



COLLEGE OF ENGINEERING  
SCHOOL OF AEROSPACE ENGINEERING

AE 6511: OPTIMAL GUIDANCE AND CONTROLS

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## HW3

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## Problem 1

Consider the problem of minimizing

$$\mathcal{J}(x) = \int_0^1 (-a^2 x^2(t) + \dot{x}^2(t)) dt$$

subject to  $x(0) = x(1) = 0$ . Show that if these boundary conditions are satisfied, then all solutions of Euler's equation are of the form

$$\begin{aligned} x(t) &= 0, & \text{if } a \neq n\pi \\ x(t) &= A \sin(n\pi t) \quad \text{or} \quad x(t) = 0, & \text{if } a = n\pi \end{aligned}$$

- (a) Show that  $\mathcal{J}(x(t)) = 0$  for all these solutions.
- (b) Do all the solutions actually minimize  $\mathcal{J}$ ? What does the Legendre condition give?
- (c) Are there some values of  $a^2$  such that  $\mathcal{J}$  can be negative? To answer this, evaluate  $\mathcal{J}(x)$  for a few choices of  $x(t)$ :

$$\begin{aligned} x(t) &= t(1-t) \\ x(t) &= t^m(1-t), \quad m > 0 \\ x(t) &= \sin(\pi t) \end{aligned}$$

### Solution:

- (a) Regardless of  $a \neq n\pi$  or  $a = n\pi$ , when  $x(t) = 0$

$$\mathcal{J}(0) = \int_0^1 (0) dt = 0 \quad \text{if } a \neq n\pi \quad \text{or} \quad a = n\pi.$$

If  $a = n\pi$  and  $x(t) = A \sin(n\pi t)$

$$\begin{aligned} &\mathcal{J}(A \sin(n\pi t)) \\ &= \int_0^1 (-a^2 A^2 \sin^2(n\pi t) + n^2 \pi^2 A^2 \cos^2(n\pi t)) dt \\ &= \int_0^1 \left[ -\frac{1}{2} a^2 A^2 (1 - \cos(2n\pi t)) + \frac{1}{2} n^2 \pi^2 A^2 (1 + \cos(2n\pi t)) \right] dt \\ &= \left[ -\frac{1}{2} a^2 A^2 \left( t - \frac{1}{2n\pi} \sin(2n\pi t) \right) + \frac{1}{2} n^2 \pi^2 A^2 \left( t + \frac{1}{2n\pi} \sin(2n\pi t) \right) \right]_0^1 \\ &= -\frac{1}{2} a^2 A^2 + \frac{1}{2} n^2 \pi^2 A^2 = -\frac{1}{2} n^2 \pi^2 A^2 + \frac{1}{2} n^2 \pi^2 A^2 = 0 \end{aligned}$$

Hence,

$$\mathcal{J}(A \sin(n\pi t)) = 0 \quad \text{if } a = n\pi.$$

(b) Let

$$F(y, r, t) = -a^2 y^2 + r^2$$

then

$$F_{rr} = \frac{\partial}{\partial r} \left( \frac{\partial F}{\partial r} \right) = 2 > 0.$$

Hence, the Legendre conditions is satisfied. Next, if we use the Jacobi equation

$$\begin{aligned} \frac{d}{dt} (F_{yr} \phi + F_{rr} \dot{\phi}) &= F_{yr} \dot{\phi} + F_{yy} \phi \\ \ddot{\phi} + a^2 \phi &= 0 \end{aligned}$$

which gives

$$\phi = A \cos(at) + B \sin(at)$$

and since  $x(0) = x(1) = 0$ ,  $A = 0$  and

$$\phi = B \sin(at).$$

Now, with this sinusoidal function which oscillates we know that

$$\begin{aligned} \phi(0) &= 0 \\ \phi\left(\frac{\pi}{a}\right) &= 0 \quad \text{if } \frac{\pi}{a} < 1 \quad \text{such as } a = n\pi. \end{aligned}$$

Thus, we know that there exists a conjugate point such that  $t_2 = \pi/a < 1 = t_1$ , and therefore, all solutions do not minimize  $\mathcal{J}$ .

(c) We evaluate  $\mathcal{J}$  for the given choices, (calculations are done with MATLAB in Problem 1: MATLAB Code)

First choice:

$$\mathcal{J}(t(1-t)) = \frac{1}{3} - \frac{1}{30}a^2$$

For this choice  $\mathcal{J} < 0$  when

$$a < -\sqrt{10}, \quad \sqrt{10} < a.$$

Second choice:

$$\mathcal{J}(t^m(1-t)) = \begin{cases} -0.0833 a^2 + \infty & \text{if } m = 0.5000 \\ \frac{a^2 - m(2a^2 - 3) + 5m^2 + 2m^3}{(2m-1)(2m+1)(2m+3)(m+1)} & \text{if } 0.5000 < m \\ -\frac{a^2 - \infty}{(2m+1)(2m+3)(m+1)} & \text{if } 0 < m < 0.5000 \end{cases}$$

For this choice (assuming that  $m$  is an integer value and  $0.5000 < m$ )  $\mathcal{J} < 0$  when

$$\begin{aligned} & a^2 - m(2a^2 - 3) + 5m^2 + 2m^3 < 0 \\ & a < -\sqrt{\frac{2m^3 + 5m^2 + 3m}{2m-1}}, \quad \sqrt{\frac{2m^3 + 5m^2 + 3m}{2m-1}} < a. \end{aligned}$$

Third choice:

$$\mathcal{J}(\sin(\pi t)) = \frac{\pi^2}{2} - \frac{a^2}{2}$$

For this choice  $\mathcal{J} < 0$  when

$$a < -\pi, \quad \pi < a.$$

Hence, there exists some values of  $a^2$  such that  $\mathcal{J}$  can be negative.

## Problem 2

Consider a particle sliding along a ramp from point  $(0, 0)$  to the point  $(a, b)$  under the force of gravity with zero initial velocity and **assuming no friction**.

- (a) Show that the trip takes  $t_f$  seconds, where

$$t_f = \sqrt{\frac{2(a^2 + b^2)}{gb}}.$$

- (b) Show that the brachistochrone solution is a cycloid given by

$$\begin{aligned} x &= \alpha + \beta(\psi + \sin(\psi)) \\ y &= \beta(1 + \cos(\psi)) \end{aligned}$$

The curve is parameterized by  $\psi$ , with constants  $\alpha$  and  $\beta$ . If  $\psi_1$  and  $\psi_2$  are the values of the parameter  $\psi$  at the initial and final points, respectively, show that the time to transverse the cycloid is

$$t_f = \sqrt{\frac{\beta}{g}}(\psi_2 - \psi_1).$$

- (c) Show that  $\psi_2 = \theta + \pi$ , where  $\theta$  satisfies

$$(1 - \cos(\theta)) - \frac{b}{a}(\theta - \sin(\theta)) = 0$$

and solve for  $\psi_1$ ,  $\alpha$ , and  $\beta$ .

- (d) How much faster than the ramp is the cycloid? Let  $a = 4$  ft and  $b = 2$  ft and compare the time difference. Where is the particle on the ramp when the particle on the cycloid finishes? Show that this distance is more pronounced for  $a \gg b$ . (Assume that the gravitational acceleration is  $g = 32$  ft/sec<sup>2</sup>.)

### Solution:

- (a) The Bernoulli's Brachistochrone Problem is visually explained in Figure 1. From Pythagorean Theorem we know that

$$\Delta s = \sqrt{\Delta x^2 + \Delta y^2} = \sqrt{1 + \left(\frac{\Delta y}{\Delta x}\right)^2} \Delta x$$

where  $\Delta x \neq 0$ . Then from the conservation of energy law

$$\begin{aligned} \frac{1}{2}mv^2 - mgy &= 0 \\ v &= \sqrt{2gy}. \end{aligned}$$

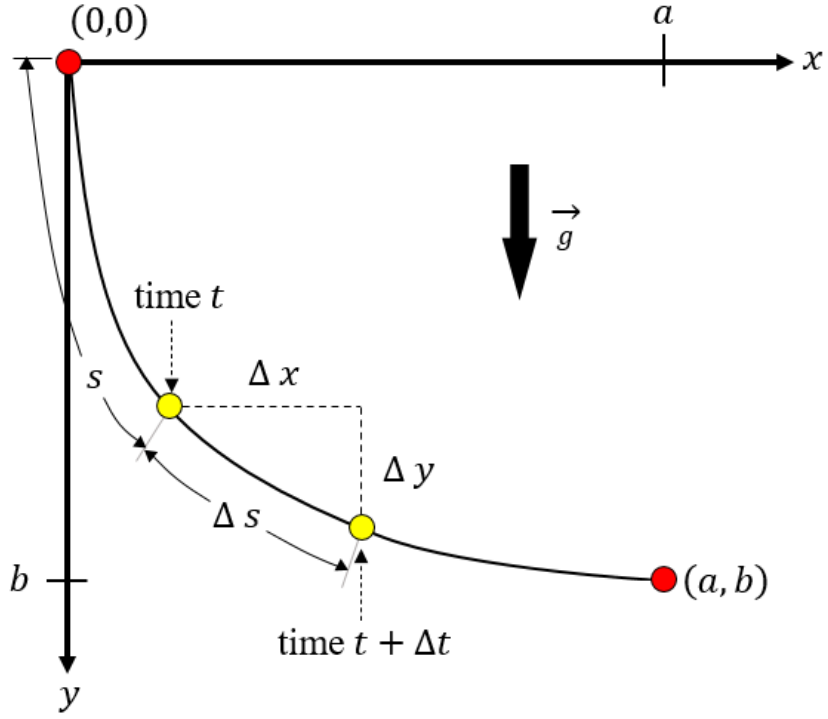


Figure 1: Bernoulli's Brachistochrone problem diagram

Next, we compute the velocity in another method

$$\begin{aligned}
 v &= \frac{\Delta s}{\Delta t} = \frac{\Delta s}{\Delta x} \frac{\Delta x}{\Delta t} \\
 &= \sqrt{1 + \left(\frac{\Delta y}{\Delta x}\right)^2} \frac{\Delta x}{\Delta t}.
 \end{aligned}$$

By equating the two velocity expressions we have

$$\begin{aligned}
 \sqrt{1 + \left(\frac{\Delta y}{\Delta x}\right)^2} \frac{\Delta x}{\Delta t} &= \sqrt{2gy} \\
 \Delta t &= \sqrt{\frac{1 + \left(\frac{\Delta y}{\Delta x}\right)^2}{2gy}} \Delta x
 \end{aligned}$$

Now if we let  $y$  be a function of  $x$ , i.e.  $y = y(x)$ , we can rewrite the above equation into

$$\Delta t = \sqrt{\frac{1 + (y')^2}{2gy}} \Delta x$$

and also let  $\Delta t = dt$  and  $\Delta x = dx$ , which brings us to

$$dt = \sqrt{\frac{1 + (y')^2}{2gy}} dx$$

Therefore,

$$T(x) = \int_0^x \sqrt{\frac{1 + (y')^2}{2gy}} dx.$$

Now, at  $(a, b)$  let  $m = b/a$ . Then,

$$\begin{aligned} t_f = T(a) &= \int_0^a \sqrt{\frac{1 + m^2}{2gmx}} dx = \sqrt{\frac{1 + m^2}{2gm}} \left[ 2\sqrt{x} \right]_0^a \\ &= \sqrt{\frac{2(1 + m^2)a}{gm}} = \sqrt{\frac{2(1 + b^2/a^2)a}{g(b/a)}} \end{aligned}$$

Hence,

$$t_f = \sqrt{\frac{2(a^2 + b^2)}{gb}}.$$

(b) For this question we use Euler-Lagrange. Let

$$F(y, y', x) = \sqrt{\frac{1 + (y')^2}{2gy}}$$

then the minimum  $T$  requires that the Euler-Lagrange equation suffice

$$\frac{d}{dx} \left( \frac{\partial F}{\partial y'} \right) - \frac{\partial F}{\partial y} = 0.$$

Then using Beltrami Identity, Euler-Lagrange equation reduces to

$$y' \frac{\partial F}{\partial y'} - F = C_1 = \text{const.}$$

which becomes

$$\begin{aligned} \frac{y'^2}{\sqrt{y(1 + y'^2)}} - \sqrt{\frac{1 + y'^2}{y}} &= C_1 \frac{1}{y(1 + y'^2)} = C_1 \\ y(1 + y'^2) &= \frac{1}{C_1^2} \\ y' &= \sqrt{\frac{C - y}{y}} \quad \because C = \frac{1}{C_1^2}. \end{aligned}$$



This becomes

$$dx = \sqrt{\frac{y}{C-y}} dy$$

$$x = \int \sqrt{\frac{y}{C-y}} dy$$

and if we use

$$\begin{cases} y = C \sin^2 \phi \\ dy = 2C \sin \phi \cos \phi d\phi \end{cases}$$

we obtain

$$\begin{aligned} x &= \int \sqrt{\frac{C \sin^2 \phi}{C \cos^2 \phi}} 2C \sin \phi \cos \phi d\phi \\ &= 2C \int \sin^2 \phi d\phi \\ &= C \int (1 - \cos 2\phi) d\phi \\ &= C \left( \phi - \frac{1}{2} \sin 2\phi \right) + D \end{aligned}$$

Let  $\theta = 2\phi$

$$x = \frac{C}{2}(\theta - \sin \theta) + D$$

then let  $\psi = \theta + \pi$

$$x = \frac{C}{2}(\psi + \sin \psi) + D - \frac{\pi C}{2}$$

Now if we let

$$\alpha = D - \frac{\pi C}{2}, \quad \beta = \frac{C}{2}$$

we obtain

$$x = \alpha + \beta(\psi + \sin \psi).$$

For y we have

$$\begin{aligned} y &= Cs \sin^2 \phi \\ &= C(1 - \cos^2 \phi) \\ &= C\left(1 - \frac{1}{2}(1 + \cos 2\phi)\right) \\ &= \frac{C}{2}(1 - \cos 2\phi) \\ &= \frac{C}{2}(1 - \cos \theta) \\ &= \frac{C}{2}(1 + \cos \psi) \end{aligned}$$

Hence we have the solution of the brachistochrone problem to be (answer for (b))

$$\begin{aligned}x &= \alpha + \beta(\psi + \sin(\psi)) \\y &= \beta(1 + \cos(\psi))\end{aligned}$$

Next we compute

$$\begin{aligned}\frac{dy}{d\psi} &= -\beta \sin(\psi) \\ \frac{dx}{d\psi} &= \beta(1 + \cos(\psi))\end{aligned}$$

and

$$\begin{aligned}\frac{1 + y'^2}{y} &= \frac{1 + \frac{\sin^2(\psi)}{(1 + \cos(\psi))^2}}{\beta(1 + \cos(\psi))} \\ &= \frac{\beta(1 + 2\cos(\psi) + \cos^2(\psi) + \sin^2(\psi))}{1 + \cos(\psi)} \\ &= 2\beta.\end{aligned}$$

Thus, when  $\psi_1$  and  $\psi_2$  corresponds to points  $x = 0$  and  $x = a$  respectively

$$\begin{aligned}t_f = T(a) &= \int_{\psi_1}^{\psi_2} \frac{2\beta}{2g} dx \\ &= \sqrt{\frac{\beta}{g}} \int_{\psi_1}^{\psi_2} \frac{dx}{d\psi} d\psi \\ &= \sqrt{\frac{\beta}{g}} [x]_{\psi_1}^{\psi_2}\end{aligned}$$

Hence,

$$t_f = \sqrt{\frac{\beta}{g}} (\psi_2 - \psi_1).$$

(c) From the boundary conditions we have

$$\begin{aligned}a &= \alpha + \beta(\psi_2 + \sin(\psi_2)) \\ b &= \beta(1 + \cos(\psi_2))\end{aligned}$$

which gives

$$\frac{b}{a} = \frac{\beta(1 + \cos(\psi_2))}{\alpha + \beta(\psi_2 + \sin(\psi_2))}.$$

We know from what we are given that

$$\frac{b}{a} = \frac{A(1 - \cos \theta)}{A(\theta - \sin \theta)}$$

where this factor of  $A \in \mathbb{R}$  is a common divider making the fraction reducible. Comparing these two we know that

$$\begin{aligned} 1 - \cos \theta &= \beta(1 + \cos \psi_2) \\ \theta - \sin \theta &= \alpha + \beta(\psi_2 + \sin \psi_2) \end{aligned}$$

Thus, if we substitute  $\psi_2 = \theta + \pi$ , the right hand side of the equations become

$$\begin{aligned} \beta(1 + \cos(\theta + \pi)) &= \beta(1 - \cos \theta) \\ \alpha + \beta(\theta + \pi + \sin(\theta + \pi)) &= \alpha + \pi\beta + \beta(\theta - \sin \theta) \end{aligned}$$

Therefore, when  $\beta = A$  and  $\alpha = -\pi A$  we have  $\psi_2 = \theta + \pi$ .

$$\alpha = -\pi A, \quad \beta = A, \quad \psi_1 = \pi \quad \text{where } A \in \mathbb{R}.$$

(d) From Problem (c) we know that the time for the cycloid is

$$t_f = \sqrt{\frac{\beta}{g}} \theta$$

then the time difference for the ramp and the cycloid becomes

$$\Delta T = \sqrt{\frac{2(a^2 + b^2)}{gb}} - \sqrt{\frac{\beta}{g}} \theta.$$

Using MATLAB we can compute this time difference as well as show the visualize the problem (refer to the code in ). To numerically compute  $\theta$ , we use the 15th order (for precision) Taylors series expansion and solve the given equation

$$(1 - \cos(\theta)) - \frac{b}{a}(\theta - \sin(\theta)) = 0.$$

Then, since we know that the cycloid is

$$\begin{aligned} x &= A(\theta - \sin \theta) \\ y &= A(1 - \cos \theta) \end{aligned}$$

and by substituting  $x = a$  and  $y = b$  we can find  $A$ . Then we can visualize the locations of the point at  $t_f$ . Thus, for  $a = 4$  and  $b = 2$  we get the time difference of

$$\Delta T = 0.7906 - 0.6202 = 0.1704 \text{ s.}$$

and the positions of each point is as follows.

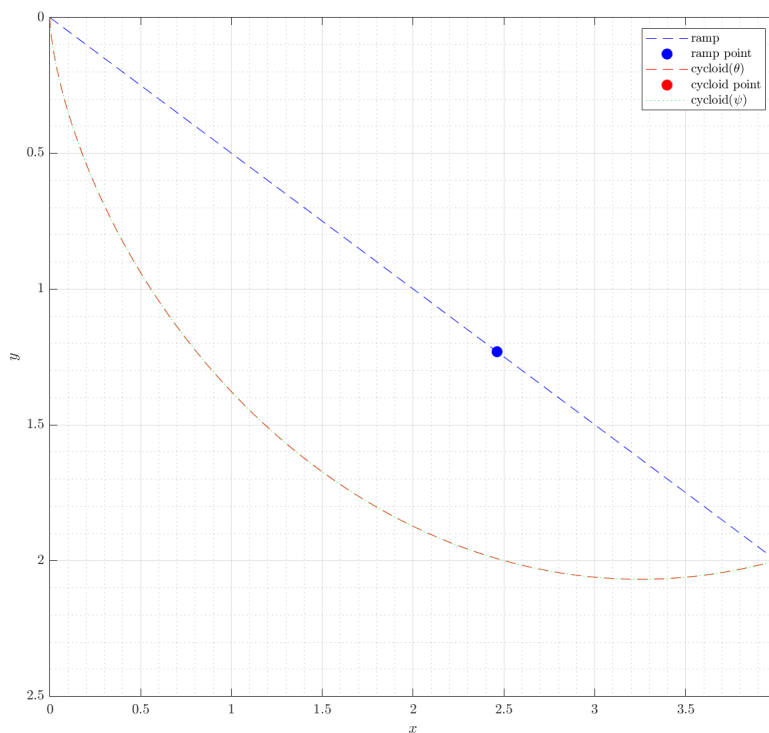


Figure 2: Position of points for ramp and brachistochrone solution ( $a = 4$ ,  $b = 2$ )

Now if we choose an  $a = 10$  and  $b = 2$  that is more  $a \gg b$  we get the following results

$$\Delta T = 0.9906 \text{ s}$$

and the plot looking as follows.

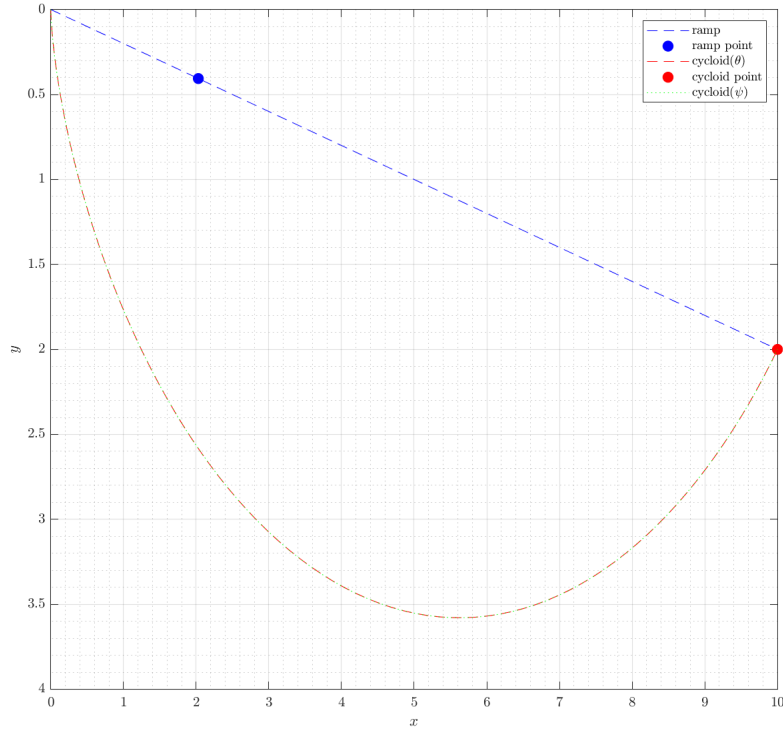


Figure 3: Position of points for ramp and brachistochrone solution ( $a = 10, b = 2$ )

Hence, we can observe that the distance difference between the ramp case and the cycloid case becomes larger for the condition of  $a \gg b$ .

### Problem 3

Analyze the following problem

$$\min \left\{ \int_0^1 (\dot{y}^2(t) + 12ty(t)) dt \right\}$$

subject to  $y(0) = y(1) = 0$ .

**Solution:**

Let

$$F(y, r, t) = r^2 + 12ty$$

then,

$$\begin{aligned} F_r &= 2r, & F_{rr} &= 2 > 0 \\ F_y &= 12t, & F_{yy} &= 0, & F_{yr} &= 0 \end{aligned}$$

From the fact that  $F_{rr} = 2 > 0$  we know that the Legendre condition is satisfied and that the problem is a regular problem. Next, with the Euler Lagrange equation we have

$$\begin{aligned} \frac{d}{dt} \left( \frac{\partial F}{\partial \dot{y}} \right) &= \frac{\partial F}{\partial y} \\ 2\ddot{y} &= 12t \\ \ddot{y} &= 6t \\ \therefore y^* &= t^2 + c_1t + c_2 \end{aligned}$$

and with the boundary conditions of  $y(0) = y(1) = 0$  we have

$$\begin{aligned} c_1 &= -1, & c_2 &= 0 \\ \therefore y^* &= t^3 - t. \end{aligned}$$

From the Weierstrass condition

$$\begin{aligned} E(t, y, r, q) &= q^2 + 12ty - r^2 - 12ty - (q - r)2r \\ &= q^2 - r^2 - (q - r)2r \\ &= (q - r)(q + r - 2r) \\ &= (q - r)^2 > 0. \end{aligned}$$

Therefore, the necessary condition for a strong local minimizer is sufficed. Lastly, we check to see if there are any conjugate points. From the Jacobi equation

$$\begin{aligned} \frac{d}{dt} (F_{yr}\phi + F_{rr}\dot{\phi}) &= F_{yr}\dot{\phi} + F_{yy}\phi \\ \frac{d}{dt} (0 + 2\dot{\phi}) &= 0 \\ 2\ddot{\phi} &= 0 \\ \phi(t) &= at + b \end{aligned}$$

and since  $\phi(0) = 0$

$$\phi(t) = at$$

from this equation there is no repeating values in the range of  $(0, 1]$ , and therefore, there are no conjugate points for this solution. Hence, the minimizer for this problem becomes

$$y^* = t^3 - t.$$

With this solution the minimum value is

$$\min \mathcal{J} = \int_0^1 (3t^2 - 1)^2 + 12t(t^3 - t)dt$$

Hence,

$$\min \mathcal{J} = -\frac{4}{5}.$$

## Problem 4

Find the extremals for the problem

$$\min\{\mathcal{J}\} = \int_{t_0}^{t_1} (3t^2x^2 + 2t^3x\dot{x})dt$$

with boundary conditions  $x(t_0) = x_0$  and  $x(t_1) = x_1$ . Calculate the optimal value of the cost  $\mathcal{J}$ .

---

**Solution:**

Let

$$F(x, r, t) = 3t^2x^2 + 2t^3xr$$

then,

$$\begin{aligned} F_r &= 2t^3x, & F_{rr} &= 0, & F_x &= 6t^2x + 2t^3r \\ F_{xr} &= 2t^3, & F_{xx} &= 6t^2. \end{aligned}$$

From the Euler-Lagrange equation we have

$$\begin{aligned} &\frac{d}{dt} \left( \frac{\partial F}{\partial \dot{x}} \right) - \frac{\partial F}{\partial x} \\ &= \frac{d}{dt} (2t^3x) - 6t^2x - 2t^3\dot{x} \\ &= 0. \end{aligned}$$

From these results we know that this problem  $J(x)$  is *singular*. Hence, from Corollary 4.3 in [1] page 23, we know that for when  $F = M(t, x) + N(t, x)\dot{x}$  the Euler Lagrange equation becomes  $N_t - M_x = 0$  for extremals. Therefore,

$$\begin{aligned} M &= 3t^2x^2, & N &= 2t^3x \\ M_x &= 6t^2x, & N_t &= 6t^2x \\ \therefore M_x &\equiv N_t \end{aligned}$$

and  $\mathcal{J}(x)$  is independent of the path  $t \in [t_0, t_1]$ . Thus, the problem with fixed end points has no relevance. Hence,

$x^* = 0.$

and

$\min \mathcal{J} = \int_{t_0}^{t_1} 0dt = 0.$



## Problem 5

Recall that the conjugate points for the problem

$$\min_{y(x)} \left\{ \int_a^b F(x, y, y') dx \right\}$$

are given by the solution  $\phi(x)$  of the Euler-Lagrange equations of the accessory minimization problem

$$\min_{\phi(x)} \left\{ \int_a^b (F_{yy}\phi^2 + 2F_{yy'}\phi\phi' + F_{y'y'}(\phi')^2) dx \right\}$$

also known as the Jacobi equation.

- (a) Show that the Jacobi equation can be written as follows

$$\left( F_{yy} - \frac{d}{dx} F_{yy'} \right) \phi - \frac{d}{dx} \left( F_{y'y'} \frac{d\phi}{dx} \right) = 0$$

where  $F_{yy}$ ,  $F_{yy'}$ ,  $F_{y'y'}$  are evaluated at the candidate weak local minimizer, say  $y^*(x)$ .

- (b) Show that the ratio  $\phi_1(x)/\phi_2(x)$  is constant for all conjugate points where  $\phi_1(x)$  and  $\phi_2(x)$  are two independent solutions of the Jacobi equation. (**Hint:** since the Jacobi equation is a second-order ordinary differential equation, its solutions are given by  $\phi(x) = c_1\phi_1(x) + c_2\phi_2(x)$  where  $c_1$  and  $c_2$  are some constants. For  $\phi(x) = 0$  for  $x = a$ , then we have  $\phi_1(a)/\phi_2(a) = -c_2/c_1$ .)

### Solution:

- (a) We know that the Jacobi equation is

$$\omega_{\phi'\phi'}\phi'' + \omega_{\phi'\phi}\phi' + \omega_{\phi x} - \omega_{\phi} = 0.$$

With some computation we know that

$$\omega_{\phi} = F_{yy}\phi + F_{yy'}\phi' \quad \text{and} \quad \omega_{\phi'} = F_{yy'}\phi + F_{y'y'}\phi'$$

and plugging this into the Euler-Lagrange equation we have

$$\frac{d}{dx} (F_{yy'}\phi + F_{y'y'}\phi') = F_{yy}\phi + F_{yy'}\phi'.$$

If we manipulate this a little bit, we have

$$\begin{aligned} F_{yy}\phi + F_{yy'}\phi' - \left( \frac{d}{dx} F_{yy'} \right) \phi - F_{yy'}\phi' - \frac{d}{dx} F_{y'y'}\phi' &= 0 \\ F_{yy}\phi - \left( \frac{d}{dx} F_{yy'} \right) \phi - \frac{d}{dx} F_{y'y'}\phi' &= 0 \end{aligned}$$

Now since,

$$\frac{d}{dx} F_{y'y'} \phi' = \frac{d}{dx} \left( F_{y'y'} \frac{d\phi}{dx} \right)$$

the equation becomes,

$$\left( F_{yy} - \frac{d}{dx} F_{yy'} \right) \phi - \frac{d}{dx} \left( F_{y'y'} \frac{d\phi}{dx} \right) = 0.$$

(b) Observing the Jacobi equation we know that it is a second-order ordinary differential equation of  $\phi$ , and for such ODE with two independent solutions we can presume the solution to be in the shape of

$$\phi(x) = C_1 \phi_1(x) + C_2 \phi_2(x)$$

where  $C_1$  and  $C_2$  are constant values. Now if at some point  $x = a$ , if this problem is  $\phi(a) = 0$ , we have a conjugate point  $x = b$  in which  $\phi(b) = 0$  is also satisfied and we know that the following holds true.

$$\begin{aligned} 0 &= C_1 \phi_1(a) + C_2 \phi_2(a) \\ \therefore \frac{\phi_1(a)}{\phi_2(a)} &= -\frac{C_2}{C_1}. \end{aligned}$$

and

$$\begin{aligned} 0 &= C_1 \phi_1(b) + C_2 \phi_2(b) \\ \therefore \frac{\phi_1(b)}{\phi_2(b)} &= -\frac{C_2}{C_1}. \end{aligned}$$

which proves that the ratio  $\phi_1(x)/\phi_2(x)$  is a constant for all conjugate points where  $\phi_1(x)$  and  $\phi_2(x)$  are two independent solutions of the Jacobi equation.

## Problem 6

Consider the problem of minimizing

$$\mathcal{J}(y) = \int_{t_0}^{t_1} (\dot{y}^2(t) - 1)^2 dt$$

- (a) Write down the Euler-Lagrange equations for this problem and show that the extremals for this problem are curves of constant slope (e.g., line segments).
  - (b) Using the Erdmann corner conditions, show that the only extremals with corners are those such that the slope is  $\pm 1$ .
  - (c) Let  $t_0 = 0$  and  $t_1 = 3$ , and assume that  $y(0) = 1$  and  $y(3) = 2$ . Find the *global* minimizer for this case.
  - (d) What about the case when  $t_0 = 0$  and  $t_1 = 1$  and  $y(0) = 0$  and  $y(1) = 2$ ?
- 

**Solution:**

(a) Let

$$F = (\dot{y}^2 - 1)^2$$

then, the Euler Lagrange equation becomes

$$\begin{aligned} \frac{d}{dt} \left( \frac{\partial F}{\partial \dot{y}} \right) &= \frac{\partial F}{\partial y} \\ \frac{d}{dt} (4(\dot{y}^2 - 1)\dot{y}) &= 0 \\ 4\ddot{y}(3\dot{y}^2 - 1) &= 0. \end{aligned}$$

Thus, from

$$\begin{cases} \dot{y}^2 = \frac{1}{3} \\ \ddot{y} = 0 \end{cases}$$

the minimizer for this problem becomes a constant slope

$$\begin{cases} y = \pm \frac{1}{\sqrt{3}}t + C \\ y = C_1 t + C_2 \end{cases}$$

Hence, the extremals for this problem are curves of constant slope.

(b) Using Erdmann corner (weak and strong) conditions, when  $F(y, r, t) = (r^2 - 1)^2$  we have

$$p(p^2 - 1) = q(q^2 - 1)$$

where  $p, q$  are the left and right derivatives. Furthermore, the strong case being

$$(p^2 - 1)(3p^2 - 1) = (q^2 - 1)(3q^2 + 1).$$

Looking at these equations, especially the first one for the weak condition, we can say that we have an infinite amount of trivial solutions for when  $p = q$ . However, for when  $p \neq q$ , we only have  $(p, q) = (1, -1), (-1, 1)$  that satisfy both equations. Thus, **the only extremals with corners are those such that the slope is  $\pm 1$ .**

(c) Applying the boundary conditions to

$$y = C_1 t + C_2 \quad y(0) = 1, \quad y(3) = 2$$

we have

$$\begin{cases} C_2 = 1 \\ C_1 = \frac{1}{3}(2 - C_2) = \frac{1}{3}. \end{cases}$$

Hence, the global minimizer for this is

$$y^* = \frac{1}{3}t + 1.$$

where

$$\min \mathcal{J} = \frac{64}{27} \approx 2.370.$$

(d) Applying the boundary conditions to

$$y = C_1 t + C_2 \quad y(0) = 0, \quad y(1) = 2$$

we have

$$\begin{cases} C_2 = 0 \\ C_1 = (2 - C_2) = 2. \end{cases}$$

Hence, the global minimizer for this is

$$y^* = 2t.$$

where

$$\min \mathcal{J} = 9.$$

## References

- [1] M. Bendersky. *The Calculus of Variations*. 2008.

## Appendix

### 8.1 Problem 1: MATLAB Code

```
1 % AE6511 Hw3 Problem 1 MATLAB code
2 % Tomoki Koike
3 clear all; close all; clc; % housekeeping commands
4 %%
5 syms a x(t)
6
7 F = -a^2 * x^2 + diff(x, t)^2
8 J = int(F, t, 0, 1)
9 %%
10 syms m
11 assume(m, {'real', 'positive'})
12 subs(J, x, t*(1-t))
13 subs(J, x, t^m * (1-t))
14 subs(J, x, sin(pi*t))
```

### 8.2 Problem 2: MATLAB Code

```
1 % AE6511 Hw3 Problem 2 MATLAB code
2 % Tomoki Koike
3 clear all; close all; clc; % housekeeping commands
4 %%
5 syms psi_1 psi_2 theta
6 assume(0 <= psi_1 & psi_1 < 2*pi);
7 assume(0 <= psi_2 & psi_2 < 2*pi);
8 assume(0 <= theta & theta < 2*pi);
9 a = 4;
10 b = 2;
11 m = b/a;
12
13 alpha = -pi;
14 beta = 1;
15 g = 32;
16
17 % 10th order Taylor series of sine and cosine to make equation solvable
18 tsin = @(x) taylor(sin(x), 'Order', 25);
19 tcos = @(x) taylor(cos(x), 'Order', 25);
20 eqn0 = 1 - tcos(theta) - m * (theta - tsin(theta)) == 0;
21 %%
```

```

22 % Brachistochrone solution
23 sol = double(solve([eqn0], [theta]))
24 theta_sol = sol(sol > 0)
25 tf_brachi = sqrt(beta/g) * theta_sol
26
27 % Ramp solution
28 tf_ramp = sqrt(2 * (a^2 + b^2)/ g / b)
29
30 % Delta T
31 DT = tf_ramp - tf_brachi
32 %%
33 % Get the exact cycloid
34 syms A
35 x_cyc = @(x) A*(x - sin(x));
36 y_cyc = @(x) A*(1 - cos(x));
37
38 A1 = vpasolve(subs(x_cyc(theta), theta, theta_sol)==a, A)
39 A2 = vpasolve(subs(y_cyc(theta), theta, theta_sol)==b, A)
40 assert(round(A1 - A2, 3) == 0, 'Precision is not acceptable for A not
    matching.')
41 A_sol = A1;
42 %%
43 % Get the location of point on ramp
44 syms X
45 x_ramp_term = solve(sqrt(2*X*(1 + m^2)/g/m) == tf_brachi, X)
46 %%
47 % Plot
48
49 % Cycloid
50 th = linspace(0, theta_sol, 100);
51 xc = A_sol * (th - sin(th));
52 yc = A_sol * (1 - cos(th));
53
54 % second cycloid for verification
55 psi = linspace(pi, pi + theta_sol, 100);
56 % beta = A_sol and alpha = - pi*A_sol
57 xc2 = -pi*A_sol + A_sol*(psi + sin(psi));
58 yc2 = A_sol * (1 + cos(psi));
59
60
61 % Ramp
62 x_ramp = linspace(0, a, 100);
63 y_ramp = x_ramp * m;
64
65 set(groot, 'defaulttextinterpreter','latex');

```

```

66 set(groot, 'defaultAxesTickLabelInterpreter','latex');
67 set(groot, 'defaultLegendInterpreter','latex');
68 fig = figure("Renderer","painters","Position",[60 60 900 800]);
69 plot(x_ramp, y_ramp, '—b', 'DisplayName', 'ramp')
70 hold on; grid on; grid minor; box on;
71 plot(x_ramp_term, x_ramp_term*m, '.b', 'MarkerSize',25, 'DisplayName','
    ramp point')
72 plot(xc, yc, '—r', 'DisplayName', 'cycloid($\theta$)')
73 plot(xc(end), yc(end), '.r', 'MarkerSize',25, 'DisplayName','cycloid
    point')
74 plot(xc2, yc2, ':g', 'DisplayName', 'cycloid($\psi$)')
75 hold off
76 xlabel('$x$')
77 ylabel('$y$')
78 legend
79 set(gca, 'Ydir','reverse');
80 % saveas(fig, 'p2_brachi_2.png')

```