

Du2025_figures

March 11, 2025

- 1 Figures for “On the anisotropic velocity distribution of newborn pickup ions in the heliosheath”
- 2 Du, Opher, & Kornbleuth, 2025, ApJL, DOI: 10.3847/2041-8213/adbc9a
- 3 Figure 1, estimated radial velocity distribution of interstellar neutrals and the neutral solar wind

```
[1]: import numpy as np
import matplotlib.pyplot as plt

import matplotlib, scipy, plasmapy, h5py
print("numpy: ", np.__version__, ", matplotlib: ", matplotlib.__version__, "\n",
      "scipy: ", scipy.__version__, ", plasmapy: ", plasmapy.__version__, "\n",
      "h5py: ", h5py.__version__)
```

```
numpy: 1.24.3 , matplotlib: 3.7.1 ,
scipy: 1.10.1 , plasmapy: 2024.7.0 ,
h5py: 3.7.0
```

```
[2]: # 1D cut through v=0 in an isotropic Maxwellian
def Maxwell(v, n0, v0, vth):
    return n0/(2*np.pi*vth**2)**(3/2)*np.exp(-0.5*(v-v0)**2/vth**2)

# Maxwellian for different vth in radial and perp, cut through vperp = 0
def Maxwell12(v, n0, v0, vthr, vthp):
    return n0/(2*np.pi)**(3/2)/(vthr*vthp*vthp)*np.exp(-0.5*(v-v0)**2/vthr**2)

# convert non-relativistic proton energy in eV to velocity in km/s
def eV2kms(eV):
    return np.sqrt(eV * 1.6e-19 * 2 / 1.67e-27)/1000
```

```
[3]: vx = np.linspace(-350, 550, 200)
nh0 = 0.1      # interstellar/H wall neutral number density (cm^-3)
vt0 = 30.0     # thermal velocity (km/s)
```

```

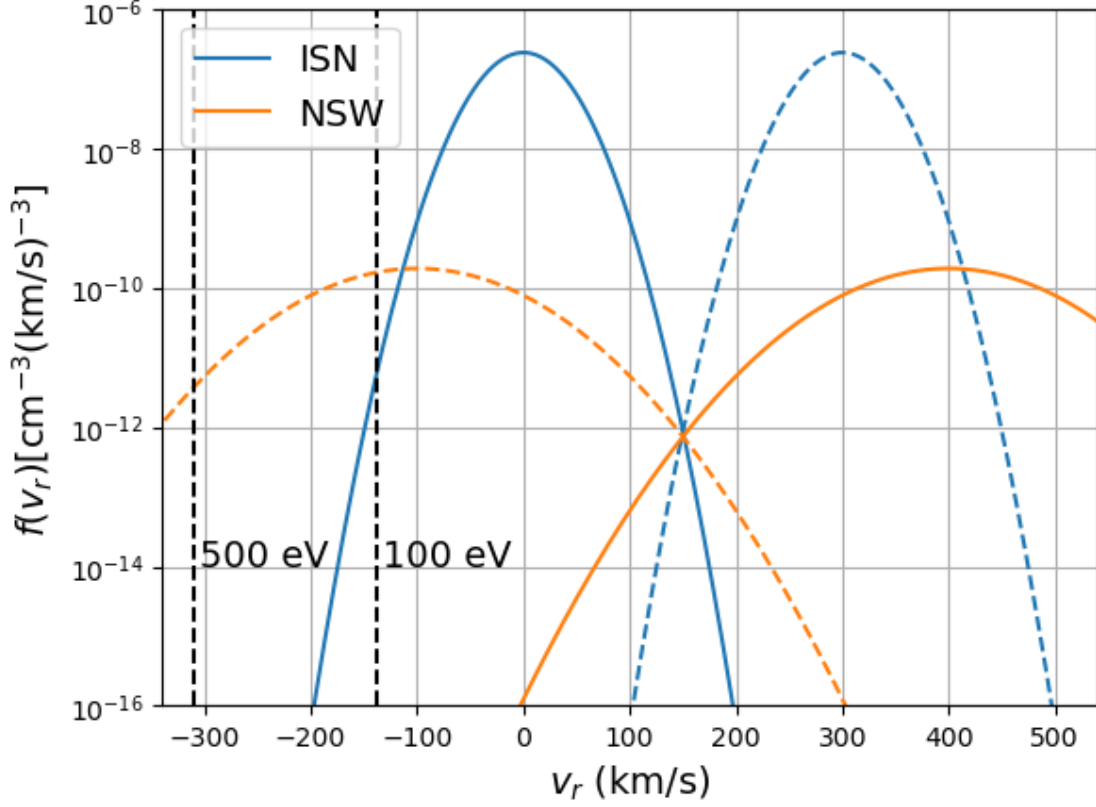
v01 = 0.0      # bulk velocity (km/s)
v02 = 300.0    # bulk velocity mirrored with respect to the heliosheath flow
                ↪ (150 km/s)
nnw = 2.0e-4   # neutral solar wind number density (cm-3)
vtw = 75.0     # NSW thermal parallel thermal velocity
vtwp = 30.0    # NSW perpendicular thermal velocity
vw1 = -100.0   # NSW bulk velocity (mirrored)
vw2 = 400.0    # NSW bulk velocity

plt.plot(vx, Maxwell(vx, nh0, v01, vt0), label='ISN')
plt.plot(vx, Maxwell(vx, nh0, v02, vt0), color='C0', ls='--')
plt.plot(vx, Maxwell2(vx, nnw, vw2, vtw, vtwp), label='NSW')
plt.plot(vx, Maxwell2(vx, nnw, vw1, vtw, vtwp), color='C1', ls='--')
plt.axvline(-eV2kms(100), ls='--', color='k')
plt.axvline(-eV2kms(500), ls='--', color='k')
plt.text(-eV2kms(100)+5, 1e-14, '100 eV', fontsize=14)
plt.text(-eV2kms(500)+5, 1e-14, '500 eV', fontsize=14)
plt.legend(fontsize=14, ncol=1, loc='upper left')
plt.yscale('log')
plt.xlabel(r'$v_r$ (km/s)', fontsize=14)
plt.ylabel(r'$f(v_r)$ [\mathrm{cm}^{-3}(\mathrm{km/s})^{-3}]$', fontsize=14)
plt.xlim(-340, 540)
plt.ylim(1e-16, 1e-6)
plt.grid(True)

# print the ratio between total distribution and ISN at ~100 eV
print((Maxwell2(-138.4, nnw, vw1, vtw, vtwp) + Maxwell(-138.4, nh0, v01, vt0)) /
      ↪ Maxwell(-138.4, nh0, v01, vt0))

```

30.354662490177816



4 **Figure 5, inverse maximum growth rate and charge-exchange timescale**

Using the approximate dispersion relation for mirror mode with Maxwellian and ring components:

$$\frac{\omega^2}{\omega_{pi}^2} = \frac{c^2 k^2}{\omega_{pi}^2} - \sum_m \frac{n_m}{n_0} \zeta_m Z(\zeta_m) \lambda_m e^{-\lambda_m} [2I_0(\lambda_m) - 2I_1(\lambda_m)] \quad (1)$$

$$+ \sum_r \frac{n_r}{n_0} \frac{2}{\sqrt{\pi} \delta v_\perp^3 \text{erfc}(-v_r/\delta v_\perp)} \int_0^\infty dv_\perp [2v_\perp(v_\perp - v_r) + \delta v_\perp^2] \left| J_1 \left(\frac{k_x v_\perp}{\Omega_0} \right) \right|^2 \exp \left[-\frac{(v_\perp - v_r)^2}{\delta v_\perp^2} \right] \quad (2)$$

$$- \sum_r \frac{n_r}{n_0} \frac{4}{\sqrt{\pi} \delta v_\perp \delta v_\parallel^2 \text{erfc}(-v_r/\delta v_\perp)} \int_0^\infty dv_\perp v_\perp^2 \left| J_1 \left(\frac{k_x v_\perp}{\Omega_0} \right) \right|^2 \exp \left[-\frac{(v_\perp - v_r)^2}{\delta v_\perp^2} \right] \left[1 + \frac{\omega}{k_z \delta v_\parallel} Z \left(\frac{\omega}{k_z \delta v_\parallel} \right) \right]$$

$$\zeta_m = \frac{\omega}{k_z v_t} = \frac{\omega}{\Omega} \frac{\Omega d_i}{k_z d_i v_t} = \frac{\omega}{\Omega} \frac{1}{k_z d_i} \frac{1}{\sqrt{\beta}}; \quad \lambda_m = \frac{k_x^2 v_t^2}{2\Omega^2} = \frac{\beta}{2} (k_x d_i)^2$$

$$\Gamma = \sqrt{\pi} \text{erfc} \left(-\frac{v_r}{\delta v_\perp} \right)$$

$$\frac{k_x v_\perp}{\Omega} = k_x d_i \frac{v_\perp}{v_A} \frac{v_A}{d_i \Omega} = k_x d_i \frac{v_\perp}{v_A}; \quad \frac{\omega}{k_z \delta v_\parallel} = \frac{\omega/\Omega}{(k_z d_i)(\delta v_\parallel/v_A)}$$

The beta parameter is not the effective plasma beta, but is defined with the thermal temperature and the total density, i.e.,

$$\beta = \frac{n_0 T}{B^2 / 8\pi}.$$

Need the special functions i0, i1, j1, erfc from scipy, the plasma dispersion function from plasmapy, and a generic integration method.

```
[4]: from scipy.special import i0, i1, j1, erfc
from scipy.integrate import quad
from plasmapy.dispersion import plasma_dispersion_func

# The integrands in the dispersion relation
def integ1(v, vr, kx, dv):
    return (2*v*(v-vr)+dv**2)*(j1(kx*v))**2*np.exp(-((v-vr)/dv)**2)

def integ2(v, vr, kx, dv):
    return v**2*(j1(kx*v))**2*np.exp(-((v-vr)/dv)**2)

# Approximate dielectric tensor for mirror mode assuming two isotropic
# Maxwellian + two rings
# Dispersion relation is obtained by solving Dyy = 0
def Dyy(kperp, kpara, omega, nr, beta, vr1, dvperp1, dvpara1, vr2, dvperp2,
        dvpara2,
        Gamma1, Gamma2, int1a, int1b, int2a, int2b):
    k2 = kperp**2 + kpara**2
    nm = 1.0 - nr      # total density fraction of Maxwellian components n_m /
    # n_total
    nm1 = 0.75*nm      # density fraction of first Maxwellian component (thermal)
    nm2 = 0.25*nm      # density fraction of second Maxwellian component (hot/
    # transmitted pickup ions)
    nr1 = 0.998*nr      # density fraction of first ring component (ISN PUI)
    nr2 = 0.002*nr      # density fraction of second ring component (NSW PUI)
    lamb1 = 0.5*beta*kperp**2      # parameter for first Maxwellian component,
    # temperature dependent
    lamb2 = lamb1*40      # parameter for second Maxwellian component,
    # temperature 40x higher
    zeta1 = omega/kpara / np.sqrt(beta)      # parameter for first Maxwellian
    # component
    zeta2 = omega/kpara / np.sqrt(beta*40)      # parameter for second Maxwellian
    # component
    Z1 = plasma_dispersion_func(zeta1)
    Z2 = plasma_dispersion_func(zeta2)
    wk1 = omega / (kpara*dvpara1)      # parameter for first ring component
    wk2 = omega / (kpara*dvpara2)      # parameter for second ring component
    Z1r = plasma_dispersion_func(wk1)
    Z2r = plasma_dispersion_func(wk2)
```

```

D0 = k2 - nm1*zeta1*Z1*lamb1*np.exp(-lamb1)*2*(i0(lamb1)-i1(lamb1)) \
      - nm2*zeta2*Z2*lamb2*np.exp(-lamb2)*2*(i0(lamb2)-i1(lamb2)) \
      + nr1*2/dvperp1**3/Gamma1*int1a + nr2*2/dvperp2**3/Gamma2*int2a \
      - nr1*4/dvperp1/dvpara1**2/Gamma1*int1b*(1 + wk1*Z1r) \
      - nr2*4/dvperp2/dvpara2**2/Gamma2*int2b*(1 + wk2*Z2r)

return D0

```

```

[5]: """
We did not automate the process of solving the dispersion relation.
To produce the figure, we manually choose a range of number density ratios, and
    ↳for each ratio,
we search in a grid of (k_parallel, k_perp) to find the wavevector that
    ↳correspond to
the largest growth rate. The results are then recorded in the next code block.
    ↳We search
twice for each ratio with a coarse grid and a refined grid to improve accuracy.
    ↳Here,
the search parameters are for the smallest density ratio. However, we recommend
    ↳searching
in reverse order starting with the highest ratio (nr = 0.1), and using the
    ↳contour plot to aid
the selection of the search parameters (e.g., kperps ~ [0.01, 0.9], kparas ~ [0.
    ↳01, 1.2],
ois ~ [-2e-3, 2e-1]).
"""

beta = 0.75
nr = 0.00843    # total density ratio for the ring components, recommend
    ↳starting with 0.1
vr1 = 3.0      # first ring velocity in V_A
vr2 = 5.0      # second ring velocity
dvperp1 = 0.6  # parallel thermal spread in the first ring in V_A
dvpara1 = 0.6  # perpendicular thermal spread in the first ring in V_A
dvperp2 = 1.5  # parallel thermal spread in the second ring in V_A
dvpara2 = 0.6  # perpendicular thermal spread in the second ring in V_A
vrd1 = vr1 / dvperp1
vrd2 = vr2 / dvperp2
Gamma1 = np.sqrt(np.pi)*erfc(-vrd1)
Gamma2 = np.sqrt(np.pi)*erfc(-vrd2)

Nper = 25
Npar = 25
# search grid in wavevector space, need to change for each ratio
kperps = np.linspace(0.001, 0.027, Nper) # for rough estimates
kparas = np.linspace(0.0001, 0.0007, Npar)
# kperps = np.linspace(0.01, 0.019, Nper) # uncomment this and the next lines
    ↳for refining

```

```

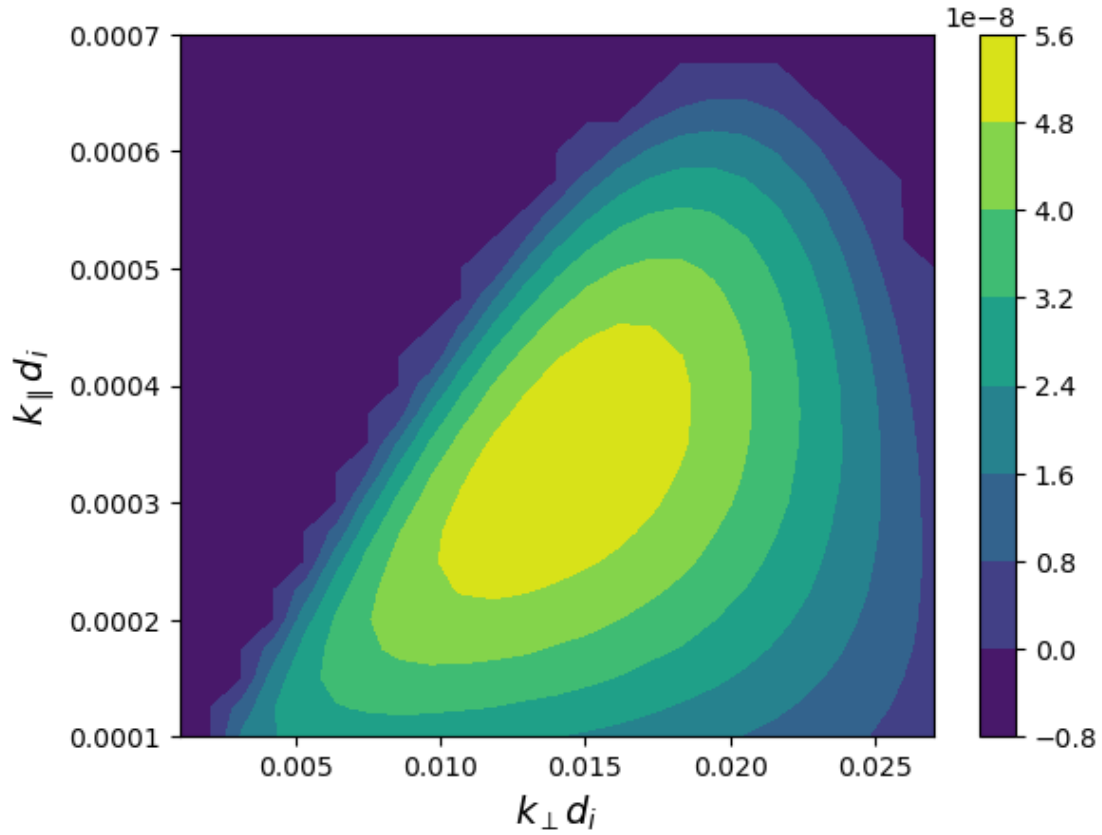
# kparas = np.linspace(0.00022, 0.00047, Npar)
om = np.zeros((Nper, Npar), dtype=complex)

for iper in range(Nper):
    kperp = kperps[iper]
    int1a = quad(integ1, 0, np.inf, args=(vr1, kperp, dvperp1))[0] # compute
    ↪the integrals
    int2a = quad(integ1, 0, np.inf, args=(vr2, kperp, dvperp2))[0]
    int1b = quad(integ2, 0, np.inf, args=(vr1, kperp, dvperp1))[0]
    int2b = quad(integ2, 0, np.inf, args=(vr2, kperp, dvperp2))[0]
    for ipar in range(Npar):
        kpara = kparas[ipar]
        Nr = 3
        Ni = 201
        ors = np.linspace(-0.1, 0.1, Nr) # not important because the real
    ↪frequency is ~ 0
        ois = np.linspace(-1.5e-11, 6.0e-8, Ni) # this determines the accuracy
    ↪in growth rate
        D0 = np.zeros((Nr, Ni), dtype=complex)
        omegas = np.zeros((Nr, Ni), dtype=complex)
        for i in range(Nr):
            for j in range(Ni):
                omega = ors[i] + 1j*ois[j]
                D0[i,j] = Dyy(kperp, kpara, omega, nr, beta, vr1, dvperp1,
    ↪dvpara1, vr2, dvperp2, dvpara2,
                                Gamma1, Gamma2, int1a, int1b, int2a, int2b)
                omegas[i,j] = omega
            omega0 = omegas.flatten()[np.argmin(abs(D0))] # find where |Dyy| is
    ↪closest to 0
            om[iper,ipar] = omega0

# plot the growth rate in k-space
plt.contourf(kperps, kparas, om.T.imag)
plt.colorbar()
plt.xlabel(r'$k_{\perp}$ d_i$', fontsize=14)
plt.ylabel(r'$k_{\parallel}$ d_i$', fontsize=14)
print(abs(om.imag).max()) # save this number for the maximum growth rate in
    ↪the next block

```

5.459865e-08



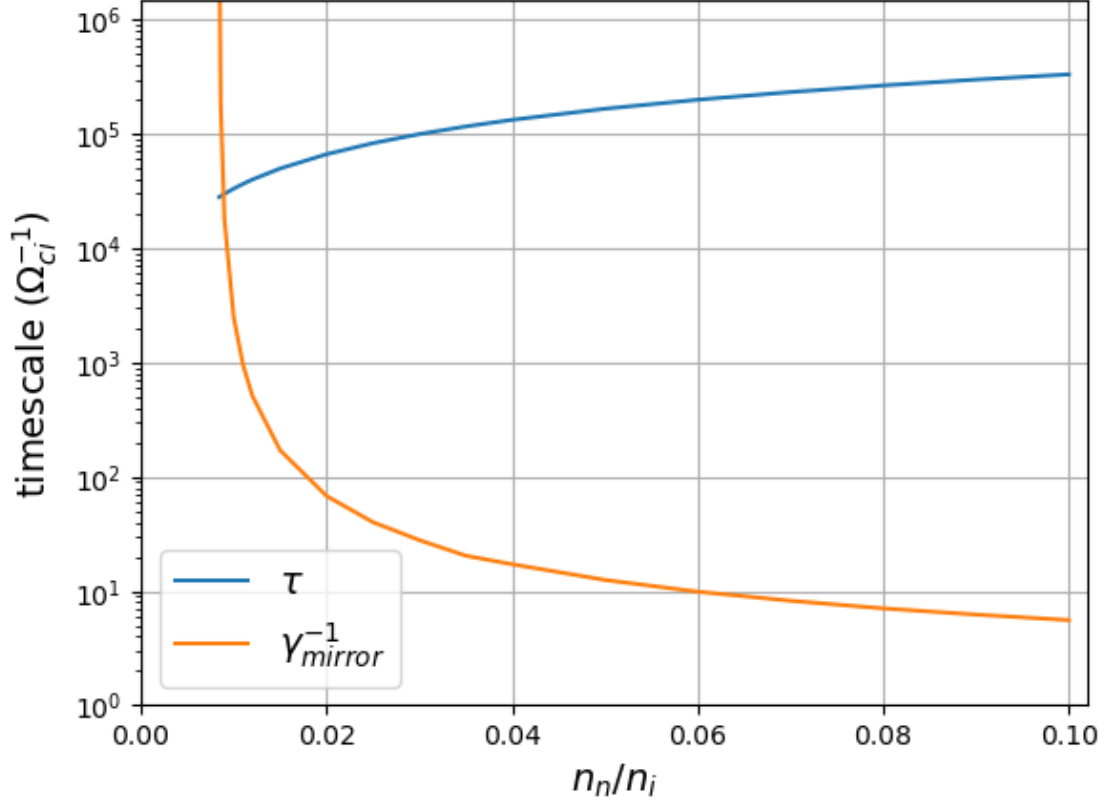
```
[6]: # Maximum growth rate from the last block
ratios = np.array([0.00843, 0.00845, 0.0086, 0.009, 0.01, 0.011, 0.
    ↪0.012, 0.015, 0.02,
    ↪0.025, 0.03, 0.035, 0.04, 0.05, 0.06, 0.07, 0.08, 0.09, 0.1])
gammas = np.array([5.46e-8, 2.37e-7, 5.67e-6, 5.572e-5, 4.025e-4, 1.044e-3, 1.
    ↪95e-3, 5.85e-3, 0.01470,
    ↪0.02492, 0.03580, 0.049, 0.05820, 0.07989, 0.1012, 0.1215, 0.1417, 0.1606, 0.1795])
# charge exchange timescale
nH = 0.1 # neutral density cm^-3
cxcs = 2.0e-15 # charge exchange cross section cm^2
usw = 1.5e7 # solar wind speed cm/s
wcis = 100.0 # Omega_ci^-1 / s
taus = ratios/(nH*cxcs*usw)/wcis

plt.plot(ratios, taus, label=r'\tau$')
plt.plot(ratios, 1/gammas, label='$\gamma_{\text{mirror}}^{-1}$')
plt.legend(fontsize=14)
```

```

plt.xlabel(r"$n_n/n_i$", fontsize=14)
plt.ylabel(r"timescale ($\Omega_{ci}^{-1}$)", fontsize=14)
plt.yscale('log')
plt.xlim(0, 0.102)
plt.ylim(1.0e0, 1.5e6)
plt.grid(True)

```



5 Figure 2, velocity distributions from the simulation

The particle velocity distribution data (“fnp_0.npy” and “fnp_8.npy”) are needed. They are macroparticle counts obtained from the full particle output of the hybrid simulation at two different times.

```

[7]: def plot_vxvz(tframe):
    fsave_name = "./fnp_%d.npy" %(tframe)
    with open(fsave_name, "rb") as file_p:
        pdist = np.load(file_p)
        vxb = np.load(file_p)
        vyb = np.load(file_p)
        vzb = np.load(file_p)

```



```

Nbs = vyb.size
fxz = 0.5*(pdist[:,Nbs//2-1,:] + pdist[:,Nbs//2,:]) # select the  $v_y \sim 0$ 
↪ slice
wcit = tframe*200 #  $\Omega_{ci} t$ 

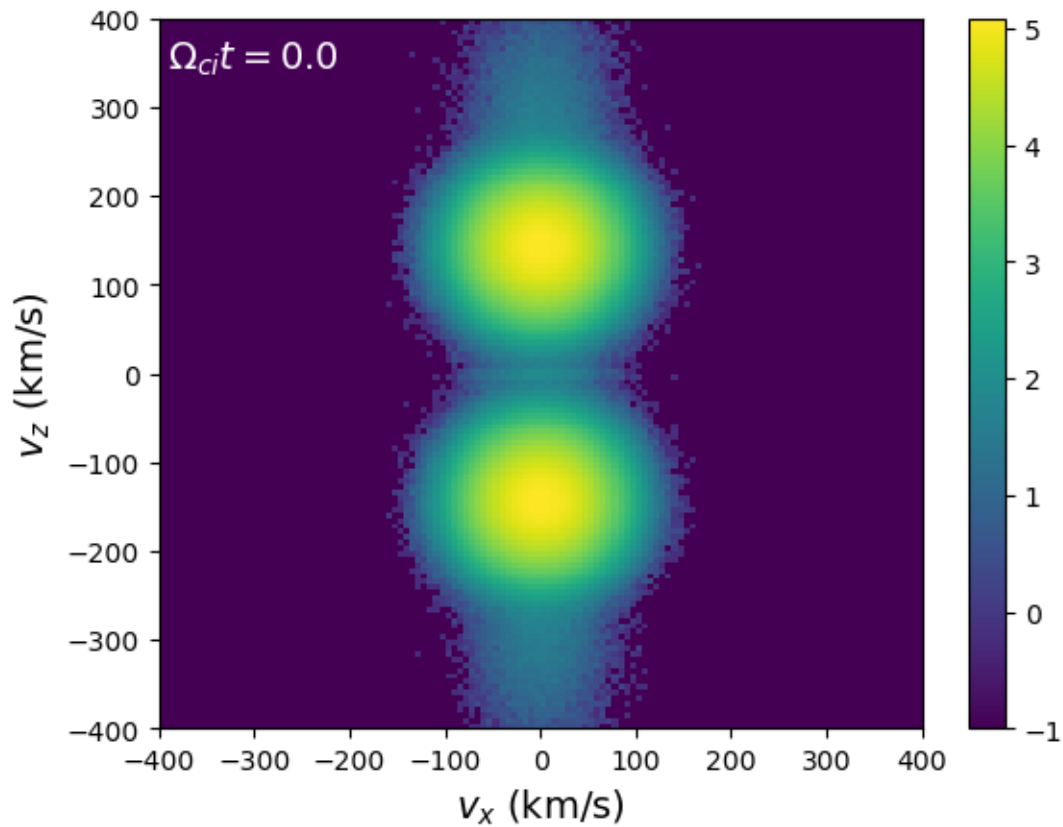
plt.pcolormesh(vxb, vzb, np.log10(fxz.T+0.1))
plt.text(vxb[0]+0.01*(vxb[-1]-vxb[0]), vzb[0]+0.93*(vzb[-1]-vzb[0]),
        r"$\Omega_{ci}t = %.1f$" %(wcit), color="w", fontsize=14)
plt.colorbar()
plt.xlabel(r"$v_x$ (km/s)", fontsize=14)
plt.ylabel(r"$v_z$ (km/s)", fontsize=14)

```

```

[8]: # Time = 0 (initial)
plot_vxvz(0)

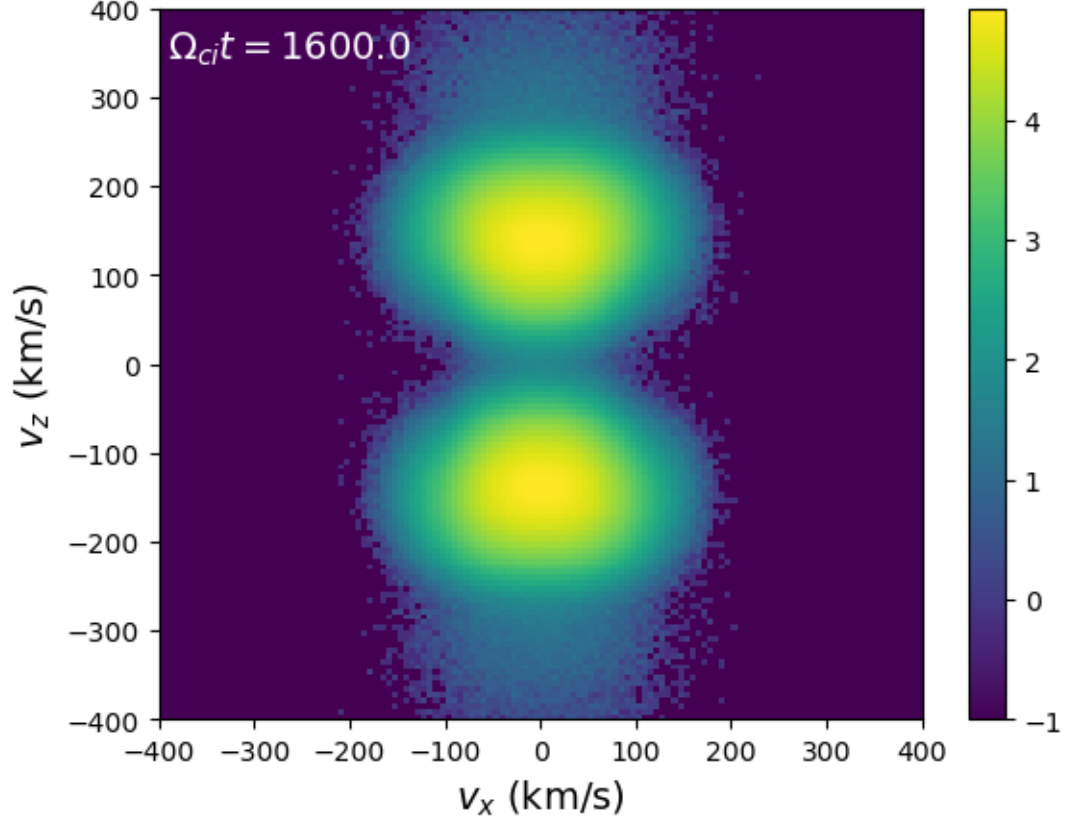
```



```

[9]: # Time  $t = 1600$  (end)
plot_vxvz(8)

```



6 Figure 3, pitch angle distribution of pickup ions

Need particle pitch angle distribution data (“TS1a/pa100_np_0”, ..., “TS1a/pa100_np_8” and “TS1b/pa100_np_0”, ..., “TS1b/pa100_np_12”) from the hybrid simulations. Each data file contains the counts of macroparticles between 280 km/s and 300 km/s in 50 bins of μ uniformly spaced between -1 and 1.

The left panel simply plots the particle counts in each bin. The top right panel plots the width of the distribution $\Delta\mu$, calculated by the full width at half maximum using the interpolation and root finding functions from scipy. The bottom right panel shows the ratio f_{ani}/f_{iso} . The numerator is calculated by the two bins around $\mu = 0$, and the denominator is calculated by the average at $t = 0$.

```
[10]: # Read data from 7% run
Nframe = 9
Nmu = 50
cpa = np.zeros((Nframe, Nmu)) # cosine pitch angle
fpa = np.zeros((Nframe, Nmu)) # distribution

for i in range(Nframe):
```

```

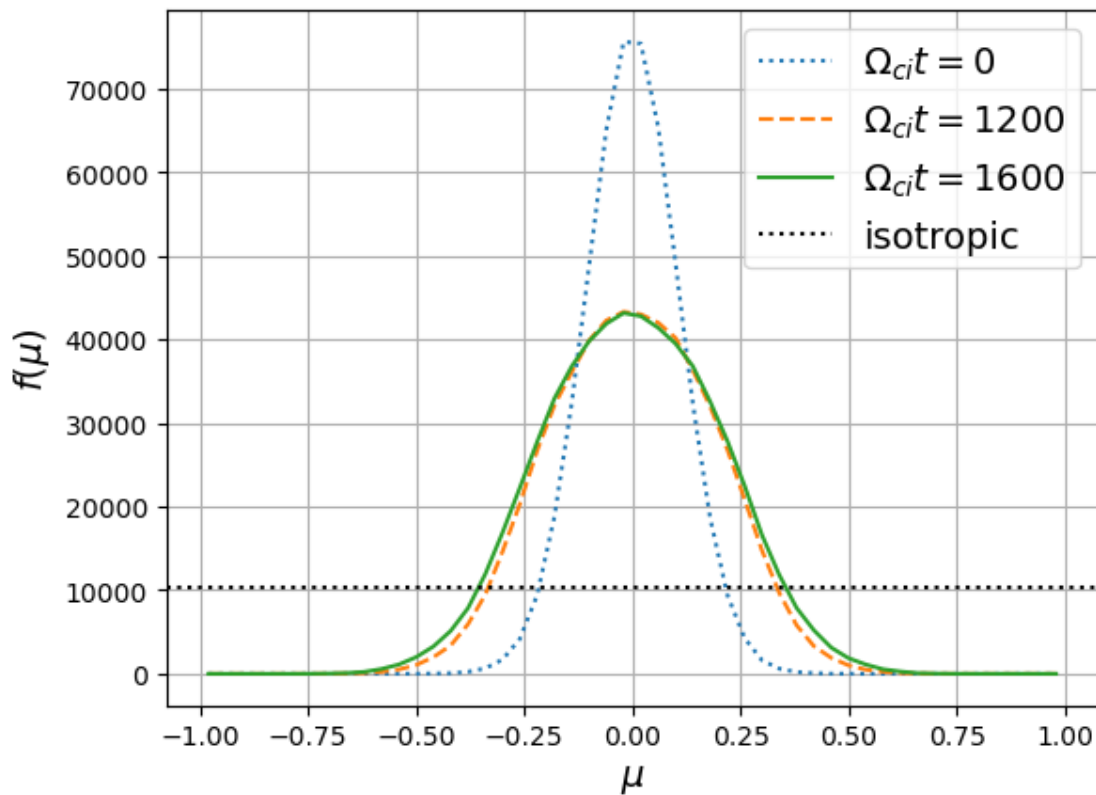
dat_dir = "./TS1a/"
fname = dat_dir + "mu100_np_%d.npy" %(i)
with open(fname, "rb") as file_p:
    fa0 = np.load(file_p)
    pa0 = np.load(file_p)
fpa[i] = fa0
cpa[i] = (pa0[1:] + pa0[:-1])/2

```

```

[11]: # Left panel of Figure 3
plt.plot(cpa[0], fpa[0], ls=':', label=r'$\Omega_{ci}t = 0$')
plt.plot(cpa[4], fpa[6], ls='--', label=r'$\Omega_{ci}t = 1200$')
plt.plot(cpa[8], fpa[8], label=r'$\Omega_{ci}t = 1600$')
plt.axhline(fpa[0].mean(), ls=':', color='k', label='isotropic')
plt.xlabel(r'$\mu$', fontsize=14)
plt.ylabel(r'$f(\mu)$', fontsize=14)
plt.legend(fontsize=14)
plt.grid()

```



```

[12]: from scipy import interpolate
      from scipy.optimize import root_scalar

```

```

# read data from 3.5% run
Nframeb = 13
fpab = np.zeros((Nframeb, Nmu))
for i in range(Nframeb):
    dat_dir = "./TS1b/"
    fname = dat_dir + "mu100_np_%d.npy" %(i)
    with open(fname, "rb") as file_p:
        fa0 = np.load(file_p)
        fpab[i] = fa0

ws = np.zeros(Nframe)
fa_fi = np.zeros(Nframe)
wsb = np.zeros(Nframeb)
fa_fib = np.zeros(Nframeb)
time = np.arange(Nframe)*200    # Particle outputs are saved every 200
    ↪  $\Omega_{ci}^{-1}$ 
timeb = np.arange(Nframeb)*200
fpamean = fpa[0].mean()
fpabmean = fpab[0].mean()

# find width and enhancement for the 7% run
for i in range(Nframe):
    fm2 = 0.5*fpa[i].max()
    f0 = interpolate.interp1d(cpa[0], fpa[i]-fm2)    # interpolate to find half
    ↪ maximum
    r0 = root_scalar(f0, bracket=[cpa[0,0], cpa[0,Nmu//2]], xtol=1.0e-5)
    r1 = root_scalar(f0, bracket=[cpa[0,Nmu//2], cpa[0,-1]], xtol=1.0e-5)
    ws[i] = r1.root - r0.root    # width of the
    ↪ distribution
    fa_fi[i] = (fpa[i, Nmu//2-1]+fpa[i, Nmu//2])/2/fpamean    # enhancement ratio

# find width and enhancement for the 3.5% run
for i in range(Nframeb):
    fm2 = 0.5*fpab[i].max()
    f0 = interpolate.interp1d(cpa[0], fpab[i]-fm2)
    r0 = root_scalar(f0, bracket=[cpa[0,0], cpa[0,Nmu//2]], xtol=1.0e-5)
    r1 = root_scalar(f0, bracket=[cpa[0,Nmu//2], cpa[0,-1]], xtol=1.0e-5)
    wsb[i] = r1.root - r0.root
    fa_fib[i] = (fpab[i, Nmu//2-1]+fpab[i, Nmu//2])/2/fpabmean

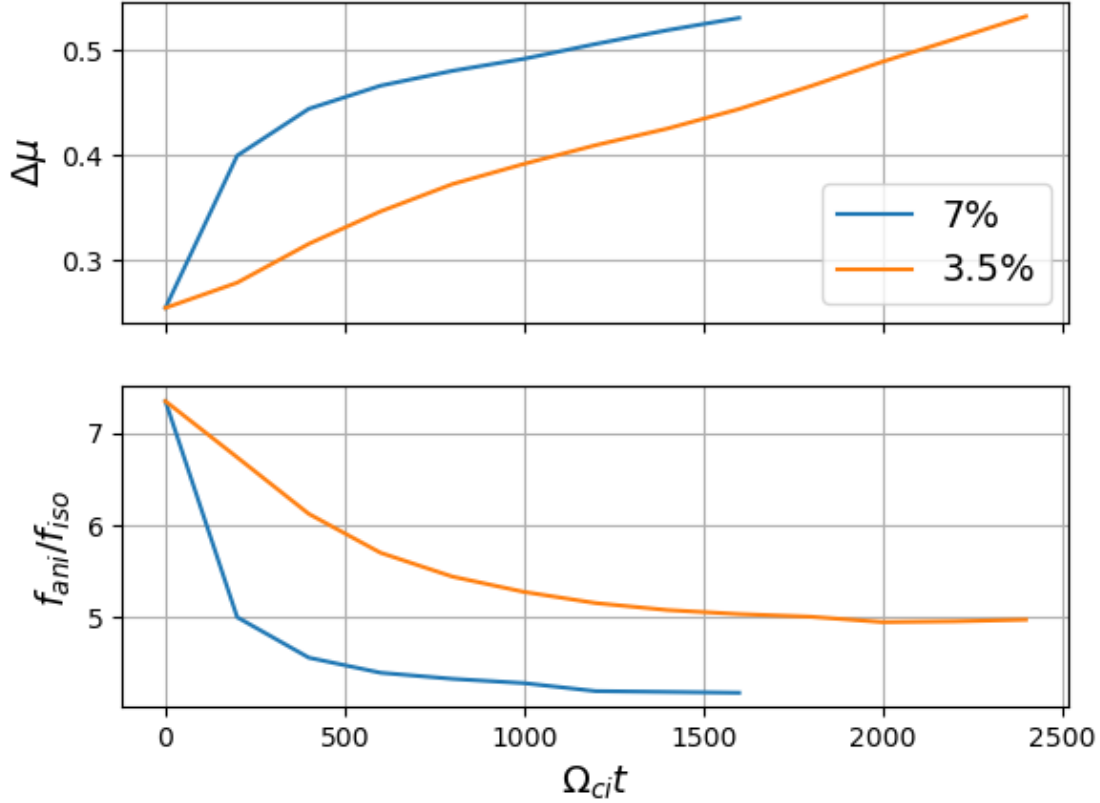
fig, ax = plt.subplots(2, 1, sharex=True)
ax[0].plot(time, ws, label="7%")
ax[0].plot(timeb, wsb, label="3.5%")
ax[1].plot(time, fa_fi)
ax[1].plot(timeb, fa_fib)
ax[1].set_xlabel(r"$\Omega_{ci}$ t$", fontsize=14)
ax[0].set_ylabel(r"$\Delta$ \mu$", fontsize=14)

```

```

ax[1].set_ylabel(r"$f_{\text{ani}}/f_{\text{iso}}$", fontsize=14)
ax[0].legend(fontsize=14)
ax[0].grid()
ax[1].grid()

```



7 Figure 4, snapshots of the simulation and frequency spectra

The snapshots need data “./hydro_hdf5/T.40000/hydro_ion_40000.h5”, “hydro_pui_40000.h5”, “hydro_npu_40000.h5”, and “field_hdf5/T.40000/fields_40000.h5”. These are outputs from the hybrid simulation for the plasma moments (including three separate ion species) and fields. The frequency spectrum needs “fields_20000.h5”, ..., “fields_122200.h5”.

Need to import h5py package to read hdf5 data.

```

[13]: import h5py

# Functions to plot magnetic field B_x and total ion density
def plot_bx(tframe):
    tindex = 200 * tframe
    wcit = tindex*0.005      # Time Omega_ci t
    lx_di = 256.0           # simulation size in x (d_i)

```

```

lz_di = 256.0          # simulation size in z (d_i)
xmin, xmax = 0, lx_di
zmin, zmax = 0, lz_di

fname = ("./field_hdf5/T." + str(tindex) +
        "/fields_" + str(tindex) + ".h5")
with h5py.File(fname, 'r') as fh:
    group = fh["Timestep_" + str(tindex)]
    Bx = group["cbx"][:,0,:]

nx = 512      # grid resolution
nz = 512
xgrid = np.linspace(xmin, xmax, nx)
zgrid = np.linspace(zmin, zmax, nz)

plt.figure()
plt.pcolormesh(xgrid, zgrid, Bx.T,
               vmin=0.9, vmax=1.1)
plt.colorbar()
plt.title(r"$B_x$, $\Omega_{ci}t = \%.2f$"%(wcit), fontsize=14)
plt.xlabel(r"$x$ (d_i)", fontsize=14)
plt.ylabel(r"$z$ (d_i)", fontsize=14)

plt.gca().set_aspect("equal")

def plot_rho(tframe):
    tindex = 200 * tframe
    wcit = tindex*0.005
    lx_di = 256.0
    lz_di = 256.0
    xmin, xmax = 0, lx_di
    zmin, zmax = 0, lz_di

    # read thermal ion data
    fname = ("./hydro_hdf5/T." + str(tindex) +
            "/hydro_ion_" + str(tindex) + ".h5")
    with h5py.File(fname, 'r') as fh:
        group = fh["Timestep_" + str(tindex)]
        rho_i = group["rho"][:,0,:]

    # read transmitted pickup ion data
    fname = ("./hydro_hdf5/T." + str(tindex) +
            "/hydro_pui_" + str(tindex) + ".h5")
    with h5py.File(fname, 'r') as fh:
        group = fh["Timestep_" + str(tindex)]
        rho_p = group["rho"][:,0,:]

```

```

# read newborn pickup ion data
fname = ("./hydro_hdf5/T." + str(tindex) +
        "/hydro_npu_" + str(tindex) + ".h5")
with h5py.File(fname, 'r') as fh:
    group = fh["Timestep_" + str(tindex)]
    rho_n = group["rho"][:,0,:]

nx = 512
nz = 512
xgrid = np.linspace(xmin, xmax, nx)
zgrid = np.linspace(zmin, zmax, nz)

plt.figure()
plt.pcolormesh(xgrid, zgrid, (rho_i+rho_p+rho_n).T,
               vmin=0.99, vmax=1.01)
plt.colorbar()
plt.title(r"$n_i$, $\Omega_{ci}t = \%.2f$"%(wcit), fontsize=14)
plt.xlabel(r"$x$ (d_i)", fontsize=14)
plt.ylabel(r"$z$ (d_i)", fontsize=14)

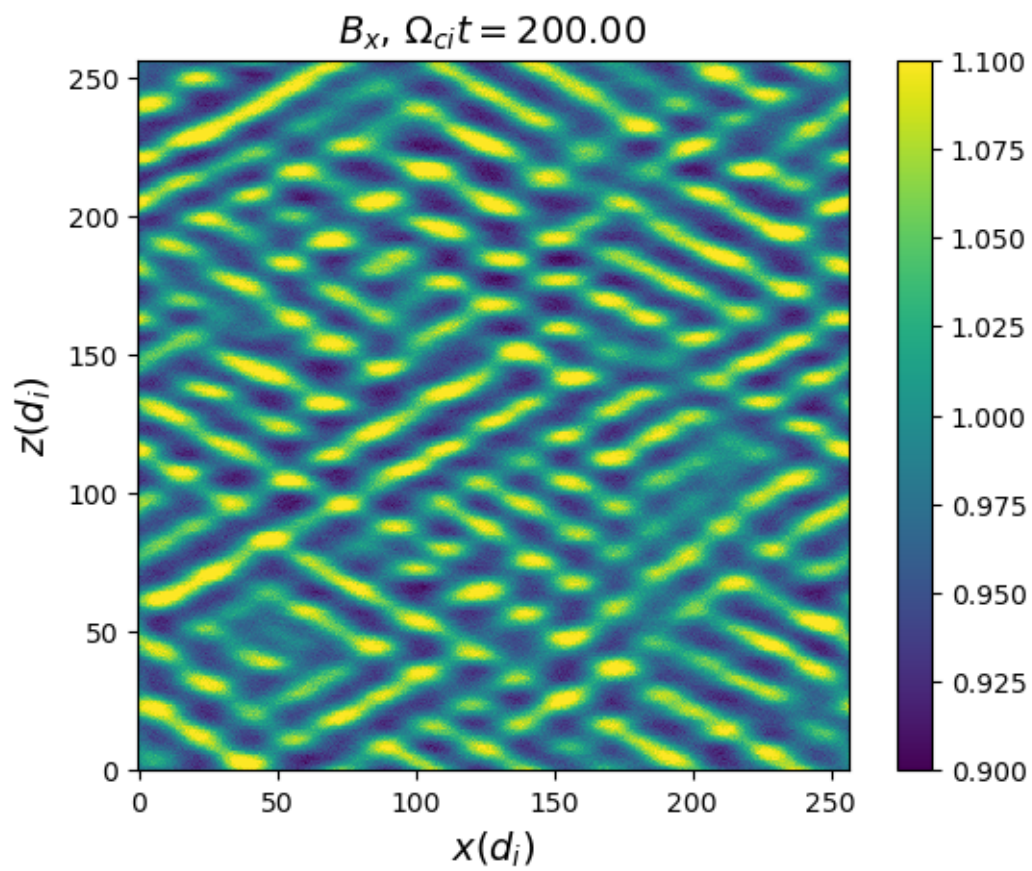
plt.gca().set_aspect("equal")

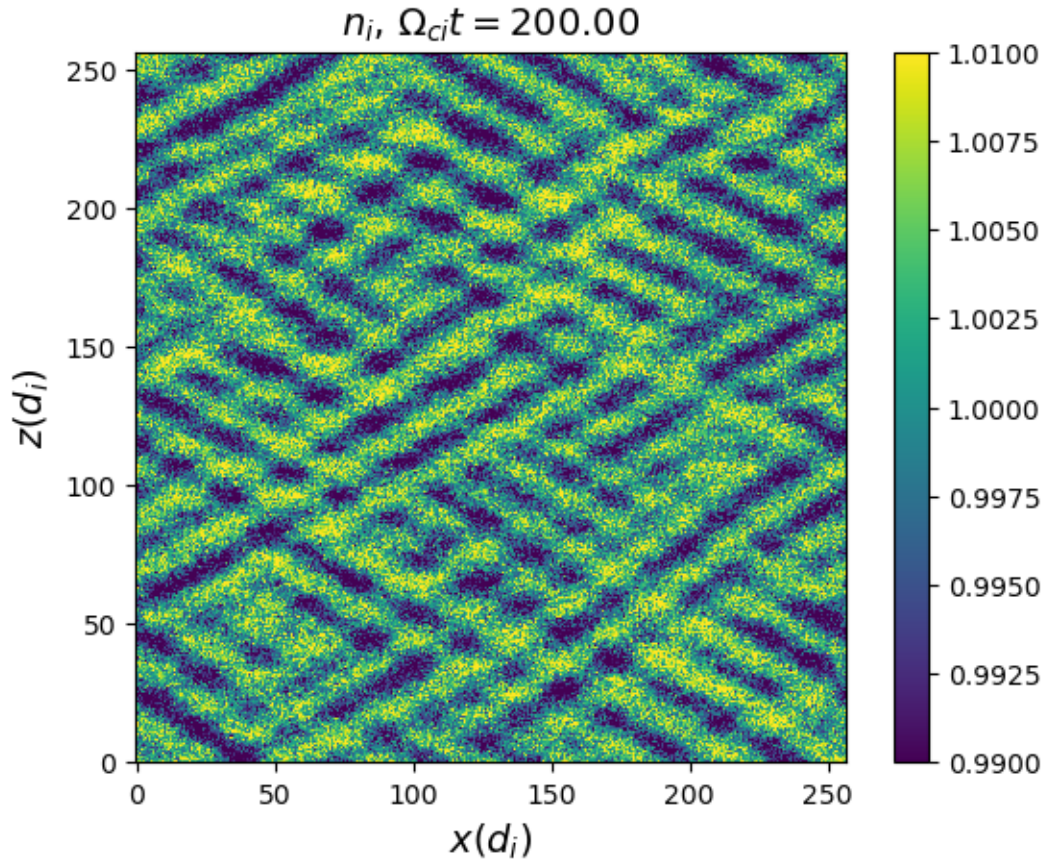
```

```

[14]: # Snapshots of B_x and rho (left and middle panels)
plot_bx(200)
plot_rho(200)

```





```
[15]: # plot frequency spectrum using data frames from nt1 to nt2
def plot_wk_b(nt1=0, nt2=1, dt=1):
    Nx = 512
    Nz = 512
    lx_di = 256.0
    lz_di = 256.0
    wcidt = 1.0

    nframes = nt2 - nt1 + 1
    Bx = np.zeros((Nx, Nz, nframes))
    for nt, tframe in enumerate(range(nt1, nt2+1, dt)):
        tindex = 200 * tframe
        fname = ("./field_hdf5/T." + str(tindex) +
                 "/fields_" + str(tindex) + ".h5")
        with h5py.File(fname, 'r') as fh:
            group = fh["Timestep_" + str(tindex)]
            Bx_ = group["cbx"][:, :, :]
            Bx[:, :, nt] = Bx_[:, 0, :] - Bx_[:, 0, :].mean()
```

```

    Bxwk = np.fft.rfftn(Bx)      # spatial-temporal Fourier transform for real_
↪data
    omegas = np.fft.rfftfreq(nframes)*2*np.pi/wcidt  # frequencies of Fourier_
↪transform
    kxs = np.fft.fftfreq(Nx)*Nx*2*np.pi/lx_di        # wavenumber array for k_x
    kzs = np.fft.fftfreq(Nz)*Nz*2*np.pi/lz_di        # wavenumber array for k_z

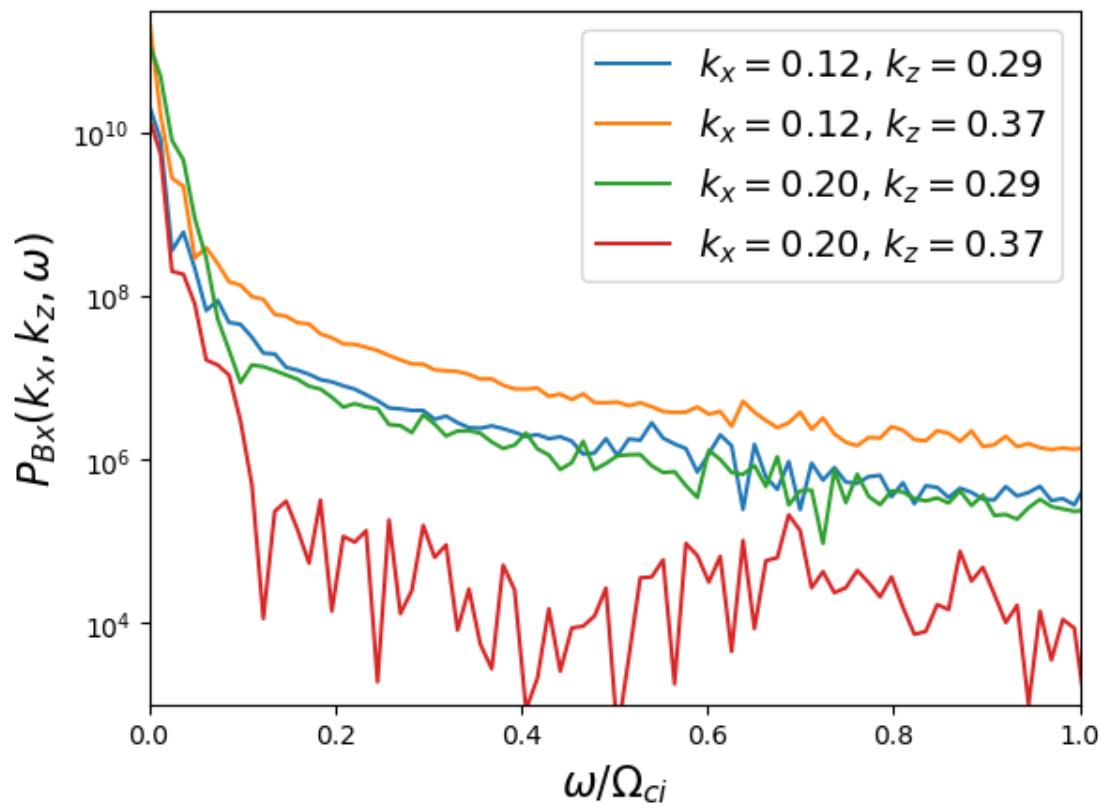
    # select a few wavenumbers and plot the frequency spectrum
    plt.plot(omegas, abs(Bxwk[5,12].T)**2,
              label=r'$k_x=%.2f, $k_z=%.2f'%(kxs[5], kzs[12]))
    plt.plot(omegas, abs(Bxwk[5,15].T)**2,
              label=r'$k_x=%.2f, $k_z=%.2f'%(kxs[5], kzs[15]))
    plt.plot(omegas, abs(Bxwk[8,12].T)**2,
              label=r'$k_x=%.2f, $k_z=%.2f'%(kxs[8], kzs[12]))
    plt.plot(omegas, abs(Bxwk[8,15].T)**2,
              label=r'$k_x=%.2f, $k_z=%.2f'%(kxs[8], kzs[15]))
    plt.yscale('log')
    plt.xlim(0, 1)
    plt.ylim(1e3, 3e11)
    plt.legend(fontsize=14)
    plt.xlabel(r"$\omega$ / $\omega_{ci}$", fontsize=16)
    plt.ylabel(r"$P_{Bx}(k_x, k_z, \omega)$", fontsize=16)

```

```

[16]: # right panel
      plot_wk_b(100, 100+511, 1)

```



[]: