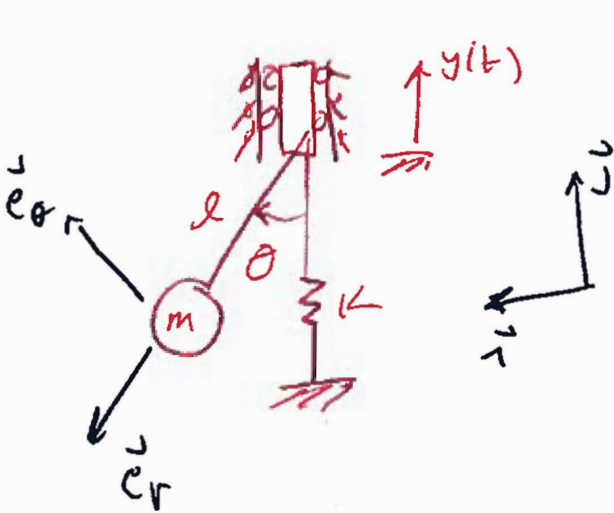


Problem 1

- Find EoMs using Lagrange's Eqs.



$$q_1 = \theta, \dot{q}_1 = \dot{\theta}$$

$$q_2 = y, \dot{q}_2 = \dot{y}$$

$$\hat{i} = \cos\theta \hat{e}_\theta + \sin\theta \hat{e}_r ; \quad \hat{j} = -\cos\theta \hat{e}_r + \sin\theta \hat{e}_\theta$$

* velocity of mass (additive)

$$\vec{v} = l\dot{\theta} \hat{e}_\theta + \dot{y} \hat{j} = (l\dot{\theta} + \dot{y} \sin\theta) \hat{e}_\theta - \dot{y} \cos\theta \hat{e}_r$$

$$\begin{aligned} \vec{v} \cdot \vec{v} &= (l\dot{\theta} + \dot{y} \sin\theta)^2 + (\dot{y} \cos\theta)^2 \\ &= \dot{y}^2 + (l\dot{\theta})^2 + 2l\dot{y}\dot{\theta} \sin\theta \end{aligned}$$

KE

$$T = \frac{1}{2} m [\dot{y}^2 + (l\dot{\theta})^2 + 2l\dot{y}\dot{\theta} \sin\theta]$$

PE

$$V = mgl(1 - \cos\theta) + mgy + \frac{1}{2} Ky^2$$

1st EQN.

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$$

0, NO NONCONSERVATIVE WORK

$$L = \frac{1}{2} m [\dot{y}^2 + (l\dot{\theta})^2 + 2l\dot{y}\dot{\theta}\sin\theta] - mgl(1 - \cos\theta) - mgy - \frac{1}{2} ky^2$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) = \frac{d}{dt} (m l^2 \ddot{\theta} + m l \dot{y} \sin\theta)$$

$$= m l^2 \ddot{\theta} + m l \ddot{y} \sin\theta + m l \dot{y} \dot{\theta} \cos\theta$$

$$\frac{\partial L}{\partial \theta} = m l \dot{y} \dot{\theta} \cos\theta - mgl \sin\theta$$

Thus,

$$m l^2 \ddot{\theta} + m l \ddot{y} \sin\theta + mgl \sin\theta = 0$$

or

$$\ddot{\theta} + \frac{\ddot{y} \sin\theta}{l} + \frac{g}{l} \sin\theta = 0$$

~~~~~  
"extra"  
term  
not seen in  
a fixed pendulum

2<sup>nd</sup> EQN.

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{y}} \right) - \frac{\partial L}{\partial y} = \cancel{Q_y} \rightarrow 0, \text{ NO NONCONSERVATIVE WORK}$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{y}} \right) = \frac{d}{dt} (m\dot{y} + ml\dot{\theta} \sin\theta)$$

$$= m\ddot{y} + ml\ddot{\theta} \sin\theta + ml\dot{\theta}^2 \cos\theta$$

$$\frac{\partial L}{\partial y} = -mg - ky$$

Thus

$$m\ddot{y} + ml\ddot{\theta} \sin\theta + ml\dot{\theta}^2 \cos\theta + ky + mg = 0$$

or

$$\ddot{y} + l\ddot{\theta} \sin\theta + l\dot{\theta}^2 \cos\theta + \frac{k}{m}y + g = 0$$

extra terms due to pendulum's swing motion

Note: I have measured  $y$  from zero spring stretch. If  $y$  is measured from static equilibrium, the term "+g" would be absent

First term is due to tangential acceleration;  
Second term due to normal acceleration

2.

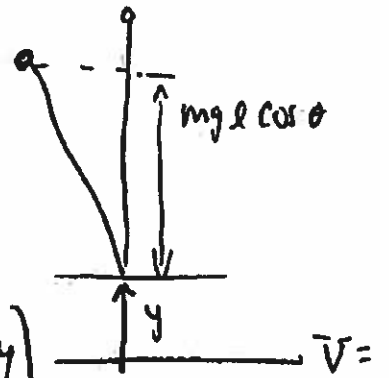
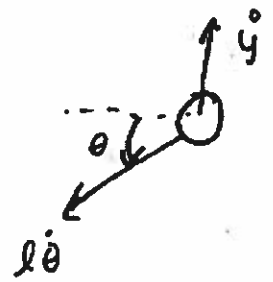
$$v^2 = (-l\dot{\theta} \cos \theta)^2 + (-l\dot{\theta} \sin \theta + \dot{y})^2$$

$$T = \frac{1}{2} m v^2$$

$$V = mgl \cos \theta + mgy$$

$$L = T - V$$

$$= \frac{1}{2} m (l^2 \dot{\theta}^2 - 2\dot{y} l \dot{\theta} \sin \theta + \dot{y}^2) - mgl (\cos \theta + y)$$



$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) = m l^2 \ddot{\theta} - m \dot{y} l \sin \theta - m \dot{y} l \dot{\theta} \cos \theta$$

$$\frac{\partial L}{\partial \theta} = -m \dot{y} l \dot{\theta} \cos \theta + mgl \sin \theta$$

$$\Downarrow \quad \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$$

$$m l^2 \ddot{\theta} - m \dot{y} l \sin \theta - mgl \sin \theta = 0$$

$$\ddot{\theta} + \left( -\frac{g}{l} - \frac{\ddot{y}}{l} \right) \sin \theta = 0$$

$$\boxed{\ddot{\theta} + \left( -\frac{g}{l} + \frac{A \Omega^2 \cos \alpha t}{l} \right) \sin \theta = 0}$$

stable when

$$\frac{A^2 \Omega^2}{2gl} > 1$$

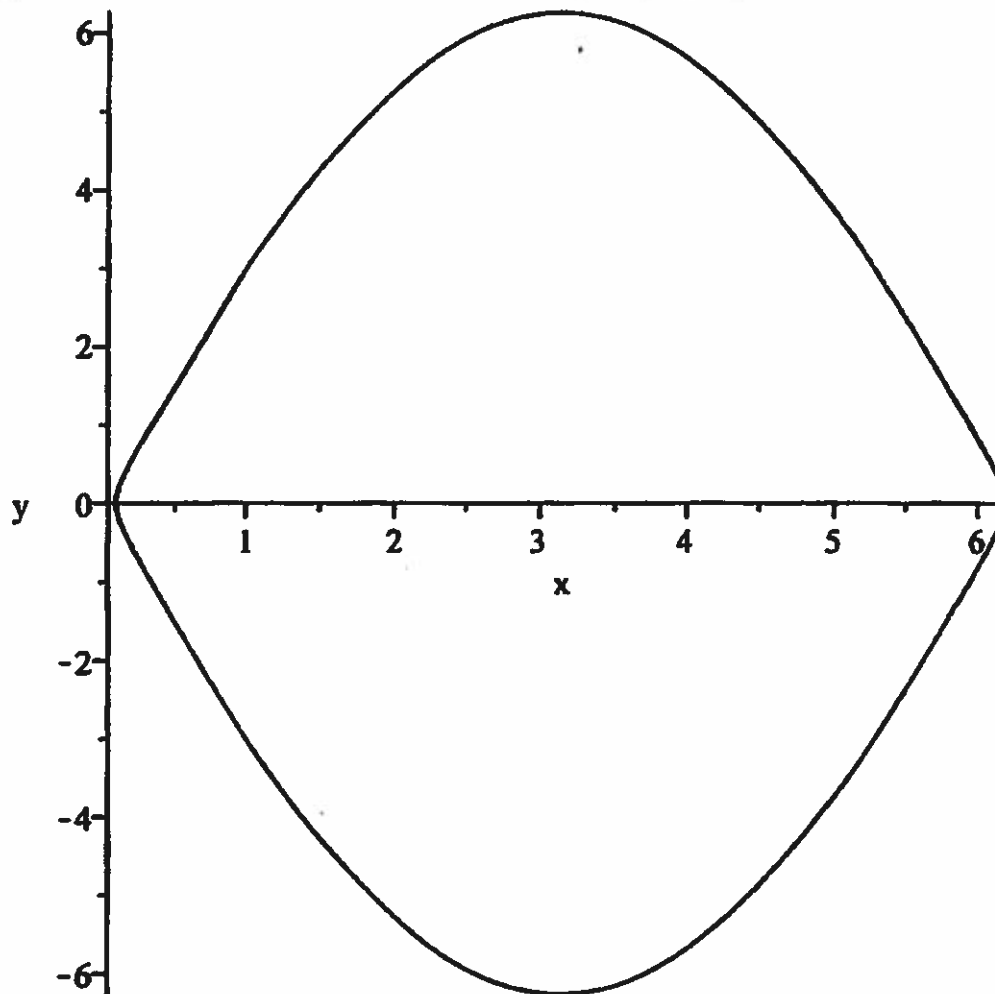
(Known Result)

```
> restart; with(DEtools):
> g:=9.81;l:=1;Omega:=0.01;A:=0.1;
      g:=9.81
      l:=1
      Ω:=0.01
      A:=0.1
```

(1)

Low frequency forcing - inverted pendulum does loop-over-loop

```
> DEplot([diff(x(t),t)=y(t),diff(y(t),t)=(g/l-A*Omega^2*sin(Omega*
t))*sin(x(t))],[x(t),y(t)],t=0..30.0,[[x(0)=0.1,y(0)=0]],
stepsize=0.001,linecolor=black,thickness=1);
```

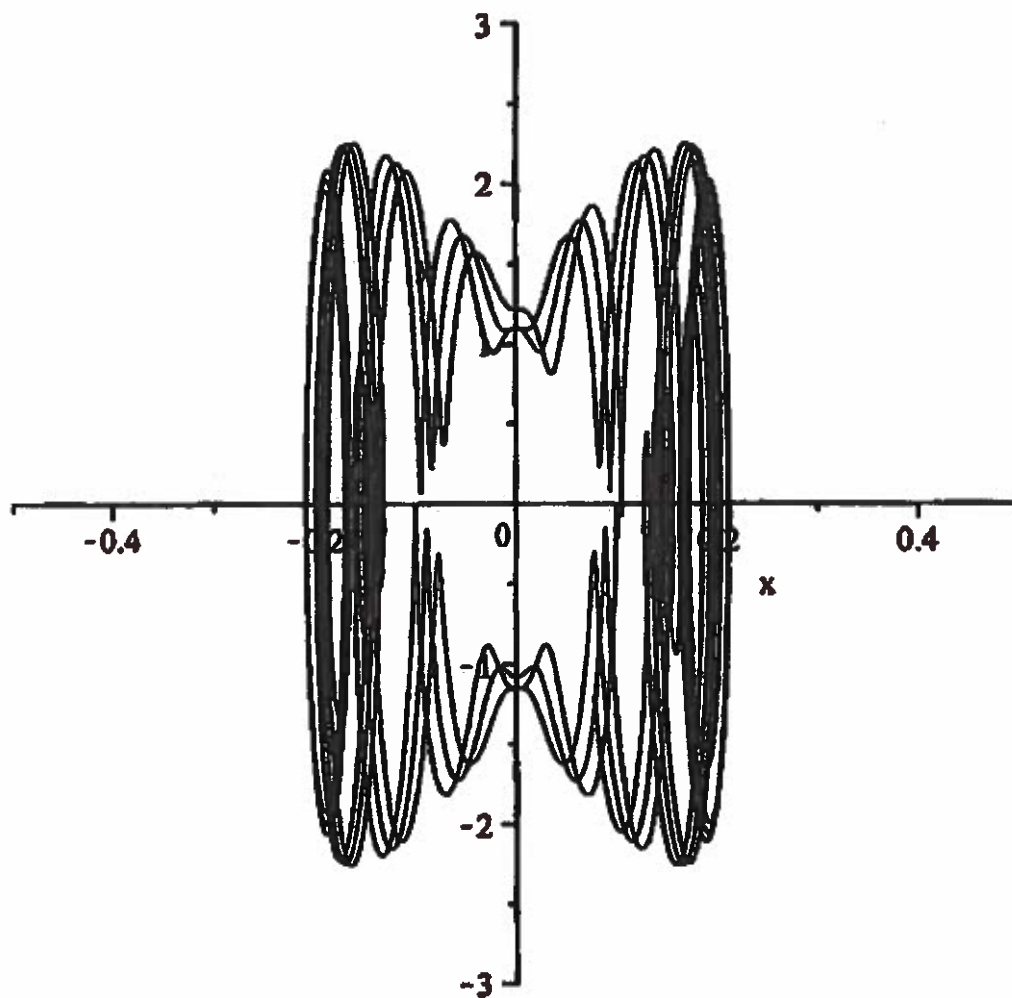


```
> g:=9.81;l:=1;Omega:=100;A:=0.1;
      g:=9.81
      l:=1
      Ω:=100
      A:=0.1
```

(2)

High frequency forcing - inverted pendulum stabilized with maximum angular displacement just over  $\pm 0.2$  radians

```
> DEplot([diff(x(t),t)=y(t),diff(y(t),t)=(g/l-A*Omega^2*sin(Omega*
t))*sin(x(t))],[x(t),y(t)],t=0..3,0,[[x(0)=0.1,y(0)=0]],x=-0.5.
.0.5,y=-3..3,stepsize=0.001,linecolor=black,thickness=1);
```



Problem 3

$$\delta L = \lim_{\epsilon \rightarrow 0} \left[ L(u + \delta u, u' + \delta u', u'' + \delta u'', \dot{u} + \delta \dot{u}) - L(u, u', u'', \dot{u}) \right]$$

where  $\delta u = \epsilon \phi(x, t)$

By a Taylor Expansion about  $\epsilon = 0$ ,

$$L(u + \delta u, u' + \delta u', u'' + \delta u'', \dot{u} + \delta \dot{u}) =$$

$$\underbrace{L(\epsilon=0)}_{L(u, u', u'', \dot{u})} + \left. \frac{\partial L}{\partial \epsilon} \right|_{\epsilon=0} \epsilon + \frac{1}{2} \left. \frac{\partial^2 L}{\partial \epsilon^2} \right|_{\epsilon=0} \epsilon^2 + \dots$$

Note: 
$$\frac{\partial L}{\partial \epsilon} = \frac{\partial L}{\partial (u + \epsilon \phi)} \frac{\partial (u + \epsilon \phi)}{\partial \epsilon} + \frac{\partial L}{\partial (u' + \epsilon \phi')} \frac{\partial (u' + \epsilon \phi')}{\partial \epsilon} + \frac{\partial L}{\partial (u'' + \epsilon \phi'')} \frac{\partial (u'' + \epsilon \phi'')}{\partial \epsilon} + \frac{\partial L}{\partial (\dot{u} + \epsilon \dot{\phi})} \frac{\partial (\dot{u} + \epsilon \dot{\phi})}{\partial \epsilon}$$

$$\left. \frac{\partial \mathcal{L}}{\partial \epsilon} \right|_{\epsilon=0} = \frac{\partial \mathcal{L}}{\partial u} \phi + \frac{\partial \mathcal{L}}{\partial u'} \phi' + \frac{\partial \mathcal{L}}{\partial u''} \phi'' + \frac{\partial \mathcal{L}}{\partial \dot{u}} \dot{\phi}$$

$\frac{\partial^2 \mathcal{L}}{\partial \epsilon^2}$  gets very messy:

$$\begin{aligned} \frac{\partial^2 \mathcal{L}}{\partial \epsilon^2} = & \frac{\partial}{\partial (u + \epsilon \phi)} \left[ \frac{\partial \mathcal{L}}{\partial (u + \epsilon \phi)} \frac{\partial (u + \epsilon \phi)}{\partial \epsilon} + \frac{\partial \mathcal{L}}{\partial (u' + \epsilon \phi')} \frac{\partial (u' + \epsilon \phi')}{\partial \epsilon} \right. \\ & + \frac{\partial \mathcal{L}}{\partial (u'' + \epsilon \phi'')} \frac{\partial (u'' + \epsilon \phi'')}{\partial \epsilon} + \left. \frac{\partial \mathcal{L}}{\partial (\dot{u} + \epsilon \dot{\phi})} \frac{\partial (\dot{u} + \epsilon \dot{\phi})}{\partial \epsilon} \right] \frac{\partial (u + \epsilon \phi)}{\partial \epsilon} \\ & + \frac{\partial}{\partial (u' + \epsilon \phi')} \left[ \text{same as above} \right] \frac{\partial (u' + \epsilon \phi')}{\partial \epsilon} \\ & + \frac{\partial}{\partial (u'' + \epsilon \phi'')} \left[ \text{same as above} \right] \frac{\partial (u'' + \epsilon \phi'')}{\partial \epsilon} \\ & + \frac{\partial}{\partial (\dot{u} + \epsilon \dot{\phi})} \left[ \text{same as above} \right] \frac{\partial (\dot{u} + \epsilon \dot{\phi})}{\partial \epsilon} \end{aligned}$$



$$\begin{aligned}
 \frac{\partial^2 L}{\partial \epsilon^2} \bigg|_{\epsilon=0} = & \left[ \frac{\partial^2 L}{\partial u^2} \phi + \frac{\partial^2 L}{\partial u \partial u'} \phi' + \frac{\partial^2 L}{\partial u \partial u''} \phi'' + \frac{\partial^2 L}{\partial u \partial \dot{u}} \dot{\phi} \right] \phi \\
 & + \left[ \frac{\partial^2 L}{\partial u' \partial u} \phi + \frac{\partial^2 L}{\partial u'^2} \phi' + \frac{\partial^2 L}{\partial u' \partial u''} \phi'' + \frac{\partial^2 L}{\partial u' \partial \dot{u}} \dot{\phi} \right] \phi' \\
 & + \left[ \frac{\partial^2 L}{\partial u'' \partial u} \phi + \frac{\partial^2 L}{\partial u'' \partial u'} \phi' + \frac{\partial^2 L}{\partial u''^2} \phi'' + \frac{\partial^2 L}{\partial u'' \partial \dot{u}} \dot{\phi} \right] \phi'' \\
 & + \left[ \frac{\partial^2 L}{\partial \dot{u} \partial u} \phi + \frac{\partial^2 L}{\partial \dot{u} \partial u'} \phi' + \frac{\partial^2 L}{\partial \dot{u} \partial u''} \phi'' + \frac{\partial^2 L}{\partial \dot{u}^2} \dot{\phi} \right] \dot{\phi}
 \end{aligned}$$

Note: used

$$\frac{\partial}{\partial (u + \epsilon \phi)} \left( \frac{\partial (u + \epsilon \phi)}{\partial \epsilon} \right) = \frac{\partial}{\partial \epsilon} \left( \frac{\partial (u + \epsilon \phi)}{\partial (u + \epsilon \phi)} \right) = 0$$

similarly

$$\begin{aligned}
 \frac{\partial}{\partial (u + \epsilon \phi)} \left( \frac{\partial (u' + \epsilon \phi')}{\partial \epsilon} \right) &= \frac{\partial}{\partial \epsilon} \left( \frac{\partial (u' + \epsilon \phi')}{\partial (u + \epsilon \phi)} \right) \\
 &= \frac{\partial}{\partial \epsilon} \left( \frac{\partial \phi'}{\partial u} \right) = 0
 \end{aligned}$$

Back to  $\delta L$

$$\delta L = \lim_{\epsilon \rightarrow 0} \left[ \frac{\partial L}{\partial u} \overbrace{\epsilon \phi}^{\delta u} + \frac{\partial L}{\partial u'} \overbrace{\epsilon \phi'}^{\delta u'} + \frac{\partial L}{\partial u''} \overbrace{\epsilon \phi''}^{\delta u''} + \frac{\partial L}{\partial \dot{u}} \overbrace{\epsilon \dot{\phi}}^{\delta \dot{u}} \right. \\ \left. + O(\epsilon^2) \right]$$

so, to first order

a.

$$\delta L = \frac{\partial L}{\partial u} \delta u + \frac{\partial L}{\partial u'} \delta u' + \frac{\partial L}{\partial u''} \delta u'' \\ + \frac{\partial L}{\partial \dot{u}} \delta \dot{u}$$

b. At second order

$$\frac{1}{2} \delta^2 L = \frac{1}{2} \left. \frac{\partial^2 L}{\partial \epsilon^2} \right|_{\epsilon=0} \epsilon^2$$

use  $\epsilon^2 \phi^2 \rightarrow \delta u^2$

$$\frac{1}{2} \delta^2 L = \frac{1}{2} \left\{ \left[ \frac{\partial^2 L}{\partial u^2} \delta u^2 + \frac{\partial^2 L}{\partial u \partial u'} \delta u \delta u' + \frac{\partial^2 L}{\partial u \partial u''} \delta u \delta u'' + \frac{\partial^2 L}{\partial u \partial \dot{u}} \delta u \delta \dot{u} + \dots \right] + \left[ \frac{\partial^2 L}{\partial \dot{u}^2} \delta \dot{\phi}^2 \right] \right\}$$

above is clearly  $\frac{1}{2} \delta(\delta L)$