

P1: B BP

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## Solving/Analyzing a spacecraft dynamics problem requires

1. Construction of a kinematical representation of the motion (numerous options available)
2. Using principles of mechanics to formulate equations of motion that govern the quantities appearing in the kinematical representation

→ critical - no use for #3 if step 2 incorrect

3. Extracting useful info from equations

Since every problem in spacecraft dynamics analysis requires the use of various kinematical relationships (some important only since the space age)

→ Here we start!!

## Rotational Kinematics

To investigate spacecraft attitude, we need effective ways to describe attitude / orientation; rigid bodies



Begin with the most basic concept

### Simple Rotation

How to describe orientation of a **rigid body** B relative to some reference orientation?

[ Note: Description is mathematical and for analysis  
NOT necessarily how the attitude is achieved ]

*ANY Change in orientation is mathematically accomplished via a simple rotation*

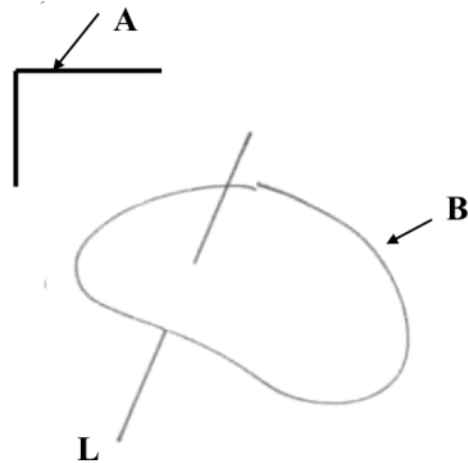
Define:

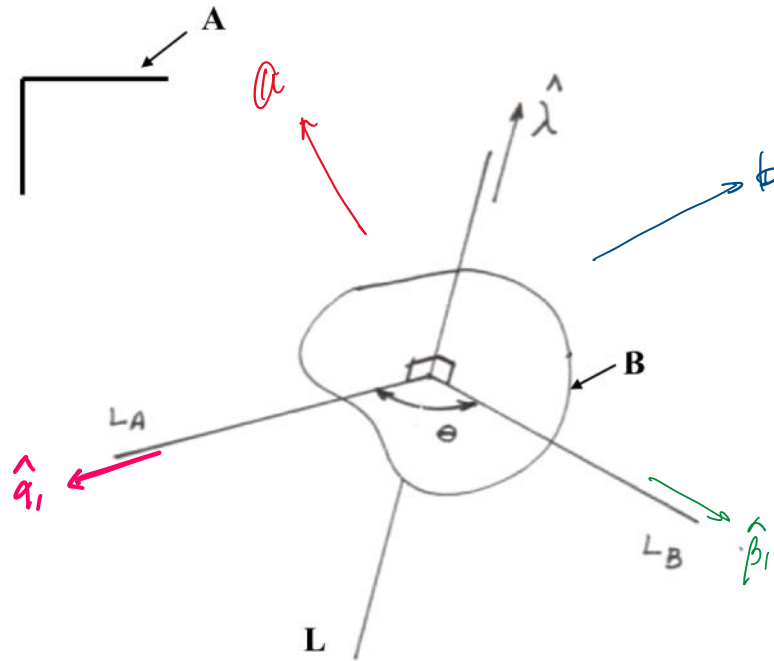
Coordinate frame A, B

[ A may be inertial  
B is body-fixed ]

Simple rotation of B in A:

There exists a line L (axis of rotation)  
whose orientation relative to both  
A and B remains fixed





To visualize the rotation

Define:  $L_A \perp L$

vector basis  $\hat{\alpha}_1, \hat{\alpha}_2, \hat{\lambda}$  fixed in A

Define:  $L_B \perp L$

vector basis  $\hat{\beta}_1, \hat{\beta}_2, \hat{\lambda}$  fixed in B

such that initially

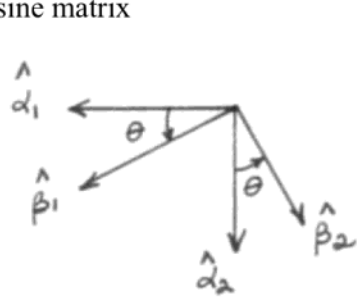
$$\begin{aligned} L_B &= L_A \\ \hat{\beta}_1 &= \hat{\alpha}_1 \\ \hat{\beta}_2 &= \hat{\alpha}_2 \end{aligned}$$

Rotate B about L ( $\hat{\lambda}$ ) through  $\theta$  to achieve the above orientation

Note:  $\hat{\lambda}$  is fixed in both A and B

Orientation of B in A is understood by knowing L ( $\hat{\lambda}$ ) and  $\theta$

Of course, can also represent information in the form of a direction cosine matrix



$$\begin{aligned}\hat{\beta}_1 &= \cos\theta \hat{\alpha}_1 + \sin\theta \hat{\alpha}_2 \\ \hat{\beta}_2 &= -\sin\theta \hat{\alpha}_1 + \cos\theta \hat{\alpha}_2\end{aligned}$$

Note: not symmetric

$$\begin{bmatrix} \hat{\beta}_1 & \hat{\beta}_2 & \hat{\lambda} \end{bmatrix} = \begin{bmatrix} \hat{\alpha}_1 & \hat{\alpha}_2 & \hat{\lambda} \end{bmatrix} \underbrace{\begin{bmatrix} \cos\theta & -\sin\theta & 0 \\ \sin\theta & \cos\theta & 0 \\ 0 & 0 & 1 \end{bmatrix}}$$

Direction Cosine Matrix  $A^B$

Add an arbitrary vector  $\bar{a}$  fixed in A; define  $\bar{b}$  such that  $\bar{b} = \bar{a}$  initially

$$\bar{a} = p \hat{\alpha}_1 + q \hat{\alpha}_2 + r \hat{\lambda}$$

From the way the unit vectors are defined

$$\bar{b} = p \hat{\beta}_1 + q \hat{\beta}_2 + r \hat{\lambda}$$



Describe  $\bar{b}$  as a function of  $\bar{a}$  (its initial value) and the "rotation parameters"  $\hat{\lambda}, \theta$

$$\bar{b} = \bar{a} \cos \theta - \bar{a} \times \hat{\lambda} \sin \theta + \bar{a} \cdot \hat{\lambda} \hat{\lambda} (1 - \cos \theta)$$

$\hat{\lambda} \hat{\lambda}$   
dyad

Proof: LHS (write  $\bar{b}$  in terms of  $\hat{\alpha}_i$ 's)

$$\bar{b} = p \hat{\beta}_1 + q \hat{\beta}_2 + r \hat{\lambda}$$

$$\begin{aligned} &= p (\cos \theta \hat{\alpha}_1 + \sin \theta \hat{\alpha}_2) + q (-\sin \theta \hat{\alpha}_1 + \cos \theta \hat{\alpha}_2) + r \hat{\lambda} \\ &= (p \cos \theta - q \sin \theta) \hat{\alpha}_1 + (p \sin \theta + q \cos \theta) \hat{\alpha}_2 + r \hat{\lambda} \end{aligned}$$

RHS (perform operations)

$$\begin{aligned} \bar{b} &= (p \hat{\alpha}_1 + q \hat{\alpha}_2 + r \hat{\lambda}) \cos \theta \\ &\quad - (p \hat{\alpha}_1 + q \hat{\alpha}_2 + r \hat{\lambda}) \times \hat{\lambda} \sin \theta \\ &\quad + (p \hat{\alpha}_1 + q \hat{\alpha}_2 + r \hat{\lambda}) \cdot \hat{\lambda} \hat{\lambda} (1 - \cos \theta) \end{aligned}$$

$$\bar{b} = (p \cos \theta - q \sin \theta) \hat{\alpha}_1 + (p \sin \theta + q \cos \theta) \hat{\alpha}_2 + r \hat{\lambda}$$

QED

$\bar{b}$  is always going to be a func. of initial cond.  $\bar{a}$  + representation of the rotation  $\hat{\lambda}, \theta$

We know that the theorem is correct but it is not a particularly convenient representation and does not generalize well

➡ More convenient to use a dyadic format

Define:  $\bar{\bar{R}} = \bar{\bar{U}} \cos \theta - \bar{\bar{U}} \times \hat{\lambda} \sin \theta + \hat{\lambda} \hat{\lambda} (1 - \cos \theta)$

simple  
rotation  
dyadic

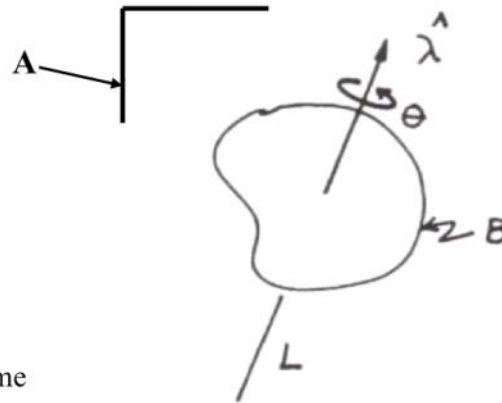
Unit or identity dyadic

Then rewrite theorem:

$$\bar{D} = a \cdot \bar{\bar{R}}$$

|||

Simple rotation is the most basic way to think of a change in orientation and easiest to visualize



But, really need to relate initial/final orientation to some reference

requires unit vectors and a set of variables that are common to all; these are variables we can use to describe orientation; variables that can be used to write EOMs for attitude motion

*diff var. each w/ version of simple rot. (SRT)*

Many variables to choose from; all have advantages / disadvantages



Examples (certainly not all-inclusive)



### 3 rotational DOF

To actually use these variables, 4 quantities must be defined ( $\lambda_i, \theta$ )

#### 1. Euler axis / angle $\hat{\lambda}, \theta$ (4)

Pro: physical meaning

Con:  $\rightarrow$  one redundant variable  
 $\rightarrow$  trig funcs.  $\left. \begin{array}{l} \text{non-linear} \\ \text{double-valued} \end{array} \right\}$

Common application: slow maneuvers

#### 2. Euler angles (3) roll, pitch, yaw (already familiar with these; in class consider 24 different sets)

Pro: no redundant parameters  
 physical interpretation in some cases.

Con:  $\rightarrow$  trig funcs  
 $\rightarrow$  singularities  
 $\rightarrow$  no convenient rule for successive rotations

Common applications: analytical studies  
 input / output  
 attitude control of 3-axis stabilized spacecraft

(Some suggest that Euler angles overemphasized in textbooks)

#### 3. Direction cosines elements $C_{ij}$ (9)

Pro: no singularities  
 no trig func.  
 convenient rule for successive approx.

Con: 6 redundant parameters

Common applications: analysis  
 vector transformations

4. Euler parameters – also called **quaternions** (4)

- Pro: *no singularities*  
*no trig*  
*convenient rule for successive approx.*
- Con: *one redundant parameter*  
*no obvious physical interpretation*

Common applications: **onboard inertial navigation**  
*attitude analysis and control*

5. Rodrigues parameters – also called Gibbs vector (3)

- Pro: *no redundant parameter*  
*no trigs*  
*convenient rule for successive approx.*
- Con: *singularities*

Common application: analytical studies

6. ....

All are ways of describing orientation; each set of variables are related to each other; each set are therefore related to  $\hat{\lambda}, \theta$

We will not actually work with  $\hat{\lambda}, \theta$  as kinematic variables; we will use  $\hat{\lambda}, \theta$  to help visualize problem; use  $\hat{\lambda}, \theta$  to relate variable sets to each other

Examine the new variable sets one at a time → start with direction cosines (already somewhat familiar)

Four important truths

- (1) A minimum of 3 coord are required to describe relative ang. orientation between two frames
- (2) Any minimal set 3 orientation

2) Any minimal set 3 orientation  
include at lease 1 geometrical attitude

3) singularity

4) regularization  $\rightarrow$  at least one redundant parameter