

Linking ocean mixing and overturning circulation

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¹⁶ ABSTRACT:

¹⁷ **Begin of info box**

¹⁸ 11th Warnemünde Turbulence Days on *Linking ocean mixing and circulation at various scales*

¹⁹ **What:** 45 participants from 10 countries met to discuss about contemporary issues of marine
²⁰ turbulence with focus on the linkage between mixing and overturning circulation on all scales
²¹ (<https://www.io-warnemuende.de/wtd-2023.html>).

²² **When:** September 17-20, 2023

²³ **Where:** University of Rostock, Germany

²⁴ **End of info box**

²⁵ Walter Munk, in his famous *abyssal recipes*, showed more than half a century ago that the
²⁶ strength of the global overturning circulation is closely linked to diapycnal mixing. Since then,
²⁷ oceanographic research has succeeded in identifying more and more processes generating the
²⁸ mixing required to drive the global overturning circulation: internal-wave mixing, boundary
²⁹ mixing, wake eddies, gravity currents, double diffusion, and many more. The same dependence
³⁰ was also observed in other marine systems at smaller scales, including marginal and semi-enclosed
³¹ seas and estuaries. Numerical models describing these mechanisms often include discretization
³² errors that become evident in particular in the form of spurious *numerical* mixing, which may
³³ trigger artificial circulation patterns at all scales.

³⁴ The Warnemünde Turbulence Days (WTD, <http://www.io-warnemuende.de/wtd.html>) have
³⁵ been established in 2003 to provide a regular international forum for discussing new developments
³⁶ in marine turbulence. Since then, the WTD have been organised every other year such that in 2023
³⁷ we organised the 11th WTD that took place during September 17-20, 2023 in Rostock. We invited
³⁸ contributions discussing all aspects of mixing in the ocean, especially, however, those that focus
³⁹ on the relation of mixing and circulation at all relevant scales.

40 **1. Context**

41 Here we report on the presentations and discussions that occurred during the WTD workshop.
42 The topic *Linking ocean mixing and circulation at various scales* provided the common thread of
43 discussions, with a special focus on the overturning circulations, i.e. circulations in the vertical
44 plane involving density transformations. Often, the notion *Mixing is the driver of overturning*
45 *circulation* was advanced, but it was in general agreed that both are interdependent, and that the
46 notion of *driver* is ambiguous and should be used with caution. Many of the discussions were
47 also guided by the scale-specific differences of ocean dynamics, specifically to find out what can
48 be learned for basin-scale dynamics from the relatively simple functioning of estuaries. In this
49 report we do therefore highlight characteristics of mixing and overturning circulation on different
50 regimes and scales.

51 The tank experiments of Sandström 1908 illustrated more than a century ago how an ocean forced
52 only by surface buoyancy fluxes would produce a very weak and shallow stratification with minute
53 overturning rates controlled by the molecular diffusion across the sharp thermocline. While the
54 applicability of these experimental results to the real ocean is debatable (Coman et al. 2006), the
55 central importance of mechanically-forced interior mixing in maintaining deep density gradients is
56 now well established (Munk 1966; Munk and Wunsch 1998). This has led some authors to assert
57 that the ocean circulation is powered by winds and tides, and that it is *not driven by the difference*
58 *in density between equator and pole* (Paparella and Young 2002). However, this seems to be at
59 odds with the simple observation that overturning circulations are greatly shaped by the surface
60 buoyancy forcing (Kuhlbrodt et al. 2007).

61 On the estuarine scale, the basic theory was already laid out by Knudsen (1900) who quantified
62 exchange flows on smaller scales. He mentioned mixing in his paper, but a quantitative connection
63 between exchange flow and mixing, based on the principles of Knudsen (1900) was only formulated
64 by MacCready et al. (2018). The beauty and simplicity of estuarine theories of mixing and exchange
65 flow is the possibility to map the dynamics into salinity coordinates.

⁶⁶ **2. Diagnostics**

⁶⁷ Approaches to the key question *Is mixing driving overturning circulation?* do strongly depend
⁶⁸ on the definitions of both *mixing* and *overturning circulation*. For each of them numerous concepts
⁶⁹ exist as shown below.

⁷⁰ *a. What is mixing?*

⁷¹ In the ocean, turbulent diffusivity (Munk 1966; St Laurent and Thurnherr 2007) or dissipation
⁷² rate are often used as measures of mixing, following the Osborn (1980) relation that links both
⁷³ together via the buoyancy frequency. Both quantities are however not suitable measures for mixing
⁷⁴ on their own, since they would typically show large values in well-mixed regions such as in the
⁷⁵ surface and bottom mixed layers (where mixing is expected to be low). Moreover, both quantities
⁷⁶ do not indicate which tracer is mixed. Diffusive tracer fluxes do not have this disadvantage, because
⁷⁷ they are specific for each individual tracer. Often, (twice) the product of tracer flux and tracer
⁷⁸ gradient is defined as *mixing*, as it quantifies the *local tracer variance decay* (Burchard et al. 2009).
⁷⁹ Under the eddy diffusivity assumption, it results in twice the eddy diffusivity times the square of the
⁸⁰ tracer gradient, i.e., it is non-negative and vanishing for zero turbulence and zero tracer gradient.
⁸¹ This definition of variance decay as mixing is also consistent with turbulence theory that assumes
⁸² stirring to increase the micro-structure tracer variance and mixing to dissipate it (Umlauf and
⁸³ Burchard 2005; Thorpe 2007; Moum 2021). Many instabilities act to produce stirring, forming a
⁸⁴ spectrum of turbulent patterns linking the mesoscale to the microscale at which diffusion operates.
⁸⁵ Some rely on rotation, such as barotropic and baroclinic instabilities, some act at finer scale such as
⁸⁶ the Kelvin-Helmholtz instability or the static instability, while others rely on mechanisms specific
⁸⁷ to seawater (e.g., double diffusion, cabbeling, thermobaricity).

⁸⁸ *b. What is overturning circulation?*

⁸⁹ The question of how to define overturning circulation is far from trivial. Like the word *overturn*
⁹⁰ implies, it should involve a flow that loops vertically over itself as opposed to a circulation that
⁹¹ would be purely horizontal. Most overturning circulations involve exchanges of heat, salt and other
⁹² biogeochemical tracers between the surface and the ocean interior, making them of relevance to
⁹³ ecosystems and the climate.

Overturning circulations typically results from the superimposition of various intricate circulation patterns generated by different processes such as wind pumping, eddy transport or boundary dynamics. The overturning streamfunction obtained by integrating bottom up the basin wide meridional transport provides a natural proxy. The vertical integration is usually done using depth coordinates, although the use of density coordinates is also popular. Interior flows being preferentially oriented along neutral surfaces, the latter choice helps to distinguish the diabatic (mixing-driven) component of the overturning. Salinity coordinates are often applied to quantify the overturning circulation and exchange flow in estuaries (Knudsen 1900; MacCready 2011, and Sec. 3c). The water-mass transformation (WMT) framework links the overturning streamfunction to dia-surface volume transports associated with diffusive fluxes across isosurfaces of water mass properties (Groeskamp et al. 2019). The original integral framework has been developed in salinity space (Walsh 1977), but can be extended to temperature-salinity space (Döös et al. 2012). Later, also the direct link of water mass transformation to mixing in terms of tracer variance decay (MacCready et al. 2018; Burchard 2020), and local formulations (Winters and D'Asaro 1996; Wang et al. 2017; Klingbeil and Henell 2023) have been presented. For streamfunctions and the WMT framework each coordinate space has its pros and cons, and it is usually through their combination that a faithful picture emerges.

3. Overturning and mixing processes in different scales and regimes

a. Large scale overturning circulation

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Large scale overturning circulations rely on a balance between several key processes, involving mixing on a wide range of spatial and temporal scales. The simplistic picture of a deep convective flux in the high-latitude being balanced by a vertical diffusive flux across a thermocline is at odds with the insufficient amount of mixing seemingly available in the ocean interior (Munk and Wunsch 1998). To resolve this apparent conundrum, two main threads have been followed: one focusing on searching the “missing mixing” that led to intense observational efforts (e.g. Polzin 2009) and a renewed understanding of internal wave dynamics (section 3d); the other highlighting the role of the Southern Ocean providing a conduit for the overturning flow (Toggweiler and Samuels 1998). A new paradigm has since emerged that is best encapsulated in the model of Nikurashin and Vallis

123 (2012), involving a balance between interior mixing, high-latitude convective mixing, together
124 with wind-driven upwelling and eddy transports in the Southern Ocean (Cessi 2019).

125 However, the key dynamics of the laterally closed parts of the overturning circulation in particular
126 in the Atlantic Ocean – the AMOC – are still not understood (Straub 1996; Brüggemann et al.
127 2011). It is clear that the zonal pressure difference at the boundaries balances the AMOC, but
128 it is also clear that the north-south pressure or density gradient should not be part of the AMOC
129 dynamics as in the model by Nikurashin and Vallis (2012) and several others. The key to understand
130 the AMOC might be the (dissipative) processes in the western boundaries, but for which simple
131 consistent models are still missing. The missing understanding of the AMOC is reflected (and
132 could be resolved) by missing consistent models of the zonally averaged properties of Atlantic and
133 Pacific ocean.

134 Furthermore, the surface buoyancy forcing leading to deep convection is not independent of
135 the ocean circulation itself, as is often assumed. This is a consequence of the nonlinear nature
136 of the equation of state, making the impact of heat fluxes on the buoyancy a function of surface
137 temperature (Roquet et al. 2022; Caneill et al. 2022). This means that freshwater and heat forcings
138 can not be treated as independent forcings separately acting on the ocean buoyancy, but should be
139 seen as coupled. This coupling is also strong during sea ice formation, as the heat loss generates a
140 salt flux in the ocean, and in the subtropics, as the evaporation generates a surface cooling.

141 Southern wind-driven upwelling has been invoked as the major driver of the return path of the deep
142 water, closing the overturning cells and explaining why the global overturning is deep (Marshall
143 and Speer 2012; Cessi 2019). However, this interpretation has recently been challenged by Klocker
144 et al. (2023) who found that a deep-reaching overturning can be generated in the absence of Ekman
145 pumping and with moderate interior mixing, as long as there is a differential heating on either
146 side of a Southern Ocean-like open channel. In parallel, Miller et al. (2020) showed that Southern
147 wind-driven upwelling alone fails to explain the mid-depth exponential ocean stratification, arguing
148 that boundary-intensified mixing is still needed to explain observations. While these studies do not
149 conflict with the idea that Southern wind-driven upwelling is an important control of the global
150 overturning circulation, it calls for a more nuanced view on its overall role.

151 The geostrophically balanced mesoscale eddies are omnipresent in the ocean and mix properties
152 along isopycnals. This eddy-induced isopycnal mixing becomes most pronounced in eddy-rich

regions, as e.g. the western boundary currents and in frontal regions such as the Antarctic Circumpolar Current (ACC). Climate models and coarse ocean models do not resolve mesoscale eddies, and consequently still rely on eddy parameterizations, usually involving a diffusion tensor consisting of an advective (antisymmetric) and a diffusive (symmetric) part. The advective part is often parameterized with the Gent–McWilliams (GM) closure (Gent et al. 1995), whereas the diffusive part of the mixing tensor can be interpreted as turbulent mixing of properties along isopycnal surfaces, and is often parameterized with the isotropic, so-called Redi scheme (Redi 1982) involving a single scalar isopycnal diffusivity. While the role of the eddy-advective part is well studied, the role of eddy-diffusive part has not garnered much attention. Recently, however, Chouksey et al. (2022) have shown the sensitivity of ocean circulation to isopycnal diffusivity using a non-eddy-resolving global ocean configuration and changing isopycnal diffusivity along with different formulations of surface forcing. The results in Chouksey et al. (2022) show that the strength of isopycnal diffusion not only controls the uptake of carbon and other passive tracers, but is also as important as the surface fluxes in shaping the circulation in the ocean. The results emphasize that both the advective and diffusive eddy effects are crucial to the large scale ocean circulation.

169 *b. Regional scale overturning circulations*

170 Ocean basin-scale observations of bottom-enhanced dissipation rates ε (Polzin et al. 1997; 171 Waterhouse et al. 2014) signaled that a one-dimension description of the open-ocean water masses 172 transformations could not explain the upwelling branch of the global overturning circulation as 173 previously thought (Munk 1966; Ferrari et al. 2016). Resolution of this conundrum requires 174 rejecting standard eddy diffusive closures for internal wave-driven boundary layers and-or properly 175 accounting for complex topography in watermass budgets (Polzin and McDougall 2022).

176 A recent observational effort within a canyon on the eastern slope of the Rockall Trough (the 177 Boundary Layer Turbulence and Abyssal Recipes Experiment) set out to directly observe the 178 near-bottom water-mass transformation using dye and tracer releases and a suite of microstructure 179 profilers and moored turbulence measurements. The dye, released ~ 10 m above the seafloor, 180 experienced rapid diapycnal upwelling equivalent to $O(100)$ m day $^{-1}$ (Wynne-Cattanach et al. 181 2024). However, tidally resolving observations from microstructure profilers do not find mixing

182 efficiencies nor stratification that decrease to zero as required by 1D models (Alford et al. 2024).
183 The 3D internal tide breaking processes, captured by moored estimates of diapycnal fluxes Polzin
184 et al. (2024), appears to be of vital importance to the observed upwelling.

185 The Samoan Passage in the tropical South Pacific is one of the major global channel flows
186 connecting ocean basins. As hypothesized for other abyssal ocean passage flows (Bryden and
187 Nurser 2003), it is a site of intense water mass transformation (Alford et al. 2013a) and thus
188 intricately links turbulent mixing processes with topography (Voet et al. 2015; Carter et al. 2019)
189 and large-scale meridional density gradients. Recent work has shown the importance of topographic
190 lee waves for the detailed dynamics of turbulent mixing via momentum redistribution at one of the
191 major Samoan Passage sills (Voet et al. 2023).

192 The earliest works on exchange flows of water masses (Knudsen 1900) as well as water mass
193 transformation and circulation (Walin 1977) were developed for the Baltic Sea, a large non-tidal
194 fjord-type estuary. The overturning circulation in the Baltic (the *Baltic Haline Conveyor Belt*;
195 Döös et al. 2004) consists of sporadic bottom salt water inflows from the North Sea into the deep
196 basins, freshwater supply due to major rivers and net precipitation, and the outflow of brackish
197 surface waters created by mixing. Local hotspots of this diahaline overturning circulation have
198 been recently investigated by Henell et al. (2023).

199 c. Estuarine circulation and river plumes

200 At the meeting, there were intense discussions about the different characteristics of mixing and
201 exchange flow between estuaries and larger ocean scales. Here, the estuarine and river plume
202 perspective is briefly given.

203 Estuarine circulation is commonly thought of as inflow of salty water and outflow of brackish
204 water into semi-enclosed coastal water bodies, generally driven by freshwater inflow from rivers
205 (MacCready and Geyer 2010; Geyer and MacCready 2014). The brackish water outflow is a product
206 of turbulent mixing between the river water and the inflowing ocean water, such that estuaries can
207 be characterised as *mixing machines* Wang et al. (2017). Estuarine exchange flow may also be
208 viewed in relation to an isohaline rather than to a fixed transect (Li et al. 2022; Reese et al. 2024),
209 a procedure that gives a horizontal rather than a vertical view.

210 Estuaries (including fjords) discharge fresher waters into their coastal ocean (Hetland 2005)
211 typically forming river plumes downstream. Mixing processes and water mass transformation
212 within river plumes control how estuarine waters ultimately enter the ocean. Wind and tides advect
213 and mix river plumes. In shallow frictional river plumes, such as the Rhine river plume, mixing
214 is further dominated by tidal straining and wind straining (de Boer et al. 2006, 2008; Fischer et al.
215 2009; Rijnsburger et al. 2016).

216 Moreover, tidal plume fronts (TPF's) generate internal waves ahead of them that can lead to
217 substantial mixing, e.g. in the near to mid-field of the Columbia river plume (Kilcher and Nash
218 2010). Recently Rijnsburger et al. (2018, 2021a,b) showed trapping of TPF's and the presence
219 of Internal Solitary Waves (ISW's) ahead of TPF's in the near to mid-field plume of the Rhine
220 river plume and their significant role in dispersion and transport of freshwater, in agreement with
221 Horner-Devine et al. (2015). It is important to understand the contribution of TPF's and ISW's to
222 mixing in river plumes.

223 *d. Interior mixing by wave-eddy and wave-wave interactions*

224 Small-scale turbulent diapycnal mixing caused by breaking internal waves contribute to main-
225 taining the large-scale overturning circulation (Talley 2013; Cimoli et al. 2023). The understanding
226 of this process breaks down to: (i) A forcing problem, detailing the large-scale sources of internal
227 wave field, and (ii) A dynamical problem, describing the energy cascades from internal waves to-
228 ward small scales until it breaks by shear or convective instabilities (Staquet and Sommeria 2002)
229 and turns into turbulent kinetic energy (TKE).

230 A major forcing is provided by the interaction of tides or non-oscillatory flows with the sea-floor
231 (Musgrave et al. 2022), generating internal tides or lee waves, respectively, with an energy input
232 of about 1 TW (Egbert and Ray 2000; Nycander 2005) or 0.2-0.75 TW (Nikurashin and Ferrari
233 2011; Wright et al. 2014). Another major forcing is the wind stress fluctuations at the ocean
234 surface (Thomas and Zhai 2022), forcing near-inertial waves with about 0.3 TW globally (Alford
235 2001; Rimac et al. 2013; Alford 2020; von Storch and Lüschow 2023). Additional mechanisms
236 include spontaneous loss of balance (Vanneste 2013; Alford et al. 2013b; Chouksey et al. 2018)
237 and shear instabilities of the balanced flow (Chouksey et al. 2022), with the global input from the
238 former estimated at 0.3 TW from idealized studies (Brüggemann and Eden 2015; Sugimoto and

239 Plougouven 2016), although a precise representation of these processes in realistic ocean models
240 remains a challenging feat and poses several open questions.

241 Recent progress on (i) involves estimating the direction of the barotropic-to-baroclinic energy
242 transfer (Pollmann and Nycander 2023) and the consistent representation of lee wave drag in ocean
243 general circulation models (Eden et al. 2021). The dynamical problem (ii) involves nonlinear wave-
244 wave interactions, described via the internal wave kinetic equation (Olbers 1976; Müller et al.
245 1986) and is the basis for the semi-empirical finescale parameterizations (Polzin et al. 2014). These
246 parameterizations have recently found solid theoretical grounding (Eden et al. 2019b; Dematteis
247 et al. 2023) and apply quite ubiquitously in the ocean interior (in an averaged sense) (Waterhouse
248 et al. 2014; Whalen et al. 2015; Pollmann et al. 2017), but are known to be incorrect in certain
249 regions (Waterman et al. 2014) and have yet to be applied in regions with strong eddies or currents
250 (Dematteis et al. 2023; Lozovatsky et al. 2023).

251 Addressing wave-eddy interactions, Müller and Briscoe (2000) and Natarov and Müller (2005)
252 proposed the Internal Wave Action Model (IWAM) based on the radiative transfer equation (Müller
253 and Olbers 1975) within the WKBJ approximation, but without practical applicability. Olbers and
254 Eden (2013) proposed IDEMIX (Internal wave Dissipation, Energy and MIXing) as a parame-
255 terized solution applicable to ocean and atmosphere (Quinn et al. 2020) including wave-mean
256 flow interactions (Olbers and Eden 2017; Eden and Olbers 2017). Building on IWAM, Sebas-
257 tia Saez et al. (2024) introduce the Internal Wave Energy Model (IWEM) to resolve the spectral
258 energy balance between waves and meso-scale eddies (with depth-varying stratification) in its
259 full form. They show a nontrivial wave energy gain outside and decay within the eddy due to
260 wave-eddy interactions, and enhanced wave dissipation by vertical refraction at eddy rims with
261 vertical diffusivities of $O(10^{-7})$ to $O(10^{-5}) \text{ m}^2/\text{s}$. Further, scale-separated stationary geostrophic
262 flow effects on small-amplitude waves are studied using a diffusion equation (Kafiabad et al. 2019;
263 Cox et al. 2023), whereas Savva et al. (2021) derived a kinetic equation using the Wigner transform
264 for resonant interactions under constant stratification. Both approaches reveal an eddy-induced
265 “diffusion of wave action” towards larger wavenumbers at constant frequency, enhancing TKE
266 production and mixing. From weak turbulence theory (e.g. Nazarenko (2011)), Eden et al. (2019a)
267 derive the kinetic equation under the weak interaction assumption (Hasselmann 1966) for mixed
268 Rossby–gravity wave–wave interactions. Recently, Olbers et al. (2024) show that energy cannot be

269 exchanged by resonant but by non-resonant wave-eddy interactions, provided the wave and eddy
270 fields show horizontal anisotropy. The question of energy exchange between internal tides and
271 mesoscale eddies, lying at the brink of the WKBJ limit, remains open.

272 **4. Discussion**

273 During the meeting, there were vivid discussions about the relation between mixing and over-
274 turning circulation. It was concluded that the question whether mixing drives the overturning
275 circulation, is that of a chicken-and-egg problem: Both are interdependent and can only exist
276 together. Small-scale tracer mixing requires the resupply of tracer gradients by the large-scale
277 circulation, and the resulting dia-surface fluxes are required to close the overturning circulation.

278 Besides giving answers, some general questions were formulated, among others:

- 279 • What are the roles of wave-wave and wave-eddy interaction and nonlinearities of the Equation
280 of State on mixing and overturning circulation?
- 281 • What is the role of small-scale mixing in the bottom boundary layer of the ocean on the
282 large-scale overturning circulation and where are the mixing hotspots?
- 283 • Where and when do the most important mixing events occur in estuaries and river plumes and
284 on which time scales does this contribute to exchange flow?
- 285 • Can we directly quantify the relation between mixing and overturning circulation on different
286 scales?

287 The presentations and discussions during the 11th Warnemünde Turbulence Days contributed
288 to all these questions and brought new insights that we have reported about in this short meet-
289 ing summary. By doing so, a wide range of scales and processes was bridged. More infor-
290 mation about the meeting including the programme, all abstracts and photos can be found at
291 <https://www.io-warnemuende.de/wtd-2023.html>.

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