

Experiments and nonlocal continuum modeling of the size-dependent fracture in elastomers^[1]

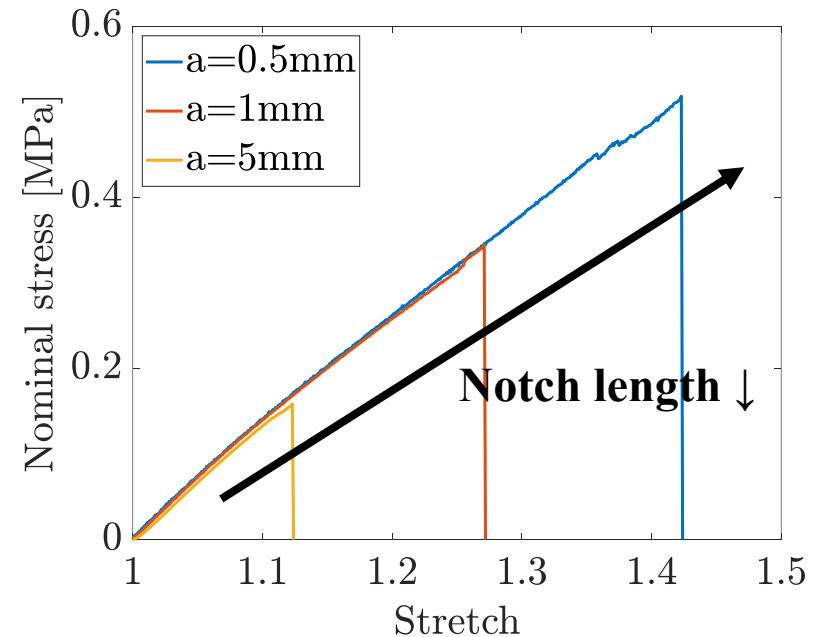
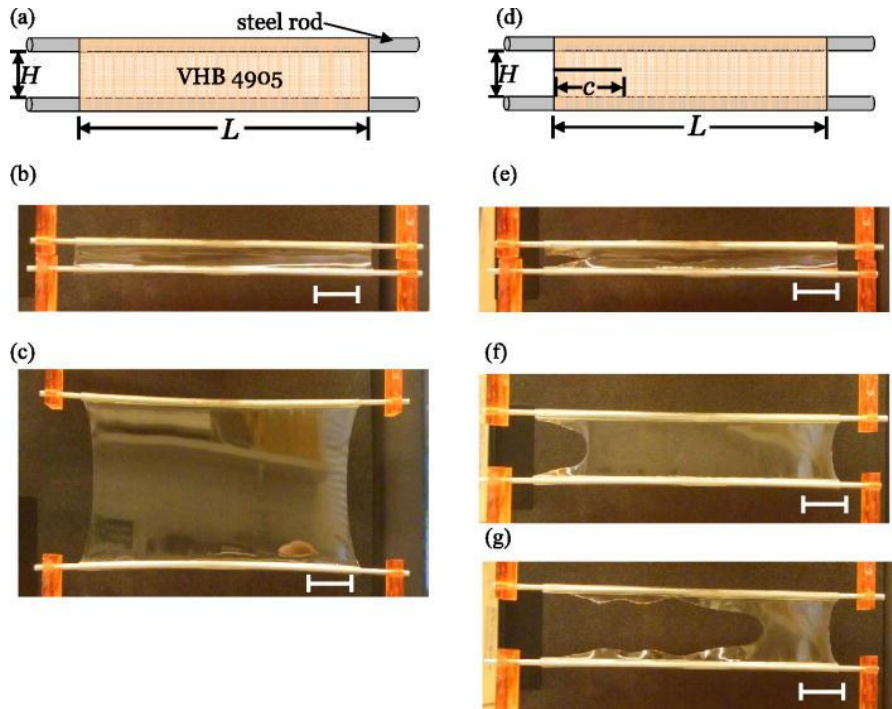
Jeongun Lee, Jaehee Lee and Hansohl Cho*

Korea Advanced Institute of Science and Technology

* hansohl@kaist.ac.kr

Fracture in elastomers

- Extreme, nonlinear deformation → nonlinear fracture
- Influenced by the size of flaws; **the size-dependent fracture**^[1,2]
 - Rupture stretch increases as the specimen size decreases

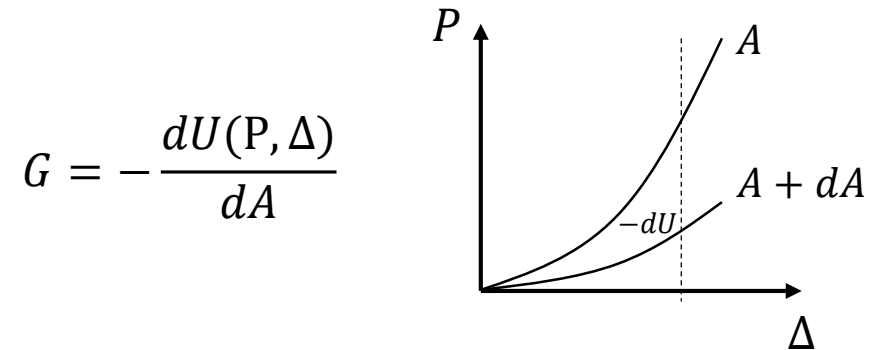


a) The presence of flaws impacts the fracture behavior^[3]

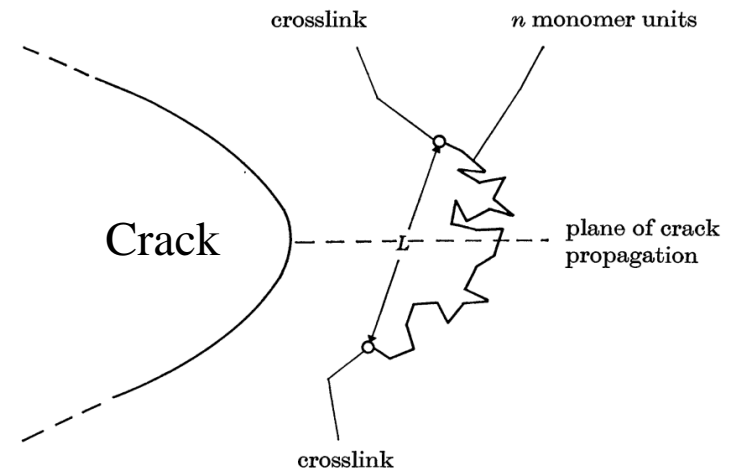
b) Size-dependent fracture in polydimethylsiloxane (PDMS) specimens

Fracture in elastomers

- Occurs when ...
 - Macroscopically, **G reaches Γ**
 - Griffith theory^[4,5]
 - G: Energy release rate
 - Γ : Fracture energy
 - Microscopically, **ε_R reaches ε_R^f**
 - Lake-Thomas theory^[6-8]
 - ε_R : Internal energy
 - ε_R^f : critical internal energy;
bond dissociation energy
- These approaches are compatible (Lake and Thomas ^[6])



$$G = -\frac{dU(P, \Delta)}{dA}$$



Objectives

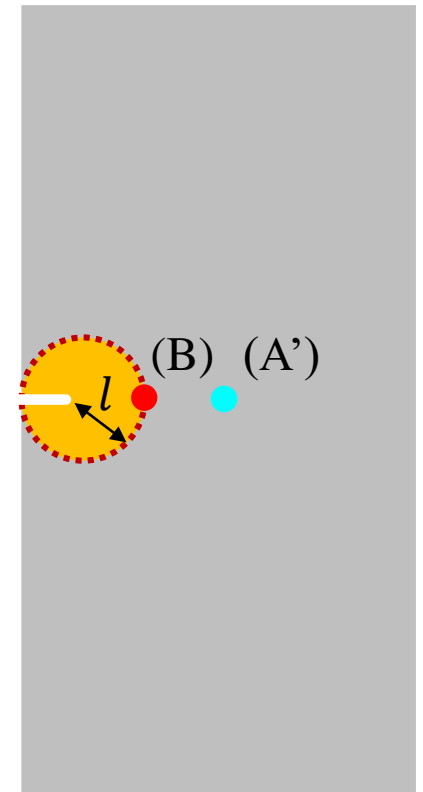
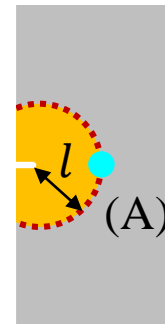
- Predicting the **size-dependent fracture** in elastomers^[1]
 - Experiments and numerical simulations were carried out
- **Internal energy-driven** fracture criterion; inspired by the Lake-Thomas model^[6-8]
- Using the **phase-field model** rooted in the gradient-damage theory^[8-12]
 - Mesh-insensitive crack propagation process
 - The internal energy-driven fracture criterion
 - Thermodynamics of the damage and fracture

Size-dependent fracture & Fracture process zone

- Fracture process zone
 - Where the polymer chains rupture = Where the dissipation mainly occurs
- Stress at point (B) is larger than those at (A) and (A')
 - $\sigma_A = \sigma_{A'} < \sigma_B$
- → Free energy at point (B) is larger than those at (A) and (A')
 - $\psi_A = \psi_{A'} < \psi_B$
- → ψ_B reaches the critical energy earlier than ψ_A
- → **The larger specimen ruptures earlier**

The size of fracture process zone^[1,2,13,14]:

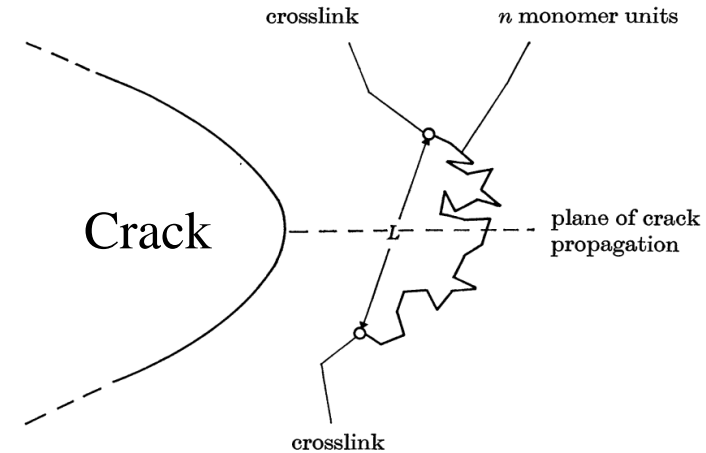
$$l = \frac{\Gamma}{W^*} = \frac{\text{Fracture energy}}{\text{Critical deformation energy}}$$



Nonlocal continuum modeling

- 1. The damage $d \in [0,1]$
 - $d=0$: intact
 - $d=1$: fully damaged
- Internal energy-driven** fracture criterion
 - Inspired by the Lake-Thomas model^[4]
 - Fracture = **Scission of polymer chains**
- Governing equations^[8]
 - Macroforce balance $\text{Div } \mathbf{T}_R = 0$
 - Microforce balance $\zeta \dot{d} = 2(1 - d)\mathcal{H}_R - \hat{\varepsilon}_R^f(d - l'^2 \Delta d)$

a)^[6]

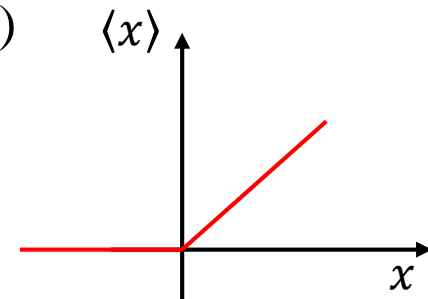


History function;
the fracture criterion

$$\mathcal{H}_R = \left\langle \varepsilon_R^0 - \varepsilon_R^f / 2 \right\rangle, \quad \text{where } \langle x \rangle = \begin{cases} x & \text{if } x > 0 \\ 0 & \text{if } x \leq 0 \end{cases}$$

Internal energy

b)



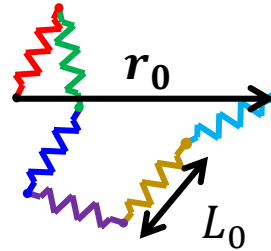
Nonlocal continuum modeling

- **Internal energy** should be considered → **Bond stretch**^[7,8]

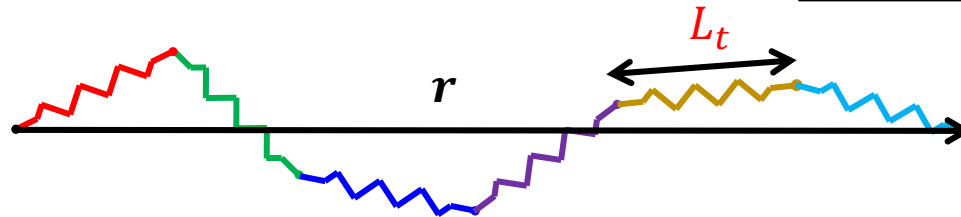
- Deformation = Chain configuration change + stretching of molecular bonds

$$\psi_R = \underbrace{(1-d)^2 \left[\frac{1}{2} N n E_b (\lambda_b - 1)^2 + \frac{1}{2} K (J - 1)^2 \right]}_{(1-d)^2 \varepsilon_R^0; \text{ Damage acts on the internal energy only}} + \underbrace{N k_b \theta n \left[\frac{\bar{\lambda} \lambda_b^{-1}}{\sqrt{n}} \beta + \ln \left(\frac{\beta}{\sinh \beta} \right) \right]}_{-\theta \eta_R; \text{ Entropic energy}} + \underbrace{\frac{1}{2} \varepsilon_R^f l^2 |\nabla d|^2}_{\text{Nonlocal energy}^{[8]}}$$

a) Reference configuration



a) Deformed configuration



$$\begin{aligned} \bar{\lambda} &= \frac{|r|}{|r_0|} \\ \lambda_b &= \frac{L_t}{L_0} \\ \mathcal{L}(x) &= \coth x - \frac{1}{x} \\ \beta &= \mathcal{L}^{-1} \left(\frac{\bar{\lambda} \lambda_b^{-1}}{\sqrt{n}} \right) \end{aligned}$$

Nonlocal continuum modeling

- 2. Phase-field model rooted in the gradient-damage theory^[8-12]

- “Diffusive damage zone”

Intrinsic length scale l'

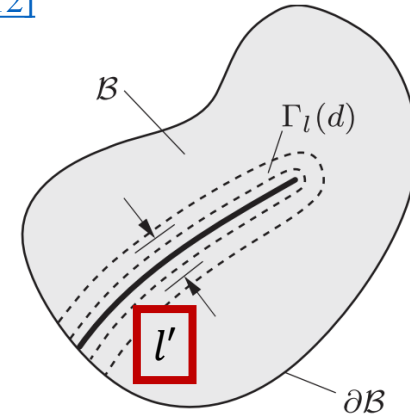
- Microforce balance $\zeta \dot{d} = 2(1 - d)\mathcal{H}_R - \hat{\varepsilon}_R^f(d - l'^2 \Delta d)$

History function;
the fracture criterion

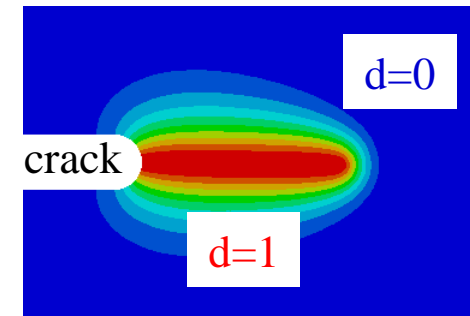
- The **intrinsic length scale l'** \rightarrow the size of diffusive damage zone

- A numerical parameter; ambiguous physical meaning

a)^[12]



b)



Crack propagation;
at reference configuration

Nonlocal continuum modeling

- Assumption^[1]: Diffusive damage zone = Fracture process zone
 - Regions of the damage evolution and the dissipation

- The size of fracture process zone

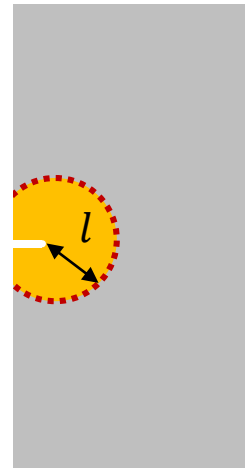
$$= \frac{\Gamma}{W^*} = \frac{\text{Fracture energy}}{\text{Critical deformation energy}} \rightarrow \text{Intrinsic length scale}$$

→ Identify the intrinsic length scale l **from experiments**

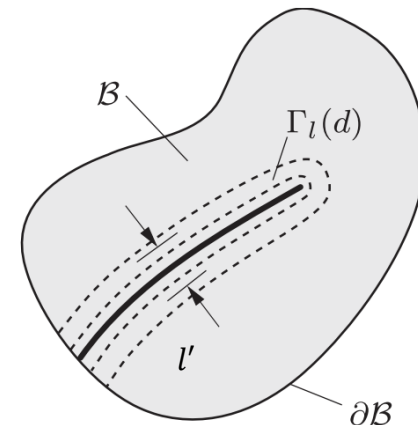
→ Apply to the phase field model

→ Predict the **size-dependent fracture** by numerical simulations^[1]

a) Fracture process zone

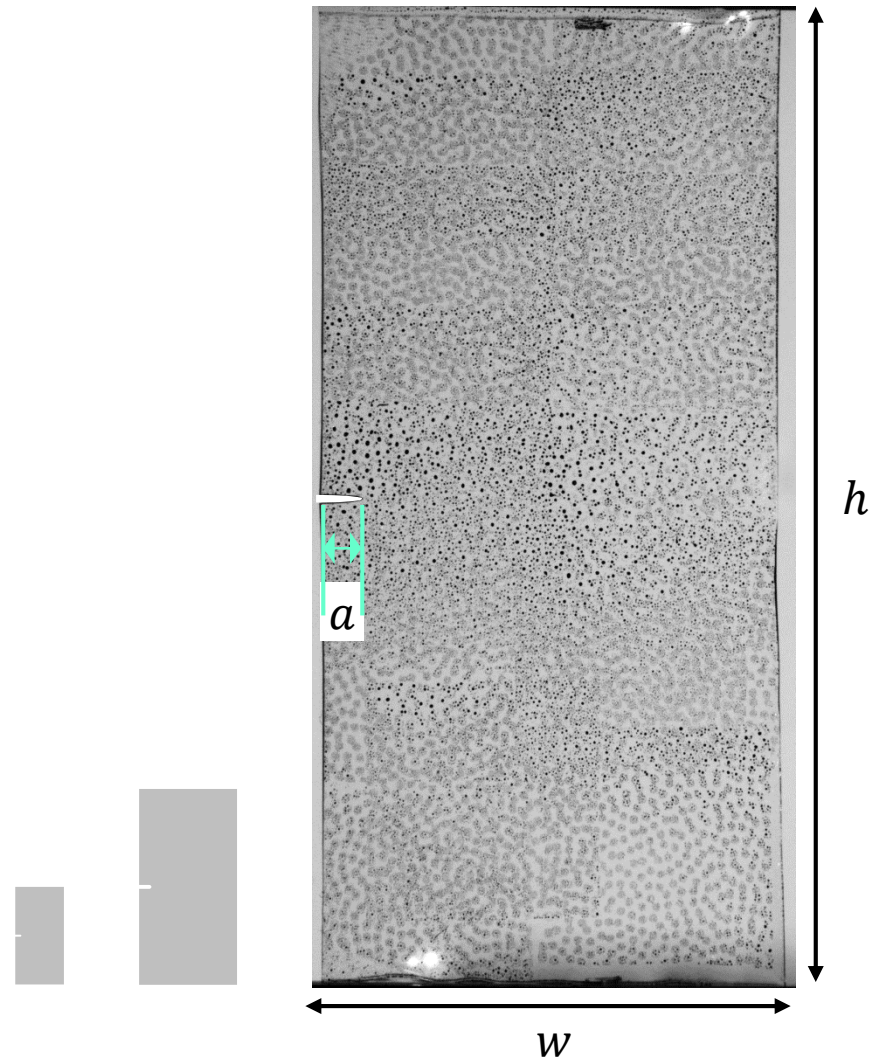


b)^[12] Diffusive damage zone



Experimental procedures^[1]

- Geometries
 - $a = \{0.5, 1, 5\}$ mm
 - $w = 10a$, $h = 20a$, specimen thickness: 0.5mm
 - $w = \{5, 10, 50\}$ mm
 - $h = \{10, 20, 100\}$ mm
- Materials
 - PDMS
 - TangoPlus (3D-printed elastomer)
- Strain rate 0.01 s^{-1} , temperature $\sim 21^\circ\text{C}$
- Digital image correlation (DIC) analysis
 - Strain fields from experiments



The intrinsic length scale l

- $l = \frac{\Gamma}{W^*} \rightarrow$ **Experimentally identified intrinsic length scale**^[1]
- Γ : Fracture energy
 - from **notched** specimens
- W^* : Critical deformation energy
 - from **unnotched** specimens

PDMS

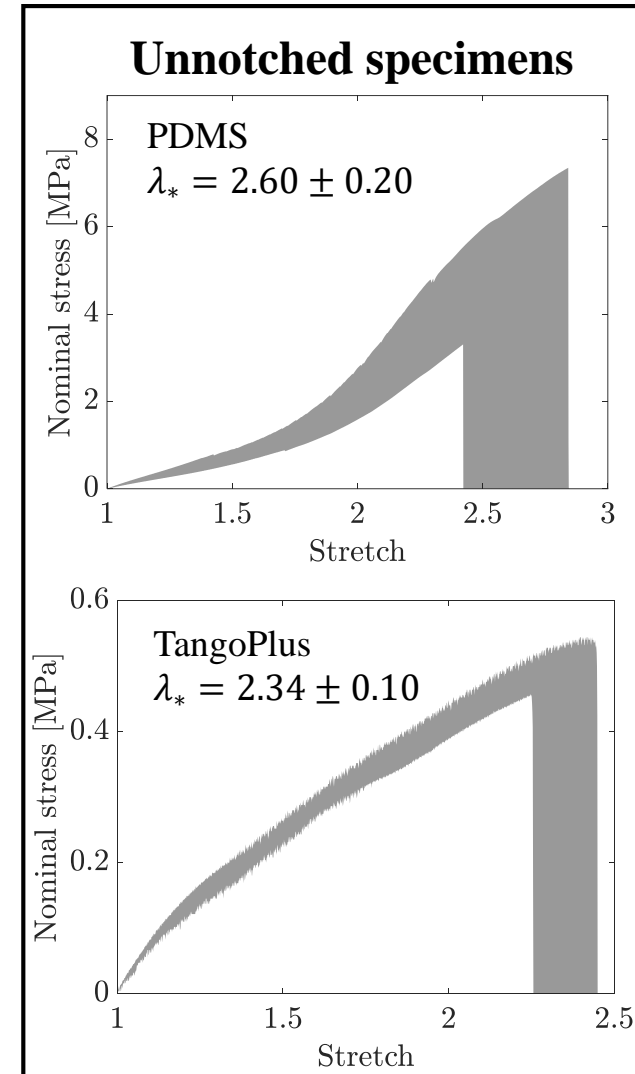
$$\Gamma \approx 0.25 \text{mJ/mm}^2, W^* \approx 2.7 \text{mJ/mm}^3$$

$$\rightarrow l \approx 0.08 \text{mm}$$

TangoPlus

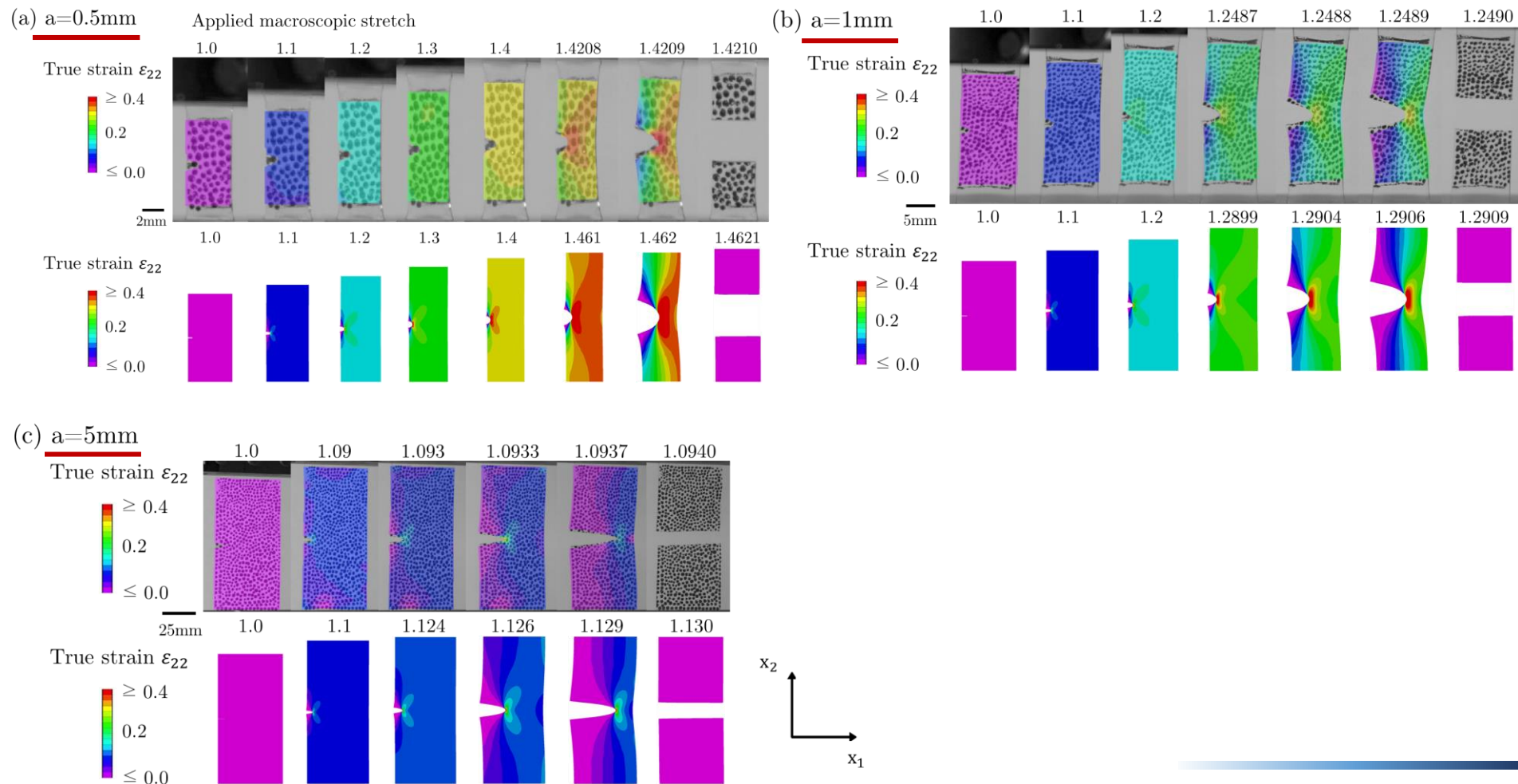
$$\Gamma = 0.5 \text{mJ/mm}^2, W^* \approx 0.45 \text{mJ/mm}^3$$

$$\rightarrow l \approx 1 \text{mm}$$



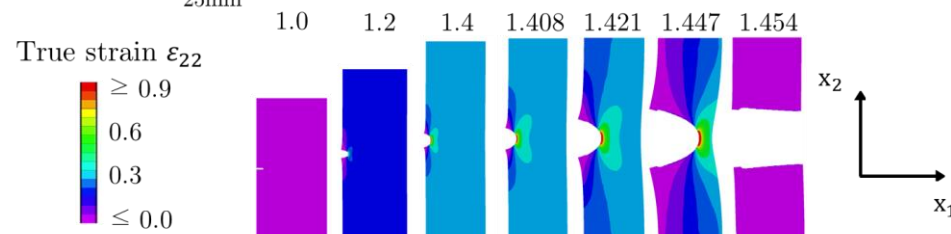
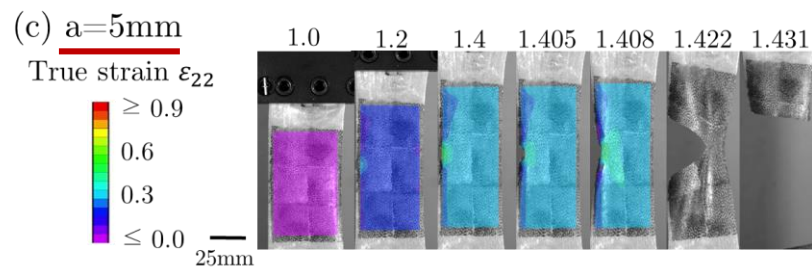
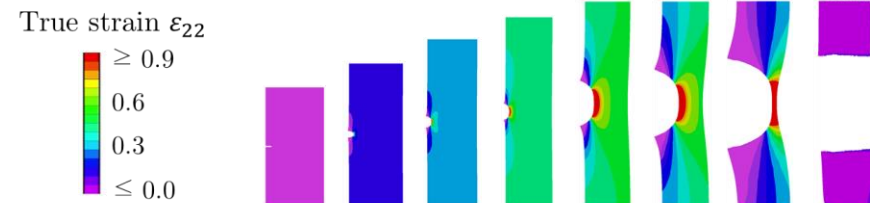
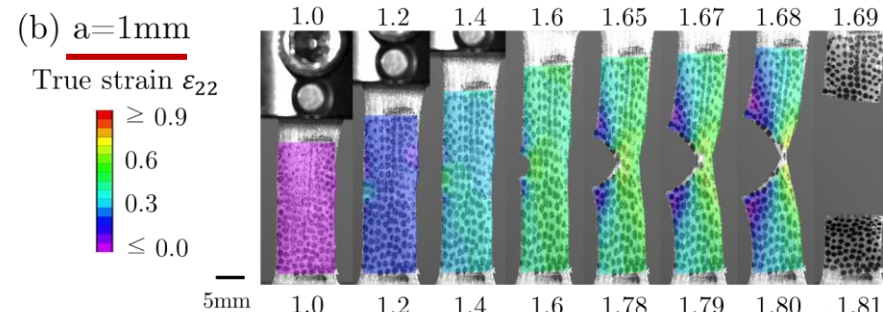
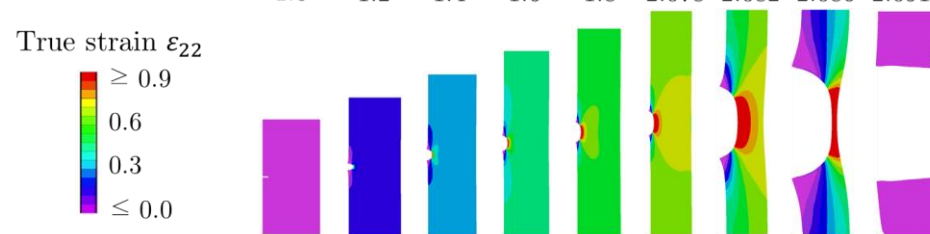
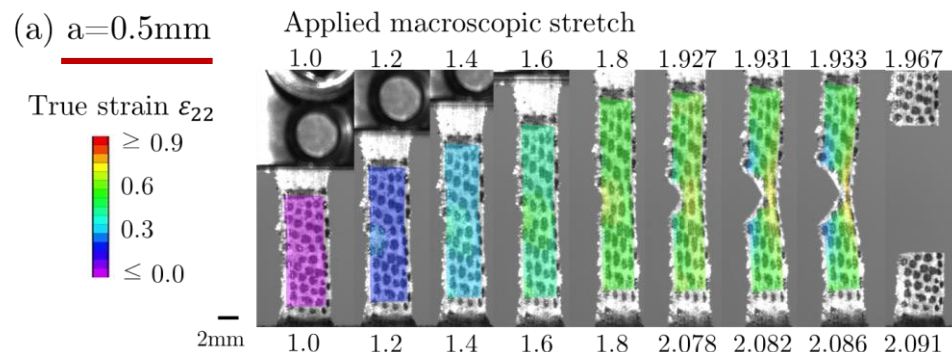
Results: Experiment vs. Numerical simulation^[1]

- Strain fields in **PDMS** specimens ($l = 0.08mm$)
 - Larger specimen ruptures earlier



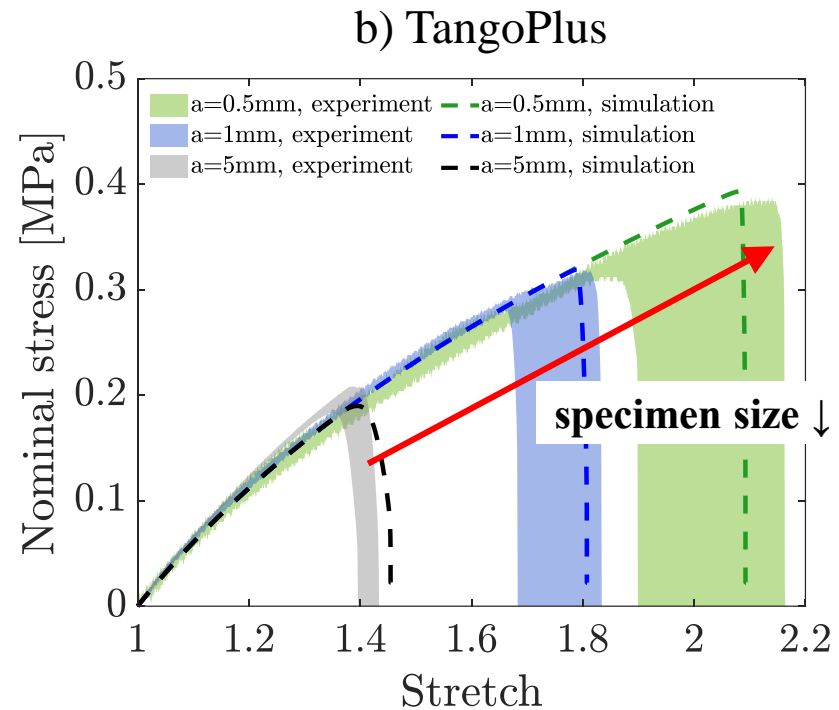
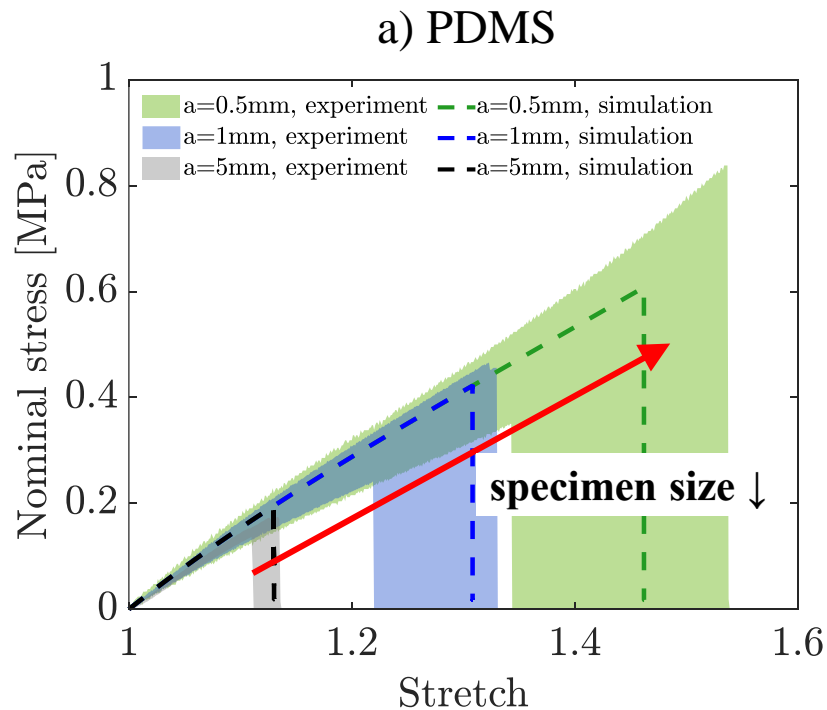
Results: Experiment vs. Numerical simulation^[1]

- Strain fields in **TangoPlus** specimens ($l = 1mm$)
 - Larger specimen ruptures earlier



Results: Experiment vs. Numerical simulation^[1]

- Notch lengths $a = \{0.5, 1, 5\}$ mm
- Geometric similarity → **Identical initial stress-stretch response**
- Smaller notch length → Higher rupture stretch



Notch-length sensitivity^[1]

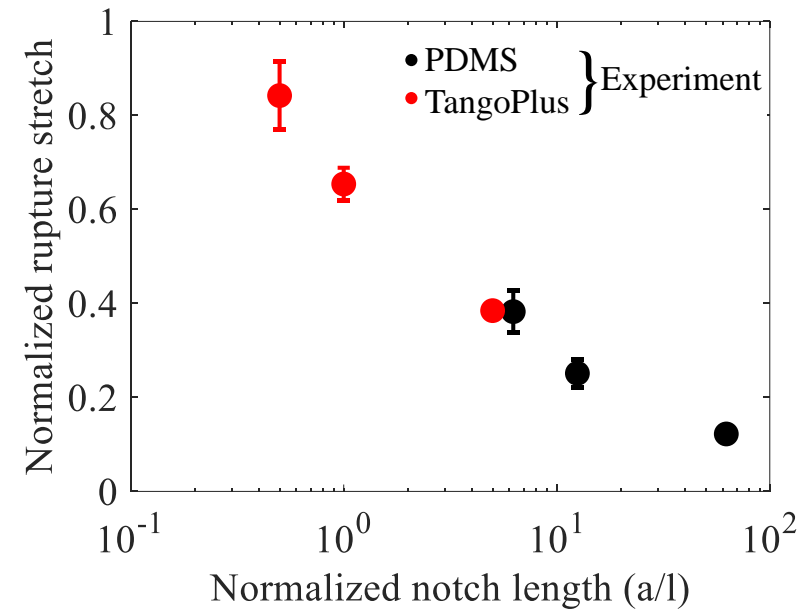
- PDMS vs. TangoPlus; same specimen sizes
 - PDMS: $l = 0.08\text{mm}$
 - TangoPlus: $l = 1\text{mm}$

More than 10 times

- Normalized rupture stretch

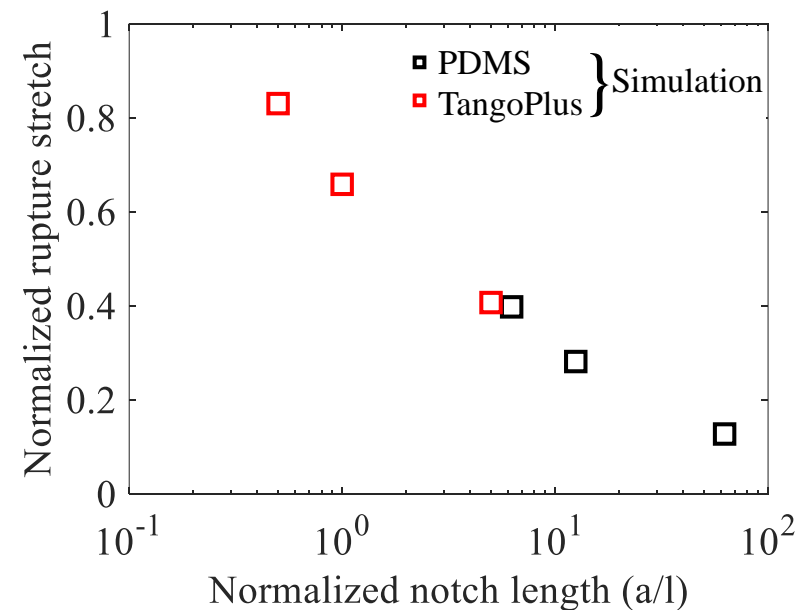
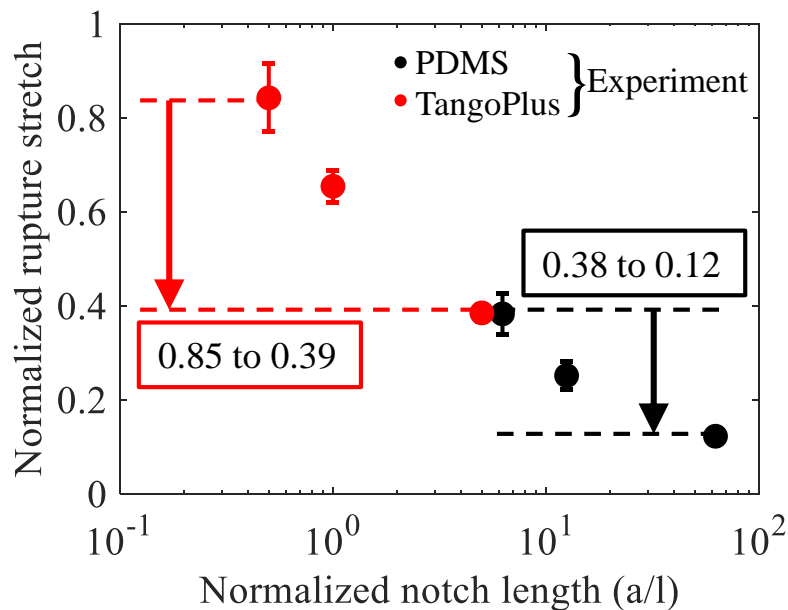
$$= \frac{\text{Rupture stretch of notched specimens}}{\text{Rupture stretch of unnotched specimens}}$$

- Normalized notch length = $\frac{\text{Notch length (a)}}{\text{Intrinsic length scale (l)}}$



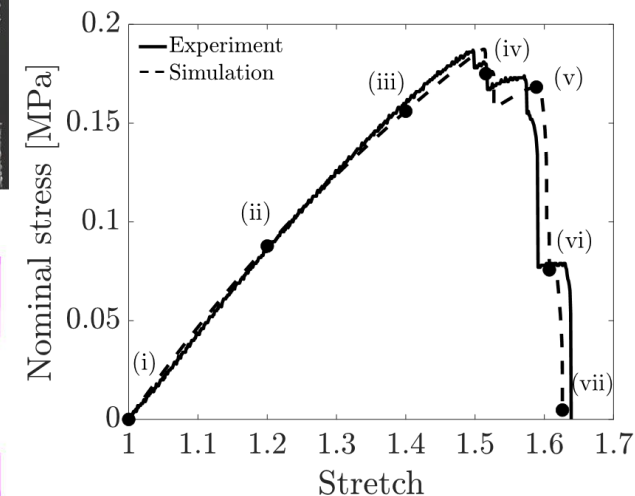
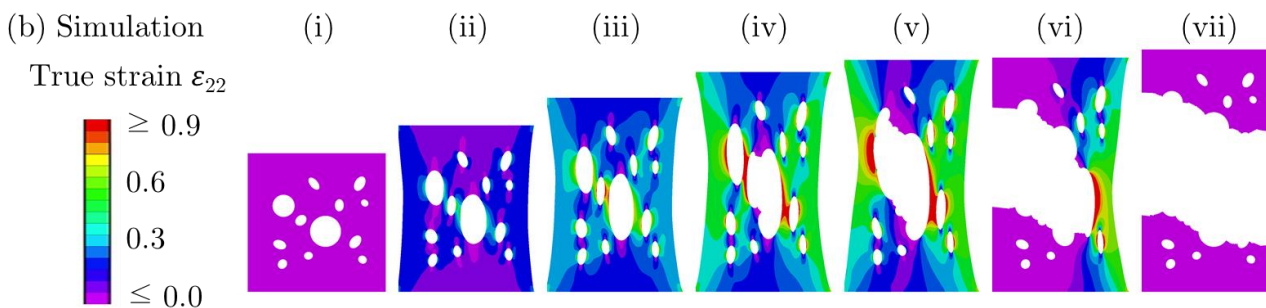
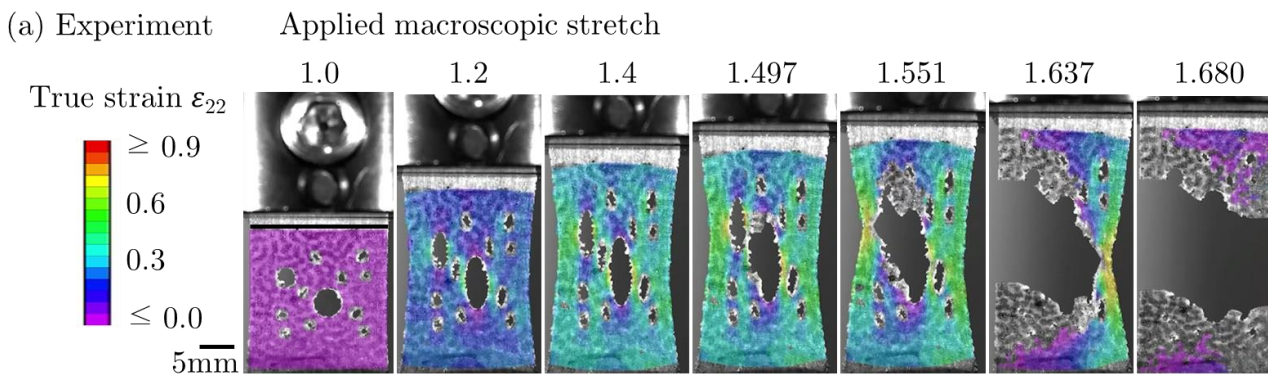
Notch-length sensitivity^[1]

- PDMS vs. TangoPlus; same specimen sizes
 - PDMS: $l = 0.08\text{mm}$
 - TangoPlus: $l = 1\text{mm}$
) **More than 10 times**
- a/l : 0.5~5 (TangoPlus; $l = 1\text{mm}$) → Highly notch length-sensitive
- a/l : 5~50 (PDMS; $l = 0.08\text{mm}$) → Less notch length-sensitive



Randomly perforated specimen (TangoPlus)^[1]

- Nicely predicted the response **without modification of parameters**
 - Progressive fracture of ligaments



Conclusion

- **Size-dependent fracture** is clearly observed in experiments
 - Rupture stretch increases as the notch length decreases
 - Size-dependence increases as the notch-root radius decreases
- **The intrinsic length scale** determines the size-dependent behavior
 - The intrinsic length scale l defines the size of diffusive damage zone / fracture process zone
 - The intrinsic length scales were identified from experiments
 - Normalized notch length (a/l) determines the size-dependence
- **Nonlocal continuum model** nicely predicted the fracture in elastomers
 - Nonlocal continuum model utilizes experimentally identified intrinsic length scales
 - The model captures the size-dependent fracture in elastomers
 - The model is capable of predicting the fracture of complex geometries

Reference

- [1] Lee et al., [*Phys. Rev. Mater.*, in revision](#).
- [2] Chen et al., [*Extreme Mech. Lett.*, 10, 2017](#).
- [3] Pharr et al., [*J Appl. Phys.*, 111, 2012](#).
- [4] Griffith, [*Philos. Trans. R. Soc. Lond. A*, 221, 1921](#).
- [5] Rivlin and Thomas, [*J. Polym. Sci.*, 10, 1955](#).
- [6] Lake and Thomas, [*Proc. R. Soc. Lond.*, 300, 1967](#).
- [7] Mao et al., [*Extreme Mech. Lett.*, 13, 2017](#).
- [8] Talamini et al., [*J. Mech. Phys. Solids*, 111, 2018](#).
- [9] Peerlings et al., [*Int. J. Numer. Methods Eng.*, 39, 1996](#).
- [10] de Borst et al., [*Eur. J. Mech. A/Solids*, 18, 1999](#).
- [11] Francfort and Marigo, [*J. Mech. Phys. Solids*, 46, 1998](#).
- [12] Miehe et al., [*Int. J. Numer. Methods Eng.*, 83, 2010](#).
- [13] Bažant, [*Int. J. Fract.*, 83, 1997](#).
- [14] Yang et al., [*J. Mech. Phys. Solids*, 131, 2019](#).