

Homework 8 (Due at class on May 5)

1 Bosonic string on a circle

1.1

The torus partition function of the bosonic string on a circle S^1 of radius R is given by

$$\begin{aligned} Z^{25} &= \text{Tr } q^{L_0-1/24} \bar{q}^{\bar{L}_0-1/24}, \\ &= |\eta(q)|^{-2} \sum_{n,w} q^{\frac{\alpha'}{4} p_R^2} \bar{q}^{\frac{\alpha'}{4} p_L^2}. \end{aligned} \quad (1.1)$$

where n are KK momenta and w are winding numbers. If we include the non-compact space $\mathbb{R}^{1,24}$, we have to multiply the partition function of the non-compact direction

$$Z^{1,24} = \text{const} \times |\eta(q)|^{-46}$$

By expanding out the Dedekind η -functions in $Z^{1,24} Z^{25}$, show that each term means the right hand sides of the mass formula and the level matching condition:

$$\begin{aligned} M^2 &= \frac{n^2}{R^2} + \frac{w^2 R^2}{\alpha'^2} + \frac{2}{\alpha'} (N + \bar{N} - 2) \\ nw &= N - \bar{N}. \end{aligned}$$

1.2

By using the Poisson resummation formula,

$$\sum_{n \in \mathbb{Z}} \exp(-\pi a n^2 + 2\pi i b n) = a^{-1/2} \sum_{m \in \mathbb{Z}} \exp\left[-\frac{\pi(m-b)^2}{a}\right],$$

show that (1.1) is modular-invariant.

1.3

Show that the partition function (1.1) of the theory at the self-dual radius $R = \sqrt{\alpha'}$ can be written as

$$Z^{25} = |\chi_1(q)|^2 + |\chi_2(q)|^2, \quad \text{where } \chi_1 = \frac{1}{\eta} \sum_n q^{n^2} \quad \chi_2 = \frac{1}{\eta} \sum_n q^{(n+1/2)^2}$$

The χ_i are the characters of the $SU(2)$ affine Lie algebra with level $k = 1$. By expanding this expression out find the massless states from above.

1.4

Show that the currents in the bosonic string theory defined by

$$j^\pm(z) = j^1(z) \pm i j^2(z) := e^{\pm 2iX^{25}(z)/\sqrt{\alpha'}} \quad j^3(z) := i \partial X^{25}(z)/\sqrt{\alpha'} ,$$

satisfy the OPEs

$$j^a(z)j^b(0) \sim \frac{\delta^{ab}}{2z^2} + \frac{i\epsilon^{abc}j^c(0)}{z} .$$

From the OPEs, show that the oscillator modes of the currents

$$j^a(z) = \sum_{m \in \mathbb{Z}} \frac{j_m^a}{z^{m+1}} ,$$

satisfy

$$[j_m^a, j_n^b] = \frac{m}{2} \delta_{m+n,0} \delta^{ab} + i\epsilon^{abc} j_{m+n}^c .$$

This infinite-dimensional algebra is called the $SU(2)$ **affine Lie algebra with level $k = 1$** . (Check that the zero modes satisfy the $SU(2)$ Lie algebra.)

2 R-R field strengths and T-duality in Type II

Let

$$\{\Gamma^\mu, \Gamma^\nu\} = 2\eta^{\mu\nu} \quad \mu = 0, \dots, 9$$

be the Clifford algebra of $SO(1,9)$ gamma matrices. The gamma matrices have the following hermiticity property,

$$(\Gamma^\mu)^\dagger = -\Gamma^0 \Gamma^\mu (\Gamma^0)^{-1} .$$

By using the chirality operator $\Gamma_{11} = \Gamma^0 \Gamma^1 \dots \Gamma^9$, chiral spinors are defined by

$$\Gamma_{11} \psi_\pm = \pm \psi_\pm . \quad (2.1)$$

Show that

$$\bar{\psi}_\pm \Gamma_{11} = \mp \bar{\psi}_\pm$$

where $\bar{\psi}_\pm = \psi_\pm^\dagger \Gamma^0$.

We define the R-R field strengths $G^{\mu_1 \dots \mu_{p+2}}$ as spinor bilinears

$$\text{IIA} : \bar{\psi}_-^L \Gamma^{\mu_1 \dots \mu_{p+2}} \psi_+^R , \quad \text{IIB} : \bar{\psi}_+^L \Gamma^{\mu_1 \dots \mu_{p+2}} \psi_+^R , \quad (2.2)$$

where ψ^R (ψ^L) comes from the right (left) movers and

$$\Gamma^{\mu_1 \dots \mu_{p+2}} = \Gamma^{[\mu_1} \dots \Gamma^{\mu_{p+2}]}$$

is the antisymmetric product of $(p+2)$ gamma matrices. Using the chirality (2.1) of the spinors, determine for which values of p the R-R field strengths (2.2) are non-zero.

In the lecture, we learn that T-duality of the 9th direction in Type II theory acts the left-moving fermion mode

$$\psi_n^9 \rightarrow -\psi_n^9, \quad n \in \mathbb{Z}.$$

The action of duality on the spinor fields is of the form

$$\bar{\psi}^L \rightarrow \bar{\psi}^L \beta_9, \quad \psi^R \rightarrow \psi^R \quad (2.3)$$

where $\beta_9 = \Gamma_{11}\Gamma^9$. Show that

$$\{\beta_9, \Gamma^9\} = 0, \quad [\beta_9, \Gamma^\mu] = 0, \quad \text{for } \mu \neq 9.$$

Using the effect of (2.3) on the R-R field strengths (2.2), show that T-duality transforms the R-R field strengths in IIA to those in IIB, and vice versa.