The Worm Algorithm physics760: Computational Physics Final Project

Ajay S. Sakthivasan Dongjin Suh

Universität Bonn

March 15, 2023

- Introduction
- 2 Theoretical Basis
 - The Ising model
 - Physical observables
- Methodology
 - Metropolis-Hastings Algorithm
 - The Worm Algorithm
 - Error Analysis
- 4 Results
 - Algorithm behaviour
 - Net Magnetisation
 - Susceptibility and Heat Capacity
 - Autocorrelation time Dynamical Exponent
- Discussion
- **6** Summary



- Introduction
- 2 Theoretical Basis
 - The Ising model
 - Physical observables
- Methodology
 - Metropolis-Hastings Algorithm
 - The Worm Algorithm
 - Error Analysis
- 4 Results
 - Algorithm behaviour
 - Net Magnetisation
 - Susceptibility and Heat Capacity
 - Autocorrelation time Dynamical Exponent
- Discussion
- **6** Summary



Introduction

- The Metropolis algorithm is a widely used Monte Carlo method for the Ising model.
- However we face the problem of critical slowing down.
- Prokof'ev and Svistunov proposed an alternative update algorithm called the Worm Algorithm (WA).
- WA preserves the local nature of the update step, but achieves a very small dynamical exponent.

- Introduction
- 2 Theoretical Basis
 - The Ising model
 - Physical observables
- Methodology
 - Metropolis-Hastings Algorithm
 - The Worm Algorithm
 - Error Analysis
- 4 Results
 - Algorithm behaviour
 - Net Magnetisation
 - Susceptibility and Heat Capacity
 - Autocorrelation time Dynamical Exponent
- Discussion
- **6** Summary



The Ising model

- A mathematical model to understand the behaviour of systems phase transitions, like ferromagnetic materials
- 2D Ising model: Magnetic system consisting of interacting spins on a two-dimensional lattice

Hamiltonian:

$$H = J \sum_{\langle i,j \rangle} s_i s_j - h \sum_i s_i$$

Partition function:

$$Z = \sum_{s} e^{-\beta H} = \sum_{s} e^{-[-\beta J \sum_{\langle i,j \rangle} s_i s_j - \beta h \sum_i s_i]}$$

Physical observables

Describe the properties of the system change at phase transition

Magnetisation per spin

$$M = \frac{1}{N} \sum_{i}^{N} \sigma_{i} \tag{1}$$

N replaced by L^2 , where L is the lattice length and σ the spin.

• Energy per site

$$E = \frac{1}{N}H\tag{2}$$

Susceptibility

$$\chi = (k_B \beta) \cdot (\langle M^2 \rangle - \langle M \rangle^2) \tag{3}$$

Specific heat

$$C = (k_B \beta)^2 \cdot (\langle E^2 \rangle - \langle E \rangle^2) \tag{4}$$

Autocorrelation time – Dynamical Exponent (z)

$$\tau \approx L^z$$
 for large J and β (5)

- Introduction
- 2 Theoretical Basis
 - The Ising model
 - Physical observables
- Methodology
 - Metropolis-Hastings Algorithm
 - The Worm Algorithm
 - Error Analysis
- 4 Results
 - Algorithm behaviour
 - Net Magnetisation
 - Susceptibility and Heat Capacity
 - Autocorrelation time Dynamical Exponent
- Discussion
- **6** Summary



Metropolis-Hastings Algorithm

- This is a Monte Carlo simulation method.
- Generate samples from a probability distribution.
- Iterative update of the system with the Accept-Reject method.
- Implement Metropolis-Hastings method.
 - **1** Random configuration of an $N \times N$ lattice.
 - Plip the spin at the site.
 - 3 Calculate the energy cost ΔE of the flip.
 - The reject/accept step.
- Repeat these steps for desired number of times and measure the observables at every iteration.

The Worm Algorithm

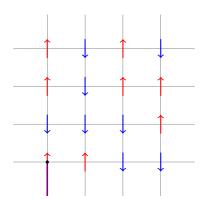
- The worm algorithm is an alternative to the standard Metropolis algorithm.
- The original implementation by Prokof'ev and Svistunov is roughly as follows:
 - Start with an arbitrary lattice configuration (with no starting) bonds between the sites).
 - 2 Select an arbitrary point as $i_1 = i_2$.
 - Grow the worm:
 - **1** When $i_1 \neq i_2$, choose to move i_1 and create or erase bond line between the sites.
 - When $i_1 = i_2$, choose with probability 0.5 to start with a new site or to move i_1 .
 - The acceptance probabilities of the moves are given by solutions to the balance equation.
 - Ollect statistics for various quantities and proceed to moving the worm.
 - Repeat the steps desired number of times.



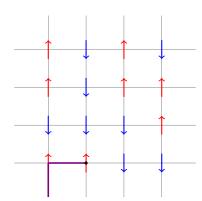
- Our first implementation of the worm algorithm is as follows:
 - Start with an arbitrary lattice configuration.¹
 - Select an arbitrary point to create the worm.
 - Grow the worm:
 - Choose a random direction to move.
 - If the new point is of the same spin as the old point, add it to the worm with probability 1.
 - If not, perform a Metropolis-like check. If a flip is favourable, flip and add
 - This way, we create a worm with equal spins.
 - Break when the worm head meets its tail or if a flip is no longer favourable.
 - Measure the observables.
 - Carry out the desired number of iterations, with a new worm every time.
- There are a few caveats.

The Worm Algorithm

- Given below is a worm at an intermediate step of its growth.
- The lattice below is a small part of the total lattice.

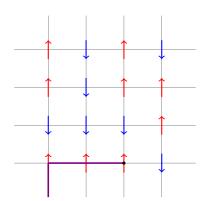


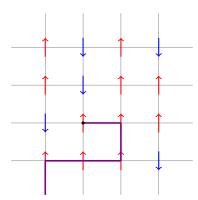
- The worm decides to move right.
- The new site has the same spin as the old site.



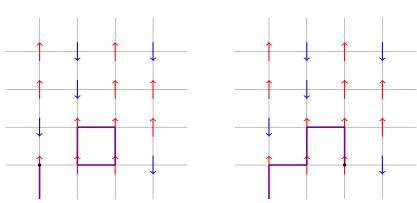
The Worm Algorithm

- The worm decides to move right.
- The new site has the opposite spin as the old site.
- Now we perform a metropolis-like check and decide whether to flip or not.





- Caveat: If the worm tries to move to a point that is already a part of the worm, we need to update the head accordingly.
- We want only an even number of bonds between the sites.
- We ensure this by choosing the new head and breaking an old bond appropriately (with equal probability).



- Caveat: If the worm tries to leave the lattice, let it choose a different direction.
- Caveat: The worm cannot be bigger than the total number of lattice sites.
- The worm dies if it hits its tail or if a new flip is no longer favourable.
- Update the observables and start with a new worm.

- Reason: Say we start with a random site. A flip is always favourable, and a worm keeps growing until breaking conditions are met.
- Attempted Solution: Try worms of alternating spins instead of same spins.
- This did solve the problem and resulted in the expected behaviour (including critical point).
- We also tried a mix of the two variations by introducing a cut-off.

How does this differ?

- We always choose to kill the worm when the head and tail meet.
- We implemented a Metropolis-like check for acceptance probabilities.
- Observables are calculated from the lattice configuration at the end of a worm's life.

- In the original implementation, this was done with a probability 0.5.
- Acceptance probabilities are calculated based on the bond configuration.
- Observables are calculated directly from the bond configuration after every time the worm moves.

Error Analysis

Error Analysis

- Bootstrap method was used for error analysis.
- The idea is roughly as follows we resample from the original sample to create multiple samples.
- For each sample, the statistic of interest is then studied.
- The standard deviation of the bootstrap distribution is an estimate of the standard error.
- This method was used to estimate the errors in the quantities in our results.

- Introduction
- 2 Theoretical Basis
 - The Ising model
 - Physical observables
- Methodology
 - Metropolis-Hastings Algorithm
 - The Worm Algorithm
 - Error Analysis
- 4 Results
 - Algorithm behaviour
 - Net Magnetisation
 - Susceptibility and Heat Capacity
 - Autocorrelation time Dynamical Exponent
- Discussion
- **6** Summary



Algorithm behaviour - Metropolis

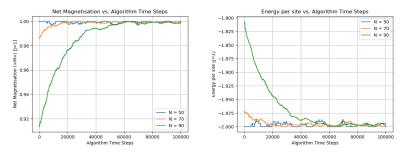


Figure 1: Metropolis behaviour studied with 100000 iterations of the algorithm with 30000 burn-in iterations, J=1 and $\beta=1$

Algorithm behaviour – Worm Algorithm

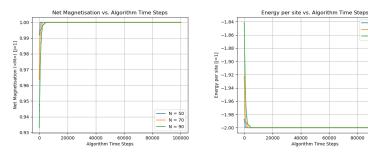


Figure 2: Worm behaviour studied with 100000 iterations of the algorithm with 30000 burn-in iterations, J=1 and $\beta=1$

80000

N = 50N = 70

100000

- ullet Both the cases were studied with the same initial configurations "almost" cold starts closer to spins = +1 and -1 respectively.
- We immediately notice that the worm algorithm equilibrates much faster.
- Lattice size doesn't significantly affect equilibration in the case of WA.

Net Magnetisation

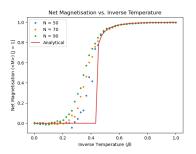


Figure 3: Metropolis algorithm: Behavior of net magnetisation

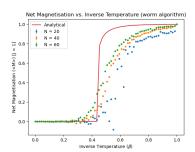


Figure 4: Worm algorithm: Behavior of net magnetisation

- We notice worm algorithm is much smoother around the critical points, even for smaller lattice sizes.
- However, it is a bit far from the analytical solution.
- As discussed, we experimented with two different implementations of worm algorithms – this plot corresponds to alternating spins.
- The other variation produced more exact result for high inverse temperature.

Susceptibility and Heat Capacity

Susceptibility and Heat Capacity - Metropolis

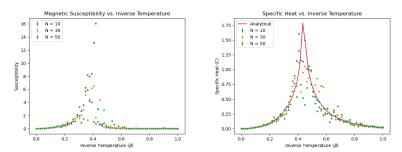


Figure 5: Metropolis Algorithm: Behaviour of Susceptibility and Heat Capacity

Susceptibility and Heat Capacity

Susceptibility and Heat Capacity – Worm Algorithm

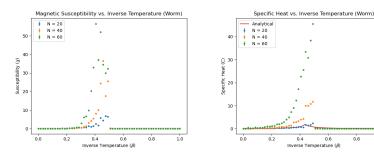


Figure 6: Worm Algorithm: Behaviour of Susceptibility and Heat Capacity

0.8

1.0

- Again, the worm algorithm produced less sporadic points compared to Metropolis.
- However, we faced an issue with normalisation.

Autocorrelation time - Dynamical Critical Exponent

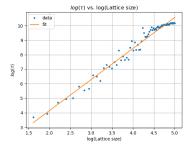


Figure 7: Metropolis Algorithm: Behaviour of Autocorrelation Time

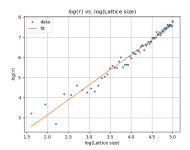


Figure 8: Worm Algorithm: Behaviour of Autocorrelation Time

- Lattice sizes from 5 to 149 were considered for autocorrelation times.
- The autocorrelation times were calculated, and then a ln-ln plot was made.
- The plots here correspond to a dynamical critical exponent of 2.13 and 1.49 for Metropolis and WA, respectively.
- 20 different runs were carried out for autocorrelation time. Dynamical exponent had a range 2.1-2.21 and 1.46-1.55 for Metropolis and WA.
- Statistical analysis was not considered, since 20 runs are not enough.

- Introduction
- 2 Theoretical Basis
 - The Ising model
 - Physical observables
- Methodology
 - Metropolis-Hastings Algorithm
 - The Worm Algorithm
 - Error Analysis
- 4 Results
 - Algorithm behaviour
 - Net Magnetisation
 - Susceptibility and Heat Capacity
 - Autocorrelation time Dynamical Exponent
- Discussion
- **6** Summary



Discussion

- The algorithm is functional to the extent that the results are not unreasonable.
- The dynamical critical exponent was reduced with WA, but still higher than the original work.
- In the original implementation, the dynamical exponent was reported as 0.25.
- We believe this is mainly due to our decision to kill the worm and using a Metropolis-like check.
- However, this hypothesis needs to tested (maybe in the future).
- It should also be noted that lattice sizes up to 512 were considered in the original work.
- Asymptotic result we might get a better result if we simulate larger lattices.



- When it comes to physical observables, the plots showed critical behaviour.
- However, they were not exact compared to analytic results.
- There is also an issue of normalisation, which, for some reason, proved to be very elusive in the case of WA.
- Primary goal for future work, replace with the variation in the original work.
- Measure observables using bond configurations.
- This should further improve our results.
- 3D Ising was not studied due to time constraints.
- Another course of work Implement the algorithm for 3D Ising model.

- Introduction
- 2 Theoretical Basis
 - The Ising model
 - Physical observables
- Methodology
 - Metropolis-Hastings Algorithm
 - The Worm Algorithm
 - Error Analysis
- 4 Results
 - Algorithm behaviour
 - Net Magnetisation
 - Susceptibility and Heat Capacity
 - Autocorrelation time Dynamical Exponent
- Discussion
- **6** Summary



Summary

- To summarise, we studied the worm algorithm in this project.
- This was then used to simulate the 2D Ising model.
- Various observables and the dynamical critical exponent were calculated.
- The corresponding results were then compared with Metropolis algorithm.
- Shortcomings and problems with the implementation were realised, and possible solutions are suggested.
- Things that were left out due to time constraints could be studied in the near future.

Thanks for your time!

References I

- ¹N. Prokof'ev and B. Svistunov, "Worm algorithms for classical statistical models", Physical Review Letters **87**, 10.1103/physrevlett.87.160601 (2001).
- ²N. Prokof'ev, B. Svistunov, and I. Tupitsyn, ""worm" algorithm in quantum monte carlo simulations", Physics Letters A 238, 253–257 (1998).
- ³L. Adzhemyan, D. Evdokimov, M. Hnatič, E. Ivanova, M. Kompaniets, A. Kudlis, and D. Zakharov, "The dynamic critical exponent z for 2d and 3d ising models from five-loop expansion", Physics Letters A 425, 127870 (2022).
- 4 L. P. Kadanoff, "Scaling laws for ising models near T_c ", Physics Physique Fizika 2, 263–272 (1966).
- ⁵W. K. Hastings, "Monte Carlo sampling methods using Markov chains and their applications", Biometrika 57, 97–109 (1970).
- ⁶B. Efron and R. Tibshirani, "Bootstrap Methods for Standard Errors, Confidence Intervals, and Other Measures of Statistical Accuracy", Statistical Science 1, 54–75 (1986).
- ⁷ Github repository physics760/final_project, http://github.com/smilex555/physics760/tree/main/final_project, Accessed: 2023-02-26.

References II

- ⁸B. Nickel, "On the singularity structure of the 2D ising model susceptibility", J. Phys. A Math. Gen. **32**, 3889–3906 (1999).
- ⁹C. Duclut and B. Delamotte, "Frequency regulators for the nonperturbative renormalization group: a general study and the model a as a benchmark", 10.48550/ARXIV.1611.07301 (2016).