

# Homework 13: Due at Dec 25

## Prob. 1 DBI action

DBI action for  $Dp$ -brane is given by

$$S_{\text{DBI}} = \int d^{p+1}\sigma \mathcal{L}_{\text{DBI}} , \quad (1)$$

$$\mathcal{L}_{\text{DBI}} = T_{Dp} - T_{Dp} \sqrt{-\det (G_{ab} + kF_{ab} - B_{ab})} , \quad (2)$$

where we assumed the world-sheet metric  $h_{ab}$  is flat (i.e.  $h_{ab} = \eta_{ab} = \text{diag}(-1, +1, \dots, +1)$ ), the first constant term is added so that the vacuum energy becomes zero,  $k$  is  $2\pi\alpha' = 2\pi l_s^2$ ,  $T_{Dp} = \frac{2\pi}{g_s(2\pi l_s)^{p+1}}$ , and

$$G_{ab} = \partial_a X^\mu \partial_b X^\nu \eta_{\mu\nu} , \quad (3)$$

namely, we assume the space-time is flat. In the following we always assume  $B_{ab} = 0$ .

Static gauge is defined by  $X^i = \sigma^i$  for  $(i = 0, 1, \dots, p)$ , and rewrite  $X^s = \phi^s(\sigma^i)$  for  $(s = p+1, \dots, 9)$ . Displacement field is defined by  $\mathbf{D} = \frac{\partial \mathcal{L}}{\partial \mathbf{E}}$ , where  $\mathbf{E} = F_{i0}$  ( $i = 1, \dots, p$ ).

### Prob. 1.1 Gauge theory on D-brane

Consider D3-brane and take the static gauge.

- Expand the Lagrangian (2) up to a quadratic order of fields and write down a canonical Lagrangian.
- Compare the coupling constant derived from the expansion and the canonical one  $g$ :

$$\mathcal{L}_{\text{gauge}} = -\frac{1}{4g^2} F_{ab} F^{ab} . \quad (4)$$

### Prob. 1.2 Features of the DBI action 1: Schwinger effect

Consider  $Dp$ -brane, take the static gauge, set  $\phi^s = 0$ , and also set the magnetic field to zero;  $\mathbf{B} = 0$ .

- Calculate the displacement field  $\mathbf{D} = \frac{\partial \mathcal{L}_{\text{DBI}}}{\partial \mathbf{E}}$ .
- Assume only  $F_{10} = E$  has non-zero value and plot  $E$  in terms of  $D$  that is a corresponding non-zero component of the displacement field. Explain physically why there are bounds for  $E$ . (This is called **Schwinger effect**.)

### Prob. 1.3 Features of the DBI action 2: Self-energy

Consider the same setting as Prob. 1.2 with  $p = 3$ .

- Calculate the self-energy  $U_Q$  of a charged particle:

$$U_Q = \int d^3\sigma \mathcal{H}(\mathbf{D}) , \quad \mathcal{H}(\mathbf{D}) = \mathbf{E} \cdot \mathbf{D} - \mathcal{L}_{\text{DBI}} , \quad (5)$$

where  $d^3\sigma$  is a spacial(not time) volume form. The charged particle is characterized by

$$\nabla \cdot \mathbf{D} = Q\delta^3(\boldsymbol{\sigma}) , \quad (6)$$

where  $\boldsymbol{\sigma}$  is a world-sheet space(not time) vector. You can use the following constant:

$$c \equiv \int_0^\infty dx \left( \sqrt{1+x^4} - x^2 \right) \simeq 1.236 . \quad (7)$$

- Compare the self-energy with that of a normal electromagnetism:

$$\mathcal{L}_{\text{EM}} = \frac{\varepsilon}{2} \mathbf{E}^2 . \quad (8)$$

Discuss also  $l_s \rightarrow 0$  limit in the self-energy.

### Prob. 1.4 Brane bending

Consider  $Dp$ -brane, take the static gauge, set  $\mathbf{B} = \mathbf{0}$ . Assume only one of  $\phi^s$  has non-zero value, say  $\phi^9 \equiv X$ . Then, the Lagrangian becomes

$$\tilde{\mathcal{L}}_{\text{DBI}} = T_{Dp} - T_{Dp} \sqrt{(1 - \boldsymbol{\mathcal{E}}^2) (1 + (\nabla X)^2) + (\boldsymbol{\mathcal{E}} \cdot \nabla X)^2 - \dot{X}^2} , \quad (9)$$

where  $\boldsymbol{\mathcal{E}} = k\mathbf{E}$ ,  $\dot{X} = \frac{\partial X}{\partial \sigma^0}$ .

- Show  $\tilde{\mathcal{L}}_{\text{DBI}} = \mathcal{L}_{\text{DBI}}$  in the case that only one of components of  $\mathbf{E}$  is non-zero, say  $E_1$ , and  $X$  only depends on  $\sigma^0$ ,  $\sigma^1$ , and  $\sigma^2$ . (This special case implies full identity  $\tilde{\mathcal{L}}_{\text{DBI}} = \mathcal{L}_{\text{DBI}}$  thanks to rotation symmetry.)

Let us look for a static solution, namely, all the fields are independent of  $\sigma^0$ . Then, we can express  $\mathbf{E}$  by  $\mathbf{E} = \nabla A_0$ .

- Derive equations of motion of  $X$  and  $A_0$ . Confirm that both equations are satisfied if

$$k\mathbf{E} = k\nabla A_0 = \pm \nabla X , \quad k\nabla^2 A_0 = \nabla^2 X = 0 . \quad (10)$$

- Show that when (10) holds we have

$$\mathbf{D} = k^2 T_{Dp} \mathbf{E} , \quad (11)$$

and the energy

$$U = \int d^p \sigma \left( \mathbf{E} \cdot \mathbf{D} - \tilde{\mathcal{L}}_{\text{DBI}} \right) = k^2 T_{Dp} \int d^p \sigma \mathbf{E}^2 . \quad (12)$$

Consider a string is ending on the  $Dp$ -brane (stretching to the 9th direction), which is a source of the electric (displacement) field

$$\nabla \cdot \mathbf{D} = \delta^p(\boldsymbol{\sigma}) . \quad (13)$$

- Derive the solution for  $\mathbf{D}$  as well as  $X$  and show that the energy  $U$  is diverging.

Let us regularize the energy by a cut off  $\delta$ :

$$U(\delta) = k^2 T_{Dp} \int_{r>\delta} d^p \sigma \mathbf{E}^2 . \quad (14)$$

- Show that  $U(\delta) = T_{\text{str}} |X(\delta)| = \frac{1}{k} |X(\delta)|$ . Draw the shape of the  $Dp$ -brane, and discuss the meaning of  $U(\delta)$ .

## Prob. 2 Bound states and supersymmetry

In type II theories we have left 16 supercharges  $Q_L$  and right 16 supercharges  $Q_R$ . Existence of string or branes break some of supersymmetry. Conditions are summarized in Table 1.

**Table 1:** Supersymmetry conditions for branes.

IIA branes	Conditions
F1-String(01)	$\Gamma^{01}Q_L = Q_L$ , $\Gamma^{01}Q_R = -Q_R$
NS5-brane(012345)	$\Gamma^{012345}Q_L = Q_L$ , $\Gamma^{012345}Q_R = Q_R$
Dp-brane(p: even)	$\Gamma^{01\dots p}Q_R = Q_L$
IIB branes	Conditions
F1-String(01)	$\Gamma^{01}Q_L = Q_L$ , $\Gamma^{01}Q_R = -Q_R$
NS5-brane(012345)	$\Gamma^{012345}Q_L = Q_L$ , $\Gamma^{012345}Q_R = -Q_R$
Dp-brane(p: odd)	$\Gamma^{01\dots p}Q_R = Q_L$

- How many supersymmetry remain if there exist Dp-brane ? (We call this Dp system and similar for multi branes.)
- How many supersymmetry remain for D0-D2 system ? D0 lies on (0) direction and D2 lies on (012) directions.
- How many supersymmetry remain for F1-D1 system ? Both string lie on (01) directions.
- How many supersymmetry remain for D0-D4 system ? D0 lies on (0) direction and D4 lies on (01234) directions.
- Show that D4-brane carries D0-brane RR-charge when there exist non-trivial instanton configuration ( $n \neq 0$ ):

$$\frac{1}{8\pi^2} \int F_{(2)} \wedge F_{(2)} = n \in \mathbb{Z} . \quad (15)$$

## Prob. 3 Branes ending on a brane in IIB superstring

Start from equations of motions of IIB SUGRA:

$$dG_{\text{odd}} = H_{(3)}G_{\text{odd}} , \quad (16)$$

$$dH_{(3)} = 0 , \quad (17)$$

$$dH_{(7)} = -\frac{1}{2} [(\mathcal{T}G_{\text{odd}})G_{\text{odd}}]_{(8)} , \quad (18)$$

introduce brane currents and construct current conservation equations. From the equations make a table of branes, which can end on some branes. Is the table compatible with S-duality ?