

Basic ODEs

Separable Equation

$$M(x) - N(y)y' = 0 \Rightarrow \int M(x)dx = \int N(y)dy.$$

Integrating Factor

$$y' + P(x)y = Q(x)$$

$$\text{Let } \mu(x) = e^{\int P(x)dx} :$$

$$\mu(x)y' + \mu'(x)y = \mu(x)Q(x)$$

$$y = \frac{\int \mu(x)Q(x)dx}{\mu(x)}$$

Bernoulli Equation

$$y' + P(x)y = Q(x)y^n$$

Let $z = y^{1-n}$, then $z' = (1-n)y^{-n}y'$:

$$y^{-n}y' + P(x)y^{1-n} = Q(x)$$

$$\frac{z'}{1-n} + P(x)z = Q(x)$$

2nd Order Equation

$$ay'' + by' + cy = r(x) \Leftrightarrow ax^2 + bx + c = 0.$$

$$\begin{cases} y = C_1e^{\lambda_1 x} + C_2e^{\lambda_2 x} & \lambda_1 \neq \lambda_2 \in \mathbb{R} \\ y = (C_1 + C_2x)e^{\lambda x} & \lambda_1 = \lambda_2 \in \mathbb{R} \\ y = e^{\alpha x}(C_1 \cos \beta x + C_2 \sin \beta x) & \lambda = \alpha \pm \beta i \in \mathbb{C} \end{cases}$$

To find the particular solution:

- If $r \in \mathbb{P}^n$, $y_p(x) \leftarrow p^n$.
- If $r(x) = g(x)e^{kx}$, $y_p(x) \leftarrow u(x)e^{kx}$.
- If $r(x) = g(x) \cos kx$ or $g(x) \sin kx$, let $z(x) \leftarrow u(x)e^{ikx}$ and take $\operatorname{Re}(z)$ or $\operatorname{Im}(z)$.

Population Models

Malthus Model

$$\frac{dN}{dt} = (B - D)N \Leftrightarrow N(t) = N_0 e^{(B-D)t}.$$

Logistic Model

$$\frac{dN}{dt} = BN - sN^2 \Leftrightarrow N(t) = \frac{B/S}{1 + e^{-Bt}(\frac{B}{N_0 s} - 1)}.$$

B/S is a stable equilibrium point.

System of Linear ODEs

General Linear ODE System

$$\frac{d}{dt} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} \Leftrightarrow \begin{cases} \frac{dx}{dt} = ax + by \\ \frac{dy}{dt} = cx + dy \end{cases}$$

B

$$\text{Eigenvalues: } r = \frac{1}{2} \left[\operatorname{Tr}(B) \pm \sqrt{\operatorname{Tr}(B)^2 - 4\det(B)} \right].$$

$$\text{Solutions: } \mathbf{u}(t) = C_+ e^{r_+ t} \mathbf{u}_+ + C_- e^{r_- t} \mathbf{u}_-$$

Nonhomogeneous System

$$\frac{d}{dt} \begin{bmatrix} x \\ y \end{bmatrix} = B \begin{bmatrix} x \\ y \end{bmatrix} + F \Rightarrow -B^{-1}F.$$

Phase Plane Classification

Both r_1 and r_2 are real:



$r_1 > 0, r_2 < 0$
saddle point



$r_1 > r_2 > 0$
Nodal source



$r_1 < r_2 < 0$
Nodal sink

Both r_1 and r_2 are complex:



$\operatorname{Re}[r] < 0$ $\operatorname{Re}[r] > 0$ $\operatorname{Re}[r] = 0$
Spiral sink Spiral source Centre

System of Non-Linear ODEs

Linearisation

$$\begin{cases} \frac{dx}{dt} = f(x, y) \\ \frac{dy}{dt} = g(x, y) \end{cases} \quad \text{with equilibrium point } (a, b).$$

By Taylor expansion,

$$\begin{cases} \frac{dx}{dt} \approx f_x(a, b)(x - a) + f_y(a, b)(y - b) \\ \frac{dy}{dt} \approx g_x(a, b)(x - a) + g_y(a, b)(y - b) \\ \dot{x} = \underbrace{\begin{bmatrix} f_x(a, b) & f_y(a, b) \end{bmatrix}}_{J(a, b)} \begin{bmatrix} x - a \\ y - b \end{bmatrix} \end{cases}$$

Lotka-Volterra Model

$$\begin{cases} \frac{dL}{dt} = uZL - D_L L & L: \text{lion} \\ \frac{dZ}{dt} = B_z Z - sLZ & Z: \text{zebra} \end{cases}$$

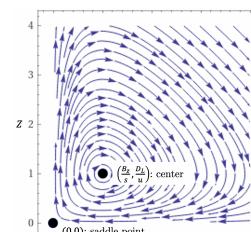
Every trajectory in the contour is periodic:

$$\left(\frac{B_z}{L} - s \right) \frac{dL}{dt} + \left(\frac{D_L}{Z} - u \right) \frac{dZ}{dt} = 0$$

$$B_z \ln L - sL + D_L \ln Z - uZ = C$$

$F(L, Z)$ has only 1 maximum \Rightarrow closed.

$$J(L, Z) = \begin{bmatrix} uZ - D_L & uL \\ -sZ & B_z - sL \end{bmatrix}.$$



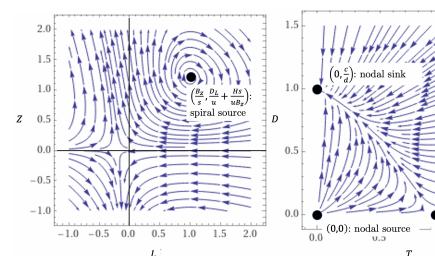
LV

Logistic LV Model

$$\begin{cases} \frac{dL}{dt} = uZL - D_L L - H \\ \frac{dZ}{dt} = (B_z Z - pZ^2) - sLZ \end{cases}$$

Lion Hunting

$$\begin{cases} \frac{dL}{dt} = uZL - D_L L - H \\ \frac{dZ}{dt} = B_z Z - sLZ \end{cases}$$



LV with Hunting

Competing Species

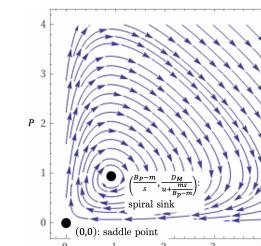
Competing Species

$$\begin{cases} \frac{dT}{dt} = (a - kD)T - bT^2 & T: \text{thylacine} \\ \frac{dD}{dt} = (c - \sigma T)D - dD^2 & D: \text{dingo.} \end{cases}$$

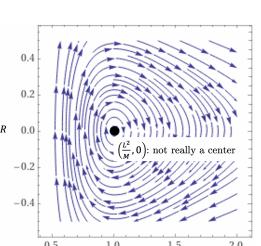
[Principle of Competitive Exclusion]: If two species are too similar, one will wipe out another.

Where Is Everybody?

$$\begin{cases} \frac{dM}{dt} = uPM - D_M M + mP & M: \text{mutant probe} \\ \frac{dP}{dt} = B_P P - sMP - mP & P: \text{normal probe.} \end{cases}$$



Mutant Probes



Planetary Orbits

Non-Linear 2nd Order ODEs

Consider the Earth moving around the Sun:

$$\ddot{r} = -\frac{M}{r^2} + \frac{L^2}{r^3} \Rightarrow \begin{cases} \dot{r} = R \\ \dot{L} = \frac{L^2}{r^2} - \frac{M}{r^3} \end{cases}$$

The equilibrium point is almost a center.

Partial Differential Equations

A PDE is an equation containing an unknown function u of 2 or more independent variables x, y, \dots and its partial derivatives with respect to them.

Separation of Variables

PDE: $u_x = f(x)g(y)u_y$. Suppose $u = X(x)Y(y)$, then $X'(x)Y(y) = f(x)g(y)X(x)Y'(y) = k$.

$$\begin{cases} X'(x) = kf(x)X(x) \\ Y'(y) = \frac{k}{g(y)}Y(y) \end{cases}$$

Wave Equation

$$c \frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 u}{\partial t^2}, \text{ where } \begin{cases} y(t, 0) = 0 \\ y(0, x) = f(x) \\ \frac{\partial y}{\partial t}(0, x) = 0 \end{cases}$$

We need 4 pieces of information for a solution.

d'Alembert's Solution

$$y(t, x) = \frac{1}{2}[f(x+ct) + f(x-ct)].$$

Separation of Variables

Let $y(t, x) = u(x)v(t)$, then $\begin{cases} u'' + \lambda u = 0 \\ v'' + \lambda c^2 v = 0 \end{cases}$

From $\begin{cases} y(t, 0) = u(0)v(t) = 0 \\ y(t, \pi) = u(\pi)v(t) = 0 \end{cases}$,

for u to cut the x -axis twice, we have $\lambda > 0$.

Let $\lambda = n^2$ and $u = C \cos(nx) + D \sin(nx)$.

Since $u(0) = 0$, $C = 0$ and $u = D \sin(nx)$.

Since $u(\pi) = 0$, $n \in \mathbb{Z}$.

Similarly, $v(t) = A \cos(nt)$.

Therefore, $y = b_n \sin(nx) \cos(nt)$ and only $y(0, x) = f(x)$ is not satisfied yet.

Fourier Series

Any odd function $f(x)$ of period 2π on $[0, \pi]$ can be expressed as $f(x) = \sum_{n=1}^{\infty} b_n \sin(nx)$, where $b_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$. So the complete solution is: $y(t, x) = \sum_{n=1}^{\infty} b_n \sin(nx) \cos(nt)$.

If we want $[0, L]$ instead of $[0, \pi]$, the Fourier formulae becomes $g(x) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right)$ and $b_n = \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi x}{L}\right) dx$. The complete solution is: $y(t, x) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right) \cos\left(\frac{n\pi ct}{L}\right)$.

Tsunami (Korteweg-de Vries)

$\partial_t \eta + \sqrt{gh} \partial_x \eta + \frac{3}{2} \sqrt{\frac{g}{h}} \eta \partial_x \eta + \frac{1}{6} h^2 \sqrt{gh} \partial_x^3 \eta = 0$, where η denotes the elevation above sea level.

Suppose $\eta = E(x-ct)$, then we can simplify to

$$-2AE' + 6BEE' + 2CE''' = 0.$$

$$-2AE + 3BE^2 + 2CE'' = 0 \text{ (integrate).}$$

$$-AE^2 + BE^3 + C(E')^2 = K \text{ (integrate w.r.t. E).}$$

Heat Equation

$$u_t = c^2 u_{xx}, \text{ where } \begin{cases} u(0, t) = u(L, t) = 0 \\ u(x, 0) = f(x) \end{cases}$$

$$\text{Solution: } u(x, t) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right) e^{-\frac{\pi^2 n^2 c^2 t}{L^2}}.$$

Heat Equation Variant

$$u_t = c^2 u_{xx}, \text{ where } \begin{cases} u(0, t) = 0 \\ u(x, 0) = f(x) \end{cases}$$

$$\text{Let } u^*(x, t) = u(x, t) - \frac{T_x}{L}.$$

Consider the Fourier series of $f(x) - \frac{T_x}{L}$, then

$$u^*(x, t) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right) e^{-\frac{\pi^2 n^2 c^2 t}{L^2}}, \text{ where } B_n = b_n - \frac{2}{L} \int_0^L \frac{T_x}{L} \sin\left(\frac{n\pi x}{L}\right) dx = b_n + \frac{2T}{\pi n} (-1)^n.$$

Fisher's Equation

$$u_t = \alpha u_{xx} + \beta u(1-u) \Rightarrow \text{heat + rumour.}$$

We seek a solution of the form $u(x, t) = U(x-ct) \equiv U(s)$, moving to the right at constant speed c , starting at $x = 0$. As $x \rightarrow \infty$, $s \rightarrow \infty$; but as $t \rightarrow \infty$, $s \rightarrow -\infty$. Now $u_{xx} = U''$ and $u_t = -cU'$. We can reduce Fisher's Equation to

$$\alpha U'' + cU' + \beta U - \beta U^2 = 0.$$

$$\begin{cases} U' = V \\ V' = -\frac{c}{\alpha}V - \frac{\beta}{\alpha}U + \frac{\beta}{\alpha}U^2 \end{cases}$$

The system has two equilibrium points $(0, 0)$ and $(1, 0)$. $(0, 0)$ is a spiral sink if $c < 2\sqrt{\alpha\beta}$, which is rejected since U cannot be negative; hence $c \geq 2\sqrt{\alpha\beta}$, where $(0, 0)$ is a nodal sink.

Diffusion of Lions (Laplace)

When $u_t = c^2(u_{xx} + u_{yy})$ and everything has settled down to a steady state ($u_t = 0$), we have:

$$u_{xx} + u_{yy} = 0, u(x, 0) = u(0, y) = u(\pi, y) = 0.$$

Suppose $0 \leq x, y \leq \pi$. Take 4 boundary conditions:

$$u(x, 0) = 0; \quad u(0, y) = 0$$

$$u(\pi, y) = 0; \quad u(x, \pi) = f(x)$$

$f(x)$ describes the density of lions along the border that has the river. Let $u(x, y) = X(x)Y(y)$, we have $X''Y + XY'' = 0 \Rightarrow -\frac{X''}{X} = \frac{Y''}{Y} = \lambda$ and $X(0) = X(\pi) = 0$. Let $\lambda = n^2$, then $X(x) = \sin(nx)$, $Y(y) = c_n \sinh(ny)$, $c_n \in \mathbb{R}$. Hence,

$$u(x, y) = \sum_{n=1}^{\infty} c_n \sin(nx) \sinh(ny).$$

Putting $u(x, \pi) = f(x)$, we have $f(x) = \sum_{n=1}^{\infty} c_n \sin(nx) \sinh(n\pi)$.

$c_n \sinh(n\pi) = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$ is the Fourier series of the odd extension of $f(x)$. Since $n \geq 1$, $\sinh(n\pi) \neq 0$. For example, if $f(x) = \sin(x) + 0.2 \sin(4x)$, then $u(x, y) = c_1 \sin(x) \sinh(y) + c_2 \sin(4x) \sinh(4y)$, where $c_1 = \frac{1}{\sinh(\pi)}$ and $c_2 = \frac{0.2}{\sinh(4\pi)}$.

Appendix

Trigonometric Identities

- sin, cos: $\sin^2 x + \cos^2 x = 1$
- tan: $\tan x = \frac{\sin x}{\cos x}$
- sec, csc: $\sec x = \frac{1}{\cos x}; \csc x = \frac{1}{\sin x}$
- cot: $\cot x = \frac{1}{\tan x} = \frac{\cos x}{\sin x}$
- $\sec^2 x - \tan^2 x = 1; \csc^2 x - \cot^2 x = 1$
- $\sin(x+y) = \sin x \cos y + \sin y \cos x$
- $\sin 2x = 2 \sin x \cos x$
- $\sin \frac{x}{2} = \pm \sqrt{\frac{1-\cos x}{2}}$
- $\cos(x+y) = \cos x \cos y - \sin x \sin y$
- $\cos 2x = \cos^2 x - \sin^2 x = 1 - 2 \sin^2 x = \cos^2 x - 1$
- $\cos \frac{x}{2} = \pm \sqrt{\frac{1+\cos x}{2}}$
- $\tan(x+y) = \frac{\tan x + \tan y}{1 - \tan x \tan y}$
- $\tan 2x = \frac{2 \tan x}{1 - \tan^2 x}$
- $\tan \frac{x}{2} = \pm \sqrt{(1 - \cos x)(1 + \cos x)}$
- $\sin x + \sin y = 2 \sin \frac{x+y}{2} \cos \frac{x-y}{2}$
- $\sin x \sin y = \frac{\cos(x+y) - \cos(x-y)}{2}$
- $\cos x + \cos y = 2 \cos \frac{x+y}{2} \cos \frac{x-y}{2}$
- $\cos x \cos y = \frac{\cos(x-y) + \cos(x+y)}{2}$
- $\sin x \cos y = \frac{\sin(x+y) + \sin(x-y)}{2}$
- sinh, cosh: $\cosh^2 x - \sinh^2 x = 1$
- $\sinh x = \frac{e^x - e^{-x}}{2}; \cosh x = \frac{e^x + e^{-x}}{2}$
- $\tanh: \tanh x = \frac{\sinh x}{\cosh x}$
- $\sech x = \frac{1}{\cosh x}$
- $\csch x = \frac{1}{\sinh x}$
- coth: $\coth x = \frac{1}{\tanh x}$
- $\tanh^2 x + \sech^2 x = 1$
- $\coth^2 x - \csch^2 x = 1$
- $\sinh(x+y) = \sinh x \cosh y + \sinh y \cosh x$
- $\sinh 2x = 2 \sinh x \cosh x$
- $\cosh(x+y) = \cosh x \cosh y + \sinh x \sinh y$
- $\cosh 2x = \cosh^2 x + \sinh^2 x$
- $\tanh 2x = \frac{2 \tanh x}{1 + \tanh^2 x}$

Fractional

- $\int \frac{1}{ax+b} dx = \frac{1}{a} \ln|ax+b| + C$
- $\int \frac{1}{a^2+x^2} dx = \frac{1}{a} \tan^{-1}\left(\frac{x}{a}\right) + C$
- $\int \frac{1}{\sqrt{a^2-x^2}} dx = \sin^{-1}\left(\frac{x}{a}\right) + C$
- $\int \frac{1}{\sqrt{a^2+x^2}} dx = \sinh^{-1}\left(\frac{x}{a}\right) + C$
- $\int \frac{1}{\sqrt{x^2-a^2}} dx = \cosh^{-1}\left(\frac{x}{a}\right) + C$
- $\int \frac{1}{x\sqrt{1-x^2}} dx = -\operatorname{sech}^{-1} x + C$
- $\int \frac{1}{|x|\sqrt{1+x^2}} dx = -\operatorname{csch}^{-1} x + C$

Logarithmic

$$\int \ln x dx = x \ln x - x + C$$

Trigonometric

- $\int \cos x dx = \sin x + C$
- $\int \sin x dx = -\cos x + C$
- $\int \tan x dx = \ln|\sec x| + C$
- $\int \sec x dx = \ln|\sec u + \tan u| + C$
- $\int \sec^2 x dx = \tan x + C$
- $\int \sec x \tan x dx = \sec x + C$
- $\int \csc x \cot x dx = -\csc x + C$
- $\int \csc^2 x dx = -\cot x + C$
- $\int \sinh x dx = \cosh x + C$
- $\int \cosh x dx = \sinh x + C$
- $\int \operatorname{sech}^2 x dx = \tanh x + C$
- $\int \operatorname{csch}^2 x dx = -\coth x + C$
- $\int \operatorname{sech} x \tanh x dx = -\operatorname{sech} x + C$
- $\int \tanh x dx = \ln|\cosh x| + C$
- $\int \coth x dx = \ln|\sinh x| + C$
- $\int \operatorname{sech} x dx = \tan^{-1}(\sinh x)$
- $\int \operatorname{csch} x dx = \ln|\tanh \frac{x}{2}| + C$

Special Integrals

- Partial fractions
- Integration by parts:
- $\int u dv = uv - \int v du$
- $\int \sin^n x \cos^m x dx$

Use trigonometric identities to convert it into $\sin^k x \cos x$ or $\cos^k x \sin x$.

Integrals

Basic

- $\int k dx = kx + C$
- $\int x^n dx = \frac{1}{n+1} x^{n+1} + C$
- $\int \frac{1}{x} dx = \ln|x| + C$
- $\int e^x dx = e^x + C$