MA3264 Mathematical Modelling

AY2022/23 Semester 1 · Prepared by Tian Xiao

Basic ODEs

Separable Equation

$$M(x) - N(y)y' = 0 \Rightarrow \int M(x)dx = \int N(y)dy.$$

Integrating Factor

$$y' + P(x)y = Q(x)$$

Let $\mu(x) = e^{\int P(x)dx}$:

$$\mu(x)y' + \mu'(x)y = \mu(x)Q(x)$$

$$y = \frac{\int \mu(x)Q(x)dx}{\mu(x)}$$

Bernoulli Equation

$$y' + P(x)y = Q(x)y^{n}$$
Let $z = y^{1-n}$, then $z' = (1-n)y^{-n}y'$:
$$y^{-n}y' + P(x)y^{1-n} = Q(x)$$

$$\frac{z'}{1-n} + P(x)z = Q(x)$$

2nd Order Equation

$$ay'' + by' + cy = r(x) \Leftrightarrow ax^2 + bx + c = 0.$$

$$\begin{cases} y = C_1 e^{\lambda_1 x} + C_2 e^{\lambda_2 x} & \lambda_1 \neq \lambda_2 \in \mathbb{R} \\ y = (C_1 + C_2 x) e^{\lambda_x} & \lambda_1 = \lambda_2 \in \mathbb{R} \\ y = e^{\alpha x} (C_1 \cos \beta x + C_2 \sin \beta x) & \lambda = \alpha \pm \beta i \in \mathbb{C} \end{cases}$$
To find the particular solution:

- If $r \in \mathcal{P}^n$, $y_p(x) \leftarrow p^n$.
- If $r(x) = g(x)e^{kx}$, $y_p(x) \leftarrow u(x)e^{kx}$.
- If $r(x) = g(x) \cos kx$ or $g(x) \sin kx$, let $z(x) \leftarrow u(x)e^{ikx}$ and take Re(z) or Im(z).

Population Models

Malthus Model

$$\frac{dN}{dt} = (B-D)N \Leftrightarrow N(t) = N_0 e^{(B-D)t}$$
.

Logistic Model

$$\frac{dN}{dt} = BN - sN^2 \Leftrightarrow N(t) = \frac{B/S}{1 + e^{-Bt}(\frac{B}{N_0s} - 1)}.$$

 B/S is a stable equilibrium point.

System of Linear ODEs

General Linear ODE System

$$\frac{d}{dt} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} \Leftrightarrow \begin{cases} \frac{dx}{dt} = ax + by \\ \frac{dy}{dt} = cx + dy \end{cases}.$$

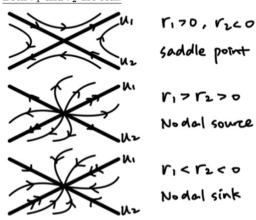
Eigenvalues: $r = \frac{1}{2} \left[\text{Tr}(B) \pm \sqrt{\text{Tr}(B)^2 - 4\text{Det}(B)} \right]$ Solutions: $\boldsymbol{u}(t) = C_+ e^{r_+ t} \boldsymbol{u}_+ + C_- e^{r_- t} \boldsymbol{u}_-$.

Nonhomogenous System

$$\frac{d}{dt} \begin{bmatrix} x \\ y \end{bmatrix} = B \begin{bmatrix} x \\ y \end{bmatrix} + F \Rightarrow -B^{-1}F.$$

Phase Plane Classification

Both r_1 and r_2 are real:



Both r_1 and r_2 are complex:



Re[r]<0 Re[r]>0 Re[r]=0 Spiral sink Spiral source Centre

System of Non-Linear ODEs

Linearisation

$$\begin{cases} \frac{dx}{dt} = f(x,y) & \text{with equilibrium point } (a,b). \\ \frac{dy}{dt} = g(x,y) & \text{By Taylor expansion,} \end{cases}$$

Lotka-Volterra Model

$$\begin{cases} \frac{dL}{dt} = uZL - D_LL \\ \frac{dZ}{dt} = B_zZ - sLZ \end{cases} L: \text{ lion; } Z: \text{ zebra.}$$

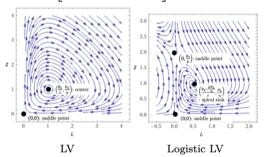
Every trajectory in the contour is periodic:

$$\left(\frac{B_z}{L} - s\right) \frac{dL}{dt} + \left(\frac{D_L}{Z} - u\right) \frac{dZ}{dt} = 0$$

$$B_z \ln L - sL + D_L \ln Z - uZ = C$$

$$F(L, Z) \text{ has only 1 maximum} \Rightarrow \text{closed.}$$

$$J(L,Z) = egin{bmatrix} uZ - D_L & uL \ -sZ & B_z - sL \end{bmatrix}.$$

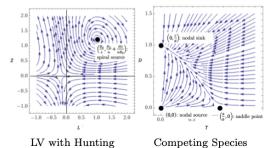


Logistic LV Model

$$\begin{cases} \frac{dL}{dt} = uZL - D_LL \\ \frac{dZ}{dt} = (B_zZ - pZ^2) - sLZ \end{cases}$$

Lion Hunting

$$\begin{cases} \frac{dL}{dt} = uZL - D_LL - H\\ \frac{dZ}{dt} = B_zZ - sLZ \end{cases}$$



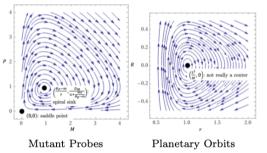
Competing Species

$$\begin{cases} \frac{dT}{dt} = (a-kD)T - bT^2 & T: \text{ thylacine;} \\ \frac{dD}{dt} = (c-\sigma T)D - dD^2 & D: \text{ dingo.} \end{cases}$$

[Principle of Competitive Exclusion]: If two species are too similar, one will wipe out another.

Where Is Everybody?

$$\begin{cases} \frac{dM}{dt} = uPM - D_MM + mP & M: \text{ mutant probe;} \\ \frac{dP}{dt} = B_PP - sMP - mP & P: \text{ normal probe.} \end{cases}$$



Non-Linear 2nd Order ODEs

Consider the Earth moving around the Sun:

$$\ddot{r} = -\frac{M}{r^2} + \frac{L^2}{r^3} \Rightarrow \begin{cases} \dot{r} = R \\ \dot{R} = -\frac{M}{r^2} + \frac{L^2}{r^3} \end{cases}$$

The equilibrium point is almost a center.

Partial Differential Equations

A PDE is an equation containing an unknown function u of 2 or more independent variables x, y, \cdots and its partial derivatives with respect to them.

Separation of Variables

PDE: $u_x = f(x)g(y)u_y$. Suppose u = X(x)Y(y), then X'(x)Y(y) = f(x)g(y)X(x)Y'(y) = k.

$$\begin{cases} X'(x) = kf(x)X(x) \\ Y'(y) = \frac{k}{g(y)}Y(y) \end{cases}$$

Wave Equation

$$c^2 \frac{\partial^2 y}{\partial x^2} = \frac{\partial^2 y}{\partial t^2}, \text{ where } \begin{array}{c} y(t,0) = 0 & y(t,\pi) = 0 \\ y(0,x) = f(x) & \frac{\delta y}{\delta t}(0,x) = 0 \end{array}$$

We need 4 pieces of information for a solution.

d'Alembert's Solution

$$y(t,x) = \frac{1}{2}[f(x+ct) + f(x-ct)].$$

Separation of Variables

Let
$$y(t,x)=u(x)v(t)$$
, then
$$\begin{cases} u''+\lambda u=0\\ v''+\lambda c^2v=0 \end{cases}$$
. From
$$\begin{cases} y(t,0)=u(0)v(t)=0\\ y(t,\pi)=u(\pi)v(t)=0 \end{cases}$$
,

for u to cut the x-axis twice, we have $\lambda > 0$.

Let
$$\lambda = n^2$$
 and $u = C \cos(nx) + D \sin(nx)$.
Since $u(0) = 0$, $C = 0$ and $u = D \sin(nx)$.

Since $u(\pi) = 0, n \in \mathbb{Z}$.

Similarly, $v(t) = A\cos(nct)$

Therefore, $y = b_n \sin(nx) \cos(nct)$ and only y(0,x) = f(x) is not satisfied yet.

Fourier Series

Any odd function f(x) of period 2π on $[0,\pi]$ can be expressed as $f(x) = \sum_{n=1}^{\infty} b_n \sin(nx)$, where $b_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$. So the complete solution is: $y(t,x) = \sum_{n=1}^{\infty} b_n \sin(nx) \cos(nct)$.

If we want [0, L] instead of $[0, \pi]$, the Fourier formulae becomes $g(x) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right)$ and $b_n = \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi x}{L}\right) dx$. The complete solution is: $y(t,x) = \sum_{n=0}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right) \cos\left(\frac{n\pi ct}{L}\right)$.

Tsunami (Korteweg-de Vries)

 $\partial_t \eta + \sqrt{gh} \partial_x \eta + \frac{3}{2} \sqrt{\frac{g}{h}} \eta \partial_x \eta + \frac{1}{6} h^2 \sqrt{gh} \partial_x^3 \eta = 0,$ where η denotes the elevation above sea level.

Suppose $\eta = E(x - ct)$, then we can simplify to -2AE' + 6BEE' + 2CE''' = 0. $-2AE + 3BE^2 + 2CE'' = 0 \text{ (integrate)}.$ $-AE^{2} + BE^{3} + C(E')^{2} = K$ (integrate w.r.t. E).

Heat Equation

$$u_t = c^2 u_{xx}$$
, where $u(0,t) = u(L,t) = 0$.
 $u(x,0) = f(x)$.
Solution: $u(x,t) = \sum_{x=1}^{\infty} b_n \sin \frac{n\pi x}{L} e^{-\frac{\pi^2 n^2 c^2}{L^2}t}$.

Heat Equation Variant

$$u_t=c^2u_{xx}, \text{ where } \dfrac{u(0,t)=0}{u(x,0)=f(x)} \qquad u(L,t)=T$$
 Let $u^*(x,t)=u(x,t)-\frac{T_x}{L}.$

Consider the Fourier series of $f(x) - \frac{Tx}{L}$, then **Appendix** $u^*(x,t) = \sum_{n=0}^{\infty} B_n \sin \frac{n\pi x}{L} e^{-\frac{\pi^2 n^2 c^2}{L^2}t}$, where $B_n = \frac{n\pi x}{L}$ $b_n - \frac{2}{L} \int_0^L \frac{Tx}{L} \sin \frac{n\pi x}{L} dx = b_n + \frac{2T}{\pi n} (-1)^n$.

Fisher's Equation

 $u_t = \alpha u_{xx} + \beta u(1-u) \Rightarrow \text{heat} + \text{rumour}.$ We seek a solution of the form $u(x,t) = U(x-ct) \equiv$ U(s), moving to the right at constant speed c, starting at x = 0. As $x \to \infty$, $s \to \infty$; but as $t \to \infty$, $s \to -\infty$. Now $u_x x = U''$ and $u_t = -cU'$. We can reduce Fisher's Equation to

$$\alpha U'' + cU' + \beta U - \beta U^2 = 0.$$

$$\begin{cases} U' = V \\ V' = -\frac{c}{\alpha}V - \frac{\beta}{\alpha}U + \frac{\beta}{\alpha}U^2 \end{cases}$$

The system has two equilibrium points (0,0) and (1,0). (0,0) is a spiral sink if $c < 2\sqrt{\alpha\beta}$, which is rejected since U cannot be negative; hence $c > 2\sqrt{\alpha\beta}$, where (0,0) is a nodal sink.

Diffusion of Lions (Laplace)

When $u_t = c^2(u_{xx} + u_{yy})$ and everything has settled down to a steady state ($u_t = 0$), we have: $u_{xx} + u_{yy} = 0$, $u(x, 0) = u(0, y) = u(\pi, y) = 0$. Suppose $0 < x, y < \pi$. Take 4 boundary conditions:

$$u(x,0) = 0;$$
 $u(0,y) = 0$
 $u(\pi,y) = 0;$ $u(x,\pi) = f(x)$

f(x) describes the density of lions along the border that has the river. Let u(x,y) = X(x)Y(y), we have $X''Y + XY'' = 0 \Rightarrow -\frac{X''}{X} = \frac{Y''}{Y} =$ λ and $X(0) = X(\pi) = 0$. Let $\lambda = n^2$, then $X(x) = \sin(nx), Y(y) = c_n \sinh(ny), c_n \in \mathbb{R}$. Hence $u(x,y) = \sum_{n=1}^{\infty} c_n \sin(nx) \sinh(ny)$. Putting $u(x,\pi) =$ f(x), we have $f(x) = \sum_{n=1}^{\infty} c_n \sin(nx) \sinh(n\pi)$. $c_n \sinh(n\pi) = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$ is the Fourier series of the odd extension of f(x). Since $n \ge - \int k \, dx = kx + C$ 1, $\sinh(n\pi) \neq 0$. For example, if $f(x) = -\int x^n dx = \frac{1}{n+1}x^{n+1} + C$ $\sin(x) + 0.2\sin(4x)$, then $u(x, y) = c_1 \sin(x)\sinh(y) + \int_{0}^{1} \frac{1}{x} dx = \ln|x| + C$ $c_2 \sin(4x)\sinh(4y)$, where $c_1 = \frac{1}{\sinh(\pi)}$ and $c_2 = \int_{0}^{1} e^x dx = e^x + C$

Trigonometric Identities

- $\sin \cos \sin^2 x + \cos^2 x = 1$ - $\tan x = \frac{\sin x}{\cos x}$ - sec, csc: $\sec x = \frac{1}{\cos x}$; $\csc x = \frac{1}{\sin x}$; - cot: $\cot x = \frac{1}{\tan x} = \frac{\cos x}{\sin x}$ - $\sec^2 x - \tan^2 x = 1$; $\csc^2 x - \cot^2 x = 1$ $-\sin(x+y) = \sin x \cos y + \sin y \cos x$
- $-\sin 2x = 2\sin x\cos x$ $-\sin\frac{x}{2} = \pm\sqrt{\frac{1-\cos x}{2}}$
- $-\cos(x+y) = \cos x \cos y \sin x \sin y$
- $-\cos 2x = \cos^2 x \sin^2 x = 1 2\sin^2 x = \cos^2 x 1$ Logarithmic $-\cos\frac{x}{2} = \pm\sqrt{\frac{1+\cos x}{2}}$
- $-\tan(x+y) = \frac{\tan x + \tan y}{1 \tan x \tan y}$
- $-\tan 2x = \frac{2\tan x}{1-\tan^2 x}$
- $-\tan\frac{x}{2} = \pm\sqrt{(1-\cos x)(1+\cos x)}$
- $-\sin x + \sin y = 2\sin \frac{x+y}{2}\cos \frac{x-y}{2}$
- $-\sin x \sin y = \frac{\cos(x+y) \cos(x-y)}{2}$
- $-\cos x + \cos y = 2\cos\frac{x+y}{2}\cos\frac{x-y}{2}$
- $-\cos x \cos y = \frac{\cos(x-y) + \cos(x+y)}{2}$
- $-\sin x \cos y = \frac{\sin(x+y) + \sin(x-y)}{2}$
- \sinh , \cosh : $\cosh^2 x \sinh^2 x = 1$ $\sinh x = \frac{e^x - e^{-x}}{2}; \cosh x = \frac{e^x + e^{-x}}{2}$
- tanh: $\tanh x = \frac{\sinh x}{\cosh x}$
- $\operatorname{sech} x = \frac{1}{\cosh x}$
- $\operatorname{csch} x = \frac{1}{\sinh x}$
- coth: $\coth x = \frac{1}{\tanh x}$ $\tanh^2 x + \operatorname{sech}^2 x = 1$
- $-\coth^2 x \operatorname{csch}^2 x = 1$
- $-\sinh(x+y) = \sinh x \cosh y + \sinh y \cosh x$
- $-\sinh 2x = 2\sinh x \cosh x$
- $-\cosh(x+y) = \cosh x \cosh y + \sinh x \sinh y$
- $-\cosh 2x = \cosh^2 x + \sinh^2 x$
- $-\tanh 2x = \frac{2\tanh x}{1+\tanh^2 x}$

$$-\int k \, dx = kx + C$$

$$-\int x^n \, dx = \frac{1}{n+1}x^{n+1} + C$$

$$-\int \frac{1}{x} \, dx = \ln|x| + C$$

$$-\int e^x \, dx = e^x + C$$

Fractional

$$-\int \frac{1}{ax+b} = \frac{1}{a} \ln|ax+b| + C$$

$$-\int \frac{1}{a^2+x^2} dx = \frac{1}{a} \tan^{-1}(\frac{x}{a}) + C$$

$$-\int \frac{1}{\sqrt{a^2-x^2}} dx = \sin^{-1}(\frac{x}{a}) + C$$

$$-\int \frac{1}{\sqrt{a^2+x^2}} dx = \sinh^{-1}(\frac{x}{a}) + C$$

$$-\int \frac{1}{\sqrt{x^2-a^2}} dx = \cosh^{-1}(\frac{x}{a}) + C$$

$$-\int \frac{1}{a^2-x^2} dx = \frac{1}{a} \tanh^{-1}(\frac{x}{a}) + C$$

$$-\int \frac{1}{x\sqrt{1-x^2}} dx = -\operatorname{sech}^{-1} x + C$$

$$-\int \frac{1}{|x|\sqrt{1+x^2}} dx = -\operatorname{csch}^{-1} x + C$$

$$-\int \ln x \, dx = x \ln x - x + C$$

Trigonometric

- $\int \cos x \, dx = \sin x + C$
- $-\int \sin x \, dx = -\cos x + C$
- $\int \tan x \, dx = \ln|\sec x| + C$
- $\int \sec x \, dx = \ln|\sec u + \tan u| + C$
- $\int \sec^2 x \, dx = \tan x + C$
- $\int \sec x \tan x \, dx = \sec x + C$
- $-\int \csc x \cot x \, dx = -\csc x + C$
- $-\int \csc^2 x \, dx = -\cot x + C$
- $\int \sinh x \, dx = \cosh x + C$
- $-\int \cosh x \, dx = \sinh x + C$
- $\int \operatorname{sech}^2 x \, dx = \tanh x + C$
- $-\int \operatorname{csch}^2 x \, dx = -\coth x + C$
- $\int \operatorname{sech} x \tanh x \, dx = \operatorname{sech} x + C$
- $\int \operatorname{csch} x \operatorname{coth} x \, dx = -\operatorname{csch} x + C$

Special Integrals

- Partial fractions
- Integration by parts:
- $\int u \, dv = uv \int v \, du$
- $\int \sin^n x \cos^m x \, dx$:

Use trigonometric identities to convert it into $\sin^k x \cos x$ or $\cos^k x \sin x$.