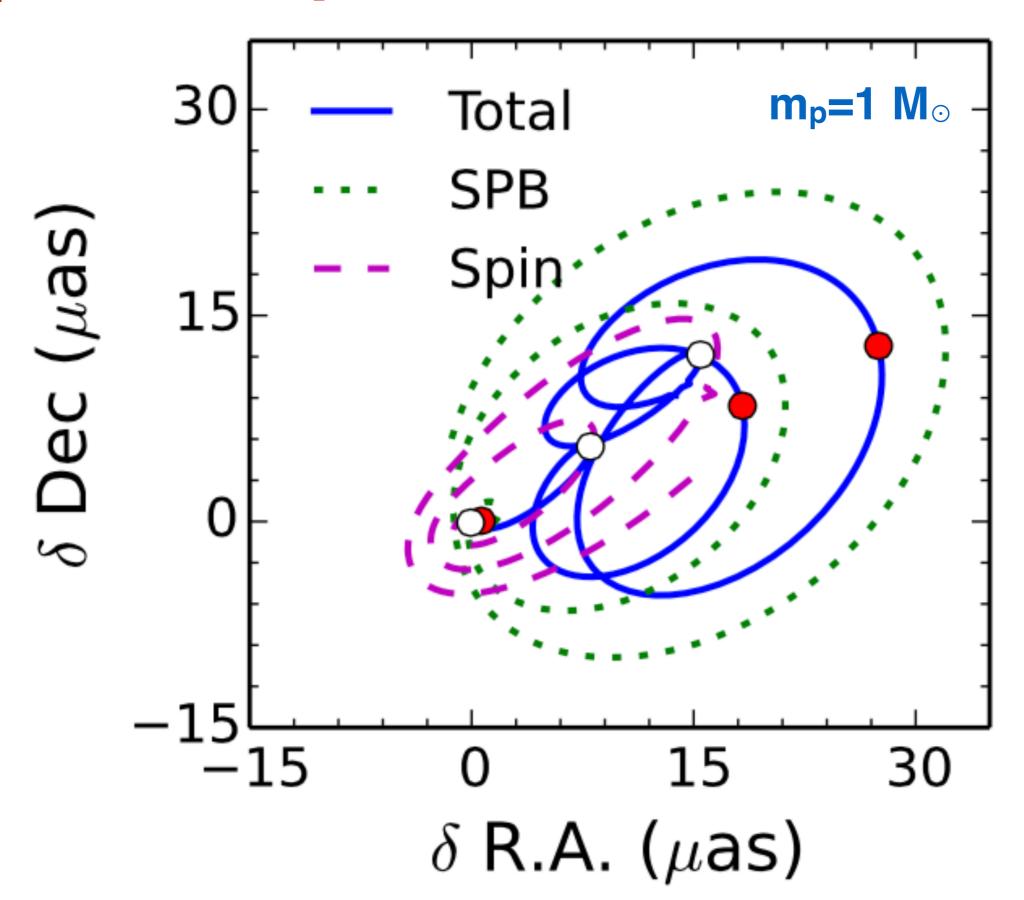
- S0-102
 - Orbital period of 11 years
 - e~0.68

How S0-102 affects the orbital motion of S2?

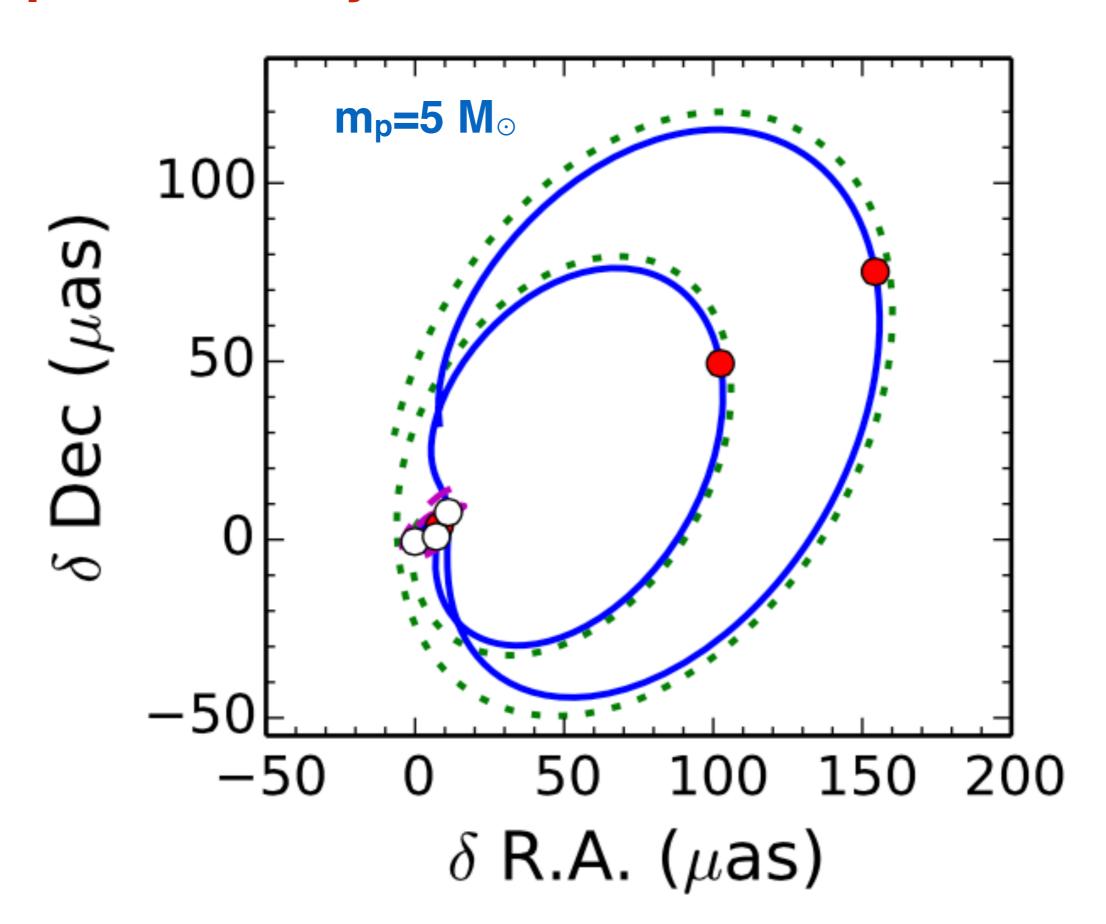


Gillessen et al. 2013

S2 perturbed by S0-102

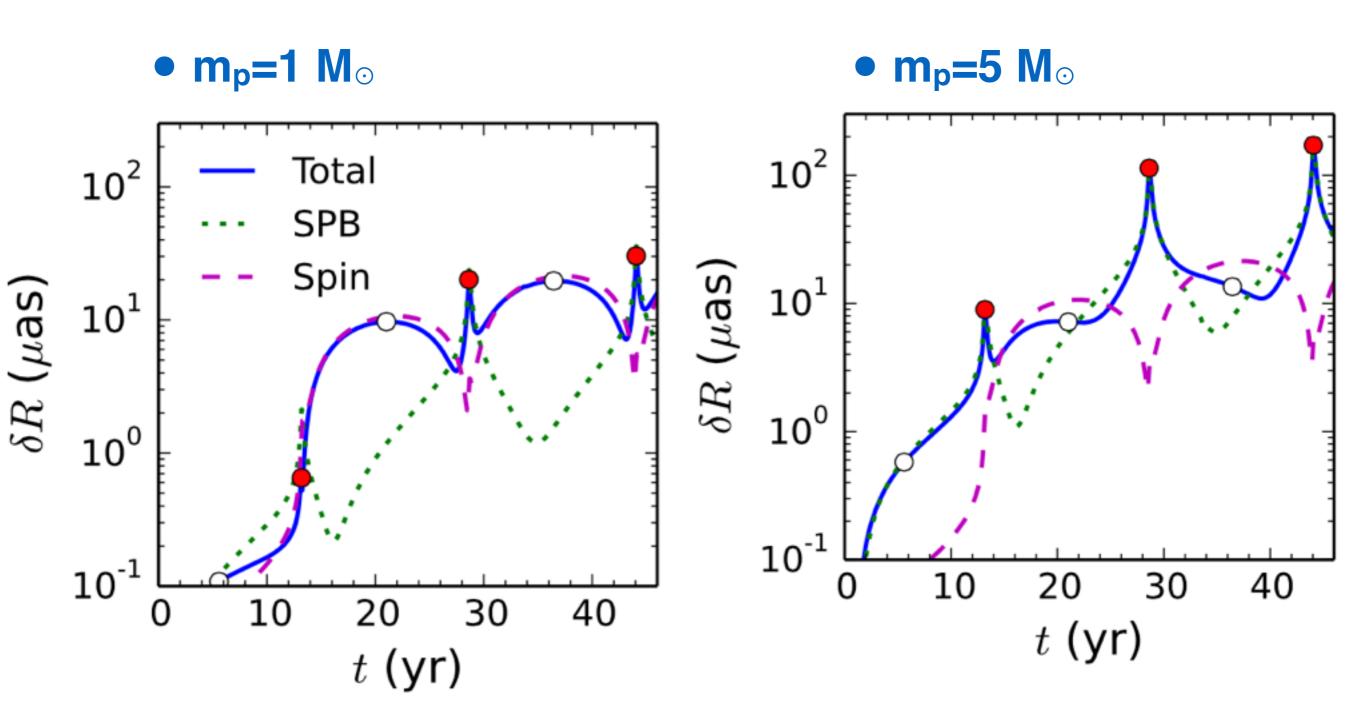


S2 perturbed by S0-102

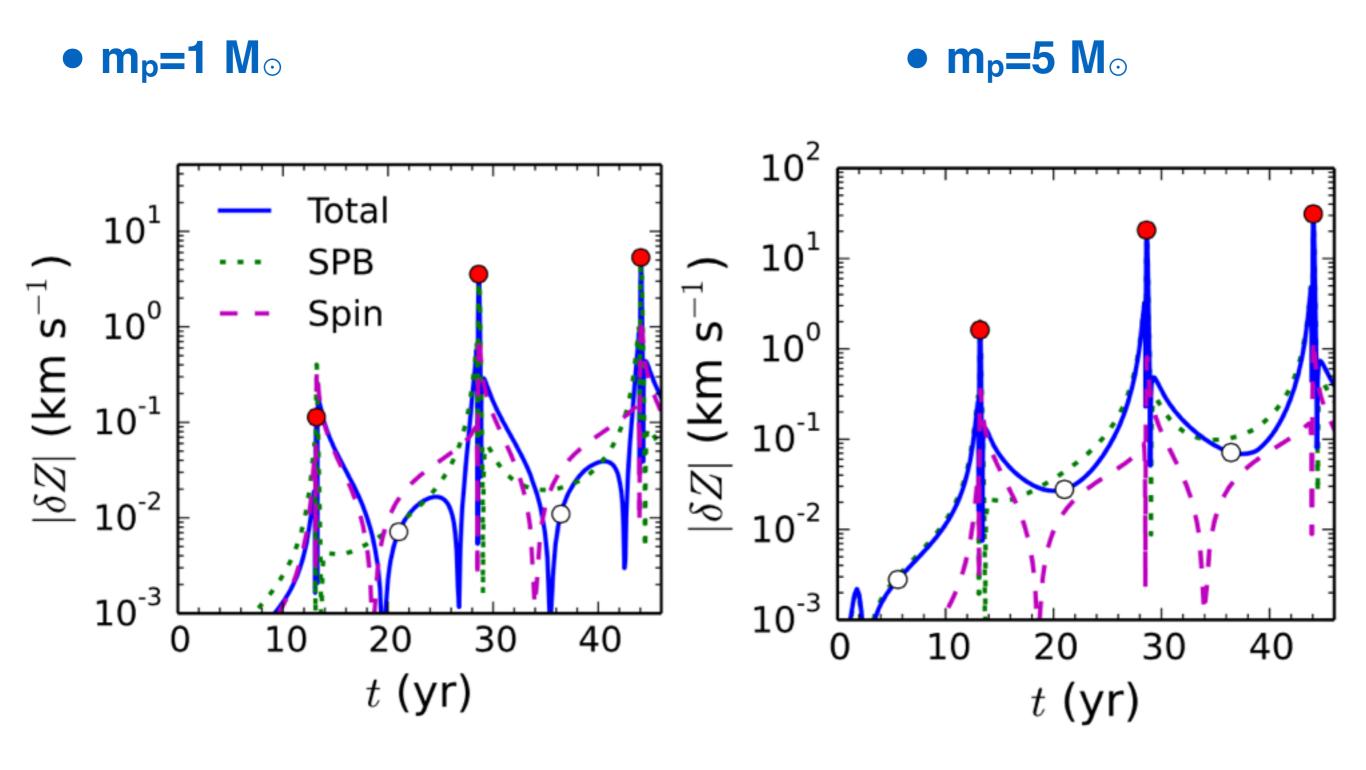


Orbits of S2 perturbed by S0-102

$$\delta R(t) = \sqrt{\delta R.A.^2 + \delta Dec^2}$$

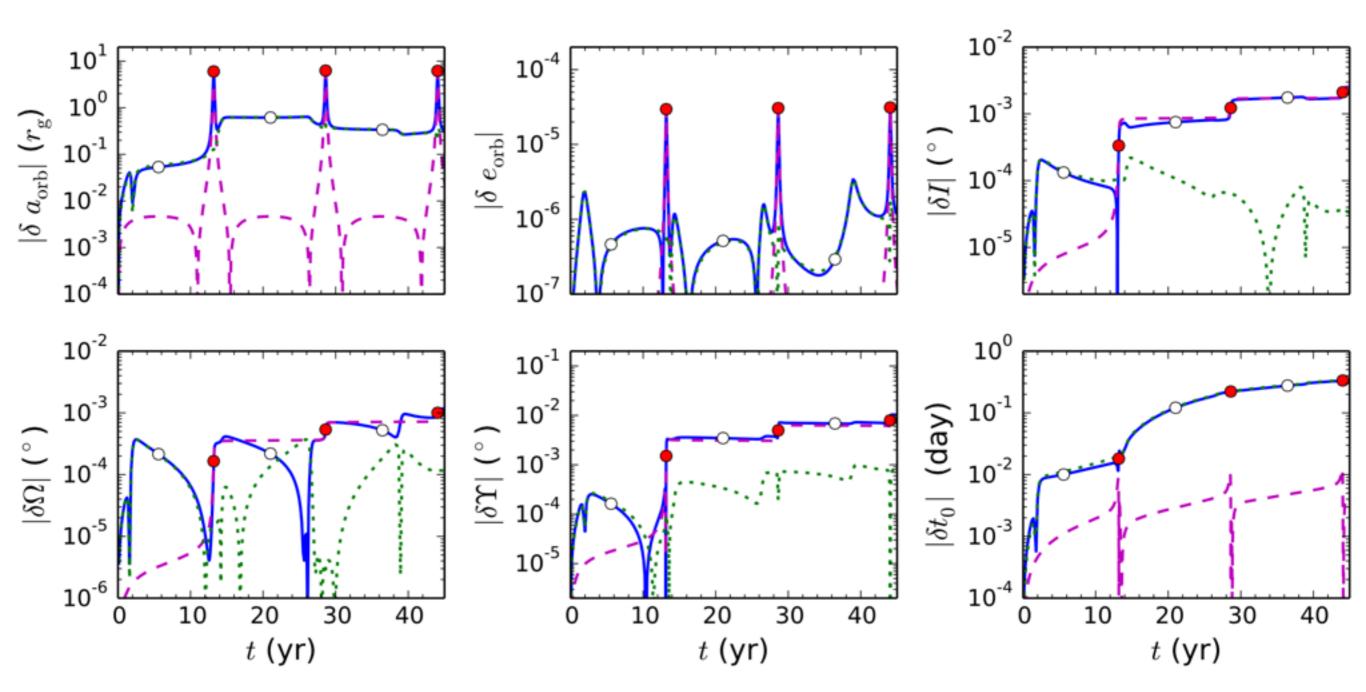


Orbits of S2 perturbed by S0-102



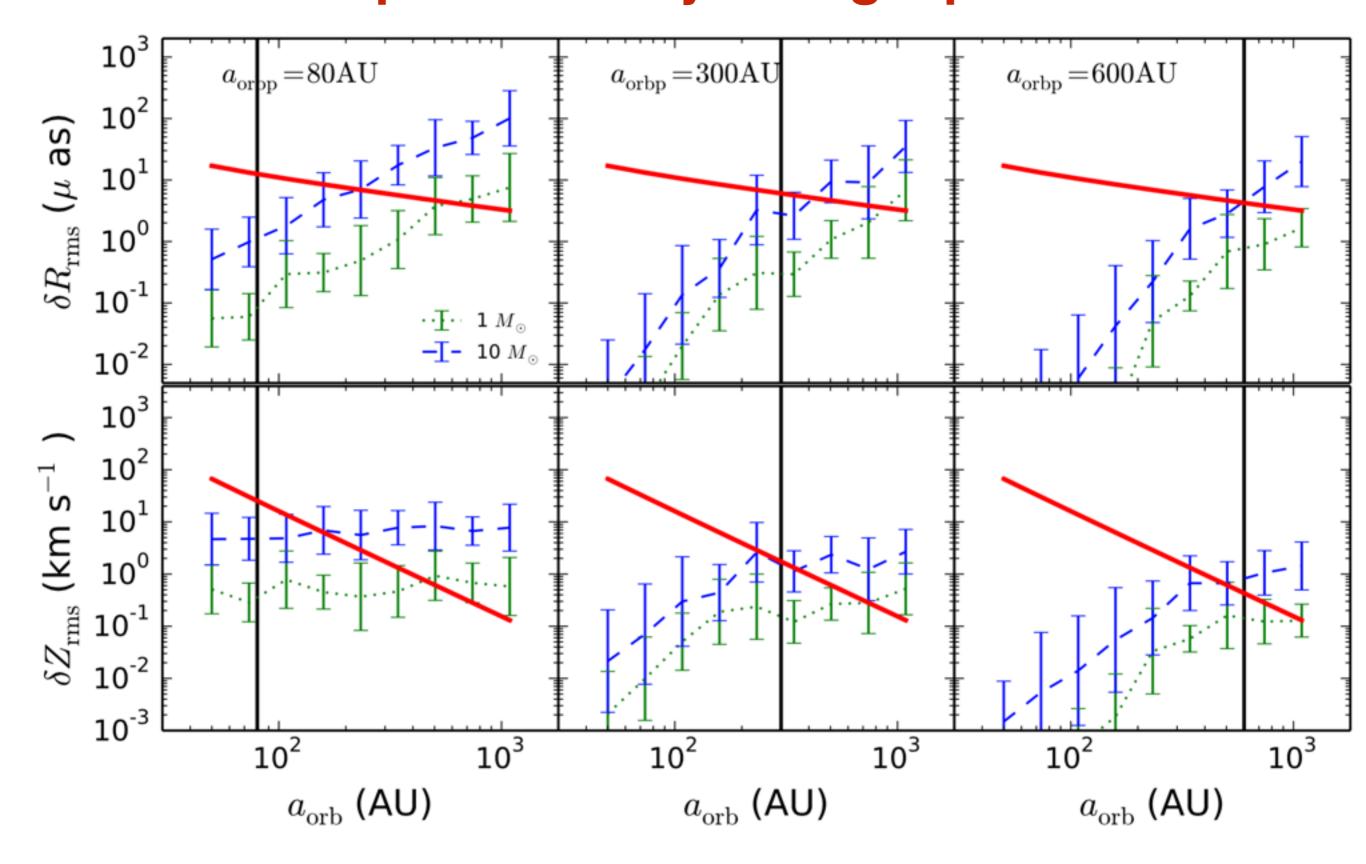
Orbits of S2 perturbed by S0-102

Variations of the orbital elements of S2



The orbital period of S2 is perturbed: ldt₀l~ 0.3 day after 45 years —>
 ~40 uas difference in sky position (> spin :10 uas)

Inner S-stars perturbed by a single perturber



The stellar perturbations are dominated by perturbers inside the target star

Results: Perturbations due to a star cluster

Star clusters

- ullet Density profile $n(r) \propto r^{-\gamma}$
 - Bahcall-Wolf Cusp (Bahcall & Wolf 1976) $\gamma = 1.75$
 - Core-like profile (Do et al. 2009) $\gamma=0.5$
- Initial conditions (Merritt et al. 2011)

$$f(a_{\text{orbp}}) \propto a_{\text{orbp}}^{2-\gamma}$$
 $f(e_{\text{orbp}}^2) \propto (1 - e_{\text{orbp}}^2)^{-\beta}$ $\beta \leq \gamma - 1/2$

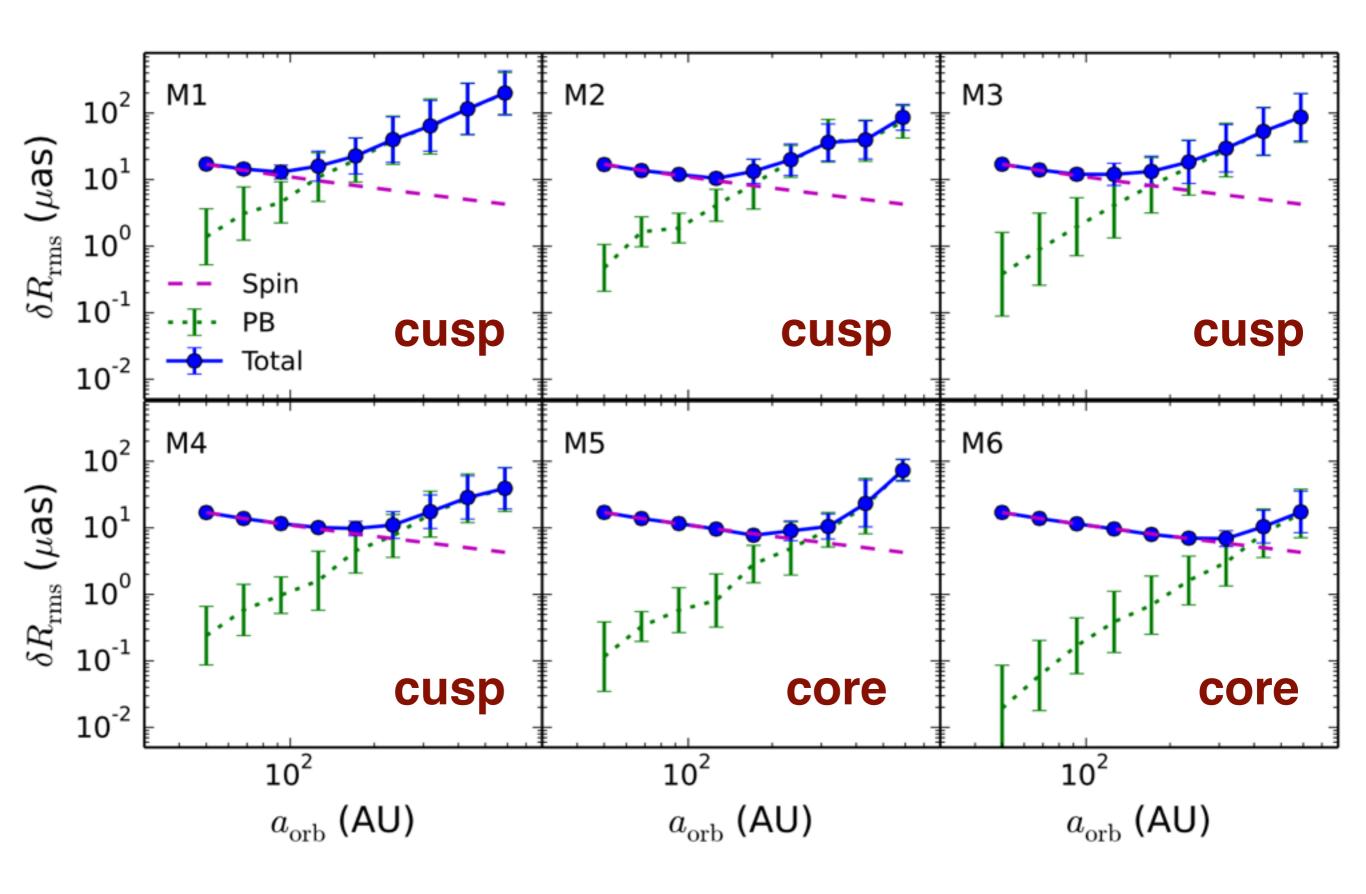
name	$M_{10}{}^{\mathrm{a}}~(M_{\odot})$	$M_1{}^{ m b}~(M_\odot)$	γ	β	$m~(M_{\odot})$	$N_p{}^{ m c}$	$N_{ m cluster}$
M1	1780	100	1.75	0.5	10	178	80
M2	1780	100	1.75	0.5	1	1780	8
M3	530	30	1.75	0.5	10	53	280
M4	530	30	1.75	0.5	1	530	28
M5	1581	5	0.5	-0.5	1	1581	10
M6	316	1	0.5	-0.5	1	316	48

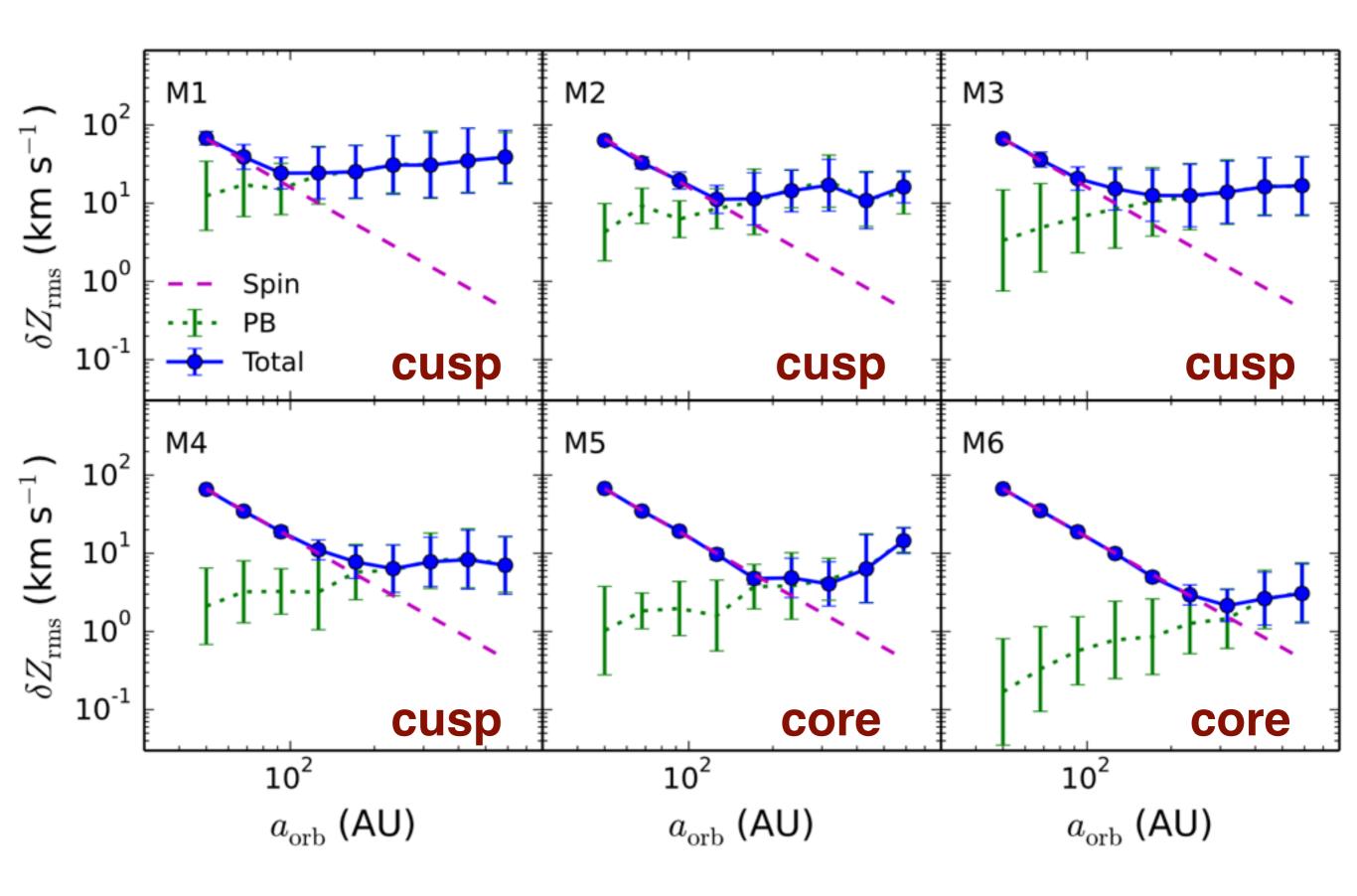
^a The total mass of the stars with $a_{\rm orb} < 10 {\rm mpc} (\sim 2062 {\rm ~AU~or} \sim 0.26 {\rm ~mas})$.

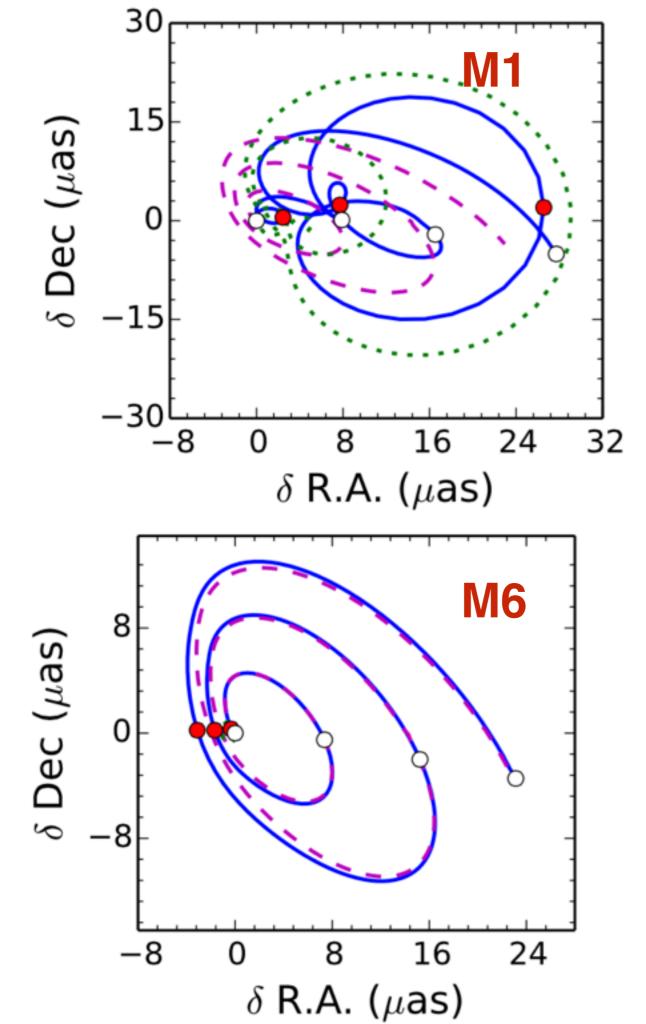
^b The total mass of the stars with $a_{\rm orb} < 1 \rm mpc$ ($\sim 206 \rm \ AU \ or \sim 0.026 \ mas$).

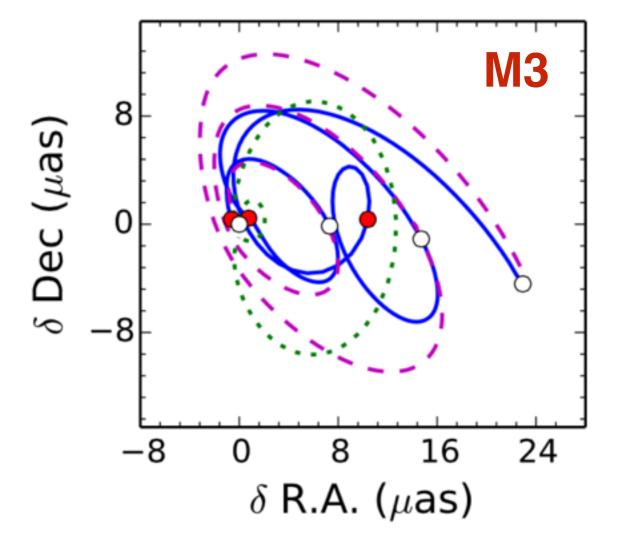
^c The total number of the stars with $a_{\rm orb} < 10 \rm mpc$.

Stellar perturbation due to a star cluster



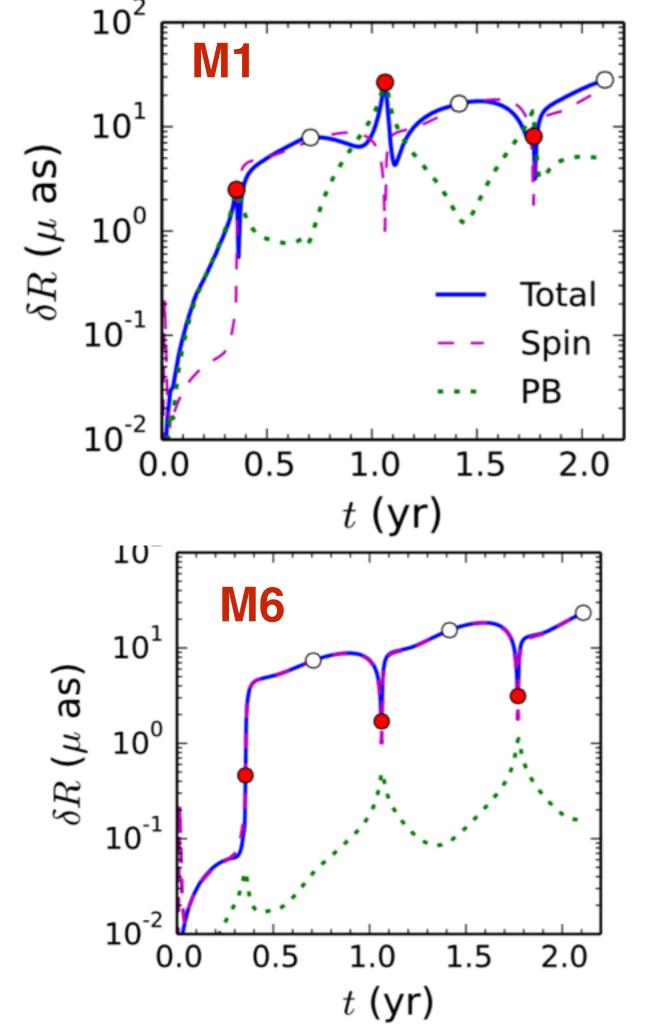


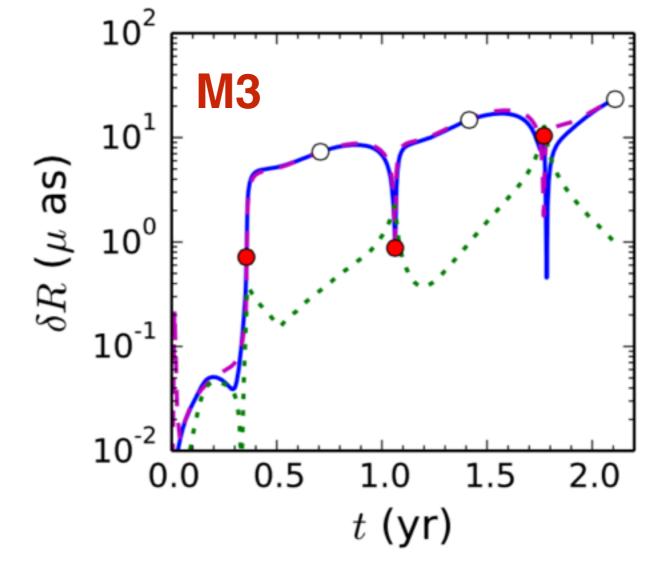




Position difference in sky position

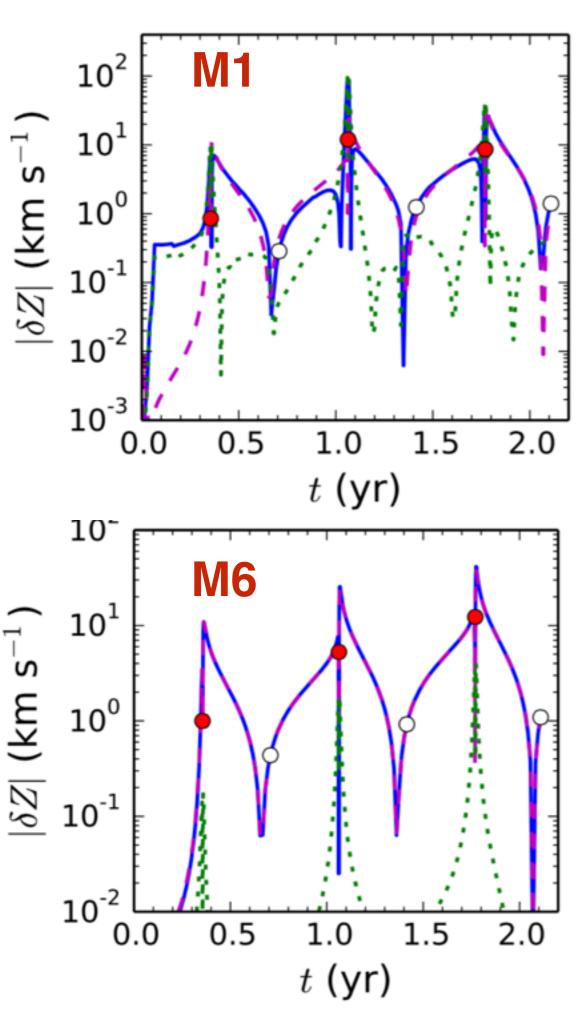
Target star a_{orb}=126AU

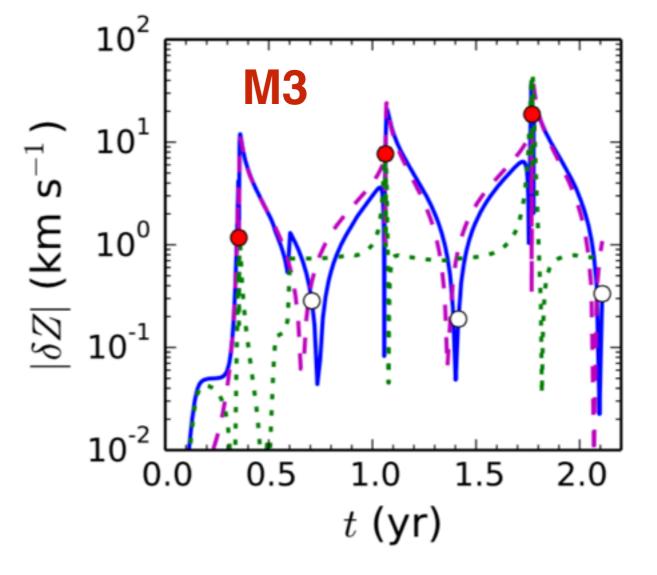




Position difference in sky plane

Target star a_{orb}=126AU





Redshift difference

Target star a_{orb}=126AU

Summary and discussion

- The spin-induced effects of S2/S0-2 are very likely obscured by the stellar perturbations from the S0-102.
- The stellar perturbations are dominated by perturbers inside the target star
- The stellar perturbations peaks around pericenter.
- Perturbed orbital period of stars
- The spin-induced effects dominates the signal for target stars inside 100-200AU if a clusters of stars exists around the MBH. But in principle the stellar perturbations are separable

Thank you!~~