

Mixed Hybrid Finite Element Method

Eddington Acceleration of Discrete Ordinates

Source Iteration

ANS Student Conference

Mathematics and Computation

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<https://github.com/smsolivier/EddingtonAcceleration.git>



NUCLEAR ENGINEERING
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1. Motivation
2. Source Iteration Background
3. Eddington Acceleration
4. Results
5. Conclusions

Motivation

Motivation

Radiation Hydrodynamics

- Propagation of thermal radiation through a fluid
- Effects of radiation on fluid momentum and energy
- Required in high energy density laboratory physics (NIF, Z Machine) and astrophysics

Spatial Discretizations

- Mixed Hybrid Finite Element Method (MHFEM) for hydrodynamics
- Linear Discontinuous Galerkin (LDG) for transport

Need hydrodynamics and transport to be consistently differenced

- Use the same method or do extra work to make differing methods agree
- Interpolating between spatial grids introduces noise
- Matching grids between methods is not always possible in higher dimensions

Problem: MHFEM and first-order form of transport are incompatible \Rightarrow can't use linear acceleration scheme

Goal

Develop a non-linear acceleration scheme that

1. Robustly reduces the number of source iterations in Discrete Ordinates calculations
2. Bridges LDG transport and MHFEM multiphysics

Test in 1D slab with lumped LDG transport

Source Iteration Background

Boltzmann Equation

Steady-state, mono-energetic, isotropically-scattering, fixed-source **Linear Boltzmann Equation** in 1D slab geometry:

$$\mu \frac{\partial \psi}{\partial x}(x, \mu) + \Sigma_t(x) \psi(x, \mu) = \frac{\Sigma_s(x)}{2} \int_{-1}^1 \psi(x, \mu') d\mu' + \frac{Q(x)}{2}$$

$\mu = \cos \theta$ the cosine of the angle of flight θ relative to the x -axis

$\Sigma_t(x)$, $\Sigma_s(x)$ total and scattering macroscopic cross sections

$Q(x)$ the isotropic fixed-source

$\psi(x, \mu)$ the angular flux

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Integro-differential equation

Discrete Ordinates (S_N) Angular Discretization

Compute angular flux on N discrete angles

$$\psi(x, \mu) \xrightarrow{S_N} \begin{cases} \psi_1(x), & \mu = \mu_1 \\ \psi_2(x), & \mu = \mu_2 \\ \vdots \\ \psi_N, & \mu = \mu_N \end{cases}$$

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Integrate order $2N - 1$ polynomials exactly with

$$\phi(x) = \int_{-1}^1 \psi(x, \mu) d\mu \xrightarrow{S_N} \sum_{n=1}^N w_n \psi_n(x)$$

S_N Equations

$$\mu_n \frac{d\psi_n}{dx}(x) + \Sigma_t(x)\psi_n(x) = \frac{\Sigma_s(x)}{2}\phi(x) + \frac{Q(x)}{2}, \quad 1 \leq n \leq N$$

$$\phi(x) = \sum_{n=1}^N w_n \psi_n(x)$$

N coupled, ordinary differential equations

All coupling in scattering term

Source Iteration

Decouple by lagging scattering term

$$\mu_n \frac{d\psi_n^{\ell+1}}{dx}(x) + \Sigma_t(x)\psi_n^{\ell+1}(x) = \frac{\Sigma_s(x)}{2}\phi^\ell(x) + \frac{Q(x)}{2}, 1 \leq n \leq N$$

N independent, first-order, ordinary differential equations

Solve each equation with well-known sweeping process

Source Iteration

1. Given previous estimate for $\phi^\ell(x)$, solve for $\psi_n^{\ell+1}$
2. Compute $\phi^{\ell+1}(x) = \sum_{n=1}^N w_n \psi_n^{\ell+1}(x)$
3. Update scattering term with $\phi^{\ell+1}(x)$ and repeat until:

$$\frac{\|\phi^{\ell+1}(x) - \phi^\ell(x)\|}{\|\phi^{\ell+1}(x)\|} < \epsilon$$

Need For Acceleration in Source Iteration

Convergence rate is linked to the number of collisions in a particle's lifetime

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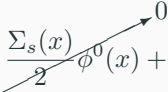
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$$\mu_n \frac{d\psi_n^1}{dx}(x) + \Sigma_t(x)\psi_n^1(x) = \frac{\Sigma_s(x)}{2}\phi^0(x) + \frac{Q(x)}{2}, 1 \leq n \leq N$$

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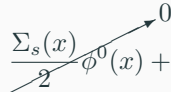
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$\phi^1(x)$ is the uncollided flux

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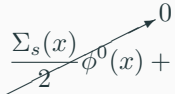
$\phi^1(x)$ is the uncollided flux

$\phi^2(x)$ is uncollided and once collided flux

Need For Acceleration in Source Iteration

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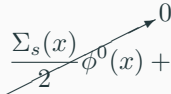
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$\phi^\ell(x)$ is the scalar flux of particles that have undergone at most $\ell - 1$ collisions

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Slow to converge in optically thick systems with minimal losses to absorption and leakage

Need For Acceleration in Source Iteration

Radiation Hydrodynamics problems often contain highly diffusive regions

S_N is too expensive in these regions

Need an **acceleration scheme** that rapidly increases the rate of convergence of source iteration

Eddington Acceleration

Zeroth Angular Moment

Boltzmann Equation

$$\mu \frac{d\psi}{dx}(x, \mu) + \Sigma_t(x)\psi(x, \mu) = \frac{\Sigma_s(x)}{2}\phi(x) + \frac{Q(x)}{2}$$

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Integrate over all angles

$$\int_{-1}^1 \mu \frac{d\psi}{dx}(x, \mu) d\mu + \int_{-1}^1 \Sigma_t(x)\psi(x, \mu) d\mu = \int_{-1}^1 \frac{\Sigma_s(x)}{2}\phi(x) d\mu + \int_{-1}^1 \frac{Q(x)}{2} d\mu$$

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Use $J(x) = \int_{-1}^1 \mu\psi(x, \mu) d\mu$, $\phi(x) = \int_{-1}^1 \psi(x, \mu) d\mu$

Zeroth Angular Moment

$$\frac{d}{dx}J(x) + \Sigma_a(x)\phi(x) = Q(x)$$

First Angular Moment

Multiply by μ and integrate

$$\int_{-1}^1 \mu^2 \frac{d\psi}{dx}(x, \mu) d\mu + \int_{-1}^1 \mu \Sigma_t(x) \psi(x, \mu) d\mu = \int_{-1}^1 \mu \frac{\Sigma_s(x)}{2} \phi(x) d\mu + \int_{-1}^1 \mu \frac{Q(x)}{2} d\mu$$

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Eddington Factor

Rearrange derivative

$$\frac{d}{dx} \int_{-1}^1 \mu^2 \psi(x, \mu) d\mu$$

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$$\frac{d}{dx} \int_{-1}^1 \mu^2 \psi(x, \mu) d\mu$$

Multiply and divide by $\int_{-1}^1 \psi(x, \mu) d\mu$

$$\frac{d}{dx} \int_{-1}^1 \psi(x, \mu) d\mu \frac{\int_{-1}^1 \mu^2 \psi(x, \mu) d\mu}{\int_{-1}^1 \psi(x, \mu) d\mu}$$

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$$\langle \mu^2 \rangle(x) = \frac{\int_{-1}^1 \mu^2 \psi(x, \mu) d\mu}{\int_{-1}^1 \psi(x, \mu) d\mu}$$

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Shape function

Moment Equations

Moment Equations

$$\frac{d}{dx} J(x) + \Sigma_a(x) \phi(x) = Q(x) \quad (\text{Zeroth Moment})$$

$$\frac{d}{dx} \langle \mu^2 \rangle(x) \phi(x) + \Sigma_t(x) J(x) = 0 \quad (\text{First Moment})$$

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Solve First Moment for $J(x)$

$$J(x) = -\frac{1}{\Sigma_t(x)} \frac{d}{dx} \langle \mu^2 \rangle(x) \phi(x)$$

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Transport information passed through $\langle \mu^2 \rangle(x)$

Just as accurate as S_N

Solving the Moment Equations requires knowledge of the angular flux (the solution)

Eddington Acceleration

Use S_N to compute $\langle \mu^2 \rangle(x)$ and Moment Equations to find $\phi(x)$

Eddington Acceleration

1. Given the previous estimate for the scalar flux, $\phi^\ell(x)$, solve for $\psi_n^{\ell+1/2}(x)$
2. Compute $\langle \mu^2 \rangle^{\ell+1/2}(x)$
3. Solve the Moment Equations for $\phi^{\ell+1}(x)$ using $\langle \mu^2 \rangle^{\ell+1/2}(x)$
4. Update the scalar flux estimate with $\phi^{\ell+1}(x)$ and repeat the iteration process until the scalar flux converges

Acceleration occurs because

1. Angular shape of the angular flux converges quickly \Rightarrow Eddington factor quickly converges
2. Moment Equations model all scattering at once \Rightarrow dependence on source iterations to introduce scattering information is reduced

Non-linear scheme \Rightarrow produces 2 solutions (S_N and Moment)

Eddington Acceleration Properties

Non-linear scheme \Rightarrow produces 2 solutions (S_N and Moment)

Relaxes consistent differencing requirements

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Solutions converge as cell spacing decreases

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Relaxes consistent differencing requirements

Solutions converge as cell spacing decreases

Benefits

1. Transport can be LDG and Moment can be MHFEM
2. Moment Equations are conservative and relatively inexpensive compared to transport sweep
3. Difference between S_N and Moment solution can be used as a measure of spatial truncation error (measure of mesh convergence)

Results

Diffusion Limit

Scale cross sections, source

$$\Sigma_t \rightarrow \Sigma_t/\epsilon$$

$$\Sigma_a \rightarrow \epsilon \Sigma_a$$

$$Q \rightarrow \epsilon Q$$

System becomes diffusive as $\epsilon \rightarrow 0$

Diffusion Limit

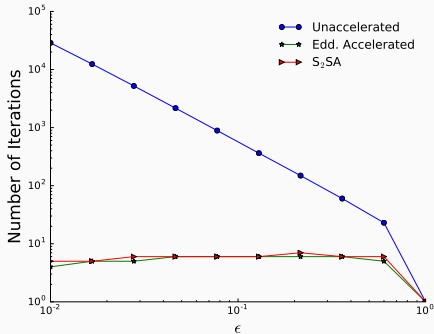
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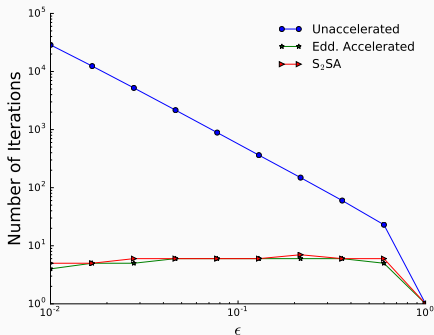
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Accelerates source iteration, survives diffusion limit

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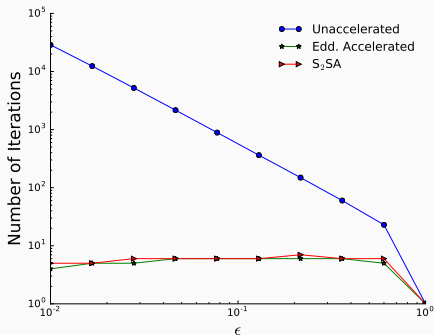
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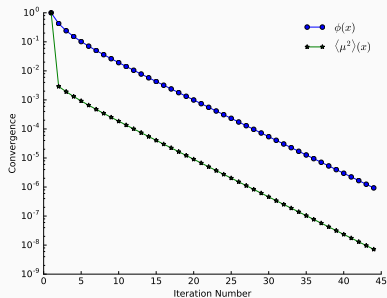


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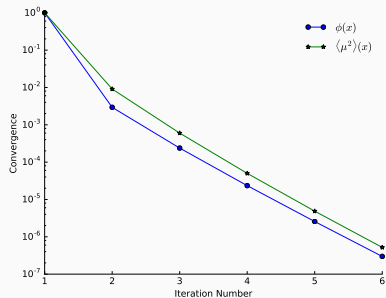
Performs similarly to consistently differenced, linear acceleration (S_2SA)

Convergence Rate Comparison

Unaccelerated

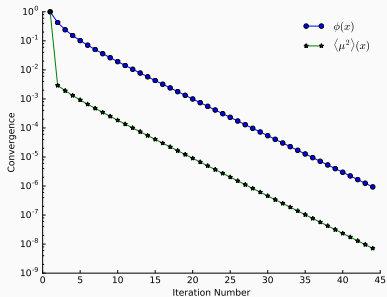


Accelerated

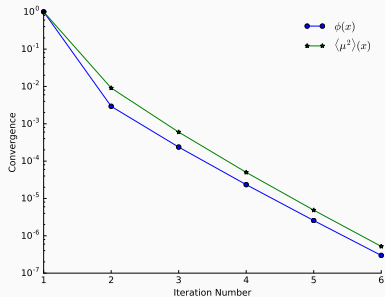


Convergence Rate Comparison

Unaccelerated



Accelerated



Fast rate of convergence of $\langle \mu^2 \rangle(x)$ is transferred to $\phi(x)$

Solution Convergence

Compare

$$\frac{\|\phi_{S_N}(x) - \phi_{\text{Moment}}(x)\|}{\|\phi_{\text{Moment}}(x)\|}$$

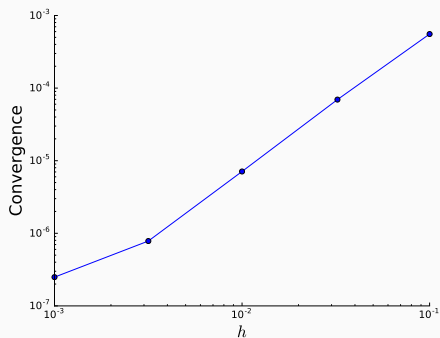
as $h \rightarrow 0$

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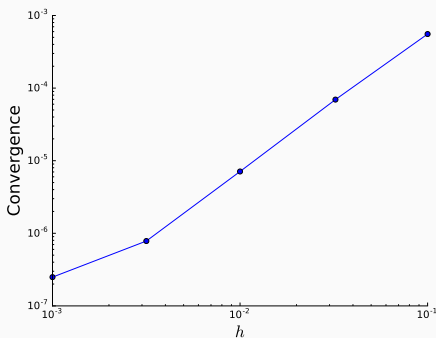


Solution Convergence

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S_N and Moment solutions converge as mesh is refined

Method of Manufactured Solutions Order of Accuracy

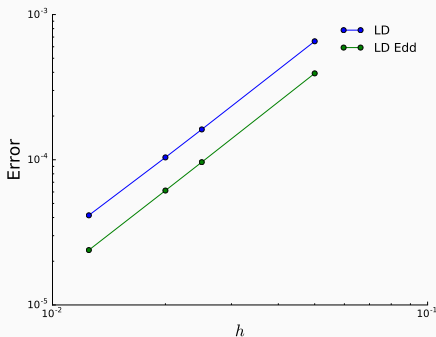
Set source term to force solution to

$$\phi(x) = \sin\left(\frac{\pi x}{x_b}\right)$$

Method of Manufactured Solutions Order of Accuracy

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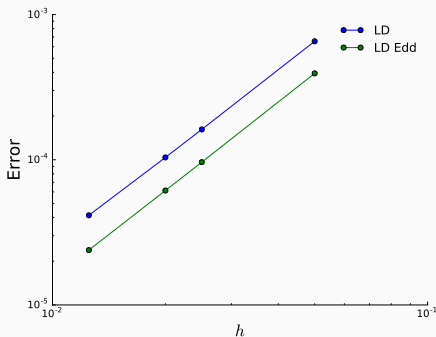
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Set source term to force solution to

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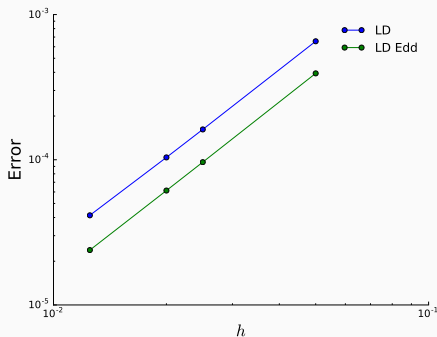


Both second order accurate

Method of Manufactured Solutions Order of Accuracy

Set source term to force solution to

$$\phi(x) = \sin\left(\frac{\pi x}{x_b}\right)$$



Both second order accurate

Eddington Acceleration did not effect the order of accuracy of lumped LDG

Conclusions

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- Scheme successfully accelerated source iteration in 1D slab geometry
- Eddington Acceleration is uniquely suited for radiation hydrodynamics
 - LDG transport
 - MHFEM hydrodynamics
 - Source iteration acceleration
 - Provides inexpensive, conservative solution
- Proved MHFEM can be used to accelerate lumped LDG transport in 1D slab geometry

Summary

Conclusions

- Scheme successfully accelerated source iteration in 1D slab geometry
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Future Work

- Develop a rad-hydro algorithm
 - Make use of inexpensive Moment solution in multiphysics iterations
- Add energy, time dependence
- Test in 3D
- Explore other multiphysics applications

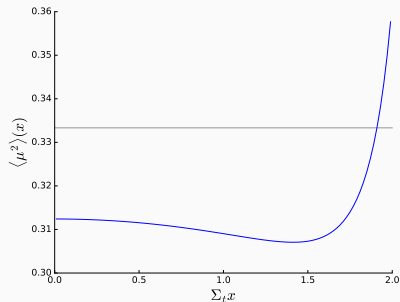
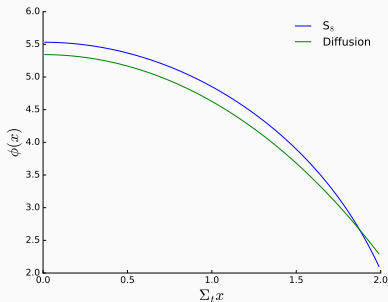
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Questions?

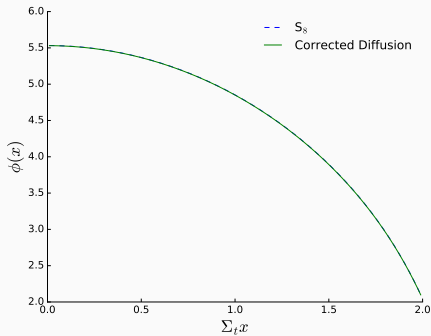
S_8 v. Diffusion

Small system \Rightarrow diffusion not expected to be accurate



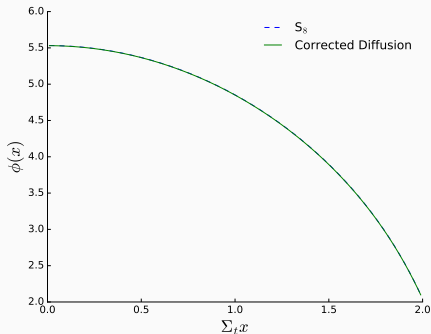
S_8 v. Drift Diffusion

Use $\langle \mu^2 \rangle(x)$ from S_8 in Moment Equations



S_8 v. Drift Diffusion

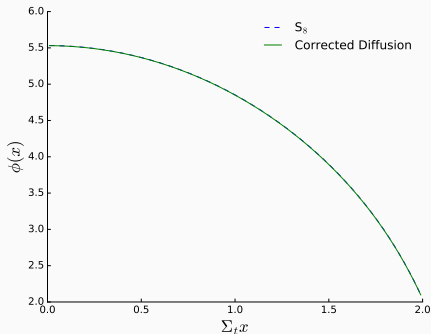
Use $\langle \mu^2 \rangle(x)$ from S_8 in Moment Equations



Moment Equations and S_N match!

S_8 v. Drift Diffusion

Use $\langle \mu^2 \rangle(x)$ from S_8 in Moment Equations



Moment Equations and S_N match!

Requires knowledge of angular flux