ISPy3 - Integrated-light Spectroscopy with Python 3

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1 Introduction

ISPy3 is a collection of Python routines that can be used to model the integrated-light spectra of stellar populations. The actual spectral synthesis and related tasks (setting up model atmospheres, etc) are done via calls to external codes. Currently, the Kurucz codes (ATLAS/SYNTHE) and MARCS/TurboSpectrum are supported, but it should be relatively straight forward to implement other similar codes. As such, ISPy3 requires these external codes as well as their input data (atomic and molecular line lists, opacity distribution files, etc) to be installed in addition to the Python files.

The original motivation for developing ISPy3 was to determine chemical abundances from integrated-light spectra of globular clusters, and the package includes high-level routines that can be used to determine the abundances of specified elements via fits to observed spectra. However, the package can also be used more generally to compute model atmospheres and synthetic spectra via Python wrapper routines for the model atmosphere and spectral synthesis codes. In addition, ISPy3 includes utility functions for a variety of other purposes, such as computing equivalent widths of individual spectral lines (using a modified version of the Kurucz WIDTH9 programme) and generating opacity distribution functions (ODFs) for ATLAS9.

The emphasis in the present document is on practical aspects of setting up and running the Python code. The ATLAS/SYNTHE and MARCS/TurboSpectrum models and codes are documented elsewhere, and that documentation should be consulted for more information. This document is *not* intended as a detailed guide to these external codes and their various features, caveats, and idiosyncracies. An earlier version of ISPy3 was introduced in Larsen et al. (2012), which may be consulted for an example of a practical use case, as well as a more in-depth discussion of the input data and a number of tests of the performance of the code. Additional references are Hernandez et al. (2017, 2018) and Larsen et al. (2014, 2017, 2022). Before proceeding it is worth repeating the warning given on the Kurucz website, which also applies here:

Neither the programs nor data are "black boxes". You should not be using them if you do not have some understanding of the physics and of the programming in the source code.

The code comes with no warranty and almost certainly contains bugs that may compromise your science. It is the responsibility of the user to carry out the necessary sanity checks to verify that the output makes sense. If you believe that you have found a bug, or have suggestions for improvements, please do not hesitate to get in touch!

2 Installation

The ISPy3 Python tools require a basic Python 3 installation, including the numpy and scipy packages. To install the Python files, simply download and untar the ispy3.tar.gz file:

>> tar xvzf ispy3.tar.gz

and make sure that the directory in which the package is installed is included in your Python search path. The Python files will not by themselves be very useful without the external executables, which ISPy3 expects to find in the directory defined by the absetup.binpath variable. Various data files are also required (line lists, ODFs, etc.), whose location is specified in the absetup.catpref variable. The module absetup.py contains a few additional top-level options for configuration of ISPy3. See Section 6.

Be warned that ISPy3 will generally trust that the ancillary data files and external executables are present in the correct format in the expected locations. If this is not the case, the code will typically crash ungracefully without necessarily producing very informative error messages. Worse still, it may run, but produce meaningless output.

3 General overview and prerequisites

Neither the Kurucz nor TurboSpectrum codes make use of modern multi-core architectures, but the problem of computing integrated-light spectra nevertheless lends itself well to making use of parallel processing capabilities. In ISPy3, an integrated-light model spectrum is computed by dividing the underlying Hertzsprung-Russell Diagram (HRD) of the population into a number of bins (typically 50-100) and computing model atmospheres and synthetic spectra for each bin. The individual model spectra are then co-added:

$$L_{\rm SSP}(\lambda) = \sum_{i=1}^{n} w_i L_i(\lambda) \tag{1}$$

where $L_{\rm SSP}(\lambda)$ is the luminosity of the resulting "simple stellar population" (SSP) at wavelength λ , $L_i(\lambda)$ the luminosity of a star in the *i*th bin, and w_i a specified set of weights. The weights are typically given by the stellar mass function $\xi(M) \equiv {\rm d}N/{\rm d}M$:

$$w_i = \xi(M_i) \Delta M_i \tag{2}$$

for a bin covering a range ΔM_i of stellar masses. As each bin is independent of the others, the individual model spectra $L_i(\lambda)$ can be computed in parallel. The main limitations on the number of parallel processes are set by the available number of CPU cores and disc space for storage of temporary files. In principle, memory might also be a limitation but in practice none of these codes are particularly memory intensive by modern standards.

In general, the Kurucz SYNTHE spectral synthesis code is used together with ATLAS model atmospheres, while TurboSpectrum is used with MARCS model atmospheres (recent versions of TurboSpectrum, 2019 and later, can also read ATLAS model atmospheres). Both combinations have pros and cons: The ATLAS atmosphere models are computed via calls to the actual ATLAS9 or ATLAS12 codes, and can thus in principle be computed for any desired set of physical parameters (composition, gravity, temperature, etc). Various details of the models can be customised, such as the number of atmospheric layers, the range of optical depths, and microturbulence. The spectral synthesis code, SYNTHE, also supports rotational broadening of the model spectra. A drawback is that the ATLAS models only support plane-parallel geometry. The MARCS models

are available for spherical geometry, but as the MARCS code is not publically available the parameter range is restricted to what can be obtained by interpolation in the precomputed model grids available via the MARCS website.

It should be emphasised that the computation of model atmospheres and the corresponding synthetic spectra are separate steps; both SYNTHE and TurboSpectrum can compute model spectra for compositions that do not necessarily match those of the atmospheres in detail. In fact, there is nothing to stop the adventurous user from computing a model spectrum with a completely different composition than that of the model atmosphere, although this is not physically consistent and the usefulness of such an experiment may be rather questionable.

Apart from these and various other technical differences in the detailed implementations of the codes, the atmosphere and spectral synthesis models supported by ISPy3 are classical models with the usual standard assumptions: the models are one-dimensional (i.e. the only spatial variable is the depth) and static, and they generally assume Local Thermodynamic Equilibrium (LTE). As any good textbook on stellar atmospheres will stress, none of these assumptions are physically realistic, but they are nevertheless very convenient computationally, and make it feasible to compute large numbers of models in a reasonable amount of time. From version 20, TurboSpectrum is able to compute spectra in NLTE. This is supported in ISPy3 at a somewhat experimental level and requires a number of additional input files, in particular a grid of pre-computed departure coefficients for the MARCS models that requires several hundred Gb of storage space.

3.1 The Kurucz codes: ATLAS, SYNTHE, etc.

As noted above, the ATLAS models come in two flavours. ATLAS9 uses pre-defined opacity distribution functions (ODF) to calculate the opacities, whereas ATLAS12 calculates opacities via direct opacity sampling (OS). The OS method allows the calculation of atmosphere models for any specified composition, but is computationally much more expensive. In addition, it can sometimes be difficult to get ATLAS12 models to converge, especially for high surface gravity models with low temperatures ($T_{\rm eff} < 4000\,{\rm K}$). The ATLAS9 models are much faster to compute and tend to converge more easily, but are restricted to compositions for which ODFs are available. In ISPy3, ODFs are included for solar-scaled and α -enhanced composition and metallicities $-3.50 \le {\rm [Fe/H]} \le +0.50$, but ODFs for other compositions can be computed with the kurucz.calcodf() function.

ISPy3 computes ATLAS models for any specified combination of physical parameters ($T_{\rm eff}$, log g, composition) by calling the corresponding Kurucz codes. An ATLAS9 model can be computed from scratch, or starting from a pre-existing model with physical parameters close to the desired ones. The default behaviour is to look for the closest matching model among those available in the pre-computed grid by Castelli and use that as a starting guess. Alternatively, an initial model can be specified directly when calling the kurucz.mkatm() function. The recommended procedure is to start from a pre-existing model, as this usually ensures faster convergence. An ATLAS12 model must always be computed starting from a pre-existing model, usually computed with ATLAS9. Again, the default behaviour is to look for a Castelli model, but a starting model can also be specified directly when calling kurucz.mkatm_a12().

The final computed models will be stored with a filename ending with .A9 or .A12. In addition, a file ending with .out is produced, containing diagnostic output from the ATLAS9 or ATLAS12 codes. At the end of this file, some extra information about each layer in the model is listed, with the last two columns (the errors on the flux and on the flux derivative) being useful for assessing whether the model has properly converged. If this information is not found at the end of the .out file, it means that the procedure crashed for some reason.

Spectral synthesis based on ATLAS models is normally done with the SYNTHE code. Assuming that an ATLAS (9 or 12) model is available, the first step is to call syntherk.synbeg(). This will define the wavelength range and spectral resolution, as well as the elements and molecules to include in the model spectrum, whether the predicted lines should be included, etc. It is possible to specify explicitly which molecules and elements to include in the spectral synthesis; the default is to include everything. Line data for the specified wavelength range will then be read and stored in temporary files. syntherk.synbeg() will produce two versions of these files, with and without the TiO lines as these can be skipped for temperatures higher than 4500 K, which makes the calculations much faster. The actual model spectrum is then calculated with a call to syntherk.synthespec() which automatically selects the appropriate input files, depending on the temperature.

3.1.1 Obtaining and installing the Kurucz codes

The Kurucz suite includes quite a large number of programmes to carry out various tasks related to the computation of model atmospheres and synthetic spectra. The primary reference for the source codes, line data, etc., is Kurucz' website at http://kurucz.harvard.edu where the original (VMS) versions of the Fortran codes can be found. Versions of the source code suitable for compilation under Linux and macOS are available at the website of F. Castelli: http://www.ser.oats.inaf.it/castelli/. The Intel Fortran compiler (ifort) is required to compile the Fortran source code and is available for free at https://www.intel.com/content/www/us/en/developer/articles/tool/oneapi-standalone-components.html#fortran. The Castelli website has versions of the codes that can be compiled with gfortran and also differ in other ways (e.g. carrying out some calculations in double instead of single precision), but I haven't tested these extensively. I also have my own modified versions of the Kurucz source codes (based more directly on the Sbordone/Castelli ifort versions) that can be compiled with gfortran and seem to give (very nearly) identical results to the ifort versions. These can be provided upon request.

Once the codes have been compiled, they should be copied to the directory absetup.binpath, where ISPy3 expects to find them. Below is an overview of the various Kurucz codes.

3.1.2 Kurucz codes for calculating model atmospheres

atlas12.exe ATLAS12 code for calculation of model atmospheres for arbitrary composition and microturbulence with the opacity sampling method. I am using a slightly modified version of the code than that distributed by Sbordone / Castelli, in which the maximum number of atmospheric layers has been increased to 99 (instead of 72) and the TCORR subroutine has been slightly modified so that it is similar to the one in atlas9mem.for. The latter

seems to alleviate convergence problems with cool, high surface gravity atmospheres in some cases.

atlas9mem.exe ATLAS9 code for calculation of model atmospheres with pre-computed Opacity Distribution Functions.

kappa9.exe Computes Rosseland mean opacities from ODFs.

kapreadts.exe Rearranges output from kappa9.exe for ATLAS9.

dfinterpbig.exe Interpolates in pre-computed ODFs to make an ODF for the appropriate metallicity.

3.1.3 Kurucz codes for spectral synthesis

rgfalllinesnew.exe Read atomic line data.

rmolecasc.exe Read molecular line data from ASCII files. The most recent version of this programme, available from the Kurucz web site, is required in order to read data for MgO, CaO, NaH, KH, CaH, and CrH. However, the Kurucz version needs to be slightly modified to write the data in a format that is compatible with the Castelli/Sbordone distributions of SYNTHE (the output must include the 'alpha' exponent for the ABO-style van der Waals damping, even though alpha is always 0 for molecular data).

rpredict.exe Read data for predicted lines.

rh2ofast.exe Read binary H₂O line data.

rschwenk.exe Read binary TiO line data (Schwenke).

xnfpelsyn.exe Precomputes continuum opacities and number densities for various elements.

synbeg.exe Initialises the spectral synthesis procedure - defines wavelength range and spectral resolution.

synthe.exe Calculates the line opacity data for each line in the spectrum.

spectrv.exe Calculates continuum opacities and final synthetic spectrum.

rotate.exe Adds rotational broadening to the computed spectrum (optional).

converfsynnmtoa.exe Converts the binary output files (from spectrv.exe or rotate.exe)
to ASCII format.

3.1.4 Kurucz codes required for calculation of equivalent widths

The WIDTH9 code is quite versatile and can be used in several different ways, depending on the input cards provided. ISPy3 uses the inpwidth.for code written by F. Castelli to produce the input files needed by WIDTH9 to compute equivalent widths, consisting of the control cards (including a model atmosphere) and line list in a modified format.

inpw_sl.exe Slightly modified version of Castelli's **inpwidth.for** code, which only needs a single line list as input and allows the wavelength matching tolerance for line identification to be specified by the user. The source code has also been modified to be compatible with **qfortran**.

width9sl.exe Modified version of the Kurucz WIDTH9 code. The modifications are as follows: 1) a curve of growth can be computed for up to 99 abundance steps (rather than the standard 9), 2) The continuum flux will be written to the output file (needed by abutils.hrd2ew()), 3) The new VOIGT subroutine by Castelli has been replaced by the original Kurucz subroutine.

3.1.5 Kurucz codes required for calculation of ODFs

The four programmes listed below are called by the calcodf() function when computing new ODFs to be used with ATLAS9. The procedure is described in Castelli (2005).

xnfdf.exe: Calculates the atomic and molecular number densities for all combinations of temperature and gas pressure for a given chemical composition.

dfsynthe.exe: This is the programme that actually computes the ODFs for a given composition. The output is a set of opacity tables for microturbulence velocities of 0, 1, 2, 4, and 8 km/s and for 57 temperatures between 1995 K and 199526 K. Each file contains opacities for a range of gas pressures (25 different values). However, the output is not yet in a format that can be directly used by ATLAS9.

dfsortp.exe, **separatedf.exe**: These two programmes rearrange and merge the output from dfsynthe.exe into one ODF file per chemical composition and microturbulence velocity that can be used as input for ATLAS9.

3.1.6 Miscellaneous

avgrot.exe: Computes up to 25 SYNTHE spectra with different $v \sin i$ values. This programme is not part of the Kurucz suite but is distributed together with the ISPy3 Python code.

3.1.7 Input files for Kurucz codes

Line lists for SYNTHE are available at the Kurucz website, where descriptions of the formats can also be found. The locations of the line lists are defined in the syntherk.py module in the variables syntherk.atomdir (for the atomic lines) and syntherk.moldir (for the molecular

lines). There is a separate directory for the predicted lines, syntherk.preddir, which are only included in the calculations if specifically requested by setting the argument predicted=True when calling syntherk.synbeg(). The line lists are, of course, a critical ingredient for the modelling of any spectrum, and the raw line lists from the Kurucz website are far from perfect.

ISPy3 expects the full atomic line list to be split up into smaller chunks. These files can be produced from the full line list with the programme splitgf100.for:

gf0150.100: Wavelengths < 1500 Å

gf0200.100: $1500 \text{ Å} < \lambda < 2000 \text{ Å}$

gf0300.100: 2000 Å $< \lambda < 3000$ Å

 $\mathbf{gf0400.100} : 3000 \,\text{Å} < \lambda < 4000 \,\text{Å}$

gf0500.100: $4000 \text{ Å} < \lambda < 5000 \text{ Å}$

gf0600.100: 5000 Å $< \lambda < 6000$ Å

gf0800.100: $6000 \,\text{Å} < \lambda < 8000 \,\text{Å}$

gf1200.100: $8000 \text{ Å} < \lambda < 12000 \text{ Å}$

gf3000.100: 12000 Å $< \lambda < 30000$ Å

SYNTHE supports the use of "ABO" (Anstee, Barklem, & O'Mara) damping constants for atomic lines, although these are not included in the standard line lists available at the Kurucz website and there is little documentation for this feature.

The molecules listed below are currently supported. The file names are defined in the list syntherk.molfiles, while the molecules are defined in syntherk.molecule_ids. Other molecules or different line lists can be included by editing these variables. For example, the Kurucz website also has a single large file that includes all diatomic molecules except TiO (diatomics.asc).

CH: chmasseron.asc (from Kurucz website, dated 2-Aug-2015)

MgH: mgh15082018.asc (from Kurucz website, dated 15-Aug-2015)

NH: nhaxCas.asc, nhcaCas.asc (the Kurucz website now has a single file, nh.asc, dated 11-Apr-2014)

OH: ohupdate.asc (from Kurucz website, dated 10-Apr-2014)

SiH: sihaxsightly12012018.asc, sihxxsightly12012018.asc (from Kurucz website, dated 12-Jan-2018)

H2: h2.asc (dated 11-Apr-2014), h2xx.asc (from Kurucz website, dated 17-Jun-2015)

C2: c2ax.asc, c2ba.asc, c2dabrookek.asc, c2ea.asc

CN: cnaxbrookek.asc, cnbxbrookek.asc, cnxx12brooke.asc

CO: coax.asc, coxx.asc

SiO: sioax.asc, sioex.asc, sioxx.asc

CaH: cah.asc

FeH: fehfx.asc

VO: voax.asc, vobx.asc, vocx.asc

TiO: tioschwenke.bin, eschwenke.bin: From Castelli website. The file with the energies, eschwenke.bin, can be produced from several ASCII files and a programme available at the Kurucz website.

H2O: h2ofastfix.bin (from Kurucz website, dated 09-feb-2012).

In addition to the line lists, the following files are needed (all directories given relative to absetup.catpref):

cats/Castelli2015/molecules.dat: List of data for molecules

cats/Castelli2015/pfiron.dat: Partition functions for iron-group elements (other
partition functions are hard-coded)

cats/Castelli2015/continua.dat: Data for continuum opacity calculations

cats/Castelli2015/linelists: Predicted lines and other line lists for ATLAS12.

cats/Castelli2015/dflines/*.bin: Line lists for computation of ODFs. These can be generated from the files in cats/Castelli2015/linelists with the programmes repack*.exe (available at the Kurucz and Castelli websites)

cats/Castelli2015/odfsl/*.bdf: ODFs computed for scaled-solar and alpha-enhanced composition in 0.25 dex steps of metallicity.

cats/Castelli2015/atlas9/*odfnew.dat: Pre-computed ATLAS9 models for scaled-solar and alpha-enhanced composition, computed by Castelli. These can be used as initial guesses when computing new ATLAS9 or ATLAS12 models.

cats/Kurucz/colors/*.dat: Theoretical colours, from Castelli & Kurucz (2003). These are used when deriving stellar parameters from photometry.

3.2 Marcs atmospheres / TurboSpectrum

ISPy3 uses the <code>interpol.modeles</code> programme (written by T. Masseron) to generate MARCS models for any specified combination of $T_{\rm eff}$, $\log g$, [Fe/H] and [α /Fe] by interpolation in the pre-computed model library available via the MARCS website (https://marcs.astro.uu.se/search.php). ISPy3 uses three grids of MARCS models: "alpha-poor" ([α /Fe] = 0 for $-2.50 < [{\rm Fe/H}] < -0.25$), "alpha-enhanced" ([α /Fe] = +0.4 for $-0.75 < [{\rm Fe/H}] < +1.0$, and "standard composition", which has scaled-solar composition at high metallicities and alpha-enhanced composition at lower metallicities. The default behaviour is to automatically select models from the appropriate grid, based on the requested composition. Spherical models are used for low surface gravities ($\log g < 3$) and plane-parallel models otherwise. The model grid covers temperatures in the range 2500 K $\leq T_{\rm eff} \leq 8000$ K, surface gravities $-0.5 \leq \log g \leq 5.0$, and metallicities $-5.0 \leq [{\rm Fe/H}] \leq +0.75$. There are nearly 3000 spherical models with a microturbulent velocity of 2 km/s and a mass of 1 M $_{\odot}$, and 2400 plane-parallel models with a microturbulence of 1 km/s. The webform on the MARCS website will not allow all of these to be downloaded at once, so it is necessary to narrow down the search to a more limited set of parameters per request.

For temperatures in the range $3800 \, \text{K} \le T_{\text{eff}} \le 7000 \, \text{K}$ an "optimised" (non-linear) interpolation scheme is used, while linear interpolation is used for temperatures outside this range. Details can be found in Masseron (2008). If a model outside the grid is requested, the closest point on the grid will be used - hence, it is not wise to request models with parameters that lie much outside those included in the grid. User discretion is advised!

For spherical geometry, the MARCS grid contains models for several different masses (0.5, 1, 2, 5, and 15 M_{\odot}); the physical radius of curvature then follows once the surface gravity is specified. ISPy3 uses the models computed for a mass of $1 M_{\odot}$ (for which the grid is most complete), and while this is probably appropriate for modelling of giants in old globular clusters, it may not be for cool supergiants in younger populations.

To compute synthetic spectra from the MARCS models, the TurboSpectrum spectral synthesis code (Alvaro & Plez 1998; Plez 2012) is used. The abundances in the model spectra are not restricted to those in the input atmosphere models. TurboSpectrum is available at Github (https://github.com/bertrandplez/) and extensive molecular line lists can be downloaded from the website of B. Plez (https://www.lupm.in2p3.fr/users/plez/). Recent versions of TurboSpectrum (since 2019) can read ATLAS model atmospheres, and NLTE spectral synthesis is possible since version 20. However, older versions (2014 or 2015) are faster. ISPy3 can be configured to use different versions of the TurboSpectrum executables, depending on whether MARCS or ATLAS models are used, or if NLTE spectral synthesis is desired.

TurboSpectrum uses a common file format for both atomic and molecular line lists, which is different than those used by the Kurucz codes. Atomic line lists in the Kurucz format can be converted to the TurboSpectrum format with the kur2bsyn.py utility and molecular line lists in the Kurucz ASCII format can be converted with the kmol2bsyn.py utility. This way, consistency between SYNTHE and TurboSpectrum can be achieved. Atomic line lists for TurboSpectrum can also be downloaded from the VALD website (http://vald.astro.uu.se/~vald/php/vald.php) and put in the right format with utilities available at the Plez website. The TiO line list is a somewhat special case, as SYNTHE uses the Schwenke line list in binary format, whereas

TurboSpectrum instead uses the line list from the VALD database (Plez 1998, updated version from January 2012).

The ISPy3 interface for TurboSpectrum uses a similar sequence of initialisation and computation as SYNTHE, although the initialisation procedure is internally simpler for TurboSpectrum. The equivalent of the synbeg() function is called turbospec.turboinit(), but does not involve execution of any external code and simply copies some files to the temporary directory. The synthetic spectra are then computed with a call to turbospec.turbospec().

3.2.1 TurboSpectrum NLTE

As of version 20, TurboSpectrum is able to carry out the spectral synthesis in NLTE for selected elements. The NLTE effects are modelled using grids of pre-computed NLTE departure coefficients. Since the departure coefficients are different for each transition and depend on the physical conditions in the atmosphere, they must be computed for each transition and for each layer in each model atmosphere. Moreover, departure coefficients are typically calculated for several values of the abundance of each element. Hence, it is easy to understand that the corresponding data files become quite large. In addition to the departure coefficients, TurboSpectrum NLTE requires the model atoms corresponding to each departure coefficient grid, as well as modified line lists that allow each transition to be identified with the corresponding states in the model atom. A special version of the interpol_modeles programme, interpol_modeles_nlte, is used to interpolate simultaneously in both the MARCS atmosphere and departure coefficient grids.

Once the required files are installed, the procedure for computing NLTE model spectra is mostly quite similar to that for regular TurboSpectrum models. Instead of calling marcs.intpatm() to interpolate in the model atmosphere grid, the function tsnlte.intpDC() is used to interpolate in both the model atmosphere and departure coefficient grids. In the call to turbospec(), NLTE is then enabled by setting the argument nlte=True.

3.2.2 External executables for MARCS/TurboSpectrum

These must be available in the directory absetup.binpath.

babsma_lu.14 Calculation of continuous absorption coefficients (2014 or 2015 version)

bsyn_lu.14 Spectral synthesis code (2014 or 2015 version)

babsma_lu.20 Calculation of continuous absorption coefficients (2020 version)

bsyn_lu.20 Spectral synthesis code that can read ATLAS models and compute NLTE spectra (2020 version)

The names of these executables are defined in turbospec.py via the variables babsma_marcs_lte, babsma_marcs_lte, babsma_marcs_lte, bsyn_marcs_lte, bsyn_marcs_lte. These variables specify which versions of the TurboSpectrum executables are to be used for LTE MARCS, ATLAS, and NLTE MARCS calculations.

interpol_modeles Code to interpolate in MARCS model grid, available from the MARCS website (https://marcs.astro.uu.se/software.php). Compiled with the optimize option set to .true. (i.e. optimised interpolation for models in the range 3800 - 7000 K).

interpol_modeles_lin Same as above, but compiled with optimize = .false. (linear interpolation for models cooler than 3800 K).

interpol_modeles_nlte Modified version of interpol_modeles that also interpolates in the departure coefficient model grids. Needed for TurboSpectrum NLTE. Available at https: //github.com/bertrandplez/Turbospectrum_NLTE.

3.2.3 Input files for MARCS/TurboSpectrum

All directories are given relative to absetup.catpref.

cats/Marcs/Alpha_enhanced/Spherical/*mod.gz: Spherical MARCS models, surface gravities $-0.5 \le \log g \le +3.5$, microturbulence of 2 km/s, $[\alpha/\text{Fe}] = +0.4$

cats/Marcs/Alpha_enhanced/Plane-parallel/*mod.gz: Plane-parallel MARCS models, surface gravities $+3.0 \le \log g \le +5.5$, microturbulence of 1 km/s, $[\alpha/\text{Fe}] = +0.4$

cats/Marcs/Alpha_poor/Spherical/*mod.gz: Spherical MARCS models, surface gravities $-0.5 \le \log g \le +3.5$, microturbulence of 2 km/s, $[\alpha/\text{Fe}] = 0.0$

cats/Marcs/Alpha_poor/Plane-parallel/*mod.gz: Plane-parallel MARCS models, surface gravities $+3.0 \le \log g \le +5.5$, microturbulence of 1 km/s, $[\alpha/\text{Fe}] = 0.0$

cats/Marcs/Standard_composition/Spherical/*mod.gz: Spherical MARCS models, surface gravities $-0.5 \le \log g \le +3.5$, microturbulence of 2 km/s.

cats/Marcs/Standard_composition/Plane-parallel/*mod.gz: Plane-parallel MARCS models, surface gravities $+3.0 \le \log g \le +5.5$, microturbulence of 1 km/s.

In turbospec.py: turbospec.atomdir and turbospec.moldir point to directories containing atomic and molecular line lists in the BSYN format. turbospec.plezdir points to a directory containing various input files for TurboSpectrum, as well as the line lists for CH and TiO. The DATA/ directory in the TurboSpectrum distribution should be copied as a subdirectory under this directory.

3.2.4 Input files for TurboSpectrum NLTE

TurboSpectrum NLTE needs additional input files, available at https://keeper.mpdl.mpg.de/d/6eaecbf95b88448f98a4/ where additional documentation can also be found. In particular, some of the departure coefficient grids are very large as they contain precomputed departure coefficients for each transition for every layer in each MARCS atmosphere for several compositions. ISPy3 expects to find the various files in the following directory structure:

cats/NLTE-TS/atmospheres/marcs_standard_comp/*.mod: The MARCS atmospheres corresponding to the NLTE departure coefficient grids. In principle they should be the same as those in the cats/Marcs/Standard_composition directory but are kept separate here to ensure consistency with the departure coefficient grids.

cats/NLTE-TS/dep-grids/<elem>/: The precomputed grids of departure coefficients for each atmosphere model. There is one subdirectory for each element. Each subdirectory contains an NLTEgrid file and a corresponding auxData file that describes the structure of the NLTEgrid file. The NLTEgrid files can be very large (tens of Gb each)!

cats/NLTE-TS/linelists/atoms/*: Line lists in the modified format required by TurboSpectrum NLTE.

cats/NLTE-TS/MODELATOMS/atom.*: Model atoms.

4 Basic examples

This section provides some examples to illustrate how ISPy3 can be used to compute stellar model atmospheres and synthetic spectra.

4.1 Computing an ATLAS9 model

To compute an ATLAS9 model for a red giant with an effective temperature of $T_{\rm eff} = 4200$ K, surface gravity $\log g = 2.0$, metallicity [m/H]=-0.5, and microturbulent velocity of 2 km/s, we use the following sequence of Python commands:

```
> from ispy3 import kurucz
> m, teff, logg, vturb = -0.5, 4200., 2.00, 2.0
> kurucz.mkodf(m, odfs='S')
> kurucz.mkatm(teff, logg, m, 't4200g200m050', vturb=vturb)
```

First, the kurucz module is imported and the stellar parameters defined. Next, kurucz.mkodf() is used to set up an ODF for the specified metallicity by interpolation in the pre-computed set of ODFs. The argument odfs='S' specifies that the solar-scaled ODFs will be used; the other option is 'A' (alpha-enhanced ODFs). The ODFs are computed for a fixed set of microturbulent velocities: 0 km/s, 1 km/s, 2 km/s, 4 km/s, and 8 km/s, all of which are stored in the working directory (odfbig0.bdf, odfbig1.bdf, ..., odfbig8.bdf). Finally, kurucz.mkatm() selects the interpolated ODF with the closest matching microturbulence and computes the model atmosphere. Since we haven't specified otherwise, the initial guess for the structure of the model will be selected among the pre-existing Castelli models, which must therefore be available. The temperature structure of the pre-existing model will then be scaled to the temperature of the requested model, and iterations started from the rescaled initial guess. The final model will be stored in an output file with the name t4200g200m050.A9, where the extension is automatically added. There will also be a diagnostic output file called t4200g200m050.out. The computation time for this model (on a 2010 Mac Pro workstation) was about 9 s, with an additional 15 s to interpolate in the ODFs (this has to be done only once for a given metallicity).

To compute a model from scratch, we set kurucz.A9_FROMSCRATCH = True. The procedure is otherwise identical:

```
> from ispy3 import kurucz
> m, teff, logg, vturb = -0.5, 4200., 2.00, 2.0
> kurucz.mkodf(m)
> kurucz.A9_FROMSCRATCH = True
> kurucz.mkatm(teff, logg, m, 't4200g200m050s', vturb=vturb)
```

To compare the two models, we can inspect the .out files. Of interest here are the two last columns in the model output at the end of these files, which list the percentage error on the flux (per layer) and the flux derivative, both of which should be small if the model has converged. An important point to be aware of is that the ATLAS code does not itself test for convergence. It simply carries out the specified number of iterations (the default is 15), and it is up to the user to verify that the model has indeed reached a satisfactory degree of convergence.



ATLAS9 model with initial guess from Castelli library:

				Č			•					
TEFF 0	4200.	LOG G 2	2.00000	[-0.5] ELECTRON] VTURB=2.0	L/H=1.25 N ROSSELAND	OVER NEW OF HEIGHT	F ROSSELAND	EDACTION	RADIATIVE	ITERATIO	N 15 ENT FLUX
U	RHOX	TEMP	PRESSURE	NUMBER	DENSITY	MEAN	(KM)	DEPTH		ACCELERATION	ERROR	DERIV
1 8.	.183E-03	2362.8	8.183E-01	5.226E+06		1.630E-05	0.000E+00	1.334E-07	0.000E+00	2.653E-03	0.004	2.905
	.087E-02	2388.7	1.087E+00	6.964E+06	6.862E-12	1.683E-05	4.528E+03	1.778E-07	0.000E+00	2.506E-03	0.004	4.353
	.426E-02	2412.0	1.426E+00	9.143E+06	8.917E-12 1.144E-11	1.816E-05	8.900E+03	2.371E-07	0.000E+00	2.349E-03	0.004	4.765
	.846E-02 .372E-02	2435.4 2459.8	1.846E+00 2.372E+00	1.186E+07 1.532E+07	1.144E-11 1.455E-11	1.950E-05 2.050E-05	1.309E+04 1.716E+04	3.162E-07 4.217E-07	0.000E+00 0.000E+00	2.186E-03 2.031E-03	0.004	3.079 2.027
	.048E-02	2485.7	3.048E+00	1.984E+07	1.850E-11	2.114E-05	2.127E+04	5.623E-07	0.000E+00	1.884E-03	0.004	1.210
	.914E-02	2512.1	3.915E+00	2.572E+07	2.351E-11	2.218E-05	2.542E+04	7.499E-07	0.000E+00	1.740E-03	0.004	0.204
	.999E-02	2538.6	4.999E+00	3.318E+07	2.972E-11	2.395E-05	2.951E+04	1.000E-06	0.000E+00	1.596E-03	0.004	-1.037
	.345E-02 .039E-02	2565.3 2592.5	6.345E+00 8.040E+00	4.262E+07 5.475E+07	3.733E-11 4.680E-11	2.552E-05 2.690E-05	3.355E+04 3.759E+04	1.334E-06 1.778E-06	0.000E+00 0.000E+00	1.472E-03 1.364E-03	0.004 0.004	-2.883 -5.161
	.019E-02	2620.5	1.019E+01	7.050E+07	5.871E-11	2.818E-05	4.169E+04	2.371E-06	0.000E+00	1.271E-03	0.004	-7.835
	288E-01	2649.3	1.288E+01	9.074E+07	7.338E-11	3.068E-05	4.578E+04	3.162E-06	0.000E+00	1.181E-03	0.004	-10.567
	.616E- 0 1	2678.0	1.617E+01	1.162E+08	9.111E-11	3.346E-05	4.979E+04	4.217E-06	0.000E+00	1.103E-03	0.004	-10.530
	.019E-01	2706.5	2.020E+01	1.482E+08	1.126E-10	3.624E-05	5.376E+04	5.623E-06	0.000E+00	1.038E-03	0.004	-10.274
	.517E-01 .132E-01	2735.0 2764.6	2.518E+01 3.133E+01	1.888E+08 2.408E+08	1.389E-10 1.711E-10	3.903E-05 4.226E-05	5.774E+04 6.172E+04	7.499E-06 1.000E-05	0.000E+00 0.000E+00	9.833E-04 9.389E-04	0.004 0.004	-11.367 -12.848
	.876E-01	2795.5	3.877E+01	3.071E+08	2.094E-10	4.734E-05	6.565E+04	1.334E-05	0.000E+00	9.009E-04	0.004	-10.547
	.764E-01	2826.0	4.765E+01	3.892E+08	2.545E-10	5.271E-05	6.949E+04	1.778E-05	0.000E+00	8.687E-04	0.004	-11.108
	.831E-01	2856.5		4.917E+08	3.082E-10	5.841E-05	7.330E+04	2.371E-05	0.000E+00	8.420E-04	0.004	-11.748
	.114E-01	2887.9	7.116E+01	6.214E+08	3.718E-10	6.483E-05	7.708E+04	3.162E-05	0.000E+00	8.226E-04	0.004	-13.220
	.636E-01 .042E+00	2921.5 2954.8	8.637E+01 1.042E+02	7.872E+08 9.918E+08	4.461E-10 5.323E-10	7.380E-05 8.363E-05	8.081E+04 8.447E+04	4.217E-05 5.623E-05	0.000E+00 0.000E+00	8.128E-04 8.039E-04	0.004	-10.598 -9.311
	251E+00	2988.4	1.251E+02	1.246E+09	6.317E-10	9.586E-05	8.807E+04	7.499E-05	0.000E+00	8.007E-04	0.004	-7.748
	.494E+00	3022.7	1.494E+02	1.562E+09	7.457E-10	1.100E-04	9.161E+04	1.000E-04	0.000E+00	8.006E-04	0.003	-10.617
	.773E+00	3057.4	1.773E+02	1.949E+09	8.745E-10	1.296E-04	9.505E+04		0.000E+00	8.122E-04	0.003	-7.380
	.089E+00	3090.8	2.089E+02	2.408E+09	1.019E-09	1.513E-04	9.840E+04	1.778E-04	0.000E+00	8.233E-04	0.003	-5.817
	.453E+00 .871E+00	3123.2 3154.8	2.452E+02 2.871E+02	2.957E+09 3.609E+09	1.184E-09 1.372E-09	1.754E-04 2.025E-04	1.017E+05 1.050E+05	2.371E-04 3.162E-04	0.000E+00 0.000E+00	8.358E-04 8.517E-04	0.002 0.002	-4.151 -2.899
	350E+00	3186.6	3.350E+02	4.392E+09	1.585E-09	2.388E-04	1.082E+05	4.217E-04	0.000E+00	8.798E-04	0.002	-1.898
	.890E+00	3217.0	3.889E+02	5.300E+09	1.823E-09	2.824E-04	1.114E+05	5.623E-04	0.000E+00	9.138E-04	0.002	-1.281
	.502E+00	3245.9	4.501E+02	6.343E+09	2.091E-09	3.306E-04	1.145E+05	7.499E-04	0.000E+00	9.504E-04	0.002	-0.825
	.202E+00 .007E+00	3273.5 3300.2	5.201E+02 6.006E+02	7.547E+09 8.941E+09	2.395E-09 2.744E-09	3.842E-04 4.441E-04	1.177E+05 1.208E+05	1.000E-03 1.334E-03	0.000E+00 0.000E+00	9.910E-04 1.037E-03	0.001 0.001	-0.521 -0.342
	.935E+00	3326.1	6.934E+02	1.056E+10	3.143E-09	5.147E-04	1.240E+05	1.778E-03		1.097E-03	0.001	-0.259
	.002E+00	3351.4	8.001E+02	1.243E+10	3.599E-09	5.971E-04	1.271E+05	2.371E-03	0.000E+00		0.001	-0.182
	.231E+00	3375.8	9.230E+02	1.458E+10	4.123E-09		1.303E+05	3.162E-03	0.000E+00	1.236E-03	0.001	-0.190
	.065E+01	3400.1	1.065E+03	1.707E+10	4.722E-09	7.981E-04	1.335E+05	4.217E-03	0.000E+00	1.326E-03	0.001	-0.195
	.228E+01 .417E+01	3423.8 3449.6	1.228E+03 1.417E+03	1.995E+10 2.342E+10	5.410E-09 6.194E-09	9.225E-04 1.069E-03		5.623E-03 7.499E-03	0.000E+00 0.000E+00	1.432E-03 1.555E-03	0.000 -0.002	-0.336 -1.790
	.634E+01	3474.7	1.634E+03	2.742E+10	7.092E-09	1.235E-03	1.433E+05	1.000E-02		1.691E-03	-0.007	-2.143
41 1.	.884E+01	3499.7	1.884E+03	3.205E+10	8.122E-09	1.430E-03	1.466E+05	1.334E-02	0.000E+00	1.851E-03	-0.012	-1.387
	.173E+01	3526.3	2.172E+03	3.755E+10	9.295E-09	1.658E-03	1.499E+05	1.778E-02	0.000E+00	2.040E-03	-0.016	-0.495
	.503E+01 .881E+01	3555.5 3588.5	2.503E+03 2.880E+03	4.423E+10 5.245E+10	1.062E-08 1.211E-08	1.931E-03 2.263E-03	1.532E+05 1.565E+05	2.371E-02 3.162E-02	0.000E+00 0.000E+00	2.266E-03	-0.017 -0.015	0.167 0.347
	.308E+01	3627.0		6.278E+10	1.376E-08	2.263E-03 2.677E-03	1.505E+05	4.217E-02	0.000E+00	2.537E-03 2.864E-03	-0.013	0.030
	.788E+01	3670.1		7.557E+10	1.557E-08	3.192E-03	1.631E+05	5.623E-02	0.000E+00	3.252E-03	-0.012	0.036
	.323E+01	3719.7		9.171E+10	1.753E-08	3.835E-03	1.664E+05	7.499E-02	0.000E+00	3.718E-03	-0.011	0.017
	.915E+01	3776.6	4.914E+03	1.121E+11	1.962E-08	4.634E-03	1.695E+05	1.000E-01	0.000E+00	4.276E-03	-0.011	0.011
	.566E+01 .281E+01	3842.5 3918.2	5.566E+03 6.280E+03	1.381E+11 1.710E+11	2.183E-08 2.414E-08	5.625E-03 6.845E-03	1.727E+05	1.334E-01 1.778E-01	0.000E+00 0.000E+00	4.925E-03 5.681E-03	-0.010 -0.009	0.004 0.000
	.064E+01	4005.1	7.063E+03	2.122E+11	2.655E-08	8.317E-03	1.789E+05	2.371E-01	0.000E+00	6.545E-03	-0.008	-0.003
	.927E+ 0 1	4104.6	7.926E+03	2.631E+11	2.906E-08	1.001E-02	1.820E+05	3.162E-01	0.000E+00	7.460E-03	-0.006	-0.006
	.889E+01	4217.9	8.888E+03	3.250E+11	3.169E-08	1.191E-02	1.852E+05	4.217E-01	0.000E+00	8.424E-03	-0.006	-0.026
	.976E+01	4346.7 4492.6	9.975E+03	3.992E+11	3.450E-08	1.394E-02	1.885E+05	5.623E-01	0.000E+00	9.389E-03	-0.001 0.002	-0.005 -0.053
	.123E+02 .268E+02	4658.1	1.123E+04 1.268E+04	4.888E+11 6.055E+11	3.756E-08 4.090E-08	1.602E-02 1.845E-02	1.919E+05 1.956E+05	7.499E-01 1.000E+00	0.000E+00 0.000E+00	1.031E-02 1.143E-02	-0.000	-0.093
	433E+02	4846.7	1.433E+04	7.931E+11	4.442E-08	2.222E-02	1.995E+05	1.334E+00	1.086E-08	1.335E-02	-0.003	-0.165
58 1.	.607E+02	5060.7	1.607E+04	1.175E+12	4.769E-08	2.988E-02	2.033E+05	1.778E+00	5.072E-06	1.755E-02	0.002	-0.107
	.768E+02	5305.5	1.767E+04	2.028E+12		4.610E-02	2.066E+05	2.371E+00	3.078E-04	2.672E-02	-0.008	0.000
	.899E+02 .997E+02	5585.5	1.899E+04	3.953E+12 8.302E+12		7.918E-02	2.091E+05	3.162E+00	3.414E-03	4.565E-02	0.013	-0.001
	.997E+02 .066E+02			8.302E+12 1.769E+13							-0.032 0.048	0.000 -0.000
	.117E+02			3.473E+13							-0.068	-0.000
	.156E+ 0 2	6895.4	2.156E+04	5.813E+13	4.683E-08	7.909E-01	2.143E+05	1.000E+01	4.847E-01	2.387E-01	0.065	0.000
	.191E+02			8.566E+13							-0.040	-0.001
	.224E+02 .258E+02			1.159E+14 1.508E+14							-0.006 0.048	0.008 -0.000
	.292E+02			1.892E+14							-0.044	-0.000
	.329E+02			2.338E+14							0.012	0.001
	.369E+02			2.828E+14							0.009	-0.000
	412E+02			3.411E+14							0.003	-0.000
/2 2.	.458E+02	8224.6	2.45/E+04	4.031E+14	4.403E-08	o.869E+00	4.211E+05	1.000E+02	9.508E-01	1.090E-01	-0.019	-0.001

ATLAS9 model computed from scratch:

TEFF	4200	100.0	00000		1 17771DD 2 6	I /II 1 2F N	OUED NEW OF	· F			TTEDATIO	N 15
TEFF 0	4200.	LOG G 2	.00000	ELECTRON] VTURB=2.0	ROSSELAND	OVEK NEW OL HEIGHT	ROSSELAND	FRACTION	RADIATIVE	ITERATION PER CE	N 15 ENT FLUX
ŭ	RHOX	TEMP	PRESSURE	NUMBER	DENSITY	MEAN	(KM)	DEPTH		ACCELERATION	ERROR	DERIV
1 9.	.894E-03	2411.4	1.082E+00	7.708E+06	6.915E-12	1.236E-05	0.000E+00	1.334E-07	0.000E+00	2.034E-03	0.009**	****
	.333E-02	2474.2	1.449E+00	1.159E+07	8.987E-12	1.198E-05	4.855E+03	1.778E-07	0.000E+00	1.761E-03		*****
	.778E-02 .310E-02	2487.3 2500.2	1.919E+00	1.472E+07 1.810E+07	1.185E-11	1.329E-05	9.530E+03 1.377E+04	2.371E-07	0.000E+00	1.686E-03		******
	.960E-02	2524.3	2.478E+00 3.144E+00	2.281E+07	1.521E-11 1.907E-11	1.516E-05 1.646E-05	1.769E+04	3.162E-07 4.217E-07	0.000E+00 0.000E+00	1.669E-03 1.606E-03		6730.372 4606.210
	.768E-02	2549.7	3.962E+00	2.889E+07	2.373E-11	1.787E-05	2.153E+04	5.623E-07	0.000E+00	1.531E-03		3376.736
	.778E-02	2581.6	4.985E+00	3.744E+07	2.941E-11	1.868E-05	2.542E+04	7.499E-07	0.000E+00	1.431E-03		3350.106
	.082E-02	2618.9	6.304E+00	4.959E+07	3.655E-11	1.926E-05	2.945E+04	1.000E-06	0.000E+00	1.327E-03		3561.478
	.743E-02 .840E-02	2656.0 2696.4	7.979E+00 1.008E+01	6.570E+07 8.804E+07	4.552E-11 5.650E-11	2.059E-05 2.179E-05	3.356E+04 3.769E+04	1.334E-06 1.778E-06	0.000E+00 0.000E+00	1.236E-03 1.142E-03		3328.005 3354.180
	.245E-01	2731.0	1.269E+01	1.160E+08	7.016E-11	2.179E-05 2.374E-05	4.185E+04	2.371E-06	0.000E+00	1.052E-03		2988.694
	.560E-01	2759.0	1.583E+01	1.484E+08	8.655E-11	2.654E-05	4.587E+04	3.162E-06	0.000E+00	9.889E-04		2352.597
	.937E-01	2787.4	1.958E+01	1.889E+08	1.058E-10	2.973E-05	4.978E+ 0 4	4.217E-06	0.000E+00	9.370E-04		1932.162
	.386E-01	2813.0	2.405E+01		1.287E-10	3.303E-05	5.361E+04	5.623E-06	0.000E+00	8.866E-04		1586.224
	.927E-01	2835.2 2855.7	2.945E+01 3.592E+01	2.934E+08 3.589E+08	1.564E-10 1.895E-10	3.651E-05 4.081E-05	5.742E+04 6.119E+04	7.499E-06 1.000E-05	0.000E+00 0.000E+00	8.437E-04 8.156E-04		1241.031 -913.117
	.347E-01	2877.1	4.361E+01	4.380E+08	2.283E-10	4.600E-05	6.488E+04	1.334E-05	0.000E+00	8.015E-04		-667.850
	.253E-01	2901.7	5.261E+01	5.370E+08	2.731E-10	5.274E-05	6.848E+04	1.778E-05	0.000E+00	7.947E-04		-499.952
19 6.	.310E-01	2925.9	6.312E+01	6.551E+08	3.248E-10	6.002E-05	7.199E+04	2.371E-05	0.000E+00	7.863E-04	-0.021 -	-377.925
	.550E-01	2950.3	7.548E+01	7.975E+08	3.852E-10	6.789E-05	7.548E+04	3.162E-05	0.000E+00	7.792E-04		-283.060
	.017E-01	2974.8 3001.2	9.009E+01 1.073E+02	9.693E+08 1.183E+09	4.561E-10 5.388E-10	7.646E-05 8.660E-05	7.896E+04 8.244E+04	4.217E-05 5.623E-05	0.000E+00 0.000E+00	7.759E-04 7.776E-04	-0.026 -	-205.891 -152.356
	.075E+00 .277E+00	3029.7	1.073E+02 1.274E+02	1.450E+09	6.336E-10	1.004E-04	8.588E+04	7.499E-05	0.000E+00	7.776E-04 7.927E-04		-132.336
	.506E+00	3065.8	1.502E+02	1.811E+09	7.380E-10	1.188E-04	8.921E+04	1.000E-04	0.000E+00	7.974E-04		-122.721
	.768E+00	3090.5	1.763E+02	2.171E+09	8.593E-10	1.368E-04	9.248E+04		0.000E+00	7.951E-04	-0.037	-87.612
	.072E+00	3116.8	2.066E+02	2.608E+09	9.984E-10	1.572E-04	9.574E+04	1.778E-04	0.000E+00	8.054E-04	-0.041	-60.832
	.425E+00	3142.8	2.417E+02	3.126E+09	1.159E-09	1.802E-04	9.901E+04	2.371E-04		8.176E-04	-0.043	-40.720
	.834E+00 .304E+00	3169.2 3196.1	2.825E+02 3.295E+02	3.746E+09 4.482E+09	1.343E-09 1.554E-09	2.075E-04 2.424E-04	1.023E+05 1.055E+05	3.162E-04 4.217E-04	0.000E+00 0.000E+00	8.362E-04 8.654E-04	-0.046 -0.048	-25.405 -14.323
	.838E+00	3223.4	3.829E+02	5.355E+09	1.790E-09	2.843E-04	1.033E+03	5.623E-04	0.000E+00	9.004E-04	-0.050	-8.624
	.448E+00	3249.9	4.438E+02	6.365E+09	2.058E-09	3.311E-04	1.119E+05	7.499E-04	0.000E+00	9.377E-04	-0.051	-4.924
32 5.	.147E+00	3275.8	5.138E+02	7.542E+09	2.364E-09	3.837E-04	1.151E+05	1.000E-03	0.000E+00	9.796E-04	-0.052	-2.700
	.954E+00	3301.4	5.945E+02	8.916E+09	2.715E-09	4.428E-04	1.182E+05	1.334E-03	0.000E+00	1.027E-03	-0.052	-1.551
	.885E+00 .955E+00	3326.6	6.876E+02	1.052E+10	3.116E-09	5.129E-04	1.214E+05	1.778E-03	0.000E+00 0.000E+00	1.084E-03	-0.053	-0.881
	.955E+00 .189E+00	3351.3 3375.3	7.947E+02 9.181E+02	1.237E+10 1.451E+10	3.575E-09 4.102E-09	5.946E-04 6.871E-04	1.246E+05 1.279E+05	2.371E-03 3.162E-03		1.151E-03 1.228E-03	-0.053 -0.053	-0.435 -0.272
	.061E+01	3399.3	1.060E+03	1.699E+10	4.705E-09	7.946E-04		4.217E-03	0.000E+00	1.318E-03	-0.054	-0.135
38 1.	.225E+01	3422.7	1.225E+03	1.986E+10	5.397E-09	9.186E-04	1.344E+05	5.623E-03	0.000E+00	1.425E-03	-0.054	-0.279
	.414E+01	3449.3	1.414E+03	2.338E+10	6.181E- 0 9		1.376E+05	7.499E-03	0.000E+00	1.547E-03	-0.057	-2.201
	.632E+01	3472.3	1.632E+03	2.722E+10	7.089E-09	1.226E-03	1.409E+05	1.000E-02		1.676E-03	-0.061	-1.250
	.885E+01	3496.1 3521.9	1.884E+03 2.175E+03	3.175E+10 3.716E+10	8.132E-09 9.321E-09	1.417E-03 1.641E-03	1.442E+05 1.476E+05	1.334E-02 1.778E-02	0.000E+00 0.000E+00	1.835E-03 2.026E-03	-0.062 -0.055	0.595 2.051
	.509E+01	3551.4	2.509E+03	4.383E+10	1.066E-08	1.914E-03	1.509E+05	2.371E-02	0.000E+00	2.258E-03	-0.041	2.517
44 2.	.889E+01	3586.6	2.889E+03	5.228E+10	1.216E-08	2.256E-03	1.543E+05	3.162E-02	0.000E+00	2.539E-03	-0.027	1.411
	.317E+01	3626.8		6.286E+10	1.380E-08	2.680E-03	1.576E+05	4.217E-02	0.000E+00	2.870E-03	-0.020	-0.012
	.796E+01	3669.5			1.561E-08	3.193E-03	1.608E+05	5.623E-02	0.000E+00	3.253E-03	-0.020	0.107
	.331E+01 .922E+01	3719.5 3776.3	4.331E+03 4.922E+03	9.178E+10 1.122E+11	1.756E-08 1.965E-08	3.838E-03 4.636E-03	1.641E+05 1.672E+05	7.499E-02 1.000E-01	0.000E+00 0.000E+00	3.721E-03 4.278E-03	-0.018 -0.017	0.036 0.035
	.573E+01	3842.3	5.573E+03	1.382E+11	2.186E-08	5.628E-03	1.704E+05	1.334E-01	0.000E+00	4.928E-03	-0.016	0.016
	.287E+01	3918.0	6.287E+03	1.711E+11	2.417E-08	6.848E-03	1.735E+05	1.778E-01	0.000E+00	5.684E-03	-0.014	0.009
	.070E+01	4005.0	7.070E+03	2.123E+11	2.657E-08	8.321E-03	1.766E+05	2.371E-01	0.000E+00	6.548E-03	-0.013	0.003
	.933E+01	4104.5	7.933E+03	2.632E+11	2.908E-08	1.002E-02	1.797E+05	3.162E-01	0.000E+00	7.464E-03	-0.011	-0.000
	.894E+01 .981E+01	4217.8 4346.6	8.894E+03 9.981E+03	3.251E+11 3.993E+11	3.172E-08 3.452E-08	1.192E-02 1.395E-02	1.828E+05 1.861E+05	4.217E-01 5.623E-01	0.000E+00 0.000E+00	8.427E-03 9.393E-03	-0.009 -0.003	-0.017 -0.002
	.123E+02	4492.6	1.123E+04	4.890E+11	3.758E-08	1.603E-02	1.896E+05	7.499E-01	0.000E+00	1.031E-02	-0.004	-0.088
	.268E+02	4658.0	1.268E+04	6.056E+11	4.091E-08	1.845E-02	1.933E+05	1.000E+00	0.000E+00	1.143E-02	-0.011	-0.078
	.433E+02	4846.4	1.433E+04	7.93 0 E+11	4.443E-08	2.222E-02	1.971E+ 0 5	1.334E+00	1.096E-08	1.335E-02	0.014	-0.060
	.607E+02	5061.2	1.607E+04	1.176E+12	4.770E-08	2.990E-02	2.009E+05	1.778E+00	5.066E-06	1.756E-02	0.001	-0.265
	.768E+02 .899E+02	5305.1 5585.7	1.768E+04 1.899E+04	2.026E+12 3.955E+12	5.005E-08	4.608E-02 7.921E-02	2.042E+05 2.068E+05	2.371E+00 3.162E+00	3.088E-04 3.419E-03	2.670E-02 4.568E-02	-0.029 0.019	0.000 0.001
								4.217E+00				-0.001
		6259.6	2.066E+04	1.771E+13	4.954E-08	2.751E-01	2.101E+05	5.623E+00	8.607E-02	1.468E-01	-0.000	-0.004
63 2.	.117E+02	6605.3	2.117E+04	3.468E+13	4.805E-08	4.941E-01	2.111E+05	7.499E+00	2.611E-01	2.133E-01	-0.068	0.005
								1.000E+01			0.189	
	.191E+02 .224E+02							1.334E+01 1.778E+01			-0.325 0.139	-0.001 0.054
	.258E+02							2.371E+01				-0.021
	.293E+02							3.162E+01			-1.301	0.001
69 2.	.329E+02	7798.1	2.329E+04	2.311E+14	4.437E-08	3.151E+00	2.158E+05	4.217E+01	9.007E-01	1.866E-01	1.304	0.004
	.369E+02							5.623E+01			0.123	-0.097
	.411E+02							7.499E+01 1.000E+02			-2.040	0.003
12 2.	JOE+WZ	0443.3	2.43/E+W4	7.W34E+14	4.403E-08	J.0/JE+WW	2.10/E+W5	1.000E+0Z	J.JUIE-UI	1.//OE-WI	3.026	0.203

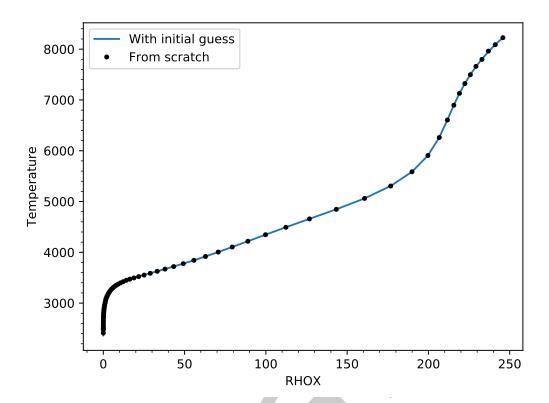


Figure 1: Comparison of two ATLAS9 models for identical physical parameters ($T_{\text{eff}} = 4200 \text{ K}$, $\log g = 2.0$, [m/H] = -0.5, but different initial conditions (initial guess vs. model computed from scratch).

The errors on the fluxes are small in both cases, but the flux derivative hasn't converged quite as well for the model that was computed from scratch, especially in the outermost layers. This is because a model calculation from scratch starts from a grey atmosphere, which will in general be a worse starting guess than a scaled pre-existing ATLAS model with parameters close to those of the requested model. Hence, more iterations are necessary. To achieve better convergence, one can increase the number of repetitions to kurucz.A9_NREP = 3. The code will then carry out $3 \times 15 = 45$ iterations. Nevertheless, the structure of the two models is already quite similar, as can be seen when plotting the temperature as a function of depth (RHOX) in Figure 1.

4.2 Computing an ATLAS12 model

The procedure for computing an ATLAS12 model is similar to that for ATLAS9, except that the step of setting up the ODF is not needed:

> from ispy3 import kurucz

```
> m, teff, logg, vturb = -0.5, 4200., 2.00, 2.0
> kurucz.mkatm_a12(teff, logg, m, 't4200g200m050', vturb=vturb)
```

A major difference is the time required to compute the model, which in this example was 50 *minutes* (as compared with 9 *seconds* for the equivalent ATLAS9 model). The computation time will vary strongly depending on the temperature and composition of the models, with cooler and more metal-rich models generally taking longer to compute as more lines contribute significantly to the opacity. The model atmosphere is now stored in t4200g200m050.A12 and the diagnostic output file is t4200g200m050b.out. The last few lines in this file are similar to those produced for the ATLAS9 models, as can be seen below. The model has converged nicely, and its structure is very similar to that of the ATLAS9 model, as expected. In the outermost and in the very deepest layers, the temperature difference between the ATLAS9 and ATLAS12 models reaches 15 K, but over most of the depth range it is less than 5 K.



ATLAS12 model:

TEFF	4200.	LOG G	2.00000	ATLAS	12						ITERATION	15
0				ELECTRON		ROSSELAND	HEIGHT	ROSSELAND	FRACTION			NT FLUX
1 2	RHOX 528E-03	TEMP 2364.5	PRESSURE 8.527E-01	NUMBER 5.431E+06	DENSITY 5.439E-12	MEAN 1 564F_05	(KM) 0.000E+00	DEPTH 1.334E-07	CONV FLUX 0.000E+00	ACCELERATION 3.264E-03	ERROR -0.000	DERIV -5.562
	129E-02	2389.6	1.129E+00	7.208E+06	7.128E-12	1.653E-05	4.483E+03	1.778E-07	0.000E+00	3.038E-03	-0.000	-2.182
	478E-02	2414.0	1.478E+00	9.459E+06	9.236E-12	1.748E-05	8.817E+03	2.371E-07	0.000E+00	2.819E-03	-0.000	-2.154
	919E-02	2438.9	1.919E+00	1.234E+07	1.187E-11		1.306E+04	3.162E-07 4.217E-07	0.000E+00	2.595E-03	-0.000	-2.980
	477E-02 185E-02	2464.5 2490.8	2.476E+00 3.184E+00	1.603E+07 2.078E+07	1.516E-11 1.929E-11	1.936E-05 2.034E-05	1.721E+04 2.134E+04	4.217E-07 5.623E-07	0.000E+00 0.000E+00	2.377E-03 2.168E-03	-0.000 -0.000	-2.472 -2.189
	083E-02	2517.7	4.083E+00	2.692E+07	2.447E-11		2.546E+04	7.499E-07	0.000E+00	1.975E-03	-0.000	-2.447
	219E-02	2545.2	5.219E+00	3.483E+07	3.094E-11	2.258E-05	2.959E+04	1.000E-06	0.000E+00	1.801E-03	-0.000	-2.403
	653E-02 457E-02	2573.1 2601.5	6.654E+00 8.457E+00	4.503E+07 5.814E+07	3.902E-11 4.906E-11	2.390E-05	3.371E+04 3.782E+04	1.334E-06 1.778E-06	0.000E+00 0.000E+00	1.646E-03 1.509E-03	-0.000 -0.000	-3.478 -4.014
	971E-02	2630.2	1.071E+01	7.497E+07	6.147E-11	2.714E-05	4.192E+04	2.371E-06	0.000E+00	1.399E-03	-0.000	-3.679
12 1.3	352E-01	2659.3	1.352E+01	9.650E+07	7.674E-11	2.915E-05	4.600E+04	3.162E-06	0.000E+00	1.288E-03	-0.000	-3.230
	700E-01 129E-01	2688.7	1.700E+01	1.240E+08	9.543E-11		5.006E+04	4.217E-06	0.000E+00	1.200E-03	-0.000	-2.292
	553E-01	2718.4 2748.5	2.129E+01 2.653E+01	1.590E+08 2.034E+08	1.182E-10 1.457E-10	3.416E-05 3.728E-05	5.409E+04 5.808E+04	5.623E-06 7.499E-06	0.000E+00 0.000E+00	1.124E-03 1.061E-03	-0.000 -0.000	-2.532 -2.326
	293E-01	2778.9	3.293E+01	2.596E+08	1.788E-10	4.093E-05	6.204E+04	1.000E-05	0.000E+00	1.007E-03	-0.000	-2.461
	967E-01	2809.9	4.067E+01	3.308E+08	2.184E-10	4.519E-05	6.595E+04	1.334E-05	0.000E+00	9.628E-04	-0.000	-2.061
	900E-01 115E-01	2841.4 2873.5	4.999E+01	4.207E+08 5.340E+08	2.655E-10 3.211E-10	5.021E-05 5.616E-05	6.982E+04 7.363E+04	1.778E-05 2.371E-05	0.000E+00 0.000E+00	9.261E-04 8.963E-04	-0.000 -0.000	-2.043 -1.831
	439E-01	2906.2	7.439E+01	6.761E+08	3.862E-10	6.327E-05	7.739E+04	3.162E-05	0.000E+00	8.733E-04	-0.000	-2.030
	901E-01	2939.3	9.001E+01	8.538E+08	4.620E-10	7.182E-05	8.108E+04	4.217E-05	0.000E+00	8.564E-04	-0.000	-1.813
	083E+00	2972.8	1.083E+02	1.074E+09	5.495E-10	8.214E-05	8.470E+04	5.623E-05	0.000E+00	8.449E-04	-0.000	-1.818
	295E+00 540E+00	3006.2 3039.8	1.295E+02 1.540E+02	1.346E+09 1.678E+09	6.498E-10 7.641E-10	9.461E-05 1.097E-04	8.825E+04 9.172E+04	7.499E-05 1.000E-04	0.000E+00 0.000E+00	8.393E-04 8.384E-04	-0.000 -0.000	-1.610 -2.170
	821E+00	3072.7	1.821E+02	2.078E+09	8.938E-10	1.277E-04		1.334E-04	0.000E+00	8.421E-04	-0.000	-1.933
	143E+00	3105.0	2.143E+02	2.556E+ 0 9	1.040E-09	1.492E-04	9.845E+04	1.778E-04	0.000E+00	8.507E-04	-0.001	-1.671
	509E+00 927E+00	3136.5 3167.1	2.509E+02 2.927E+02	3.123E+09 3.791E+09	1.206E-09 1.393E-09	1.746E-04 2.044E-04	1.017E+05 1.049E+05	2.371E-04 3.162E-04	0.000E+00 0.000E+00	8.642E-04 8.828E-04	-0.001 -0.001	-1.294 -0.892
	403E+00	3196.8	3.403E+02	4.573E+09	1.604E-09	2.391E-04	1.045E+05	4.217E-04	0.000E+00	9.063E-04	-0.001	-0.609
	945E+00	3225.5	3.945E+02	5.483E+09	1.844E-09	2.794E-04	1.113E+05	5.623E-04	0.000E+00	9.352E-04	-0.001	-0.399
	565E+00	3253.4	4.565E+02	6.540E+09	2.115E-09	3.260E-04	1.144E+05	7.499E-04	0.000E+00	9.698E-04	-0.001	-0.268
	274E+00 987E+00	3280.3 3306.3	5.274E+02 6.087E+02	7.762E+09 9.176E+09	2.424E-09 2.775E-09	3.795E-04 4.410E-04	1.175E+05 1.207E+05	1.000E-03 1.334E-03	0.000E+00 0.000E+00	1.011E-03 1.059E-03	-0.001 -0.001	-0.196 -0.141
	021E+00	3331.6	7.021E+02	1.081E+10	3.177E-09	5.114E-04	1.238E+05	1.778E-03	0.000E+00	1.115E-03	-0.001	-0.125
	096E+00	3356.3	8.096E+02	1.269E+10	3.637E-09	5.921E-04	1.270E+05	2.371E-03		1.180E-03	-0.001	-0.089
	335E+00 076E+01	3380.5	9.334E+02	1.487E+10	4.163E-09	6.848E-04	1.301E+05	3.162E-03	0.000E+00	1.256E-03	-0.001	-0.094
	241E+01	3404.5 3428.3	1.076E+03 1.241E+03	1.739E+10 2.032E+10	4.767E-09 5.460E-09	7.916E-04 9.143E-04	1.334E+05 1.366E+05	4.217E-03 5.623E-03	0.000E+00 0.000E+00	1.343E-03 1.445E-03	-0.001 -0.001	-0.089 -0.118
	431E+01	3453.0	1.431E+03	2.378E+10	6.252E-09		1.398E+05	7.499E-03	0.000E+00	1.565E-03	-0.002	-0.566
	651E+01	3477.7	1.651E+03	2.780E+10	7.160E-09	1.224E-03	1.431E+05	1.000E-02	0.000E+00	1.701E-03	-0.003	-0.476
	903E+01 194E+01	3503.2 3530.4	1.903E+03 2.194E+03	3.253E+10 3.817E+10	8.196E-09 9.375E-09	1.418E-03 1.647E-03	1.464E+05 1.497E+05	1.334E-02 1.778E-02	0.000E+00 0.000E+00	1.861E-03 2.051E-03	-0.004 -0.004	-0.194 0.057
	526E+01	3560.2	2.526E+03	4.500E+10	1.071E-08	1.923E-03	1.530E+05	2.371E-02	0.000E+00	2.278E-03	-0.004	0.110
	905E+01	3593.6	2.905E+03	5.337E+10		2.259E-03	1.563E+05	3.162E-02	0.000E+00	2.547E-03	-0.003	0.011
	333E+01	3631.0	3.333E+03	6.368E+10	1.385E-08	2.669E-03	1.596E+05	4.217E-02	0.000E+00	2.868E-03	-0.003	-0.032 -0.008
	815E+01 353E+01	3673.7 3722.8	3.815E+03 4.353E+03	7.652E+10 9.274E+10	1.566E-08 1.763E-08	3.177E-03 3.812E-03	1.629E+05 1.661E+05	5.623E-02 7.499E-02	0.000E+00 0.000E+00	3.253E-03 3.714E-03	-0.003 -0.003	-0.023
	948E+01	3779.4	4.948E+03	1.133E+11	1.974E-08	4.604E-03		1.000E-01	0.000E+00	4.263E-03	-0.003	-0.016
	504E+01	3845.0	5.604E+03	1.394E+11	2.196E-08	5.593E-03	1.725E+05	1.334E-01	0.000E+00	4.914E-03	-0.003	-0.021
	322E+01 110E+01	3920.5 4007.4	6.322E+03 7.109E+03	1.724E+11 2.139E+11		6.808E-03 8.271E-03	1.756E+05 1.787E+05	1.778E-01 2.371E-01	0.000E+00 0.000E+00	5.668E-03 6.520E-03	-0.002 -0.001	-0.014 -0.025
	977E+01	4106.9	7.977E+03	2.651E+11		9.968E-03	1.818E+05	3.162E-01	0.000E+00	7.442E-03	0.001	-0.037
	944E+01	4220.5	8.943E+03	3.274E+11		1.186E-02	1.849E+05	4.217E-01	0.000E+00	8.397E-03	0.004	-0.050
	004E+02	4349.3	1.004E+04	4.020E+11	3.469E-08 3.776E-08	1.387E-02	1.882E+05	5.623E-01 7.499E-01	0.000E+00 0.000E+00	9.344E-03	0.002 -0.000	-0.106 -0.088
	129E+02 275E+02	4495.1 4660.7	1.129E+04 1.275E+04	4.920E+11 6.094E+11	4.109E-08	1.598E-02 1.842E-02	1.917E+05 1.954E+05	1.000E+00	0.000E+00	1.029E-02 1.142E-02	0.009	-0.081
	440E+02	4849.2	1.440E+04	7.984E+11	4.461E-08	2.22E-02	1.992E+05	1.334E+00	1.292E-08	1.334E-02	0.011	-0.168
	514E+02	5063.3	1.613E+04	1.183E+12	4.787E-08	2.994E-02	2.030E+05	1.778E+00	5.575E-06	1.758E-02	-0.003	-0.135
	774E+02 905E+02	5307.8 5588.6	1.774E+04 1.904E+04	2.042E+12 3.988F+12	5.019E-08 5.117E-08	4.623E-02 7.954E-02	2.063E+05 2.088E+05	2.371E+00 3.162E+00	3.137E-04 3.118E-03	2.681E-02 4.590E-02	0.007 0.014	0.012 -0.002
	002E+02					1.460E-01					-0.025	
62 2.0	971E+02	6262.7	2.071E+04	1.784E+13	4.963E-08	2.768E-01	2.121E+05	5.623E+00	9.177E-02	1.468E-01	0.023	0.000
	122E+02					4.927E-01					-0.017	
	161E+02 196E+02					7.898E-01 1.123E+00					-0.039 0.064	
	230E+02	7306.3	2.229E+04	1.135E+14	4.559E-08	1.508E+00	2.155E+05	1.778E+01	7.584E-01	2.136E-01	-0.067	-0.001
	264E+02					1.980E+00					-0.030	
	299E+02 337E+02					2.525E+00 3.156E+00					0.031 -0.189	-0.001 0.001
						3.902E+00						-0.000
71 2.4	420E+02	8085.4	2.419E+04	3.387E+14	4.422E-08	4.821E+00	2.197E+05	7.499E+01	9.412E-01	1.653E-01	-0.391	0.000
72 2.4	466E+ 0 2	8218.4	2.465E+04	4.009E+14	4.422E-08	5.842E+00	2.208E+05	1.000E+02	9.509E-01	1.688E-01	0.621	0.041

4.3 Calculating a synthetic spectrum with SYNTHE

Once an ATLAS9 or ATLAS12 model has been computed, a synthetic spectrum can be produced with SYNTHE for any specified set of abundances. In the example below we compute a model spectrum that includes the subordinate Na I lines near 5680 Å, with Na enhanced by +0.4 dex relative to the solar-scaled abundance patterns, a microturbulent velocity of 2 km/s, and no rotational broadening.

```
> from ispy3 import syntherk
>
> atoms, abun = ['Na'], [+0.4]
> vturb, vrot = 2.0, 0.0
>
> syntherk.synbeg(5675., 5690., 0.01, initdir='/scratch/soeren')
> syntherk.synthespec(-0.5, 't4200g200m050.A9', 't4200g200m050.asc',
> initdir='/scratch/soeren', atoms=atoms, abun=abun,
> vturb=vturb, vrot=vrot)
```

The procedure has two steps: first, synbeg() must be called to define the wavelength range and spectral resolution. This will also copy the relevant data from the line lists (atomic and molecular) and stage the data in initdir. The second step is the actual call to synthespec(), in which the logarithmic baseline scaling of the abundances is the first argument (-0.5) and any offsets in the abundances of specific elements are given in the arrays atoms and abun. Hence, in this example most elements will have [X/H] = -0.5, except sodium with [Na/H] = -0.5 + 0.4 = -0.1, or [Na/Fe] = +0.4. The model atmosphere is read from t4200g200m050.A9 and the synthetic spectrum will be stored in t4200g200m050.asc. The first few lines of the output file are:

```
# ISPy3 0.90.0
# SYNTHE
# atomdir = /Users/soeren/cats/linelists/atoms/ssl/SYNTHE/
# moldir = /Users/soeren/cats/Kurucz/molecules/20apr2016/
\# atmname = t4200g200m050.A9
# TEFF
        4200.GRAVITY 2.000
          [-0.5] VTURB=2.0 L/H=1.25 NOVER NEW ODF
# TITLE
5675.0073 0.9792 4.9447e+05
5675.0171 0.9752
                  4.9247e+05
5675.0269 0.9692
                  4.8941e+05
5675.0371 0.9612
                  4.8537e+05
5675.0469 0.9528
                  4.8115e+05
```

After the ISPy3 version number, the second line indicates the code used to calculate the spectrum (here SYNTHE). The atomdir and moldir lines give the directories from which the atomic and molecular line data were read, as specified by the variables with the same names in syntherk.py.

Α

The actual model spectrum has three columns: The air wavelength (in Å), followed by the normalised spectrum, and finally the flux in physical units.

4.4 Interpolating in the MARCS grid

From the user's perspective, the procedure for obtaining a model by interpolation in the MARCS grid is (by design) closely analogous to that for computing an ATLAS model:

```
> from ispy3 import marcs
>
> m, teff, logg = -0.5, 4200., 2.00
> mass = 1.0
> marcs.intpatm(teff, logg, m, mass, 't4200g200m050')
```

The intpatm() function requires the mass as a fourth argument. However, this is currently ignored and a mass of $1\,M_\odot$ will be assumed for the spherical models no matter what is specified. The function also accepts the vturb argument, but this is also currently ignored and a microturbulent velocity of 2 km/s is assumed for spherical models while 1 km/s is assumed for plane-parallel models. Since no new model is computed, the call to intpatm() is very fast. The interpolated model is stored as t4200g200m050.marcs. The header indicates that the model was interpolated from the grid of spherical models, followed by the 56 layers of the atmosphere with a depth scale defined at 5000 Å. The last few lines list the models that were selected from the grid for the interpolation. Since there are three parameters, eight models are needed.

Interpolated MARCS model:

```
'sphINTERPOL' 56
-3.9868 2800.20
                    -3.9141
                               1.5463
                                         2.0000
                                                     0.116831E+13 -5.0000
-3.8625
          2825.70
                    -3.8243
                                1.6180
                                          2.0000
                                                     0.116800E+13
-3.7439
          2853.77
                    -3.7290
                                1.6921
                                         2.0000
                                                     0.116766E+13
                                                                    -4.7455
                                          2.0000
                                                     0.116732E+13
-3.5204
          2913.77
                    -3.5300
                               1.8447
                                         2.0000
                                                     0.116696E+13
                                                                    -4.4909
-3.4142
          2944.50
                     -3.4295
                                1.9216
                                          2.0000
                                                     0.116660E+13
                                                                     -4.3636
          2975.55
                                          2.0000
                                                     0.116624E+13
-3.3112
                    -3.3295
                               1.9979
                                                                     -4.2364
          3006.66
                     -3.2306
                                2.0732
                                          2.0000
                                                     0.116588E+13
                                         2.0000
                                                     0.116553E+13
-3.1123
          3037.69
                    -3.1333
                               2.1472
                                                                     -3.9818
-3.0156
          3068.53
                    -3.0379
                               2.2198
                                          2.0000
                                                     0.116517E+13
                                                                     -3.8545
                                          2.0000
                                                     0.116482E+13
-2.9201
          3099.15
                    -2.9444
                               2.2910
                                                                     -3.7273
-2.8253
          3129.45
                    -2.8531
                               2.3608
                                          2 0000
                                                     0.116447E+13
                                                                     -3.6000
                                          2.0000
                                                     0.116413E+13
-2.7309
          3159.28
                    -2.7640
                               2.4294
                                                                     -3.4727
-2.6365
-2.5419
                                         2.0000
          3188.70
                    -2.6769
                               2.4969
                                                     0.116379E+13
                                                                     -3.3455
          3217.41
                    -2.5922
                               2.5635
                                                     0.116345E+13
                                                                     -3.2182
-2.4465
          3244.94
                    -2.5103
                               2.6293
                                         2 0000
                                                     0.116311E+13
                                                                     -3 0909
          3272.35
-2.3505
                    -2.4294
                               2.6945
                                          2.0000
                                                     0.116278E+13
                                                                     -2.9636
-2.2537
-2.1556
                    -2.3498
-2.2725
                               2.7593
2.8236
          3299.59
                                          2.0000
                                                     0.116245E+13
                                                                     -2.8364
                                          2.0000
                                                     0.116212E+13
                                                                     -2.7091
          3325.73
          3351.27
3376.32
                               2.8878
2.9519
-2.0562
                     -2.1966
                                          2.0000
                                                     0.116179E+13
                                                                     -2.5818
-1.9555
                    -2.1219
                                          2.0000
                                                     0.116146E+13
                                                                     -2.4545
-1.8534
          3401.07
3425.70
                    -2.0482
                               3.0159
                                         2 0000
                                                     0.116112E+13
                                                                     -2 3273
                    -1.9750
                                3.0799
                                                     0.116079E+13
-1.7499
                                          2.0000
                                                                     -2.2000
-1.6451
          3450.68
                    -1.9019
                               3.1438
                                         2.0000
                                                     0.116045E+13
                                                                     -2.0727
-1.5391
          3476.47
                               3.2077
                                          2.0000
                                                     0.116011E+13
                                                                     -1.9455
                    -1.8282
-1 4320
          3503 59
                    -1 7534
                               3 2715
                                         2 0000
                                                     0 115977F+13
                                                                     -1 8182
-1.3241
                    -1.6768
                               3.3348
                                                     0.115942E+13
          3532.95
                                          2.0000
                                                                     -1.6909
-1.2154
-1.1064
          3565.25
                    -1 5976
                                3 3977
                                          2 0000
                                                     0 115908F+13
                                                                     -1 5636
          3601.11
                    -1.5155
                                3.4597
                                                     0.115874E+13
-0.9974
          3642.45
3689.59
                    -1.4290
                               3.5207
                                         2.0000
                                                     0.115840E+13
                                                                     -1.3091
-0.8886
                    -1.3386
                                3.5803
                                          2.0000
                                                     0.115807E+13
                                                                     -1.1818
-0.7804
          3744.03
                    -1.2435
                               3.6383
                                         2.0000
                                                     0.115773E+13
                                                                     -1.0545
-0.6729
          3806.63
                    -1.1443
                                3.6946
                                          2.0000
                                                     0.115741E+13
                                                                     -0.9273
-0.5666
          3879.03
                    -1.0411
                               3.7490
                                         2.0000
                                                     0.115709E+13
                                                                     -0.8000
-0.4615
          3962.30
                     -0.9348
                                3.8017
                                          2.0000
                                                     0.115677E+13
                                                                     -0.6727
-0.3575
          4058.02
                    -0.8266
                               3.8530
                                         2.0000
                                                     0.115645E+13
                                                                     -0.5455
 -0.2545
          4167.54
                     -0.7179
                                3.9034
                                          2.0000
                                                     0.115613E+13
-0.1521
          4292.54
                    -0.6097
                               3.9537
                                         2.0000
                                                     0.115581E+13
                                                                     -0.2909
 -0.0498
          4435.26
                     -0.5022
                                4.0045
                                          2.0000
                                                     0.115547E+13
                                                                     -0.1636
                               4.0565
                                          2.0000
                                                     0.115511E+13
 0.0527
          4598.03
                    -0.3930
                                                                     -0.0364
                                                                     0.0909
 0.1557
          4784.16
                                4.1096
                                          2.0000
                                                     0.115473E+13
                                         2.0000
                                                     0.115434E+13
 0.2593
          4996.95
                    -0.1083
                               4.1614
 0.3645
          5240.27
                     0.1182
                               4.2074
                                         2.0000
                                                     0.115398E+13
0.115368E+13
                                                                     0.3455
                               4.2439
 0.4717
          5518.41
                     0.4101
                                                                      0.4727
 0.5812
          5836.58
                     0.7493
                               4.2704
                                          2.0000
                                                     0.115344E+13
                                                                      0.6000
                                                                     0.7273
 0.6920
          6202.24
                     1.1171
                               4.2885
                                          2.0000
                                                     0.115328E+13
                     1.4700
                               4.3005
                                         2.0000
                                                                     0.8545
0.9818
 0.8028
          6586.93
                                                     0.115316E+13
 0.9128
          6908.37
                                                     0.115307E+13
 1.0228
          7150.16
                     1.9280
                               4.3164
                                         2.0000
                                                     0.115298E+13
                                                                      1.1091
                     2.0716
                                         2.0000
                                                     0.115291E+13
 1.1333
          7345.00
                               4.3231
                                                                      1.2364
 1.2443
1.3561
          7521.30
7681.17
                     2.1963
                               4.3298
4.3368
                                         2.0000
                                                     0.115283E+13
                                                                      1.3636
1.4909
                                                     0.115274E+13
 1.4689
1.5826
                     2.4022
2.4920
                               4.3440
4.3517
                                                     0.115266E+13
0.115256E+13
          7828.77
                                         2.0000
                                                                      1.6182
                                          2.0000
          7969.60
                                                                     1.7455
 1.6973
          8105.58
                     2.5762
                               4.3600
                                          2 0000
                                                     0.115246E+13
                                                                     1.8727
                     2.6565
                                          2.0000
                                                                     2.0000
                               4.3689
                                                     0.115234E+13
 1.8128
          8238.64
s4000\_g + 2.0\_m1.0\_t02\_st\_z - 0.50\_a + 0.20\_c + 0.00\_n + 0.00\_o + 0.20\_r + 0.00\_s + 0.00\_mod
s4000_g+2.0_m1.0_t02_st_z-0.50_a+0.20_c+0.00_n+0.00_o+0.20_r+0.00_s+0.00.mod
$4000 g+2.5 m1.0 t02 st z-0.50 a+0.20 c+0.00 n+0.00 o+0.20 r+0.00 s+0.00.mod
s4000 g+2.5 m1.0 t02 st z-0.50 a+0.20 c+0.00 n+0.00 o+0.20 r+0.00 s+0.00.mod
$4250 g+2.0 m1.0 t02 st z-0.50 a+0.20 c+0.00 n+0.00 o+0.20 r+0.00 s+0.00 mod
s4250_q+2.0_m1.0_t02_st_z-0.50_a+0.20_c+0.00_n+0.00_o+0.20_r+0.00_s+0.00.mod
s4250 g+2.5 m1.0 t02 st z-0.50 a+0.20 c+0.00 n+0.00 o+0.20 r+0.00 s+0.00.mod
```

```
Interpolation point : Teff= \, 4200. logg= 2.00 z= -0.50 Optimized interpolation applied for standard composition models
```

4.5 Calculating a synthetic spectrum with TurboSpectrum

To compute a model spectrum with TurboSpectrum for the MARCS model obtained above, the procedure is again largely equivalent to that for SYNTHE:

```
> from ispy3 import turbospec
>
> atoms, abun = ['Na'], [+0.4]
> vturb = 2.0
>
> turbospec.turboinit(5675., 5690., 0.01, initdir='/scratch/soeren')
> turbospec.turbospec(-0.5, 't4200g200m050.marcs', 't4200g200m050MT.asc',
> atoms=atoms, abun=abun, vturb=vturb, initdir='/scratch/soeren')
```

Recall, however, that TurboSpectrum has its own set of line lists and other input data. One limitation is that turboinit() does not support the elements and molecules arguments that can be specified when calling synbeg() to include specific elements and molecules in the spectral synthesis. Another is that TurboSpectrum does not support rotational broadening. For compatibility reasons, turbospec() will accept the vrot argument, but will proceed to ignore it

4.6 Calculating a synthetic spectrum with TurboSpectrum NLTE

The TurboSpectrum NLTE calculations require a line list with additional information, which must be selected by setting the variables atomdir and atomlines in turbospec.py accordingly. The file nlte_ges_linelist_jmg17feb2022_I_II, available at https://keeper.mpdl.mpg.de/d/6eaecbf95b88448f98a4/?p=%2Flinelist&mode=list, covers the range 4200 Å-9200 Å.

The first step is then to interpolate in the MARCS model atmosphere and departure coefficient grids with tsnlte.intpDC(). This replaces the call to marcs.intpatm() in the previous examples. Once this is done, turbospec.turbospec() is called as before, switching on the NLTE option by setting nlte=True:

```
> from ispy3 import tsnlte
> from ispy3 import turbospec
>
> atoms, abun = ['Na'], [+0.4]
> vturb = 2.0
> m, teff, logg = -0.5, 4200., 2.00
```

```
> mass = 1.0
>
> tsnlte.intpDC(teff, logg, m, mass, 't4200g200m050', atoms=atoms, abun=abun)
>
> turbospec.turboinit(5675., 5690., 0.01, initdir='/scratch/soeren')
> turbospec.turbospec(-0.5, 't4200g200m050.marcs', 't4200g200m050MTN.asc',
> atoms=atoms, abun=abun, vturb=vturb, nlte=True, initdir='/scratch/soeren')
```

Figure 2 shows model spectra computed in LTE and NLTE with TurboSpectrum for the region around the Na I lines at 5683 and 5688 Å. The same line list was used for both spectra. There are noticeable differences in the strengths of several lines, with the Na lines in particular becoming stronger in NLTE, hence the Na abundances derived from fitting NLTE spectra will typically be lower than in LTE.

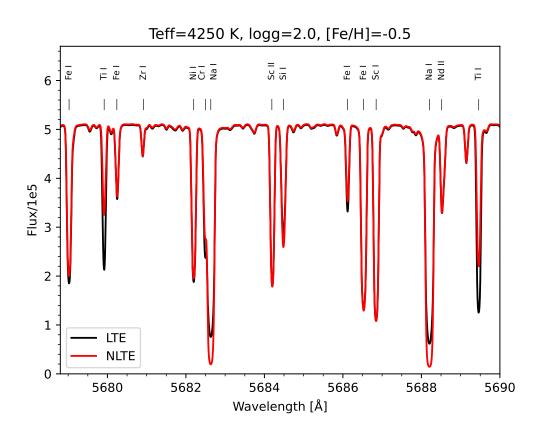


Figure 2: TurboSpectrum LTE and NLTE spectra for the spectral region around the Na I lines near 5680 Å.

4.7 Calculating equivalent widths with WIDTH9

The Kurucz WIDTH9 code is quite quirky and can be used in many different ways. In ISPy3, the function kurucz.calcew() uses a modified version of WIDTH9 to compute the equivalent width of a specified spectral line for a range of input abundances, which are given relative to the abundance of the corresponding element as specified in the header of the ATLAS9 model. Note that WIDTH9 (and hence kurucz.calcew()) will not work with ATLAS12 models.

Naturally, an input line list is required. The format is the same as for the Kurucz atomic line list, but it is not necessary to include the full Kurucz list. A point worth noting is that while the wavelengths are specified in nm in the line list, the wavelength of the line must be given in Å when calling calcew() (as for other ISPy3 functions). If we are interested in the Na I lines near 5680 Å, we might select the following entries:

When calling calcew(), the element is specified according to the same notation used in the Kurucz list, i.e. ZZ.ii where ZZ is the atomic number and ii is the degree of ionisation. Hence, for neutral sodium this would be 11.00. To calculate the equivalent width of the Na 5682.633 Å line for five abundances for the ATLAS9 model computed above, calcew() would then be called as follows:

The results are returned in a dictionary structure with the keys result['abund'] (the abundances for which equivalent widths are computed), result['ew'] (the equivalent widths), result['logew'] (the log of the equivalent widths), result['depth'] (the log(RHOX) corresponding to an optical depth of unity for the line), result['resid'] (the relative residual flux at the line centre) and result['contin'] (the continuum flux).

In the example, the equivalent widths are computed for five abundances, starting from -1.0 dex below the Na abundance in the ATLAS9 model in steps of 0.5 dex. Since the model was computed for [m/H]=-0.5 and the reference solar abundance of Na (defined in absetup.py) is $\log n(\text{Na}) = -5.71$, the equivalent widths are then computed for $\log n(\text{Na}) = -5.71 - 0.5 - 1.0, -5.71-0.5-0.5, \dots, -5.71-0.5+1.0 = -7.21, -6.71, \dots, -5.21$. These values are confirmed by printing out result['abund'], and the corresponding equivalent widths are 68.73 mÅ, 108.80 mÅ, ..., 249.8 mÅ. Kurucz's version of WIDTH9 supports a maximum of 9 individual abundances, but this has been increased to 99 in the modified version used by ISPy3.

4.8 Assigning a τ_{500} depth scale to ATLAS models

In the ATLAS models, the depth is parameterised as a column density RHOX, while the MARCS models use the continuum optical depth at 500 nm, τ_{500} . The function kurucz.tau500() can be used to generate a table with τ_{500} vs. RHOX for an ATLAS model. This is done by running the ATLAS9 code with the line opacity disabled, stopping after one iteration so that the continuum opacities are recomputed but no modifications made to the model. The procedure amounts to a single call to tau500().

```
> from ispy3 import kurucz
> kurucz.tau500('t4200g200m050.A9','t4200g200m050.tau500')
```

This is very fast, since no further iterations are needed and no line opacities are included in the calculation. A few lines of the output file are listed below. The TAUNU column contains the optical depth at 500 nm (as indicated in the header).

```
4200. GRAVITY 2.00000 LTE
WAVELENGTH 500.000 FREQUENCY 5.995850E+14
                               ABTOT
                                                                    SNU
 1 1 037F-02 4 839F-06
                           4 668F-04
                                        9 935F-01 1 937F-08
                                                                1 570F-06
                                                                             1 581F-06
                                                                                        1 092F-08
                                                                                                     9. 971F-07
                                                                                                                 6 116F-01 -3 387F+12
                            4.677E-04
4.687E-04
4.699E-04
                                                    2.196E-08
                                                                             1.581E-06 1.252E-08
1.581E-06 1.547E-08
1.581E-06 1.912E-08
    1.369E-02
                6.391E-06
                                        9.919E-01
                                                                1.569E-06
                                                                                                     9.971E-07
                                                                                                                 7.024E-01 -4.202E+12
   1.782E-02 8.324E-06
                                        9.900E-01
                                                    2.481E-08
                                                                1.566E-06
                                                                                                     9.971E-07 8.696E-01 -5.161E+12
                                                     2.810E-08
                                                                 1.562E-06
   2.970E-02 1.391E-05
                            4.712E-04
                                        9.849E-01
                                                    3.197E-08
                                                                1.558E-06 1.581E-06
                                                                                         2.346E-08
                                                                                                     9.971E-07 1.326E+00 -7.687E+12
   2.460E+02 5.170E+01 2.955E+00 1.857E-03 9.303E-05
                                                                9.303E-05 9.302E-05
                                                                                        -2.457E-09
                                                                                                    1.455E-07 -8.708E+02 -3.688E+16
```

The ATLAS and MARCS models can now be compared directly, as shown in Fig. 3 for the same physical parameters as in Fig. 1. As can be seen, the temperature structures of the models are quite similar, although the ATLAS9 model extends to lower optical depths.

4.9 Writing output to a log file

Many of the functions in ISPy3 can use the Python logging module to write diagnostic output to a log file. To make use of this, the log needs to be configured with a call to basicConfig():

This enables logging, and output will automatically be written to (in this case) ispy3.log. ISPy3 makes use of two levels of logging messages, INFO and DEBUG, where the latter produces more detailed information.

5 Determining abundances from spectral fitting

While the functions discussed above can be called separately by the user, they are also used by higher-level functions in ISPy3 to derive abundances from fits to observed spectra of simple stellar populations (SSPs). As for single stars, the procedure involves two steps: First, the model atmospheres must be defined, and then the spectral fitting can be done.

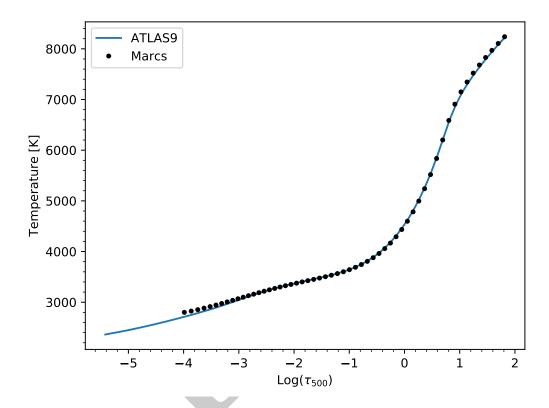


Figure 3: Comparison of ATLAS9 and Marcs models for identical physical parameters ($T_{\rm eff}$ = 4200 K, $\log g$ = 2.0, [m/H]= -0.5.

5.1 Setting up the model atmospheres

To model the integrated-light spectrum of a simple stellar population, ISPy3 needs a list of stellar parameters for the individual spectra that will be co-added. Usually this is stored in a regular text file, which can be read into a data structure in the format used by other tasks with the abutils.rdphys() utility function:

```
> from ispy3 import abutils
> stelpar = abutils.rdphys('hrd.txt')
```

The first few lines of the file hrd.txt might look as follows:

MASS	TEFF	LOGG	RSTAR	WEIGHT	LOGVT	VROT	SYNT	ID
0.513	3950.0	4.7355	0.510	1.625e+04	-0.301	0.0	MT	ISO
0.547	4091.7	4.7010	0.548	1.396e+ 0 4	-0.301	0.0	A9S	ISO
0.585	4288.4	4.6705	0.586	1.324e+04	-0.301	0.0	A9S	ISO
0.622	4536.3	4.6504	0.619	1.151e+04	-0.301	0.0	A9S	ISO

The first line is a header that allows rdphys() to parse the remaining lines in the file. Only the columns TEFF (effective temperature), LOGG (surface gravity), RSTAR (radius of the star in solar radii), and WEIGHT (the weight of the bin) are required. The logarithm of the microturbulent velocity (in km/s) can be specified in the LOGVT column; if this column is not included then rdphys() will assign the value None and the microturbulence must then be specified in some other way (for example, some functions allow a global value for LOGVT to be specified or fitted). The VROT column specifies the rotational velocity (in km/s). If nothing is specified here, no rotational broadening will be applied to the spectra. The SYNT column specifies the combination of model atmospheres and spectral synthesis codes to be used for the modelling, with the following allowed values:

```
A9S = ATLAS9+SYNTHE (default)
A12S = ATLAS12+SYNTHE
MT = MARCS+Turbospectrum
MTN = MARCS+Turbospectrum NLTE
A9T = ATLAS9+Turbospectrum,
```

The A9T combination requires the 2019 or more recent version of TurboSpectrum, and MTN requires TurboSpectrum v 20 or more recent. Note that SYNT, like the other parameters, is specified individually for each bin, so that it is straight forward to model different bins with different model atmospheres and spectral synthesis codes. The MASS column contains the mass, which is not used by rdphys(). Any other column headings are ignored and any value can be listed in the corresponding column. In the example above, the ID column is used to indicate that the bins come from an isochrone.

The rdphys() utility has an option teffsort whose default value is True. This means that the entries read from the input file will be sorted according to their temperature, and the

atmospheres and spectra computed in that order. Since the cooler models and synthetic spectra usually take longer to compute, sorting the input list in this way makes the computation more efficient. The user should be aware, however, that this means that the ordering of the models stored on disc will usually be different from that in which they appear in the input list.

The model atmospheres can now be computed with a call to abutils.hrd2atm():

In this example, the model atmospheres will be computed for a baseline scaling of the abundances of [m/H]=-0.8 relative to the solar composition. Abundances of individual elements are here specified with the atoms and abun variables. For ATLAS models, these are passed on to the ATLAS9 or ATLAS12 codes, and for MARCS they are used to select the appropriate grid. In the latter case, the selection is based on the median abundance ratio of the specified α -elements (O, Mg, Si, S, Ca, Ti). The odfs='A' option is passed on to mkodf() and specifies that α -enhanced ODFs will be used for the ATLAS9 models. The hrd2atm() routine will run the atmosphere calculations for the different HRD bins in parallel and at the end store the models in files named star0000, star0001, etc, with a filename extension according to the type of atmosphere (e.g. star*.marcs or *.A9, etc.). For TurboSpectrum NLTE calculations, files containing interpolated departure coefficients will also be created, with one file for each element per atmosphere model (e.g. star0000_Mg_coef.dat, star0000_Fe_coef.dat, etc). The various temporary files that are generated while computing each model are stored in subdirectories under tmproot (here /scratch/soeren), which must exist before starting the procedure; for each model a separate temporary subdirectory with a unique name will be created and removed again once the calculation is complete, and only the final models are kept.

5.2 Carrying out the fitting

Chemical abundances can now be determined from an integrated-light spectrum by letting ISPy3 vary the input abundances used in the spectral modelling until the best match to the observations is obtained. The fitabun.py module contains two functions for this purpose: fitabun.fit1par() is used to fit for the abundance of one element, while fitabun.fitnpar() can fit for the abundances of multiple elements simultaneously. The abfit.py module contains a utility method fit1p(), which uses fitabun.fit1par() to fit for the abundance of a single element but in addition allows more efficient managing of the input/output data, especially when multiple fits are to be carried out.

5.2.1 General aspects of the fitting procedure

To determine how well the SSP model spectrum matches the data, the continuum levels of the two spectra must be matched. If the data were properly flux calibrated, the matching would amount to a simple scaling of either the data or the model. In practice, achieving a sufficiently accurate flux calibration is often difficult, particularly for echelle spectra. Instead of a single

value, ISPy3 therefore applies a wavelength-dependent scaling $S(\lambda)$, which is defined by fitting a polynomial or a spline function to the ratio of the observed spectrum $F_{\text{obs}}(\lambda)$ and the SSP model $F_{\text{syn}}(\lambda)$.

$$S(\lambda) = \text{fit}(F_{\text{obs}}/F_{\text{svn}}) \tag{3}$$

The observed spectrum and the error spectrum $E_{\rm obs}$ are then scaled by the fit,

$$F_{\text{obs.scl}} = F_{\text{obs}}/S \tag{4}$$

$$E_{\text{obs,scl}} = E_{\text{obs}}/S \tag{5}$$

and the χ^2 of the fit is evaluated as

$$\chi^2 = \sum_{i=1}^n w_i \left(\frac{F_{\text{obs,scl}}(i) - F_{\text{syn}}(\lambda_i)}{E_{\text{obs,scl}}(i)} \right)^2$$
 (6)

where the sum is over n pixels in the observed spectrum with wavelengths λ_i and the w_i are user-specified weights.

The procedure can be customised via several variables defined in fitabun.py. The choice of continuum scaling function is controlled by fitabun.CFITFNC, which can be either 'poly' (polynomial) or 'spline'. The order of the scaling function is defined by CFITORD in both cases, with a value of 3 corresponding to a cubic spline. The choice fitabun.CFITFNC='poly' and fitabun.CFITORD=0 would correspond to a scaling by a single numerical value, CFITORD=1 would allow a linear dependence on wavelength, etc. For CFITFNC='spline', the variable NKNOTS further defines the number of interior knots for the spline.

The pixels to include in the continuum scaling can be specified in various ways:

Leave out pixels that lie below some fraction of the continuum level: The variable CFRAC specifies that only pixels with values above a certain fraction of the "true" continuum level should be included when fitting the scaling function. How to define the "true" continuum is, however, a non trivial matter. In ISPy3 this is done by finding the maximum pixel value within a wavelength range of CDLAM Å, centred on each pixel. This can be defined based on either the model spectrum (CSEL='model') or the data (CSEL='data'). If the 'data' option is selected, then there is a risk that outlying pixel values will artificially boost the continuum level, but if one has a very good quality, high S/N spectrum this option may still be useful. Using the 'model' option avoids problems associated with noise, but is susceptible to imperfections in the synthetic spectra (such as missing lines). The default, CFRAC=0, implies that this scheme is disabled altogether.

Rejecting pixels where the fit is poor: A second scheme relies on rejection of pixel values that have a poor fit. After scaling the data, the standard deviation of the data-model difference is computed, and pixels that deviate by more than CSIGP standard deviations in the positive direction and more than CSIGM standard deviations in the negative direction are rejected. The continuum scaling is then recomputed, taking into account only the remaining pixels, and the procedure is repeated CITER times. The default, CITER=1, means no iterations are performed

and this scheme is disabled. The asymmetric rejection limits are allowed on the grounds that outliers may be more likely in one direction than the other (for example, a significant number of lines may be missing from the line list). On the other hand, one has to pay careful attention to systematic effects if the rejection limits are very asymmetric.

Specifying the continuum regions explicitly: The third option is to provide a list of continuum flags explicitly when calling the fit1par() or fitnpar() functions, described below. The cont option is given in the form [[lam1, lam2, ...], [c1, c2, ...]] where a value < 0.5 for each flag c1, c2, ... specifies that the corresponding wavelength is excluded, while a value of > 0.5 specifies that the corresponding wavelength is included when fitting the continuum. ISPy3 will interpolate in these arrays to determine the flags at the wavelengths of the input data, hence the wavelengths of the cont array do not need to match those of the data.

5.2.2 Fitting a single element with fit1par()

The fit1par() task uses a simple Golden Section search to search for the best-fitting abundance of a single element, within some specified range. The uncertainty range is found by varying the abundance until the χ^2 of the fit has increased by unity relative to the χ^2 of the best fit.

The input to fit1par() is an observed spectrum which must be shifted to the rest frame, as well as the same set of stellar parameters used to set up the model atmospheres. The function abutils.hrd2spec() is used to compute model spectra for each HRD bin and co-adding them to produce the SSP models.

To determine the abundance of sodium from a fit to a spectrum stored in the text file spectrum.txt, relative to a baseline scaling of -0.8 dex relative to the solar abundances, we can call fit1par() as follows:

The first argument to fit1par() is a list of [[lambda1, flux1, err1, weight1], [lambda2, flux2, err2, weight2], ..] values for each pixel, which are typically read in from a text file. The wavelengths should be specified in Å measured in air. In the example we use np.loadtxt() to read the data from the file spectrum.txt, assuming that the radial velocity correction has already been applied. A typical use of the weight column is to exclude certain pixels from the fit by setting weight=0. The same could, of course, be achieved by setting the err to some very large value. The data list is followed by the wavelength window over which the fit should be carried out ([5677., 5695.]). The synthetic spectra will initially

be computed at a default resolving power of $\lambda/\Delta\lambda = 500\,000$, as defined by the global variable fitabun.RESOLP. They will then be smoothed as specified by the sigsm argument and rebinned to the resolution of the observations.

Next follow the stellar parameters. It is critically important that these are *exactly* the same, and appear in the same order, as those that were used to compute the model atmospheres, as fit1par() will use the corresponding already existing model atmospheres (and NLTE departure coefficient files, if applicable) for the spectral synthesis.

The fourth argument specifies the element to be fitted (here 'Na'), followed by the range of abundances to search, relative to the baseline scaling (LogZ). Here, we specify a range of [-1,+1] dex with respect to the scaled baseline abundances, [m/H] = -0.8. The pfix argument specifies which parameters are to be kept fixed in the fit and vfix their corresponding values. In addition to the abundances of specific elements, these parameters can also include LogVT (the microturbulent velocity) and LogZ (the baseline scaling of the abundances relative to solar). In the example, LogVT is set to USER, meaning that the microturbulent velocities specified in the stelpar array will be used. A single numerical value can also be given, which will then be used for the modelling of all the spectra. LogZ is set to -0.8, so the baseline scaling of the abundances is -0.8, while the α -elements are enhanced by 0.3 dex.

The sigsm parameter specifies the dispersion of the Gaussian kernel used to smooth the model spectrum (in Å) and calcerr indicates whether the one-sigma errors should be calculated. If sigsm is set to a range, fit1par() will search for the best-fitting smoothing. The value of sigsm applies to the midpoint of the wavelength range over which the fit is done. Depending on the value of the variable fitabun.SIGCONST, the actual smoothing will be scaled according to the wavelength so as to correspond to a constant velocity (for fitabun.SIGCONST = 'VEL') or kept constant (for fitabun.SIGCONST = 'LAM').

The cont argument can be used to specify the sampling points for the continuum scaling. These are defined separately from the weights and errors defined in the input spectrum, the idea here being that one might want to force exclusion of certain regions of the spectrum when matching the continuum scaling. This could be the case, for example, if the line list is known to be inadequate in some parts of the spectrum.

The fit1par() routine returns the best-fit value of the fitted parameter (vfit), the best smoothing (sigbst), the positive and negative errors (vp, vm), and the reduced χ^2 of the fit (chsq). At the end, we can inspect the results:

```
> print(vfit)
> 0.38699100866857605
> print(sigbst)
> 0.23
> print(vp, vm)
> 0.0107421875 -0.015625
> print(chsq)
> 3.013505275006322
```

where we see that the best-fitting sodium abundance is increased by 0.39 dex with respect to the baseline (LogZ= -0.8), i.e. [Na/m]= $+0.387^{+0.011}_{-0.016}$. Since sigsm was specified as a single value, that same value is returned.

In addition to the elements, fit1par() can also fit for the baseline LogZ, which is treated as a special "element". To fit for the baseline and the smoothing, one would specify:

This will then solve for a scaling of the reference abundances in the range [-3.0, 1.0] with respect to solar composition, with the α -elements enhanced by 0.3 dex, and for the best-fitting Gaussian smoothing with a dispersion between 0.01 and 1.0 Å. Clearly, it would be non-sensical to specify LogZ (or any other element) both in the pfix array and as a parameter to be fitted, but fit1par() has sufficient faith in the user that it does not explicitly check this.

In addition to the return values, fit1par() produces a file with the spectrum and the best fit, with the default name fitabun.txt. This can then be inspected to check the quality of the fit. This file will be overwritten every time fit1par() is called. A different name can be specified with the argument output='filename.txt'.

Similarly to the calculation of model atmospheres, the intermediary files for the spectral synthesis for each HRD bin are stored in temporary directories. These will be created as sub-directories relative to absetup. TMPROOT, which must exist before the calculation is started.

5.2.3 Fitting multiple elements with fitnpar()

The fitnpar() function can fit for abundances of multiple elements. Starting from an initial estimate of the abundances, the downhill simplex method of Nelder & Mead will be employed to search for the best fit. An example call to fit for the abundances of Fe and Na might look as follows:

```
> pfix=pfix, vfix=vfix, sigsm=0.23)
> print(vfit)
```

Here, the elements to be fitted are specified in the fourth argument (['Fe', 'Na']) and the initial guesses are specified in the fifth argument. The other variables are the same as in the fit1par() example. Note that in this example we have also enabled logging to the file fitn.log, so that progress of the fitting can be followed by inspecting the log file. At the end, the last few lines in the log file will contain information about the best fit:

```
FITNPAR (5677.000 - 5695.000):
    Final Fit:
        Fe : +0.008
        Na : +0.386
        Broadening = 0.230
        Reduced chi-square = 2.987
        Number of iterations = 21
```

This shows the wavelength range, the best-fit abundances and smoothing, as well as the reduced χ^2 of the best fit and the number of iterations. The best-fit values are also returned by fitnpar():

```
> print(vfit)
[0.0078125, 0.3859374999999997]
```

where we see that [Fe/m]= 0.008 and [Na/m]= +0.386. Reassuringly, the best-fit [Na/m] is very similar to the value obtained from the fit1par() fit above, and only a small adjustment of the Fe abundance relative to the baseline is required, so that we find [Na/Fe]=+0.378. While fitnpar() can fit for the broadening by giving a range (as for fit1par()), the best-fit value is not returned and must therefore be looked up in the log file.

fitnpar() allows several elements to be fit together as a group. To this end, a list of elements can be specified instead of a single element. To fit for Fe and the α -elements, one could specify

Again, one has to be careful to avoid elements being specified both as parameters to be fit and to be kept fixed.

A limitation of fitnpar() is that it does not provide a way to estimate the errors on the fitted parameters, so it is up to the user to do this in some other way.

5.2.4 Finding and applying the radial velocity shift

There are various ways in which the radial velocity of a spectrum may be measured and applied. ISPy3 includes a function called deltarv() which will search for the radial velocity offset in a given interval that gives the best match between a model and observed spectrum. This usually works best if a rough correction (to within a few km/s) has already been applied, for example based on identification of a few prominent lines (such as H β or the Mgb triplet in the optical range).

Suppose we have estimated the radial velocity shift of a spectrum to be about -10 km/s. We can then use abutils.rvcorrect() to apply this correction to the wavelength scale of the spectrum, stored in the same format used by fit1par() and fitnpar():

```
> import numpy as np
> from ispy3 import abutils
> rv = -10
> data = np.loadtxt('n0104_1d_redu.txt',usecols=(0,1,2,3))
> data2 = abutils.rvcorrect(data, rv)
```

To refine the radial velocity estimate, we can then use deltarv() to search for the best match in a range of ± 10 km/s:

As with fit1par() and fitnpar(), the model atmospheres must already be available and must match the stellar parameters specified in stelpar. In general, it is wise to repeat the procedure for several wavelength windows and use an average as the final estimate of the radial velocity shift. The deltarv() function is a fairly "quick-and-dirty" hack and you might well be able to code something much better yourself!

5.2.5 Integrated-light equivalent widths

The function abutils.hrd2ew() can be used to compute equivalent widths for lines in an integrated-light spectrum. Currently this only works with ATLAS model atmospheres and SYNTHE (the A9S and A12S SYNT options). Similarly to abutils.hrd2spec(), the model atmospheres must be computed before the call to abutils.hrd2ew(). To calculate the equivalent width for the Na I 5682 Å line, the function would then be called as follows:

```
> ew5682 = abutils.hrd2ew(stelpar, 11.00, 5682.633, lines='na.lst',
    minlog=-1.0, dablog=0.5, nablog=5, logvt='USER')
```

As in the calcew() example, the equivalent widths are computed for five abundances (nablog=5), starting from -1.0 dex with respect to the abundance specified in the model atmospheres (minlog=-1.0) and in steps of 0.5 dex (dablog=0.5). The output is an array with the equivalent widths for each abundance value.

5.3 The abfit.py module

The abfit.py module provides a higher-level interface to the fit1par() function, intended to simplify the managing of the input/output data and streamline the procedure for carrying out multiple fits to a given spectrum. The module contains a single class, abfit.setup(), which initialises the setup for a sequence of fits for a single element and opens an output file for the results. Upon initialisation, the class returns an object with the associated method fit1p(), which can then be used to carry out multiple fits for the abundance of the element in different wavelength regions and neatly write the results to the output file. At the end, the output file is closed with a call to the close() method. A slightly modified version of the fit1p() method, called fitZ(), is optimised for the global metallicity fits. Information that is not specific to a particular element, such as the name of the input files, spectral resolution, radial velocity, the continuum definition file, etc., is defined separately via functions that are passed on to abfit.setup() at the time of initialisation. The idea behind this approach is to avoid, as much as possible, that the same information has to be specified in multiple files.

Here's a basic example to fit for the abundance of iron in two spectral bins:

The initialisation method, abfit.setup(), requires several arguments to be specified in order to define the details of the fit: First, the element to be fitted must be given (here 'Fe'), followed by the arrays pfix and vfix that define the parameters to keep fixed and their corresponding values. These will be passed on to fit1par(). The stellar parameters are given in the usual format, and are then followed by the name of the output file for the results ('fitFe.out'). After this follow a number of functions: setfname(key) should return the name of the file containing the input spectrum, as a function of a key is passed as the first argument when calling the fitting method fit1p(). In the example here, the spectra are stored in two files, n1041.txt and n104u.txt, so the key can be 'l' or 'u' and setfname() will construct the full file name accordingly. This is just a way of keeping the calls to fit1p() as concise as possible. The next argument, calcrv(lambda) should be a function that returns the radial velocity correction to be applied to the spectrum, which can be a function of the wavelength. Similarly, calcsigsm(lambda) should return the Gaussian broadening to be applied to the spectrum, again as a function of wavelength. Finally, fncont() should return the name of the file containing the continuum definition data. These functions are generally common to all elements that

one might want to fit, and can therefore conveniently be defined in a separate file, here called setup.py:

```
from ispy3 import abutils

def getsp():
    stelpar = abutils.rdphys('hrd_t10m070p04_hb0104.txt')
    return stelpar

def calcrv(lamc):
    return -19.300

def calcsigsm(lamc):
    return lamc * 4.0e-5

def setfname(ordername):
    return 'n104'+ordername+'.txt'

def fncont():
    return 'cont.dat'
```

Note that a function to read the HRD data has also been defined here. In addition to the arguments discussed above, the abfit.setup method accepts the following additional parameters: vac2air=True/False specifies whether the input wavelengths are given in vacuum (and should thus be converted to air wavelengths). The parameter sigvel=True/False specifies whether the Gaussian broadening of the synthetic spectra should be specified in units of velocity (True) or in units of wavelength (False).

After initialisation, abfit.setup() returns an object that includes the fit1p() method. The first argument to this method is the key that defines the input file, followed by the wavelength range, the name of an output file for the spectral fit (which will be renamed from the fitabun.txt file produced by fit1par()) and the parameters defining the continuum scaling function. This can be either poly or spline, followed by the order of the fitting function (3 for a cubic spline) and, when using a spline, the number of internal knots. In the example, the abundance of iron is fitted in two spectral windows. At the end, the output file is closed with the close() method. This produces the output file fitFe.out with the results of the fit:

```
# ISPy3 0.90.1
# soeren@soeren12.astro.ru.nl /Users/soeren/BTSync2/projects/Abundance/doc/n0104
# Python 3.7.3 Darwin 17.7.0 2020-02-06 11:28:32.050801
# Fitting: Fe
# BestFit
                 Lam1
                         Lam2
                                Result err+
                                                                (sig)
                                                        Chisq
                5400.00 5420.00 -0.079 +0.006 -0.008
fitFe540.txt
                                                        3.859
fitFe542.txt
                5420.00 5460.00 -0.122 +0.005 -0.005
                                                        3.434
# Done at 2020-02-06 12:15:34.133749
```

A second method, fitZ() is very similar to fit1p(), but optimised to fit parameters over a broader wavelength range. The default values for the arguments to the method differ somewhat: the default range of the fitrng argument is [-4.0, 0.5] (appropriate for the LogZ fits), and the calcerr argument is set to False as default. As one typical use of fitZ() is to determine the best-fitting broadening (by setting the sigsm argument to a range), the broadening applied to the spectrum will be written to the output file as an extra column (whether it was fitted or not).

6 User configurable variables

Some of the modules in ISPy3 contain a variety of variables that can be modified by the user. Some of these may be set automatically from higher level routines.

absetup.py

absetup.binpath: A string pointing to the directory in which the external binaries are located.

absetup.catpref: This variable should point to the root directory for the various data files needed by the model atmosphere and spectral synthesis codes. More details are given in the descriptions of ATLAS/SYNTHE (sec. 3.1) and Marcs/TurboSpectrum (sec. 3.2).

absetup. TMPROOT is the name of a directory for storage of temporary files. The SYNTHE spectral synthesis code in particular relies heavily on reading and writing large amounts of temporary data, so for optimum performance it is important that this directory resides on a local disc that allows fast access to read/write operations.

absetup.MAXPROC is the maximum number of processes to be run in parallel.

absetup.stdabun contains the reference abundance scale. The default is Grevesse & Sauval (1998), but a few other options are available. Negative numbers are logarithms of fractional number densities relative to the *total*, following the convention in the ATLAS models.

fitabun.py

fitabun.RESOLP = 5e5 # Default resolving power for synthetic spectra, defining the wavelength steps ($\Delta \lambda = \lambda / \text{fitabun.RESOLP}$)

fitabun. EXPAND = $1.0 \, \text{#}$ When computing the synthetic spectra, add EXPAND extra Å at both ends of the fitted wavelength range. This ensures that wings of lines beyond the fitted range are included in the modelling. It should hardly ever be necessary to change this.

fitabun. CFITORD = 3 # Order of continuum fitting function. Usually this will be set when calling fit1p().

fitabun.CFITFNC = 'spline' # Type of function used to fit the continuum shape. Can be 'spline' or 'poly'. Will usually be set when calling fit1p().

fitabun.NKNOTS = 5 # Number of knots when CFITFNC = 'spline'.

fitabun.CFRAC = 0.0 # Only use pixels with $F(\lambda)/F_{\rm max}$ > CFRAC when matching the continuum scaling of the spectra. $F_{\rm max}$ is evaluated over a wavelength range from λ - CDLAM/2 to λ + CDLAM/2. If CFRAC = 0 (default) then all pixels are used.

fitabun.CDLAM = 5.0 # Width of wavelength range for computing F_{max} .

fitabun.CITER = 1 # Number of iterations for continuum fitting

fitabun.CSIGP = 2.0 # Upper clipping threshold for continuum fitting

fitabun.CSIGM = 1.5 # Lower clippling threshold for continuum fitting

fitabun.CSEL = 'MODEL' # Select continuum points based on 'MODEL' or 'DATA'

fitabun.NPINT = 1 # Number of interpolation points per pixel in the observed spectrum. The default value of 1 means that the synthetic spectrum is simply interpolated at the central wavelength of each pixel *after* smoothing to account for velocity broadening and instrumental resolution. If the observed spectra are severely undersampled, NPINT can be increased in order to sample the synthetic spectrum at more points within each pixel before rebinning to the resolution of the data.

fitabun. SMOOTHFNC = 'gaussian' # Type of profile used to smoothen the model spectra. Current options are 'gaussian' or 'uniform'.

fitabun.SIGCONST = 'VEL' # Keep Gaussian sigma constant in VELocity/LAMbda space

fitabun.NROT = 1 + Number of individual sin i values over which to average when synthesising spectra of rotating stars. The population of rotators is assumed to have an isotropic distribution of orientations.*Note:*this only works with SYNTHE, not with TurboSpectrum.

kurucz.py

kurucz.LOGTAU0 = -6.875 # Logarithm of Rosseland mean optical depth of outermost layer for ATLAS models (both ATLAS9 and ATLAS12)

kurucz.DLOGTAU = 0.125 # Logarithmic step in Rosseland depth per layer (both ATLAS9 and ATLAS12)

kurucz.NDEPTH = 72 # Number of layers in atmosphere model (both ATLAS9 and ATLAS12). The default values of these variables will thus produce a model with a maximum Rosseland mean depth of 100.

kurucz.AFE = 'A' # Use alpha-enhanced ('A') or solar-scaled ('S') initial models when starting from the pre-existing set. This will be set automatically when model atmospheres are computed via abutils.hrd2atm(). At any rate, the value is usually not critically important since this only concerns the starting model for the iterations.

kurucz.A9_NREP = 1 # Number of times to repeat 15 extra ATLAS9 iterations. A model calculation always starts with 15 iterations, hence the default value implies a total of 30 iterations.

kurucz.A9_INIT_EXT = None # Filename extension for ATLAS9 initial models

kurucz.A9_FROMSCRATCH = False # Calculate ATLAS9 models from scratch instead of starting from a pre-existing model

kurucz.A12_NREP = 2 # Number of times to repeat 15 extra ATLAS12 iterations. A model calculation always starts with 15 iterations, hence the default value implies a total of 45 iterations.

kurucz.A12_INIT_EXT = None # Filename extension for ATLAS12 initial models. A typical use of this option would be to first compute a set of ATLAS9 models via a call to abutils.hrd2atm()), and then use these as starting points for ATLAS12 models. In this case, one would set kurucz.A12_INIT_EXT = '.A9'. The default value, None, implies that the initial model is selected from the pre-existing set.

syntherk.py

```
syntherk.atomdir # Directory with atomic line lists for SYNTHE
syntherk.moldir # Directory with molecular line lists for SYNTHE
syntherk.preddir # Directory with predicted lines for SYNTHE
```

turbospec.py

```
babsma_marcs_lte = 'babsma_lu.14'  # babsma binary for LTE MARCS models
bsyn_marcs_lte = 'bsyn_lu.14'  # bsyn binary for LTE MARCS models
babsma_nlte = 'babsma_lu.20.1'  # babsma binary for NLTE MARCS models
bsyn_nlte = 'bsyn_lu.20.1'  # bsyn binary for NLTE MARCS models
babsma_atlas = 'babsma_lu.20.1'  # babsma binary for ATLAS models
bsyn_atlas = 'bsyn_lu.20.1'  # bsyn binary for ATLAS models
```

tsnlte.py

```
NLTEelem = [
    [1, 'H', 'nlte', 'atom.h20', 'auxData_H_MARCS_May-10-2021.txt', 'NLTEgrid_H_MARCS_May-10-2021.bin'],
    [11, 'Na', 'nlte', 'atom.na102', 'auxData_Na_MARCS_Feb-20-2022.dat', 'NLTEgrid4TS_NA_MARCS_Feb-20-2022.bin'],
    [26, 'Fe', 'nlte', 'atom.fe607a', 'auxData_Fe_MARCS_May-07-2021.dat', 'NLTEgrid4TS_Fe_MARCS_May-07-2021.bin'],
```

List of elements to be treated in NLTE. For each element included in the list, the corresponding model atoms, auxData*, and NLTEgrid* files must exist in the cats/NLTE-TS subdirectory (Sec. 3.2.4).

7 Functions

Module kurucz.py

mkodf(m,atoms=[],abun=[], odfs='S'): Interpolate in the library of pre-computed opacity distribution functions (ODFs) to define an ODF for a specified metallicity m. Parameters:

```
m: The metallicity of the desired ODF in logarithmic units, relative to the Sun (i.e. m = 0 is solar metallicity, m = -2 means 1/100 solar metallicity, etc) odfs='S': Composition of the ODF. Can be 'S' (solar-scaled) or 'A' (alpha-enhanced) atoms=[], abun=[]: ignored.
```

mkatm(teff, logg, m, atmname, workdir='.', wait=True, atoms=[], abun=[],
nlte=False, vturb=0.0, firstguess=None): Calculate an ATLAS9 model atmosphere for specified physical parameters. Before calling mkatm(), an interpolated ODF
must be generated via a call to mkodf().

Parameters:

teff: The effective temperature of the model (in K)

logg: The logarithm of the surface gravity of the model (in cgs units)

m: The metallicity of the model in logarithmic units relative to the Sun. More precisely stated, m is the logarithm of the default scaling factor for the abundances in the reference (solar) abundance scale specified in absetup.stdabun.

atmname: A string specifying the filename for the computed model. In fact, two output files will be produced, namely the model atmosphere itself (with an extension '.A9' appended to atmname) and a diagnostic output file produced by the atlas9mem.exe code with more details about the calculation (with extension '.out' appended to atmname) that can be useful for checking convergence of the models.

workdir='.': The directory in which the actual calculation is carried out. At the end of the calculation, only the atmosphere and diagnostic output file are copied back to the current directory. When multiple atmospheres are being computed in parallel by ISPy3, temporary directories will automatically be set up for each model to keep the files from being overwritten.

wait=True: Wait for the calculation to be completed? The default, True, means that the function will wait until the atmosphere has been computed. For wait=False the model calculation will be launched as a background process and mkatm will return a process ID that allows ISPy3 to keep track of calculations running in parallel.

atoms=[]: An array of elements for which the composition is specified

abun=[]: The abundances for the elements specified in atoms, relative to the scaling factor m. While the line opacity is computed via the ODFs by ATLAS9, modifying the abundances of individual elements can still have a significant effect on an ATLAS9

model. This is the case, for example, for elements such as Mg and Si that are significant electron donors and therefore affect the H⁻ opacity, which is a dominant continuum opacity source in solar-type and cooler stars.

nlte=False: Include NLTE effect in the model atmosphere calculation? Default is False, since the NLTE treatment in ATLAS9 is anyway not very realistic.

vturb=0.0: The microturbulent velocity, in km/sec. This is used to select the appropriate ODF. ODFs are computed for 0, 1, 2, 4, and 8 km/sec, of which the closest match will be selected.

firstguess=None: Explicitly specify an existing model to use as an initial guess. If None, a model from the pre-existing set will be selected.

mkatm_a12(teff, logg, m, atmname, workdir='.', wait=True, atoms=[], abun=[], nlte=False, vturb=2.0, firstguess=None): Calculate an ATLAS12 model atmosphere for specified physical parameters. Since ATLAS12 uses opacity sampling, it is not necessary to first call mkodf(). The parameters have the same meaning as for the ATLAS9 models (mkatm()), with a few differences in the detailed behaviour:

- The output model has extension '.A12' and the diagnostic output file is named by adding 'b.out' to atmname.
- Any changes in the abundance patterns specified with atoms=[] and abun[] are fully accounted for in the model.
- The specified microturbulent velocity is taken into account self-consistently by the atlas12.exe code in the opacity calculations.

tau500(atmname, tauname): Use atlas9mem.exe to generate a table with the continuum optical depth at 500 nm (τ_{500}) for an existing ATLAS9 or ATLAS12 model.

calcodf(m, odfname, atoms=[], abun=[]):: Use the dfsynthe.exe code to calculate ODFs for a specified composition. The m, atoms, and abun variables have the same meaning as described above. The ODFs will be computed for microturbulent velocities of 0, 1, 2, 4, and 8 km/sec for temperatures between 1995 K and 199526 K.

Module fitabun.py

```
fit1par(specobs, specwin, stelpar, pfit, prange, pfix=[], vfix=[], tol=0.001,
calcerr=True, sigsm=[0,1.0], cont=None):
```

```
fitnpar(specobs, specwin, stelpar, pfit, pinit, pfix, vfix, tol=0.001, sigsm=[0,1.0], cont=None):
```

Module abutils.py

hrd2atm(stelpar, mh, atoms=[], abun=[], tmproot='.', nlte=False, odfs='S'): Compute stellar atmosphere models for a given set of stellar parameters. For a more detailed description of this function, see Sec. 5.1.

stelpar: List of stellar parameters defining the input population (see Sec. 5.1).

mh: Log of the overall abundance scaling relative to solar composition.

atoms and abun: List of individual elements and their abundances, relative to mh

tmproot='.': Directory under which temporary files will be stored during the calculations. Specifically, a unique temporary subdirectory will created under tmproot for each individual entry in stelpar and will be deleted again when the calculations are complete.

nlte=False: Compute NLTE models? (currently ignored)

odfs='S': Composition of opacity distribution functions (ODFs) used when computing ATLAS9 models. Valid options are 'S' (scaled-solar) and 'A' (alpha-enhanced).

hrd2spec(stelpar, outspec, synlimits, mh, atoms=[], abun=[], logvt='USER', nrot=1, tmproot='.', initialize=True, initdir='.'): Calculate an integrated-light model spectrum for a stellar population. The model atmospheres are assumed to exist already. In general, a call to hrd2spec() must therefore be preceded by a call to hrd2atm() to set up the model atmospheres, using the exact same stellar parameters (stelpar). Be warned that hrd2spec() does not check that any pre-existing model atmospheres actually do correspond to the stellar parameters specified in stelpar - it is up to the user (you!) to make sure this is the case. If a set of models exist, but do not match the parameters in stelpar, non-sensical results are the most likely outcome. If some models are missing the procedure will most likely crash. It is also a good idea to make sure the model atmospheres are computed for the same chemical composition as that assumed in the spectral synthesis, but again hrd2spec() will happily use any set of model atmospheres provided, as long as they are of the correct flavour (ATLAS/MARCS) (although you may see a warning in the log file if there is a big difference in overall metallicity).

stelpar: List of stellar parameters defining the input population (see Sec. 5.1).

outspec: Name of output file in which the model spectrum will be stored.

synlimits: A list containing the wavelength limits and steps, e.g. synlimits=[5000,5100,0.01] to compute a model spectrum between 5000 Å and 5100 Å in steps of 0.01 Å.

mh: Log of the overall abundance scaling relative to solar composition.

atoms and abun: List of individual elements and their abundances, relative to mh

logvt='USER': Log of the microturbulent velocity (in km/sec). Can be a numerical value or logvt='USER' in which case the values specified in stelpar will be used (usually this will be the preferred approach).

nrot=1 Number of individual sin *i* values over which to average when synthesising spectra of rotating stars. The default nrot=1, means that the values specified in the stelpar list will be used directly. Otherwise the population of rotators is assumed to have an isotropic distribution of orientations and spectra computed for the nrot values of sin *i* will be weighted accordingly. *Note:* this only works with SYNTHE, not with TurboSpectrum

tmproot='.': Directory under which temporary files will be stored during the calculations. Specifically, a unique temporary subdirectory will created under tmproot for each individual entry in stelpar and will be deleted again when the calculations are complete.

initdir='.': Directory containing initialisation files needed for the spectral synthesis calculations.

initialize=True: Generate the initialisation files for the spectral synthesis? For spectral synthesis with SYNTHE and TurboSpectrum, calls to syntherk.synbeg() and turbospec.turboinit() will be made, respectively. This step can be omitted (setting initialize=False) if these files have already been generated through a previous call to hrd2spec() with the same stellar parameters and wavelength limits. This will save a lot of time (especially for SYNTHE).

hrd2ew(stelpar, elem, lam, nablog=1, minlog=0, dablog=0, logvt='USER', lamtol=0.001, lines=None): Calculate integrated-light equivalent widths for a spectral line. As for hrd2spec(), the model atmospheres must be set up with a call to hrd2atm() before invoking hrd2ew(). This function uses syntherk.calcew() to compute the equivalent widths via a modified version of the Kurucz WIDTH9 code and currently does *not* support MARCS atmospheres. The output from hrd2ew() is a list of the equivalent widths corresponding to each abundance step as specified by minlog, nablog, and dablog.

stelpar: List of stellar parameters defining the input population (see Sec. 5.1).

elem: Element code in the format ZZ.ii for atomic number ZZ and ionisation ii, e.g. 11.00 for neutral sodium. See Sec 4.7.

lam: Wavelength of the line, in Å.

nablog=1: Number of abundance steps for which to calculate equivalent widths (up to a maximum of 99).

minlog=0: Minimum abundance of the element, relative to the composition specified in the input model atmospheres.

dablog=0: Step between the abundances for which equivalent widths are computed (see example in Sec. 5.2.5).

logvt='USER': Log of the microturbulent velocity (in km/s). Can be a numerical value or logvt='USER' in which case the values specified in stelpar will be used.

lamtol=0.001: Wavelength tolerance when matching the specified wavelength to the line list.

lines=None: File with line data in the Kurucz format. The (not very sensible) default is to use the Castelli list, assumed to be located in absetup.catpref+'/cats/Castelli/atoms_SLr.

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A ISPy3 installation on the coma cluster at Radboud University

A general guide to the coma computing cluster can be found on the internal Wiki of the astro department, https://astro.ru.nl/wiki/internal/computing/cluster (you will need to log in with your science username and password).

What follows here is a brief guide that should help local users getting started with ISPy3.

The ISPy3 python modules are available in the directory

/vol/astro-constant/slarsen/ISPy3/python/ispy3/

on the coma cluster. To use these, make a python directory in your home directory (if you do not have one already) and then copy the ISPy3 python files to a subdirectory python/ispy3. Make sure you have your \$PYTHONPATH environment variable set to point to the python directory (usually by having a line such as

```
export PYTHONPATH=${HOME}/python
```

in your .profile file). You should now be able to import ISPy3 modules into python. As of this writing (Sep 2023), the default python version on coma is still 2.7.*, so python code using ISPy3 should generally be called explicitly with

```
> python3 mycode.py
```

Remember that the head node should never be used for anything requiring significant computing power. Large jobs should be submitted via the slurm scheduler (typically with the sbatch command). Smaller jobs can be run interactively on coma01. See the Wiki for details. Note that the network filesystems (/vol/astro*) are mounted read-only on the coma?? compute nodes. Any output data (including temporary data produced by the various ISPy3 functions) should be written to the local /scratch disc on each node and copied back afterwards.

The binaries and source for the Kurucz codes and TurboSpectrum, as well as the various data files needed by these codes, can be found in the directory /vol/astro-constant/slarsen/. The absetup.py file should contain the following two lines:

```
binpath = '/vol/astro-constant/slarsen/ISPy3/gbin/'
catpref = '/vol/astro-constant/slarsen'
```

which tell ISPy3 where to find the binaries and data files. There are in fact two versions of the binaries: the example above points to binaries compiled with gfortran, there are also binaries compiled with ifort in /vol/astro-constant/slarsen/ISPy3/ibin/. The latter must be used when running ISPy3 on the Linux PC workstations (where, apparently, some of the dynamically linked libraries needed by the gfortran binaries are missing).

If you wish to compile and install the binaries yourself, you can find the source codes in the subdirectories under /vol/astro-constant/slarsen/ISPy3/src/:

```
/vol/astro-constant/slarsen/ISPy3/src/kurucz:
```

The ifort versions of the Kurucz codes. To compile, simply execute the script ./compile.csh.

You will need a correctly configured Intel Fortran compiler on your system.

/vol/astro-constant/slarsen/ISPy3/src/kurucz_gfortran:

Versions of the Kurucz codes that I have modified to make them compatible with gfortran. Note that these are *not* the same versions that can be downloaded from the Castelli site. To compile, use ./gcompile.csh which will put the executables in the gbin/ directory. You will need a working gfortran compiler on your system.

/vol/astro-constant/slarsen/ISPy3/src/Turbospectrum:

The 2014 and 2019 versions of TurboSpectrum. ISPy3 uses the 2014 version with Marcs models (as it is faster), but the 2019 version is needed if you want to use TurboSpectrum with ATLAS models. The 2014 versions of the binaries should be renamed to bsyn_lu.14 and babsma_lu.14 and the 2019 versions to bsyn_lu.19 and babsma_lu.19.

/vol/astro-constant/slarsen/ISPy3/src/interpol_marcs:

The interpol_marcs code for interpolation in the MARCS atmosphere grid. Note that there are two slightly different versions, interpol_modeles_lin_sl.f and interpol_modeles_sl.f, for linear and optimal interpolation, respectively. ISPy3 will figure out which version to use, depending on the temperature. To compile, run ./compile.csh (to compile with ifort) or ./gcompile.csh (gfortran).