

## SPEECH PROCESSING

### Properties of Digital Filters

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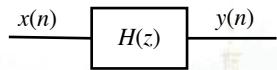
## Aims of this module

- This lecture reviews some well known facts about filters and introduces some less known ones that will be needed later on.
  - Derive the power response of first order FIR and IIR filters and relate this to the geometry of the pole-zero diagram.
  - Relate the bandwidth of a 2nd-order resonance to the geometry of the pole-zero diagram.
  - Describe the bandwidth expansion transformation of a filter.
  - Describe the effect of reversing the coefficients of a filter.
  - Derive expressions for the log frequency response and its average value.

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## 1<sup>st</sup> order FIR filter

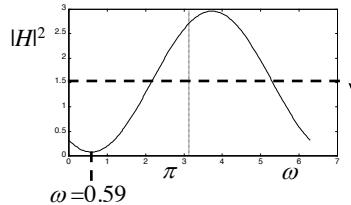


$$H(z) = 1 - az^{-1} \Leftrightarrow y(n) = x(n) - ax(n-1)$$

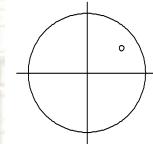
- Filter has a single zero at  $z = a = re^{j\theta}$

• Frequency response of filter is given by:  $H(e^{j\omega}) = 1 - ae^{-j\omega}$

- Power response of filter is given by:  $|H(e^{j\omega})|^2 = H(e^{j\omega})H^*(e^{j\omega})$
  - Example:  $a = 0.6 + 0.4j = 0.72e^{0.59j}$
- $$|H(e^{j\omega})|^2 = 1.52 - 1.44 \cos(\omega - 0.59)$$
- $$= (1 - ae^{-j\omega})(1 - a^* e^{+j\omega})$$
- $$= 1 + r^2 - 2r \cos(\omega - \theta)$$



Mean power gain =  $\frac{1}{1+r^2}$



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## Log Frequency Response for $|a| < 1$

- Calculate the log response of the filter  $\log(H(e^{j\omega})) = \log(1 - ae^{-j\omega})$
- If  $|a| < 1$  then  $|ae^{-j\omega}| < 1$  and we can expand the log as a power series using

$$\log(1 - d) = -\left(d + \frac{d^2}{2} + \frac{d^3}{3} + \dots\right) \quad \text{for } |d| < 1$$

- Hence  $\log(H(e^{j\omega})) = -\sum_{n=1}^{\infty} \frac{a^n}{n} e^{-jn\omega}$

• Note that  $\log(re^{j\phi}) = \log(r) + j\phi \Rightarrow \log(|x|) = \Re(\log(x))$

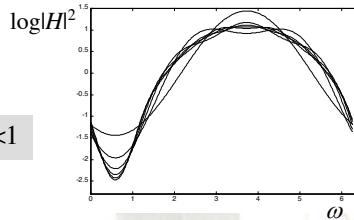
- Hence  $\log(|H(e^{j\omega})|^2) = -2 \sum_{n=1}^{\infty} \frac{r^n}{n} \cos(n(\omega - \theta))$  where  $a = re^{j\theta}$

Log power spectrum

First six terms in the summation for:

$$a = 0.6 + 0.4j$$

The average of  $\log|H|^2$  is always zero if  $|a| < 1$



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## Log Frequency Response for $|a|>1$

- If  $|a|>1$ , we can rearrange the formula in terms of  $a^{-1}$

$$\log(H(e^{j\omega})) = \log(-ae^{-j\omega}(1-a^{-1}e^{j\omega})) = \log(-ae^{-j\omega}) + \log(1-a^{-1}e^{j\omega})$$

- Since  $|a^{-1}|<1$ , we can expand the log as before to obtain

$$\log(H(e^{j\omega}))^2 = 2\log|a| - 2 \sum_{n=1}^{\infty} \frac{r^{-n}}{n} \cos(n(\omega - \theta)) \quad \text{where } a = re^{j\theta}$$

- The average of  $\log|H|^2$  is  $2\log|a|$  if  $|a|>1$

- The log response of an arbitrary filter is just the sum of the log responses of each pole or zero. For a stable filter, all the poles must be within the unit circle. Hence ...

- Given a stable filter  $H(z) = \frac{b_0 + b_1 z^{-1} + b_2 z^{-2} + \dots}{a_0 + a_1 z^{-1} + a_2 z^{-2} + \dots}$

- Then the average value of  $\log|H|^2$  is given by  $2\log\left(\frac{|b_0|}{|a_0|}\right) + 2 \sum_{\text{zeros with } |z_i|>1} \log(|z_i|)$

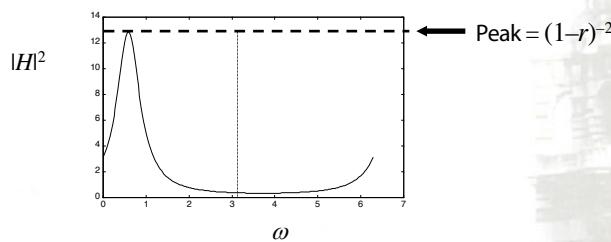
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## Single Pole Filter

$$H(z) = \frac{1}{1 - az^{-1}} \Leftrightarrow y(n) = x(n) + ay(n-1)$$

- Filter has a single pole at  $z = a = re^{j\theta}$
- Power response of filter is given by  $|H(e^{j\omega})|^2 = \frac{1}{1 + r^2 - 2r \cos(\omega - \theta)}$



- Note: response is no longer a simple cosine wave

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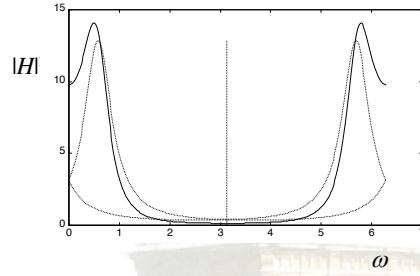
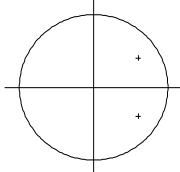
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## Pole Pairs

- If the filter coefficients are real, any complex zeros or poles will always occur in conjugate pairs
- The response of the filter is the product of the responses of the individual poles
  - Conjugate pole/zero pairs ensure a symmetric response
- Example: Poles at  $0.6 \pm 0.4j = 0.72e^{\pm 0.59j} = re^{\pm j\theta}$

$$H(z) = \frac{1}{1 - 2r\cos\theta z^{-1} + r^2 z^{-2}}$$

$$= \frac{1}{1 - 1.2z^{-1} + 0.52z^{-2}}$$

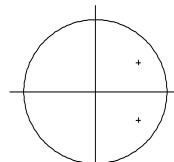


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## Geometrical Interpretation

$$H(z) = \frac{1}{(1 - az^{-1})(1 - a^* z^{-1})}$$

$$|H(z)| = \frac{1}{|1 - az^{-1}| |1 - a^* z^{-1}|}$$



- But since  $|z|=1$ , we have  $|1 - az^{-1}| = |z^{-1}| |z - a| = |z - a|$

– This is just the distance between  $z$  and  $a$

- The magnitude response of the filter

$$H(z) = \frac{b_0 + b_1 z^{-1} + b_2 z^{-2} + \dots}{a_0 + a_1 z^{-1} + a_2 z^{-2} + \dots} = \frac{b_0}{a_0} \times \frac{\prod (1 - x_i z^{-1})}{\prod (1 - y_i z^{-1})}$$

at a frequency  $\omega$  is proportional to the product of the distance from the point  $e^{j\omega}$  to all the zeros divided by the product of the distance to all the poles.

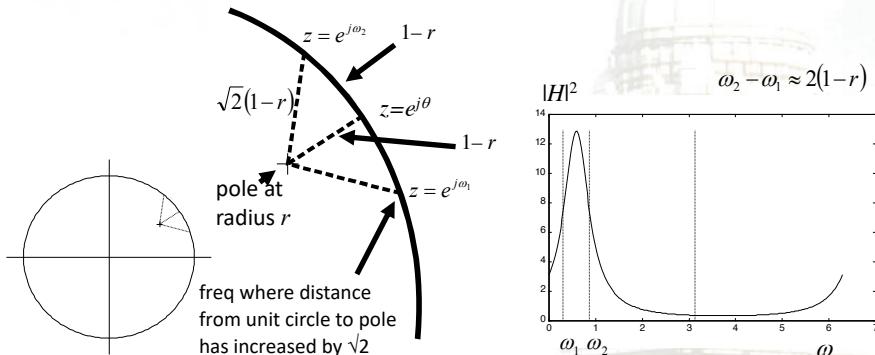
– The constant of proportionality is  $b_0/a_0$ .

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## Bandwidth of a Resonance Peak

- The bandwidth of a resonance peak is the frequency range at which the magnitude response has decreased by  $\sqrt{2}$  compared to the peak value.
- For poles near the unit circle this is approximately  $2(1-r)$  rad/s =  $(1-r)/\pi$  Hz (normalized).



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## Bandwidth Expansion

- Given a filter  $H(z) = \frac{b_0 + b_1 z^{-1} + b_2 z^{-2} + \dots}{a_0 + a_1 z^{-1} + a_2 z^{-2} + \dots}$

form a new filter by multiplying coefficients  $a_i$  and  $b_i$  by  $k^i$  for some  $k < 1$

$$G(z) = H(z/k) = \frac{b_0 + b_1 k z^{-1} + b_2 k^2 z^{-2} + \dots}{a_0 + a_1 k z^{-1} + a_2 k^2 z^{-2} + \dots}$$

- If  $H(z)$  has a pole/zero at  $z_0$ , then  $G(z)$  will have one at  $kz_0$ 
  - All poles and zero will be moved towards the origin by a factor  $k$
- If the bandwidth of a pole of  $H(z)$  is  $b=2(1-r)$ , then the bandwidth of the corresponding pole in  $G(z)$  will be expanded to:

$$2(1-kr) = b + 2r(1-k)$$

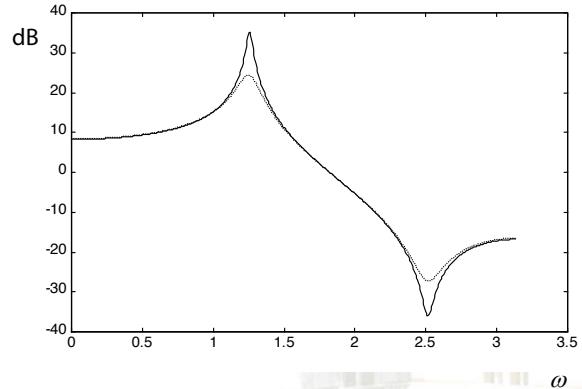
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## Example

$$H(z) = \frac{1 + 1.57z^{-1} + 0.94z^{-2}}{1 - 0.67z^{-1} + 0.96z^{-2}}$$

$k = 0.95$



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## Coefficient Reversal

- Given a filter  $H(z) = b_0 + b_1 z^{-1} + b_2 z^{-2} + \dots + b_p z^{-p}$   
form a new filter by conjugating the coefficients and putting them in reverse order:

$$G(z) = b_p^* + b_{p-1}^* z^{-1} + b_{p-2}^* z^{-2} + \dots + b_0^* z^{-p} = z^{-p} H^*(z^{*-1})$$

- If  $z_0$  is a zero of  $H(z)$  then  $z_0^{*-1}$  is a zero of  $G(z)$ 
  - This is called a *reflection* in the unit circle.
- The frequency response of  $G(z)$  is given by:

$$G(e^{j\omega}) = e^{-jp\omega} H^*(e^{j\omega})$$

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- Hence  $G(z)$  has the same magnitude response as  $H(z)$  but a different phase response:

$$|G(e^{j\omega})| = |H(e^{j\omega})| \quad \text{Arg}(G(e^{j\omega})) = -\text{Arg}(H(e^{j\omega})) - p\omega$$

